

Condensed Matter Field Theory notes

Taper

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Abstract

Notes of book [AS10].

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1 pp.33 eq.1.43

sec:Table

In page 33 of [AS10], the author derives a difference of action, when we have a symmetry transformation paraterized by ω_a :

$$x_\mu \rightarrow x'_\mu = x_\mu + \frac{\partial x_\mu}{\omega_a} |_{\omega=0} \omega_a(x) \quad (1.0.1)$$

$$\phi^i(x) \rightarrow \phi^i(x') = \phi^i(x) + \omega_a(x) F_a^i[\phi] \quad (1.0.2)$$

We have:

$$\mathcal{L} = \mathcal{L}(\phi^i(x), \partial_{x_\mu} \phi^i(x)) \quad (1.0.3)$$

$$\mathcal{L}' = \mathcal{L}'(\phi'^i(x'), \partial_{x'_\mu} \phi'^i(x')) \quad (1.0.4)$$

$$= \mathcal{L}(\phi^i + F_a^i \omega_a, (\delta_{\mu\nu} - \partial_{x_\mu}(\omega_a \partial_{\omega_a} x_\mu)) \partial_{x_\nu}(\phi^i + F_a^i \omega_a)) \quad (1.0.5)$$

And

$$\Delta S = \int d^m x' \mathcal{L}' - \int d^m x \mathcal{L} \quad (1.0.6)$$

eq:dS-integrand

$$\begin{aligned} &= \int d^m x (1 + \partial_{x_\mu}(\omega_a \partial_{\omega_a} x_\mu)) \\ &\times \mathcal{L}(\phi^i + F_a^i \omega_a, (\delta_{\mu\nu} - \partial_{x_\mu}(\omega_a \partial_{\omega_a} x_\mu)) \partial_{x_\nu}(\phi^i + F_a^i \omega_a)) \\ &- \int d^m x \mathcal{L}(\phi^i(x), \partial_{x_\mu} \phi^i(x)) \end{aligned} \quad (1.0.7)$$

Then he argues that, "for constant parameters ω_a the action difference Δa vanishes". Therefore "the leading contribution to the action difference of a symmetry transformation must be linear in the derivative $\partial_{x_\mu} \omega_a$ ".

Then he writes that "A straightforward expansion of the formula above for ΔS shows that these terms are given by"

$$\Delta S = - \int d^m x j_\mu^a(x) \partial_{x_\mu} \omega_a \quad (1.0.8)$$

where j_μ^a is:

$$j_\mu^a = \left(\frac{\partial \mathcal{L}}{\partial(\partial_{x_\mu} \phi^i)} \partial_{x_\nu} \phi^i - \mathcal{L} \delta_{\mu\nu} \right) \frac{\partial x_\nu}{\partial \omega_a} - \frac{\partial \mathcal{L}}{\partial(\partial_{x_\mu} \phi^i)} F_a^i \quad (1.0.9)$$

I am partially confused about how to do the "straightforward expansion". I guess I should do $\frac{\partial}{\partial(\partial_{x_\mu} \omega_a)}$ to the integrand inside expression for ΔS , though I don't really understand the reason. Even so, the integrand contains terms like $\partial_{x_\mu} \partial_{\omega_a} x_\mu$, which I don't know how to deal with.

Solution. The reality is a bit more complicated. We first do a first order expansion to get the infinitesimal difference:

$$\mathcal{L}' - \mathcal{L} \quad (1.0.10)$$

$$\approx \frac{\partial \mathcal{L}}{\partial \phi^i} F_a^i \omega_a + \frac{\partial \mathcal{L}}{\partial(\partial_{x_\mu} \phi^i)} \left[\partial_\mu (F_a^i \omega_a) - \partial_\mu \left(\omega_a \frac{\partial x_\nu}{\partial \omega_a} \right) \partial_\nu (\phi^i + F_a^i \omega_a) \right] \quad (1.0.11)$$

