# HOMEWORK 1: Representations of SU(2) and SU(3)

Course: Physics 262, Group Theory for Physicists (Fall 2019)

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Due by: Be ready to discuss on Friday, Feb 8

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# 1 The spin-1 representation of SU(2)

The spin-1 representation of SU(2) is the three dimensional representation with highest weight j = 1 and states  $|1\rangle$ ,  $|0\rangle$ , and  $|-1\rangle$ . A vector in this space is:

$$\begin{pmatrix} a \\ b \\ c \end{pmatrix} = a |1\rangle + b |0\rangle + c |-1\rangle .$$
 (1.1)

#### 1.1 Normalizations

We chose a normalization of our states so that:

$$T^{-}|m\rangle = N_m|m-1\rangle \tag{1.2}$$

$$T^{+}\left|m-1\right\rangle = N_{m}\left|m\right\rangle , \qquad (1.3)$$

where  $|n\rangle$  is orthonormal. Write out  $N_m$  for each of the weights  $m=\pm 1,0$ .

#### 1.2 Raising, lowering, Cartan

Write out the explicit form of the generators in this representation:  $d(T^+)$ ,  $d(T^-)$ ,  $d(T^3)$ .

#### 1.3 Generators in the 'usual' basis

Write out the explicit form of the generators in this representation:  $d(T^1)$ ,  $d(T^2)$ ,  $d(T^3)$ . Recall that:

$$d(T^{\pm}) = d(T^{1}) \pm id(T^{2}) . \tag{1.4}$$

# 2 The adjoint representation of SU(2)

[Flip: Updated 1/24: typo on  $\varepsilon^{123}$ , (-i) on  $\mathrm{ad}(T^a)^{bc}=-if^{abc}$ . Version below is corrected.]

The **adjoint** representation is one where the generators themselves are states. The action of a generator (as a matrix) on a generator (as a state) is given by the **structure constant**,  $f^{abc}$ . Recall that the structure constant is defined by

$$[T^a, T^b] = if^{abc}T^c . (2.1)$$

For SU(2)  $f^{abc} = \varepsilon^{abc}$ , the totally antisymmetric tensor with  $\varepsilon^{123} = 1$ . The matrices of the adjoint representation are:

$$ad(T^a)^{bc} = -if^{abc} . (2.2)$$

In other words, the action of a generator in the adjoint representation  $\operatorname{ad}(T^a)$  on a state  $|T^b\rangle$  is

$$\operatorname{ad}(T^a) \left| T^b \right\rangle = -i f^{abc} \left| T^c \right\rangle . \tag{2.3}$$

[Flip: Update: removed irrelevant information.]

### 2.1 Dimension of the adjoint representation

What is the dimension of the adjoint representation? (How many basis states are there?) Answer: three. (You may want to write one sentence explaining why the answer is three. If the answer is more than one sentence, then it's probably wrong.)

## 2.2 Dimension of the adjoint representation

[Flip: Update: clarified the basis. Why is the  $|T^{\pm}\rangle$ ,  $|T^{3}\rangle$  basis a useful one to use instead of  $|T^{1,2}\rangle$ ,  $|T^{3}\rangle$ ? (Think about the weight of the state.)].

Write out the explicit matrix form of  $\operatorname{ad}(T^3)$  acting on a basis of states  $|T^{\pm}\rangle$ ,  $|T^3\rangle$ . Confirm that it matches  $d(T^3)$  in the spin-1 representation acting on states  $|m=\pm 1\rangle$ ,  $|m=0\rangle$ . Write out the explicit matrix forms of  $\operatorname{ad}(T^1)$  and  $\operatorname{ad}(T^2)$  acting on the basis  $|T^{\pm}\rangle$ ,  $|T^3\rangle$ . Observe that this is *not* quite the same as  $d(T^1)$  and  $d(T^2)$  of the spin-1 representation.

### 2.3 What gives?

Confirm that  $ad(T^1)$ ,  $ad(T^2)$ , and  $ad(T^3)$  satisfy the SU(2) commutation relations:

$$[\operatorname{ad}(T^a), \operatorname{ad}(T^b)] = i\varepsilon^{abc}\operatorname{ad}(T^c) . (2.4)$$

Confirm that if some set of matrices d(T) are a representation—that is, they satisfy the algebra's commutation relations—then another set of matrices that differ by a unitary transformation,  $\tilde{d}(T) \equiv U d(T) U^{\dagger}$ , is also a representation. Write down the matrix U that transforms the spin-1 representation into the adjoint representation. This proves that the spin-1 and adjoint representations are, in fact, the same. The difference between them is purely "cosmetic." Hint: https://physics.stackexchange.com/q/279880/166736. The difference has to do with the spherical basis of rotations.

[Flip: Remark: the basis  $T^{\pm}$ ,  $T^3$  forms an algebra that has different structure constants than  $\varepsilon^{abc}$ . That's okay. This basis is useful for building the representations of the group and understanding how they transform, but  $T^{\pm}$  is not a Hermitian matrix.]

## 3 Properties of Algebras

Show that:

- If a matrix Lie group is defined to be **special** (unit determinant), then the algebra is made up of traceless matrices.
- If a matrix Lie group is defined to be **unitary**, then the algebra is made up of Hermitian matrices.

## 4 Adjoint representation of a group, algebra

[Flip: Update: Clarification and discussion.] Let G be a matrix Lie group and call the generators  $T^a$ . We can write elements of G that are "sufficiently close to the origin" as  $g = \exp(-i\theta^a T^a)$ . The adjoint representation of the group is the action of G on elements of its algebra:

$$Ad(g)|X\rangle = gXg^{-1}, \qquad (4.1)$$

where  $X = x^a T^a$  is an element of the algebra. Show that if  $g = \exp(-i\epsilon^a T^a)$  for small angles  $\epsilon^a$ , the adjoint action of the *group* reduces to the adjoint action of the *algebra*:

$$Ad(g)|X\rangle = (1 + c^a \operatorname{ad}(T^a) + \mathcal{O}(\epsilon^2)) X.$$
(4.2)

Determine what  $c^a$  is. Observe that the action of  $\operatorname{ad}(T^a)$  on X is a commutator:

$$\operatorname{ad}(T^{a})|T^{b}\rangle = -if^{acb}|T^{c}\rangle = |[T^{a}, T^{b}]\rangle . \tag{4.3}$$

[Flip: Correction: 1/30: minus sign on (4.3), see Georgi (6.8)] HINT: See Gutowski sections 2.14 - 2.26.

# Extra Credit

These problems are for your own edification. You are encouraged to explore them according to your own personal and research interests. Relevant: https://youtu.be/OobMRztklqU.

### **BCH**

Derive the Baker-Campbell-Hausdorff formula. Start by looking up what the Baker-Campbell-Hausdorff formula is. I don't have anything deep to say about this, but going through the derivation once gives you a feel for how to think of the group and algebra as a manifold and tangent space.