

Dirac Hamiltonian

Taper

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Abstract

This aims for recapitulate important information about Dirac Hamiltonian as is written in Sakurai's text [SN11], chapter 8.¹ This is intended to learn to help me deal with a Lattice Dirac Hamiltonian.

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The starting point is still the Schrodinger equation:

$$i\hbar \frac{\partial}{\partial t} |\psi(t)\rangle = H |\psi(t)\rangle \quad (0.0.1)$$

We will be using **Natural Units** from now on.

1 Klein-Gorden Equations

The Klein-Gorden equation starts with the the relativistic energy (of a particle with momentum \mathbf{p} and mass m):

$$E_p = +\sqrt{p^2 + m^2} \quad (1.0.2)$$

The plus sign here hints that we may have a negative energy solution.

We do not quantize this equation directly, because the square root is hard to represent in operators. Instead, Klein-Gorden equation starts with H^2 :

$$H^2 = p^2 + m^2 \xrightarrow{\text{quantized as}} -\nabla^2 + m^2 \quad (1.0.3)$$

On the other hand, the Schrodinger equation applying twice gives:

$$-\frac{\partial^2}{\partial t^2} = H^2 \quad (1.0.4)$$

¹Be careful that some version of this book is missing chapter 8... I don't understand why...

Hence we have the Klein-Gorden equation:

$$\left(\frac{\partial^2}{\partial t^2} - \nabla^2 + m^2\right) \psi(x, t) = 0 \quad (1.0.5) \quad \text{eq:kg-eq}$$

It can be simplified in 4-vector notation as:

$$(\partial_\mu \partial^\mu + m^2) \psi(x, t) = 0 \quad (1.0.6) \quad \text{eq:kg-eq-4vec}$$

or most succinctly as,

$$(\partial^2 + m^2) \psi(x, t) = 0 \quad (1.0.7)$$

The author looked at this equation from three different aspects.

Producing the free particle solution It is first checked that the free-particle state:

$$\psi \propto \exp(-i(Et - \mathbf{p} \cdot \mathbf{x})) = e^{-ip^\mu x_\mu} \quad (1.0.8)$$

is a solution to the Klein-Gorden equation, provided that:

$$E^2 = \mathbf{p}^2 + m^2 \quad (1.0.9)$$

Probability density By analogy with non-relativistic case, the four-vector current is defined as:

$$j^\mu = \frac{i}{2m} [\psi^* \partial^\mu \psi - (\partial^\mu \psi)^* \psi] \quad (1.0.10)$$

It is easy to show $\partial_\mu j^\mu = 0$ ². Then, this is a conserved current. The time-component of this is the density:

$$\rho \equiv j^0 = \frac{i}{2m} [\psi^* \partial_t \psi - (\partial_t \psi)^* \psi] \quad (1.0.11)$$

If one check this equation, one find that ρ is proportionally to the imaginary of some number which, one has no clear reason to say whether it is positive or negative. Therefore, the Klein-Gorden equation leaves one in a bad position to interprate the probability nature of quantum mechanics (what is a "negative probability" anyway?).

Coupling to electromagnetic field The author discusses a little bit about why, when there is electromagnetic field, we have the replacement

$$p \rightarrow p + eA \quad (1.0.12)$$

in page 490, footnote (be careful that the author uses the notation $e = |e|$, whereas I used the notation $e = -|e|$, i.e. e is just q , the charge).

This coupling effectively gives the replacement:

$$\partial_\mu \rightarrow D_\mu \quad (1.0.13)$$

where $D_\mu \equiv \partial_\mu - ieA_\mu$.

The Klein-Gorden equation now becomes:

$$[D_\mu D^\mu + m^2] \psi(x, t) = 0 \quad (1.0.14) \quad \text{eq:kg-emfield}$$

This equation reveals a connection between $e \rightarrow -e$ and $\psi \rightarrow \psi^*$ ³.

²Just note that normally we have $A^\mu B_\mu = A_\mu B^\mu$

³You will find this if you complex conjugate the Klein-Gorden equation.

Two component interpretation The author also mentions a way to split the equation into two equations. If we define $\phi(x, t)$ and $\chi(x, t)$ as in eq.(8.1.15) in p.491, then it is straightforward to check that, the Klein-Gorden equation is equivalent to:

$$iD_t \Upsilon = \left[-\frac{1}{2m} \mathbf{D}^2 (\tau_3 + i\tau_2) + m\tau_3 \right] \Upsilon \quad (1.0.15)$$

where $(D_t, \mathbf{D}) = D_\mu$, τ_i are pauli matrices, and Υ is the two component wave function:

$$\Upsilon \equiv \begin{pmatrix} \phi(x, t) \\ \chi(x, t) \end{pmatrix} \quad (1.0.16)$$

This formulation enables a mysterious interpretation of negative energy, negative probability, and particles and antiparticles. For more information, please look at pp.491-494 of [SN11].

