# Quantum Field Theory in Condensed Matter Physics

Taper

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#### Abstract

This is my study notes of various books, listed in the reference.

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#### Second Quantization 1

The [Nag99] introduces second quantization in a heuristic, non-rigorous way. It starts from conjecturing that

•  $N_n \approx \hat{N}_n$ ,

sec:Second-Quantization

 $N_n$  represents  $N_n$  times represent of the experiments done in on a single particle.  $\hat{N}_n$  represents the results from one experiments done on a system of  $\hat{N}_n$  non-interacting particles.

This equality means that we expect these two kinds of experiments to be approximately equivalent.

With such spirit, we first examine the single-particle state and tries to promote something to many-body picture. In single particle state, we have

$$\psi(\mathbf{r},t) = \sum_{n} a_n(t)\phi_n(\mathbf{r})$$
 (1.0.1)

as decomposing any wave function into orthonormal basis and concentrate the dynamics property on the coefficients  $a_n(t)$ . And we could found <sup>1</sup>

$$\frac{\mathrm{d}a_n(t)}{\mathrm{d}t} = \frac{\partial \langle \hat{H} \rangle}{\partial (i\hbar a_n^*)} \tag{1.0.2}$$

$$\frac{\mathrm{d}a_n(t)}{\mathrm{d}t} = \frac{\partial \langle \hat{H} \rangle}{\partial (i\hbar a_n^*)}$$

$$\frac{\mathrm{d}(i\hbar a_n^*(t))}{\mathrm{d}t} = -\frac{\partial \langle \hat{H} \rangle}{\partial a_n}$$
(1.0.2)

 $<sup>^{1}</sup>$ Eq (1.2.5) to eq(1.2.10) of [Nag99]

where  $\langle H \rangle = \langle \psi | H | \psi \rangle$ . These equations are analogous to Hamilton's canonical equations with  $a_n \Leftrightarrow x$ , and  $i\hbar a_n^* \Leftrightarrow p$ . Then propose that we can promote  $a_n$  and  $a_n^*$  as operators. The promotion is not analysed but we can give a quick analogy as:

- $N_n \equiv N|a_n|^2$  is promoted to  $\hat{N_n}$ , which is an observable that counts the total number of particles in state n.
- $\sqrt{N}a_n \to \hat{A}_n$ ,  $\sqrt{N}a_n^* \to \hat{A}_n^{\dagger}$ . So that  $\hat{N}_n = \hat{A}_n^{\dagger}\hat{A}_n$ .

And

$$\begin{cases} [\hat{A}_{n}, \hat{A}_{m}^{\dagger}] = \delta_{n,m}, [\hat{A}_{n}, \hat{A}_{m}] = [\hat{A}_{n}^{\dagger}, \hat{A}_{n}^{\dagger}] = 0, & \text{for bosons} \\ \{\hat{A}_{n}, \hat{A}_{n}^{\dagger}\} = \delta_{n,m}, \{\hat{A}_{n}, \hat{A}_{m}\} = \{\hat{A}_{n}^{\dagger}, \hat{A}_{n}^{\dagger}\} = 0, & \text{for fermions} \end{cases}$$
(1.0.4)

and many other usual commutative/anti-commutative relations.

The wave picture starts from promoting  $\psi$  and  $\psi^*$  into field operators :

$$\hat{\psi}(\mathbf{r}) \equiv \sum_{n} \hat{A}_{n} \phi_{n}(\mathbf{r}) \tag{1.0.5}$$

$$\hat{\psi}^{\dagger}(\mathbf{r}) \equiv \sum_{n} \hat{A}_{n}^{\dagger} \phi_{n}^{*}(\mathbf{r}) \tag{1.0.6}$$

This definition is again analogous to the single particle picture. Calculation  $^2$  shows for bosons

$$[\hat{\psi}(\mathbf{r}), \hat{\psi}^{\dagger}(\mathbf{r}')] = \delta(\mathbf{r} - \mathbf{r}') \tag{1.0.7}$$

$$[\hat{\psi}(\mathbf{r}), \hat{\psi}(\mathbf{r}')] = [\hat{\psi}^{\dagger}(\mathbf{r}), \hat{\psi}^{\dagger}(\mathbf{r}')] = 0$$
 (1.0.8)

with commutators [] replaced with anti-commutators {} for fermions. The **particle density operator** is defined as  $n(\mathbf{r}) = \hat{\psi}^{\dagger}(\mathbf{r})\hat{\psi}(\mathbf{r})$ . and the rest are old things.

The operator new to me is the **phase operator**  $\hat{\theta}_n$  for bosons <sup>3</sup>. It has the property that:

$$\hat{A}_n^{\dagger} = \sqrt{\hat{N}_n} e^{-i\hat{\theta}_n/\hbar} \tag{1.0.9}$$

$$\hat{A}_n = \sqrt{\hat{N}_n} e^{i\hat{\theta}_n/\hbar} \tag{1.0.10}$$

$$[\hat{N}_n,\hat{\theta}_n]=i\hbar \tag{1.0.11}$$
 eq:Ntheta

It is shown that equation 1.0.11 leads to  $[\hat{A}_n, \hat{A}_n^{\dagger}] = 1$ . It is shown also that

$$\exp\left(\frac{i}{\hbar}\hat{\theta}_n\right)(\hat{N}_n)^m \exp\left(-\frac{i}{\hbar}\hat{\theta}_n\right) = (\hat{N}_n + 1)^m \tag{1.0.12}$$

$$\exp\left(\frac{i}{\hbar}\hat{\theta}_n\right)|N_n\rangle \propto |N-1\rangle$$
 (1.0.13) eq:eithetaN

It is concluded, because of the above equation 1.0.13 and that there is not state  $|-1\rangle$ , the  $\hat{\theta}_n$  is not Hermitian. <sup>4</sup>

<sup>&</sup>lt;sup>2</sup>Page 15 of [Nag99]

<sup>&</sup>lt;sup>3</sup>The phase for fermions is complicated and is delegated to the task of chapter 5 of [Nag99].

<sup>&</sup>lt;sup>4</sup>See page 17 of [Nag99]

## References

[Nag99] Naoto Nagaosa. Review of Quantum Mechanics and Basic Principles of Field Theory. Springer Berlin Heidelberg, Berlin, Heidelberg, 1999.

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