

Solutions to Pathria's Statistical Mechanics

Chapter 7

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May 8, 2016

Problem 7.1

Problem 7.4

We can deduce from 7.1.30 and 7.1.26a, when p is a const:

$$g_{\frac{5}{2}} \propto T^{-\frac{5}{2}}$$

So we have:

$$\left(\frac{\partial}{\partial T} g_{\frac{5}{2}}(z)\right)_p = -\frac{5}{2T} g_{\frac{5}{2}}(z)$$

According to the D.10 from appendix D:

$$z \frac{\partial}{\partial z} g_{\frac{5}{2}}(z) = g_{\frac{3}{2}}(z)$$

Combine the two equation:

$$\frac{1}{z} \left(\frac{\partial z}{\partial T}\right)_p = -\frac{5}{2T} \frac{g_{\frac{5}{2}}(z)}{g_{\frac{3}{2}}(z)}$$

If you compare with 7.1.36 you can get:

$$\gamma = \frac{5}{3} \frac{g_{\frac{5}{2}}(z) g_{\frac{1}{2}}(z)}{g_{\frac{3}{2}}(z)^2}$$

Problem 7.11

(a)

At the critical point of B-E condensation, $N_0 \ll N, z = 1$.

$$\frac{N}{V} = \frac{1}{\lambda^3} \frac{1}{\Gamma(\frac{3}{2})} \left(\int_0^\infty \frac{x^{\frac{1}{2}} dx}{e^x + 1} + \int_0^\infty \frac{x^{\frac{1}{2}} dx}{e^{x+\beta\epsilon_1} + 1} \right) \quad (1)$$

Because $\beta\epsilon_1 \gg 1$, so

$$\int_0^\infty \frac{x^{\frac{1}{2}} dx}{e^{x+\beta\epsilon_1} + 1} = e^{-\beta\epsilon_1} \int_0^\infty e^{-x} x^{\frac{1}{2}} dx = \Gamma(\frac{3}{2}) e^{-\beta\epsilon_1} \ll 1 \quad (2)$$

So, we think $\Delta T = T_c - T_{c0} \ll 1$.

$$\frac{N}{V} = \frac{1}{\lambda^3} \left(\xi\left(\frac{3}{2}\right) + e^{-\beta\epsilon_1} \right) \approx \frac{1}{\lambda_0^3} \left(1 + \frac{3\Delta T}{2T_{c0}} \right) \left(\xi\left(\frac{3}{2}\right) + e^{\frac{\epsilon_1}{kT_{c0}}} \right) \quad (3)$$

And we know

$$\frac{N}{V} = \frac{1}{\lambda_0^3} \xi\left(\frac{3}{2}\right) \quad (4)$$

So,

$$\frac{3\Delta T}{2T_{c0}} \xi\left(\frac{3}{2}\right) + e^{\frac{\epsilon_1}{kT_{c0}}} = 0 \quad (5)$$

$$\frac{T_c}{T_{c0}} = 1 - \frac{\frac{2}{3}}{\xi\left(\frac{3}{2}\right)} e^{\frac{\epsilon_1}{kT_{c0}}} \quad (6)$$

(b)

$$\frac{N}{V} = \frac{1}{\lambda^3} \frac{1}{\Gamma\left(\frac{3}{2}\right)} \left(\int_0^\infty \frac{x^{\frac{1}{2}} dx}{e^x + 1} + \int_0^\infty \frac{x^{\frac{1}{2}} dx}{e^{x+\beta\epsilon_1} + 1} \right) = \frac{1}{\lambda^3} (g_{3/2}(1) + g_{3/2}(e^{-\beta\epsilon_1})) \quad (7)$$

As $\beta\epsilon_1 \ll 1$, we can expand $g_{3/2}(e^{-\beta\epsilon_1})$ and keep the first two terms.

$$\frac{N}{V} = \frac{1}{\lambda^3} \left(2\xi\left(\frac{3}{2}\right) + \Gamma\left(-\frac{1}{2}\right)(\beta\epsilon_1)^{\frac{1}{2}} \right) = \frac{1}{\lambda_0^3} \xi\left(\frac{3}{2}\right) \quad (8)$$

