

# Physics 262: SU(2) Conventions

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## Abstract

Some conventions for group theory. Corrections and additions welcome. We roughly follow chapter 3.2 of Georgi with slightly different conventions.

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## 1 Basis of Generators

The SU(2) algebra is

$$[T^a, T^b] = i\epsilon^{abc}T^c . \quad (1.1)$$

Let us decompose these generators into a pair of raising/lowering operators  $T^\pm$  and the diagonal generator,  $T^3$ . We choose to define the raising/lowering operators as:

$$T^\pm = T^1 \pm iT^2 \quad (1.2)$$

note that Georgi's definition differs by a factor of  $1/\sqrt{2}$ . These satisfy the commutation relations:

$$[T^3, T^\pm] = \pm T^\pm \quad (1.3)$$

$$[T^+, T^-] = 2T^3 . \quad (1.4)$$

## 2 Weights

The **weight**,  $m$ , of a state is the  $T^3$  eigenvalue:

$$T^3 |m\rangle = m |m\rangle . \quad (2.1)$$

The weight is an index for the different components of the representation.

### 3 Raising/Lowering Weights

The  $T^\pm$  operators increment the weight of a state by  $\pm 1$ . We can see this by applying  $T^3$  to the state  $T^\pm |m\rangle$ :

$$T^3 T^\pm |m\rangle = ([T^3, T^\pm] + T^\pm T^3) |m\rangle = (m \pm 1) T^\pm |m\rangle . \quad (3.1)$$

The result is that acting on a state with  $T^\pm$  moves you to a new

### 4 Finite Dimensional Representations

Let us examine representations of  $SU(2)$  that have a finite number of weights—that is, a finite number of components. The  $T^\pm$  operators mean that given a state with some state  $|m\rangle$ , you can “rotate” to a different state  $|m \pm 1\rangle$ . We are interested in finite dimensional representations—objects with only a finite number of components (weights) that rotate into each other. Thus this ladder must truncate for some maximum and minimum weights. Let  $j$  be the **maximum weight** of the representation:

$$T^+ |j\rangle = 0 . \quad (4.1)$$

Let us sequentially derive the states of the representation by applying the lowering operator. What we want to find is the value of the *minimum weight* so that we will know that we have found all of the states.

We can assume that  $|j\rangle$  is normalized. However, we have no guarantee that  $T^- |j\rangle$  is normalized. Thus let us define a normalization

$$T^- |j\rangle = N_j |j-1\rangle , \quad (4.2)$$

where  $|j-1\rangle$  is normalized. Let’s look at this first state below the maximum weight state.

#### 4.1 Normalization of $|j-1\rangle$

We determine the normalization  $N_j$  by calculating the norm  $\|T^- |j\rangle\|$ :

$$\langle j | (T^-)^\dagger T^- |j\rangle = \langle j | T^+ T^- |j\rangle \quad (4.3)$$

$$= \langle j | ([T^+, T^-] + T^- T^+) |j\rangle \quad (4.4)$$

$$= \langle j | 2T^3 |j\rangle \quad (4.5)$$

$$= 2j . \quad (4.6)$$

We have used  $T^+ |j\rangle = 0$ . Using the definition of the normalized state  $|j-1\rangle$ , this must equal:

$$2j = |N_j|^2 \langle j-1 | j+1 \rangle , \quad (4.7)$$

so that we have found the normalization

$$N_j = \sqrt{2j} . \quad (4.8)$$

Next we should check the normalization of the  $T^+ |j-1\rangle$ . To do this, we use the one strategy available: *use the commutation relations to simplify the expression*. Use the facts that

1.  $T^+$  annihilates  $|j\rangle$
2.  $T^3|m\rangle = m$ .

