# Dirac Hamiltonian

### Taper

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#### Abstract

This aims for recapitulate important information about Dirac Hamiltonian as is written in Sakurai's text [SN11], chapter 8. This is intended to learn to help me deal with a Lattice Dirac Hamiltonian.

## Contents

sec:Klien-Gorden Equations

1 Klein-Gorden Equations 1
2 Dirac Equation 3
3 License 4
The starting point is still the Schrodinger equation:

$$i\hbar \frac{\partial}{\partial t} |\psi(t)\rangle = H |\psi(t)\rangle$$
 (0.0.1)

We will be using Natural Units from now on.

# 1 Klein-Gorden Equations

The Klein-Gorden equation starts with the the relativistic energy (of a particle with momentum  $\mathbf{p}$  and mass m):

$$E_p = +\sqrt{p^2 + m^2} (1.0.2)$$

The plus sign here hints that we may have a negative energy solution.

We do not quantize this equation directly, because the square root is hard to represent in operators. Instead, Klein-Gorden equation starts with  $H^2$ :

$$H^2 = p^2 + m^2 \xrightarrow{\text{quantized as}} -\nabla^2 + m^2$$
 (1.0.3)

On the other hand, the Schrodinger equation applying twice gives:

$$-\frac{\partial^2}{\partial t^2} = H^2 \tag{1.0.4}$$

1

 $<sup>^1\</sup>mathrm{Be}$  careful that some version of this book is missing chapter 8... I don't understand why...

Hence we have the Klein-Gorden equation:

$$\left(\frac{\partial^2}{\partial t^2} - \boldsymbol{\nabla}^2 + \boldsymbol{m}^2\right) \psi(\boldsymbol{x}, t) = 0 \tag{1.0.5} \label{eq:loss}$$

It can be simplified in 4-vector notation as:

$$\left(\partial_{\mu}\partial^{\mu} + m^{2}\right)\psi(x,t) = 0 \tag{1.0.6}$$
 eq:kg-eq-4vec

or most succinctly as,

$$(\partial^2 + m^2) \,\psi(x,t) = 0 \tag{1.0.7}$$

The author looked at this equation from three different aspects.

Producing the free particle solution It is first checked that the free-particle state:

$$\psi \propto \exp(-i(Et - \mathbf{p} \cdot \mathbf{x})) = e^{-ip^{\mu}x_{\mu}}$$
 (1.0.8)

is a solution to the Klein-Gorden equation, provided that:

$$E^2 = \mathbf{p}^2 + m^2 \tag{1.0.9}$$

**Probability density** By analogy with non-relativistic case, the four-vector current is defined as:

$$j^{\mu} = \frac{i}{2m} [\psi^* \partial^{\mu} \psi - (\partial^{\mu} \psi)^* \psi]$$
 (1.0.10)

It is easy to show  $\partial_{\mu}j^{\mu}=0^2$ . Then, this is a conserved current. The time-component of this is the density:

$$\rho \equiv j^0 = \frac{i}{2m} \left[ \psi^* \partial_t \psi - (\partial_t \psi)^* \psi \right]$$
 (1.0.11)

If one check this equation, one find that  $\rho$  is proportionally to the imaginary of some number which, one has no clear reason to say whether it is positive or negative. Therefore, the Klein-Gorden equation leaves one in a bad position to interprate the probability nature of quantum mechanics (what is a "negative probability" anyway?).

Coupling to electromagnetic field The author discusses a little bit about why, when there is electromagnetic field, we have the replacement

$$p \to p + eA \tag{1.0.12}$$

in page 490, footnote (be careful that the author uses the notation e = |e|, whereas I used the notation e = -|e|, i.e. e is just q, the charge).

This coupling effectively gives the replacement:

$$\partial_{\mu} \rightarrow D_{\mu} \tag{1.0.13}$$

where  $D_{\mu} \equiv \partial_{\mu} - ieA_{\mu}$ .

The Klein-Gorden equation now becomes:

$$[D_{\mu}D^{\mu} + m^{2}]\psi(x,t) = 0 \tag{1.0.14}$$

eq:kg-emfield

This equation revels a connection between  $e \to -e$  and  $\psi \to \psi^{*3}$ .

