Solutions to Pathria's Statistical Mechanics Chapter 4

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Problem 4.1

Problem 4.4

The probability of a state with energy E_r and particle number N is

$$p_{r,N} = \frac{e^{-\beta E_{r,N} + \beta \mu N}}{\mathcal{Q}(\mu, V, \beta)} \tag{1}$$

in which $Q = \sum_{r,N} e^{-\beta E_{r,N} + \beta \mu N}$ is the grand canonical partition function. Since we have define that $z = e^{\beta \mu}$, the probability can be written as:

$$p_{r,N} = \frac{z^N e^{-\beta E_{r,N}}}{\mathcal{Q}(z, V, \beta)} \tag{2}$$

So the probability that has exactly N particles will be:

$$p_{N} = \sum_{r} p_{r,N} = \frac{z^{N} \sum_{r} e^{-\beta E_{r,N}}}{Q(z, V, \beta)}$$
(3)

easily we can find the summation in the numerator is the canonical partition function of system with V, N and β :

$$Q_N(V,\beta) = \sum_r e^{-\beta E_{r,N}} \tag{4}$$

Thus Eq.(??) will become:

$$p_N = \frac{z^N Q_N(V, \beta)}{Q(z, V, \beta)} \tag{5}$$

For ideal classical gas, the canonical partition function is:

$$Q_N(V,T) = \frac{V^N}{N!} \left(\frac{2\pi mkT}{h^2}\right)^{3N/2} \tag{6}$$

and the grand partition function is

$$Q(z, V, \beta) = \sum_{N=0}^{\infty} z^N Q_N(V, T) = \exp\left[zV \left(\frac{2\pi mkT}{h^2}\right)^{3/2}\right]$$
 (7)

Clearly the probability distribution of particle number is

$$p_N = \frac{1}{N!} \frac{(zV\lambda_T^{-3})^N}{e^{zV\lambda_T^{-3}}} \tag{8}$$

It is obvious that this distribution is a Poison distribution. From the knowledge of Poison distribution, we know the root-mean-square value of (ΔN) is

$$\Delta N = \sqrt{zV\lambda_T^{-3}} = \sqrt{e^{\beta\mu}V\left(\frac{2\pi mkT}{h^2}\right)^{3/2}} \tag{9}$$

We can also get this result from the formula of grand canonical ensemble:

$$\Delta N = kT \sqrt{\left(\frac{\partial^2 \ln \mathcal{Q}}{\partial \mu^2}\right)_{T,V}}$$

$$= \sqrt{e^{\beta \mu} V \left(\frac{2\pi m kT}{h^2}\right)^{3/2}}$$
(10)

And this result is consistent with the one we get by Poison distribution.

Problem 4.5

We could know from 4.3.20:

$$S = kT(\frac{\partial q}{\partial T})_{z,V} - Nkln(z) + kq$$

We can know partial differential:

$$(\frac{\partial q}{\partial T})_{z,V} - (\frac{\partial q}{\partial T})_{\mu,V} = (\frac{\partial q}{\partial z})_{T,V} (\frac{\partial z}{\partial T})_{\mu,V}$$
$$(\frac{\partial q}{\partial z})_{T,V} = \frac{N}{z}$$

So we can infer that:

$$S=k[\frac{\partial (Tq)}{\partial T}]_{V,\mu}$$

Problem 4.14

The ClausiusClapeyron equation is

$$\frac{dP_{\sigma}}{dT} = \frac{L}{T\Delta v}$$

Since the volume of liquid is negligible compared to that of gas, we can alternate Δv by $v_g = kT/P_{\sigma}$. Put all of these into the Clausius Clapeyron equation, we can get a differential equation:

$$\frac{dP_{\sigma}}{P_{\sigma}} = \frac{L}{R} \frac{dT}{T^2} \tag{11}$$

so the solution to the differential equation will be:

$$P_{\sigma}(T) = P_0 \exp\left[\frac{L}{R} \left(\frac{1}{T_0} - \frac{1}{T}\right)\right] \tag{12}$$

From the problem we know that L = 2260 kJ/kg = 40680 J/mol, $T_0 = 373 \text{K}$ and $P_0 = 101 \text{kPa}$. Then we put all these numbers into Eq.(??) and we get the equilibrium vapor pressure is

$$P_{\sigma}(473\text{K}) = 1619\text{kPa}$$

Experiment result is $P_{\sigma} \sim 1500 \mathrm{kPa}$, and our calculation is approximately correct.

Problem 4.15

According to Clausius-Clapeyron equation. And ignore the volume of solid phase.

$$\frac{dP_{\sigma}}{dT} = \frac{L}{TV}$$

Use the gas equation.

$$ln(p) = -\frac{L}{kT} + A$$

Use the triple point parameter.

$$ln(p) = -\frac{L}{kT} + 6.6 \times 10^{26}$$