Thus the raising of  $|j-1\rangle$  gives:

$$T^+|j-1\rangle = T^+ \left( \frac{1}{N_j} T^-|j\rangle \right) \quad (4.9)$$

$$= \frac{1}{N_j} ([T^+, T^-] + T^- T^+) |j\rangle \quad (4.10)$$

$$= \frac{2j}{N_j} |j\rangle \quad (4.11)$$

$$= N_j |j\rangle . \quad (4.12)$$

In the last line we explicitly use  $N_j = \sqrt{2j}$ . We observe that raising  $|j-1\rangle$  and lowering  $|j\rangle$  produce the same normalization:

$$T^-|j\rangle = N_j |j-1\rangle \quad (4.13)$$

$$T^+|j-1\rangle = N_j |j\rangle . \quad (4.14)$$

## 4.2 Recursion relation

Now suppose that for some general weight  $|m\rangle$  in our representation, we know that:

$$T^-|m+1\rangle = N_{m+1}|m\rangle \quad (4.15)$$

$$T^+|m\rangle = N_{m+1}|m+1\rangle , \quad (4.16)$$

where  $N_{m+1}$  is known and the states  $|m\rangle, |m+1\rangle$  are normalized. This is the case with the ‘top of the ladder,’  $(m+1) = 1$ . We now derive  $N_m$  and show that the same relation holds for  $m \rightarrow (m-1)$ .  $N_m$  is defined by

$$T^-|m\rangle = N_m|m-1\rangle . \quad (4.17)$$

$N_m$  is fixed by requiring that  $|m-1\rangle$  is normalized,  $\langle m-1|m-1\rangle = 1$ . Then we have:

$$N_m^2 = \langle m|T^+T^-|m\rangle \quad (4.18)$$

$$= \langle m|([T^+, T^-] + T^- T^+) |m\rangle \quad (4.19)$$

$$= 2m + N_{m+1}^2 . \quad (4.20)$$

We thus have a recursion relation for  $N_m$  in terms of  $N_{m+1}$ .

To complete the recursion, we must prove that

$$T^+|m-1\rangle = N_m|m\rangle . \quad (4.21)$$

The left-hand side is

$$T^+ |m-1\rangle = T^+ \left( \frac{1}{N_m} T^- |m\rangle \right) \quad (4.22)$$

$$= \frac{1}{N_m} ([T^+, T^-] + T^- T^+) |m\rangle \quad (4.23)$$

$$= \frac{1}{N_m} (2m |m\rangle + N_{m+1} T^- |m+1\rangle) \quad (4.24)$$

$$= \frac{1}{N_m} (2m |m\rangle + N_{m+1}^2 |m\rangle) \quad (4.25)$$

$$= N_m |m\rangle , \quad (4.26)$$

where we have used the explicit result, (4.20).

### 4.3 Closed expression for the normalization

We can imagine writing down the relations for the squared normalizations:

$$\begin{aligned} N_j^2 &= 2j \\ N_{j-1}^2 - N_j^2 &= 2(j-1) \\ N_{j-2}^2 - N_{j-1}^2 &= 2(j-2) \\ &\vdots \\ N_{j-k}^2 - N_{j-k-1}^2 &= 2(j-k) \end{aligned} \quad (4.27)$$

Sum both sides of the equations and recognizing that there are  $(k+1)$  terms,

$$N_{j-k}^2 = 2(k-1)j - 2 \frac{k(k+1)}{2} \quad (4.28)$$

$$= (k+1)(2j-k) . \quad (4.29)$$

Rather than using “ $k$  steps below the highest weight, ” it is more tidy to use the normalization of the  $m^{\text{th}}$  rung of the ladder via  $k = j - m$ :

$$N_m = \sqrt{(j-m+1)(j+m)} . \quad (4.30)$$

## 5 Bottom of the ladder

In order for the representation to be finite, we need the ladder to truncate at the bottom. There is thus some a limit to the number of times you can hit the highest weight state  $|j\rangle$  with  $T^-$  before it annihilates. Let this number be  $\ell$ . Then we want to use

$$T^- |j-\ell\rangle = 0 \quad (5.1)$$

with the normalization  $N_{j-\ell}$  to determine what this  $\ell$  is.

$$N_{j-\ell} = \sqrt{(\ell+1)(2j-\ell)} . \quad (5.2)$$

Observe that  $(\ell+1) > 0$  since  $\ell$  is a positive natural number. Then the requirement that the ladder terminates,  $N_{j-\ell} = 0$ , requires

$$\ell = 2j . \quad (5.3)$$

From this we observe:

1. The highest weight  $m_{\max} = j$  is half integer. It can be written as  $\ell/2$  for some positive integer  $\ell$ .
2. The lowest weight is simply  $m_{\min} = -j$ .