2 Dirac Equation

sec:Dirac Equation

The Dirac equation aims to change the Klein-Gorden equation into a first order PDE. There are two approaches to derive the equation. The first can be found in section 8.2 of [SN11], and section II.1 of [Zee10]. The second approach is found in page 99 (chapter 2) of [Gre97], and is preferred in my opinion. After all, they lead to the same equation:

$$(i\gamma^\mu \partial_\mu - m)\psi(x, t) = 0 \quad (2.0.17)$$

eq:dirac-eq

Where γ^μ are matrices satisfying:

$$\frac{1}{2} \{\gamma^\mu, \gamma^\nu\} = \eta^{\mu\nu} \quad (2.0.18)$$

In other words, γ^μ is a representation of the Clifford algebra $\text{Cl}_{1,3}(\mathbb{R})$. We can cast this equation into another form:

$$i \frac{\partial}{\partial t} \text{ket}\psi = H |\psi\rangle = \boldsymbol{\alpha} \cdot \mathbf{p} + \beta m \quad (2.0.19)$$

where \mathbf{p} is the momentum, $\boldsymbol{\alpha}$ and β satisfy a similar relation:

$$\frac{1}{2} \{\alpha_i, \alpha_j\} = \delta_{ij} \quad (2.0.20)$$

$$\{\alpha_i, \beta\} = 0 \quad (2.0.21)$$

$$\alpha_i^2 = \beta^2 = 0 \quad (2.0.22)$$

They are related to γ^μ matrices by:

$$\alpha_i = \gamma^0 \gamma^i \quad \text{and} \quad \beta = \gamma^0 \quad (2.0.23)$$

A detailed formula for γ matrices can be found in p.94 of [Zee10] or in the section . A detailed formula for α_i, β matrices can be found in p.496 of [SN11].

The density problem The dirac formulation has a appearent advantage. The current defined by

$$j^\mu \equiv \bar{\psi} \gamma^\mu \psi \quad (2.0.24)$$

where

$$\bar{\psi} \equiv \psi^\dagger \beta \quad (2.0.25)$$

is conserved. And the probability density ρ is

$$\rho \equiv j^0 = \psi^\dagger \beta \gamma^0 \psi = \psi^\dagger \psi \quad (2.0.26)$$

It is positive-definite.

Coupled with a charge When coupled with electromagnetic field, and with a the vector potential $\mathbf{A} = 0$, it becomes

$$H = \boldsymbol{\alpha} \cdot \mathbf{p} + \beta m - e\Phi \quad (2.0.27)$$

where Φ is the scalar potential.

3 Gamma matrices

sec:Gamma-matrices

It could be quite easy to find representations of Clifford Algebra in any dimension. First I note that the size of the Gamma matrices of $4n \times 4n$, i.e., the dimension could/must ⁴ be $4n$. So we could have a very easy construction by repeatedly using the Pauli matrices. The Pauli matrices follow:

$$\{\sigma_i, \sigma_j\} = 2\delta_{ij} \quad (3.0.28)$$

So the following 4×4 matrices are good Gamma matrices:

$$\gamma_{1/2/3} = \begin{pmatrix} & \sigma_{1/2/3} \\ -\sigma_{1/2/3} & \end{pmatrix}, \quad \gamma_0 = \begin{pmatrix} \mathbb{1} & \\ & -\mathbb{1} \end{pmatrix} \quad (3.0.29)$$

Note the minus sign in the $-\mathbb{1}$ of γ_0 is important to produce the correct anticommutation relations.

Next, for higher dimension 8, we can set:

$$\gamma'_{1/2/3} = \begin{pmatrix} & \gamma_{1/2/3} \\ -\gamma_{1/2/3} & \end{pmatrix}, \quad \gamma'_0 = \begin{pmatrix} \mathbb{1} & \\ & -\mathbb{1} \end{pmatrix} \quad (3.0.30)$$

And repeating the above process, we can get the Gamma matrices in any $4n$ dimension.

⁴I am not sure about this. In this [1] says it must be $4n$, the [2] says there could be Gamma matrix in odd dimension

References

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- [SN11] J. J. (Jun John) Sakurai and Jim. Napolitano. *Modern quantum mechanics*. Addison-Wesley, 2011.
- [Zee10] A. Zee. *Quantum Field Theory in a Nutshell: (Second Edition)*. In a Nutshell. Princeton University Press, 2010. URL: <https://books.google.co.jp/books?id=n8Mmbjtco78C>.

References

- [1] <https://physics.stackexchange.com/questions/53318/dimension-of-dirac-gamma-matrices>
- [2] <http://demonstrations.wolfram.com/DiracMatricesInHigherDimensions/>

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