If $\beta\epsilon_1 = 0$,

$$\lambda' = 2^{\frac{1}{3}} \lambda_0 \quad (9)$$

$$T'_c = \left(\frac{1}{2}\right)^{\frac{2}{3}} T_{c0} \quad (10)$$

And $T'_c - T_c \ll 1$.

$$\frac{N}{V} = \frac{2}{\lambda'^3} \left(1 + \frac{3\Delta T}{2T'_c} \right) \left(\xi\left(\frac{3}{2}\right) - \sqrt{\pi \left(\frac{\epsilon_1}{kT'_c} \right)} \right) = \frac{1}{\lambda_0^3} \xi\left(\frac{3}{2}\right) \quad (11)$$

$$\frac{3\Delta T}{2T'_c} \xi\left(\frac{3}{2}\right) - \sqrt{\frac{\pi\epsilon_1}{kT'_c}} = 0 \quad (12)$$

$$\frac{\Delta T}{T'_c} = \frac{\frac{2}{3}}{\xi\left(\frac{3}{2}\right)} \sqrt{\frac{\pi\epsilon_1}{kT'_c}} = \frac{2^{\frac{4}{3}}}{3\xi\left(\frac{3}{2}\right)} \sqrt{\frac{\pi\epsilon_1}{kT_{c0}}} \quad (13)$$

$$\frac{T_c}{T_{c0}} = \left(\frac{1}{2}\right)^{\frac{2}{3}} \left(1 + \frac{2^{\frac{4}{3}}}{3\xi\left(\frac{3}{2}\right)} \sqrt{\frac{\pi\epsilon_1}{kT_{c0}}} \right) \quad (14)$$

Problem 7.14

From the definition of p and U:

$$p = \frac{kT}{\lambda^n} g_{\frac{5}{2}}(z)$$

$$U = \frac{nkTV}{s\lambda^n} g_{\frac{5}{2}}(z)$$

So we get:

$$p = \frac{sU}{nV}$$

And when $T \rightarrow \infty$, we could use the ideal gas equation.

$$pV = nRT$$

So

$$C_V = \frac{n}{s} Nk, C_p = \left(\frac{n}{s} + 1\right) Nk$$

Problem 7.21

$$\frac{N}{V} = \frac{1}{h^3} \int_0^\infty \frac{8\pi \frac{\epsilon^2}{c^3} d\epsilon}{e^{\beta\epsilon} - 1} = \frac{8\pi}{\beta^3 h^3 c^3} \int_0^\infty \frac{x^2 dx}{e^x - 1} = \frac{8\pi}{\beta^3 h^3 c^3} \Gamma(3) \xi(3) \quad (15)$$

$$\frac{E}{V} = \frac{1}{h^3} \int_0^\infty \frac{8\pi \frac{\epsilon^3}{c^3} d\epsilon}{e^{\beta\epsilon} - 1} = \frac{8\pi}{\beta^4 h^3 c^3} \int_0^\infty \frac{x^3 dx}{e^x - 1} = \frac{8\pi}{\beta^4 h^3 c^3} \Gamma(4) \xi(4) \quad (16)$$

$$\frac{E}{N} = \frac{3\xi(4)}{\xi(3)} kT \approx 2.7kT \quad (17)$$

Problem 7.24

We can deduce from 7.3.12 and 7.3.19 and 7.3.23.

$$u = 4.16 * 10^{-14}$$

$$s = 2.03 * 10^{-14}$$

$$n = 4.09 * 10^8$$

Problem 7.31

The state density for transverse phonon is

$$g(\omega) = \frac{6N}{\omega_{DT}^3} \omega^2 \quad (18)$$

The state density for longitudinal phonon is

$$g(\omega) = \frac{3N}{\omega_{DL}^3} \omega^2 \quad (19)$$

And

$$C_V = C_{VT} + C_{VL} \quad (20)$$

Compared with equation (7.4.17), it is easy to find verify that

$$C_V = Nk(2D(x_{0,T}) + D(x_{0,L})) \quad (21)$$

The nature of the error in equation (7.4.17) comes from the fact that the longitudinal and the transverse modes of the solid should have their own cutoff frequencies, ω_{DL} and ω_{DT} rather than a common cutoff at ω_D .