<sup>&</sup>lt;sup>2</sup>Just note that normally we have  $A^{\mu}B_{\mu}=A_{\mu}B^{\mu}$ 

<sup>&</sup>lt;sup>3</sup>You will find this if you complex conjugate the Klein-Gorden equation.

Two component interpretation The author also mentions a way to split the equation into two equations. If we define  $\phi(x,t)$  and  $\chi(x,t)$  as in eq.(8.1.15) in p.491, then it is straightforward to check that, the Klein-Gorden equation is equivalent to:

$$iD_t\Upsilon = \left[ -\frac{1}{2m}\mathbf{D}^2(\tau_3 + i\tau_2) + m\tau_3 \right]\Upsilon \tag{1.0.15}$$

where  $(D_t, \mathbf{D}) = D_{\mu}$ ,  $\tau_i$  are pauli matrices, and  $\Upsilon$  is the two component wave function:

$$\Upsilon \equiv \begin{pmatrix} \phi(x,t) \\ \chi(x,t) \end{pmatrix} \tag{1.0.16}$$

This formulation enables a mysterious interpretation of negative energy, negative probability, and particles and antiparticles. For more information, please look at pp.491 494 of [SN11].

# 2 Dirac Equation

The Dirac equation aims to change the Klein-Gorden equation into a first order PDE. There are two approaches to derive the equation. The first can be found in section 8.2 of [SN11], and section II.1 of [Zee10]. The second approach is found in page 99 (chapter 2) of [Gre97], and is prefered in my opinion. After all, they lead to the same equation:

$$(i\gamma^{\mu}\partial_{\mu} - m)\psi(x,t) = 0 \tag{2.0.17}$$

eq:dirac-eq

Where  $\gamma^{\mu}$  are matrices satisfying:

$$\frac{1}{2} \{ \gamma^{\mu}, \gamma^{\nu} \} = \eta^{\mu\nu} \tag{2.0.18}$$

In other words,  $\gamma^{\mu}$  is a representation of the Clifford algebra  $\text{Cl}_{1,3}(\mathbb{R})$ . We can cast this equation into another form:

$$i\frac{\partial}{\partial t}ket\psi = H|\psi\rangle = \boldsymbol{\alpha} \cdot \mathbf{p} + \beta m$$
 (2.0.19)

where **p** is the momentum,  $\alpha$  and  $\beta$  satisfy a similar relation:

$$\frac{1}{2}\{\alpha_i, \alpha_j\} = \delta_{ij} \tag{2.0.20}$$

$$\{\alpha_i, \beta\} = 0 \tag{2.0.21}$$

$$\alpha_i^2 = \beta^2 = 0 \tag{2.0.22}$$

They are related to  $\gamma^{\mu}$  matrices by:

$$\alpha_i = \gamma^0 \gamma^i \quad \text{and} \quad \beta = \gamma^0$$
 (2.0.23)

A detailed formula for  $\gamma$  matrices can be found in p.94 of [Zee10]. A detailed formula for  $\alpha_i, \beta$  matrices can be found in p.496 of [SN11].

The density problem The dirac formulation has a appearent advantage. The current defined by

$$j^{\mu} \equiv \bar{\psi}\gamma^{\mu}\psi \tag{2.0.24}$$

where

$$\bar{\psi} \equiv \psi^{\dagger} \beta \tag{2.0.25}$$

is conserved. And the probability density  $\rho$  is

$$\rho \equiv j^0 = \psi^{\dagger} \beta \gamma^0 \psi = \psi^{\dagger} \psi \tag{2.0.26}$$

It is positive-definite.

Coupled with a charge When coupled with electromagnetic field, and with a the vector potential  $\mathbf{A} = 0$ , it becomes

$$H = \alpha \cdot \mathbf{p} + \beta m - e\Phi \tag{2.0.27}$$

where  $\Phi$  is the scalar potential.

## References

- [Gre97] Professor Dr. Walter Greiner. Relativistic Quantum Mechanics. Springer Berlin Heidelberg, 1997. URL: https://link.springer.com/book/10.1007{%}2F978-3-662-03425-5, doi:10.1007/978-3-662-03425-5.
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- [Zee10] A. Zee. Quantum Field Theory in a Nutshell: (Second Edition). In a Nutshell. Princeton University Press, 2010. URL: https://books.google.co.jp/books?id=n8Mmbjtco78C.

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