

(Tentative) Calculating Green Function

Taper

September 4, 2016

Abstract

This is a note for reading the paper [1], and for understanding the code produced from that file.

Contents

1	Chapter 1 - Introductions	1
2	Chapter 2 - 2 Landauer-Büttiker formalism	1
3	Chapter 3 Tight-binding model	2
3.1	3.1 Spin-degenerate system	3
3.2	3.2 Including spin degrees of freedom	3
4	Anchor	5
5	License	5

1 Chapter 1 - Introductions

This chapter is really a nice introduction to the current fields of mesoscopic physics. The writing is clear and it traces the development of this field. It gives me a lucid and holistic historical account of both the important discoveries and motives behind them. I should find those marked regions on pdf inside this part very useful.

2 Chapter 2 - 2 Landauer-Büttiker formalism

This chapter introduces the Landauer-Büttiker formalism for calculating the transport properties. The typical setup is illustrated below:

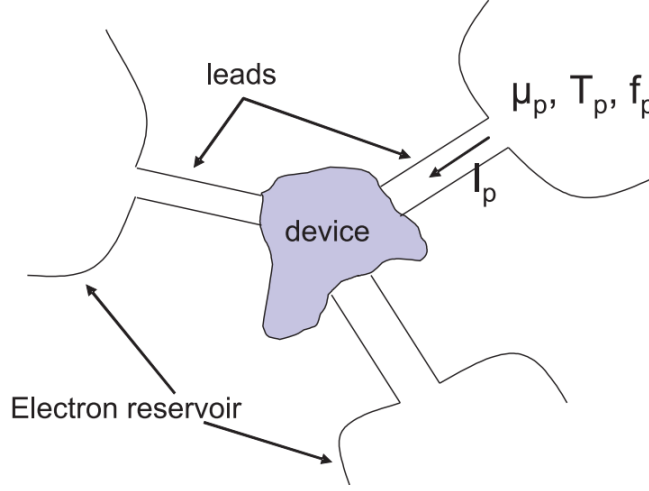


Figure 1: Setup for the Landauer-Büttiker formalism

In that formalism, the currents following through the leads have the following expression:

$$I_p = \frac{-e}{h} \sum_q \int T_{qp}(E) (f_p(E) - f_q(E)) dE \quad (2.0.1)$$

where T_{pq} is the transmission coefficients for electrons to go from lead q to lead p . This formula can be simplified/linearized into:

$$I_p = \frac{e^2}{h} \sum_q T_{pq}(E_F) (V_p - V_q) \quad (2.0.2)$$

An obvious advantage of Landauer-Büttiker formalism is that it makes the dependence of I_p on experimental setup explicit in the formula.

This chapter continue to discuss some time reversal symmetry (TR) properties of this formula, centring/centering around the coefficient T_{pq} . But I am perplexed by that he, while discussing TR, mentions the magnetic field B and formulae like:

$$T_{12}(+B) = T_{12}(-B) \quad (2.0.3)$$

3 Chapter 3 Tight-binding model

Here in this chapter the author presents the fundamental Hamiltonian of the system under consideration.

The process to obtain the Hamiltonian is discretization the Hamiltonian in continuous case, quite the reverse of the first few chapters of A. Zee's QFT in a Nutshell. It should be noted that "a site may represent a region containing many atoms", although "this region should be small compared to physically relevant quantities such as the Fermi wavelength".

Given the formula mentioned in previous chapter, the Hamiltonian presented here seemed extraneous. Maybe the author wants to examine the physics in more detail.

3.1 3.1 Spin-degenerate system

Here the general Hamiltonian, using tight-binding model, is mentioned:

$$H = \sum_{n,m} (t_{nm}^x |n+1, m\rangle \langle n, m| + t_{nm}^y |n, m+1\rangle \langle n, m| + h.c.) + \sum_{n,m} \epsilon_{nm} |n, m\rangle \langle n, m| \quad (3.1.1)$$

The important thing is to determine the coefficient t . In the absence of magnetic field, t is given by:

$$t_{nm}^x = t_{nm}^y = -t = -\frac{\hbar^2}{2m^*a^2} \quad (3.1.2)$$

When magnetic field is present, we effect a so called Peierls substitution to get t . The result is, using Landau gauge in one dimensional under a homogeneous magnetic field.

