

# HOMEWORK 1: Representations of SU(2) and SU(3)

COURSE: Physics 262, *Group Theory for Physicists* (Fall 2019)  
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DUE BY: Be ready to discuss on Friday, Feb 8  
UPDATED: Jan 24, 11:30am

## 1 The spin-1 representation of SU(2)

The spin-1 representation of SU(2) is the three dimensional representation with highest weight  $j = 1$  and states  $|1\rangle$ ,  $|0\rangle$ , and  $|-1\rangle$ . A vector in this space is:

$$\begin{pmatrix} a \\ b \\ c \end{pmatrix} = a|1\rangle + b|0\rangle + c|-1\rangle . \quad (1.1)$$

### 1.1 Normalizations

We chose a normalization of our states so that:

$$T^- |m\rangle = N_m |m-1\rangle \quad (1.2)$$

$$T^+ |m-1\rangle = N_m |m\rangle , \quad (1.3)$$

where  $|n\rangle$  is orthonormal. Write out  $N_m$  for each of the weights  $m = \pm 1, 0$ .

### 1.2 Raising, lowering, Cartan

Write out the explicit form of the generators in this representation:  $d(T^+)$ ,  $d(T^-)$ ,  $d(T^3)$ .

### 1.3 Generators in the ‘usual’ basis

Write out the explicit form of the generators in this representation:  $d(T^1)$ ,  $d(T^2)$ ,  $d(T^3)$ . Recall that:

$$d(T^\pm) = d(T^1) \pm id(T^2) . \quad (1.4)$$

## 2 The adjoint representation of SU(2)

[**Flip**: Updated 1/24: typo on  $\varepsilon^{123}$ ,  $(-i)$  on  $\text{ad}(T^a)^{bc} = -if^{abc}$ . Version below is corrected.]

The **adjoint** representation is one where the generators themselves are states. The action of a generator (as a matrix) on a generator (as a state) is given by the **structure constant**,  $f^{abc}$ . Recall that the structure constant is defined by

$$[T^a, T^b] = if^{abc}T^c . \quad (2.1)$$

For  $SU(2)$   $f^{abc} = \varepsilon^{abc}$ , the totally antisymmetric tensor with  $\varepsilon^{123} = 1$ . The matrices of the adjoint representation are:

$$\text{ad}(T^a)^{bc} = -if^{abc} . \quad (2.2)$$

In other words, the action of a generator in the adjoint representation  $\text{ad}(T^a)$  on a state  $|T^b\rangle$  is

$$\text{ad}(T^a) |T^b\rangle = -if^{abc} |T^c\rangle . \quad (2.3)$$

**[Flip: Update: removed irrelevant information.]**

## 2.1 Dimension of the adjoint representation

What is the dimension of the adjoint representation? (How many basis states are there?) ANSWER: three. (You may want to write one sentence explaining why the answer is three. If the answer is more than one sentence, then it's probably wrong.)

## 2.2 Dimension of the adjoint representation

**[Flip: Update: clarified the basis. Why is the  $|T^\pm\rangle, |T^3\rangle$  basis a useful one to use instead of  $|T^{1,2}\rangle, |T^3\rangle$ ? (Think about the weight of the state.)]**

Write out the explicit matrix form of  $\text{ad}(T^3)$  acting on a basis of states  $|T^\pm\rangle, |T^3\rangle$ . Confirm that it matches  $d(T^3)$  in the spin-1 representation acting on states  $|m = \pm 1\rangle, |m = 0\rangle$ . Write out the explicit matrix forms of  $\text{ad}(T^1)$  and  $\text{ad}(T^2)$  acting on the basis  $|T^\pm\rangle, |T^3\rangle$ . Observe that this is *not* quite the same as  $d(T^1)$  and  $d(T^2)$  of the spin-1 representation.

## 2.3 What gives?

Confirm that  $\text{ad}(T^1)$ ,  $\text{ad}(T^2)$ , and  $\text{ad}(T^3)$  satisfy the  $SU(2)$  commutation relations:

$$[\text{ad}(T^a), \text{ad}(T^b)] = i\varepsilon^{abc} \text{ad}(T^c) . \quad (2.4)$$

Confirm that if some set of matrices  $d(T)$  are a representation—that is, they satisfy the algebra's commutation relations—then another set of matrices that differ by a unitary transformation,  $\tilde{d}(T) \equiv U d(T) U^\dagger$ , is also a representation. Write down the matrix  $U$  that transforms the spin-1 representation into the adjoint representation. This proves that the spin-1 and adjoint representations are, in fact, the same. The difference between them is purely “cosmetic.” Hint: <https://physics.stackexchange.com/q/279880/166736>. The difference has to do with the spherical basis of rotations.

**[Flip: Remark: the basis  $T^\pm, T^3$  forms an algebra that has different structure constants than  $\varepsilon^{abc}$ . That's okay. This basis is useful for building the representations of the group and understanding how they transform, but  $T^\pm$  is not a Hermitian matrix.]**

### 3 Properties of Algebras

Show that:

- If a matrix Lie group is defined to be **special** (unit determinant), then the algebra is made up of traceless matrices.
- If a matrix Lie group is defined to be **unitary**, then the algebra is made up of Hermitian matrices.

### 4 Adjoint representation of a group, algebra

**[Flip: Update: Clarification and discussion.]** Let  $G$  be a matrix Lie group and call the generators  $T^a$ . We can write elements of  $G$  that are “sufficiently close to the origin” as  $g = \exp(-i\theta^a T^a)$ . The adjoint representation of the *group* is the action of  $G$  on elements of its algebra:

$$\text{Ad}(g) |X\rangle = gXg^{-1} , \quad (4.1)$$

where  $X = x^a T^a$  is an element of the algebra. Show that if  $g = \exp(-i\epsilon^a T^a)$  for small angles  $\epsilon^a$ , the adjoint action of the *group* reduces to the adjoint action of the *algebra*:

$$\text{Ad}(g) |X\rangle = (1 + \epsilon^a \text{ad}(T^a) + \mathcal{O}(\epsilon^2)) X . \quad (4.2)$$

Determine what  $c^a$  is. Observe that the action of  $\text{ad}(T^a)$  on  $X$  is a commutator:

$$\text{ad}(T^a) |T^b\rangle = -if^{abc} |T^c\rangle = -|[T^a, T^b]\rangle . \quad (4.3)$$

HINT: See Gutowski sections 2.14 - 2.26.

## Extra Credit

These problems are for your own edification. You are encouraged to explore them according to your own personal and research interests. **Relevant:** <https://youtu.be/0obMRztklqU>.

## BCH

Derive the Baker-Campbell-Hausdorff formula. Start by looking up what the Baker-Campbell-Hausdorff formula is. I don't have anything deep to say about this, but going through the derivation once gives you a feel for how to think of the group and algebra as a manifold and tangent space.