eq:1-1-omega

$$= \omega_a \left[\frac{\partial \mathcal{L}}{\partial \phi^i} F_a^i + \frac{\partial \mathcal{L}}{\partial(\partial_{x_\mu} \phi^i)} \left(\partial_\mu F_a^i - \partial_\mu \left(\frac{\partial x_\nu}{\partial \omega_a} \right) \partial_\nu (\phi^i + F_a^i \omega_a) \right) \right] \quad (1.0.12)$$

eq:1-1-pmu-omega

We also discover the integrand in Eq.1.0.6 to be

$$\left(1 + \partial_\mu \left(\omega_a \frac{\partial x_\mu}{\partial \omega_a} \right) \right) \mathcal{L}' - \mathcal{L} \quad (1.0.13)$$

$$= \left(1 + \partial_\mu \left(\omega_a \frac{\partial x_\mu}{\partial \omega_a} \right) \right) (\mathcal{L}' - \mathcal{L}) + \left(\partial_\mu \left(\omega_a \frac{\partial x_\mu}{\partial \omega_a} \right) \right) \mathcal{L} \quad (1.0.14)$$

eq:integrand-1-density

For the first term $\left(1 + \partial_\mu \left(\omega_a \frac{\partial x_\mu}{\partial \omega_a} \right) \right) (\mathcal{L}' - \mathcal{L})$, the $(\mathcal{L}' - \mathcal{L})$ already has terms of first order of ω_a and of first order of $\partial_\nu \omega_a$. For our purpose, the second order terms $(\partial_\nu (F_a^i \omega_a))$ from item 1.0.11 and item 1.0.12 can be ignored. Also, the item $(\partial_\mu \left(\omega_a \frac{\partial x_\mu}{\partial \omega_a} \right)) (\mathcal{L}' - \mathcal{L})$ in eq.1.0.14 can also be ignored.

Therefore the integrand in Eq.1.0.6 becomes

$$(\mathcal{L}' - \mathcal{L}) + \left(\partial_\mu \left(\omega_a \frac{\partial x_\mu}{\partial \omega_a} \right) \right) \mathcal{L} \quad (1.0.15)$$

$$= \omega_a \left[\frac{\partial \mathcal{L}}{\partial \phi^i} F_a^i + \frac{\partial \mathcal{L}}{\partial(\partial_{x_\mu} \phi^i)} \left(\partial_\mu F_a^i - \left(\partial_\mu \frac{\partial x_\nu}{\partial \omega_a} \right) \partial_\nu (\phi^i + F_a^i \omega_a) \right) + \left(\partial_\nu \frac{\partial x_\mu}{\partial \omega_a} \right) \mathcal{L} \right] \quad (1.0.16)$$

$$+ \partial_\mu \omega_a \left[\frac{\partial \mathcal{L}}{\partial(\partial_{x_\mu} \phi^i)} \left(F_a^i - \frac{\partial x_\nu}{\partial \omega_a} \partial_\nu (\phi^i + F_a^i \omega_a) \right) + \frac{\partial x_\mu}{\partial \omega_a} \mathcal{L} \right] \quad (1.0.17)$$

Therefore, the term we seek, i.e. the coefficient of $\partial_\mu \omega_a$ is

$$\frac{\partial \mathcal{L}}{\partial(\partial_{x_\mu} \phi^i)} \left(F_a^i - \frac{\partial x_\nu}{\partial \omega_a} \partial_\nu (\phi^i + F_a^i \omega_a) \right) + \frac{\partial x_\mu}{\partial \omega_a} \mathcal{L} \quad (1.0.18)$$

$$= \left(\mathcal{L} \delta_{\mu\nu} - \frac{\partial \mathcal{L}}{\partial(\partial_{x_\mu} \phi^i)} \partial_\nu \phi^i \right) \frac{\partial x_\nu}{\partial \omega_a} + \frac{\partial \mathcal{L}}{\partial(\partial_{x_\mu} \phi^i)} F_a^i \quad (1.0.19)$$

which is what we expect in equation 1.43 of [AS10].

Question: as for why we should ignore the term with ω_a , there are two posts ([1], [2]) might be useful for a thought.

References

[AS10] Alexander. Altland and Ben BD Ben Simons. *Condensed matter field theory*. Cambridge University Press, 2010.

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