$$t_{nm}^x = -te^{i2\pi(m-1)\Phi/\Phi_0} \quad (3.1.3)$$

$$t_{nm}^y = -t \quad (3.1.4)$$

When the magnetic field is inhomogeneous, it is generally difficult to choose a gauge to calculate analytically. The author uses a very intuitive discretization process to approach this problem. The process is best illustrated by just the following picture:

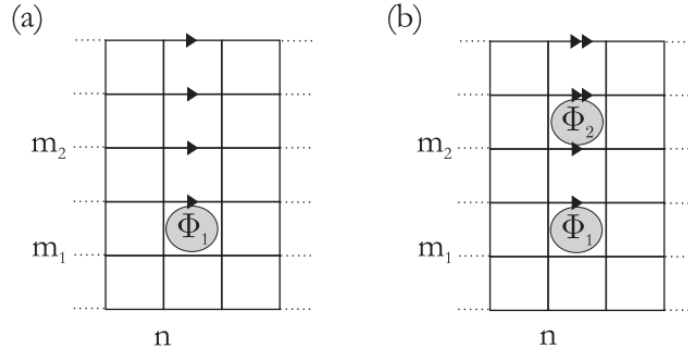


Figure 2: Not uniform magnetic field

Therefore, the result is something like:

$$t_{nm}^x = -te^{i2\pi \sum_{m' < m} \Phi_{nm'}/\Phi_0} \quad (3.1.5)$$

3.2 3.2 Including spin degrees of freedom

When the spin is taken into consideration, the formulation should be modified accordingly.

The author points to a paper, saying that it contains a "lucid discussion on the physics" of Peierls substitution. This might be something worth reading.

If Zeeman/exchange splitting is considered, then:

$$H_S = -\frac{1}{2}g^* \mu_B \sum_{nm} |n, m\rangle \langle n, m| \otimes (B_{nm}^{eff} \cdot \sigma) \quad (3.2.1)$$

Not that the magnetic field strength here is only "effective".

If Spin-orbit coupling is taken into consideration,¹ the following Hamiltonian should be considered:

$$H_{SO} = \lambda P \cdot (\nabla V \times \sigma) \quad (3.2.2)$$

Here P is the mechanical momentum operator, σ is the three pauli spin matrix ($\sigma_x, \sigma_y, \sigma_z$).

The tight-binding version of this is too complicated to be presented here, it is on page 16, equation (3.20).

One additional consideration is the Rashba spin-orbit coupling. This is the peculiar result of electrons trapped in a two dimensional surface. In the Z direction, with Z perpendicular to the 2 dimensional plane, the potential looks like something below, called a triangular potential:

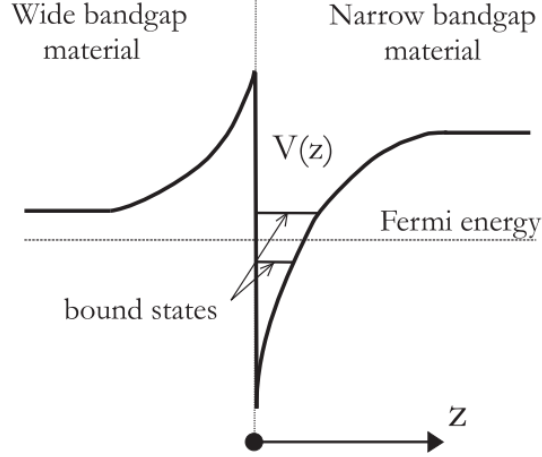


Figure 3: Conduction band at the interface of a semiconductor heterostructure. Band bending creates a potential well $V(z)$ confining the electrons to the XY plane. The asymmetry of this well leads to Rashba spin-orbit coupling.

It is pretty obvious, by equation 3.2.2, that this potential is going into our Hamiltonian. The result in tight-binding model is yet another complicated Hamiltonian, not to be presented here, numbered equation (3.23) on page 18.

¹ In my opinion, this is essentially all about Spin-Magnetic field coupling, not just Spin-Orbital coupling.

4 Anchor

References

- [1] Electronic Transport in Mesoscopic Systems, by von Georgo Metalidis.
(Link found via Google)

5 License

The entire content of this work (including the source code for TeX files and the generated PDF documents) by Hongxiang Chen (nicknamed we.taper, or just Taper) is licensed under a Creative Commons Attribution-NonCommercial-ShareAlike 4.0 International License. Permissions beyond the scope of this license may be available at `mailto:we.taper[at]gmail[dot]com`.