

Quantum Computing & Applications for Engineering



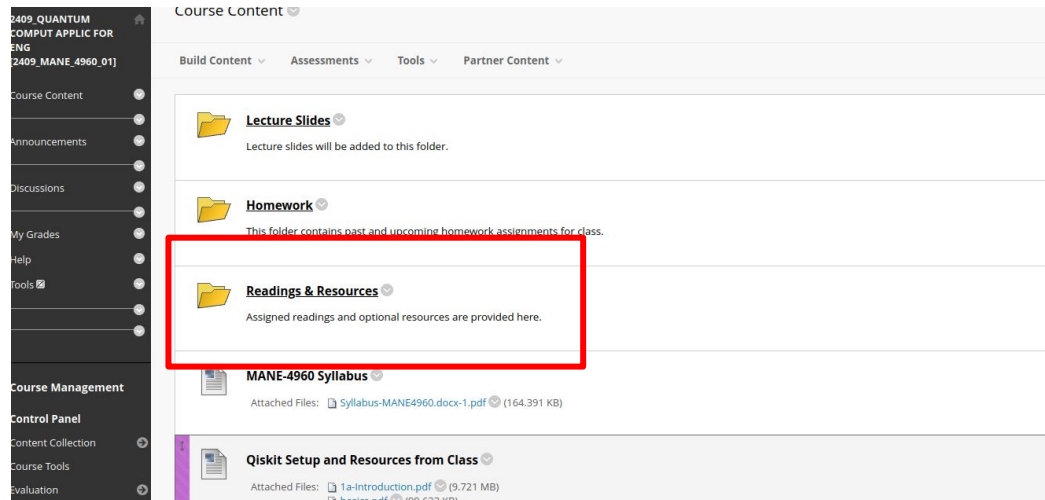
10/1/2024

Today's goals

- Qubit Errors and Noise
- Measuring Decoherence Processes
- Noise models
- Example quantum algorithm workflow
- Discuss HW

Announcements

- New resources on LMS
 - Textbook by Tom Wong
 - Recommended Reading, Websites
- Bonus quantum computing survey
 - Complete by 10/11 for +10 points

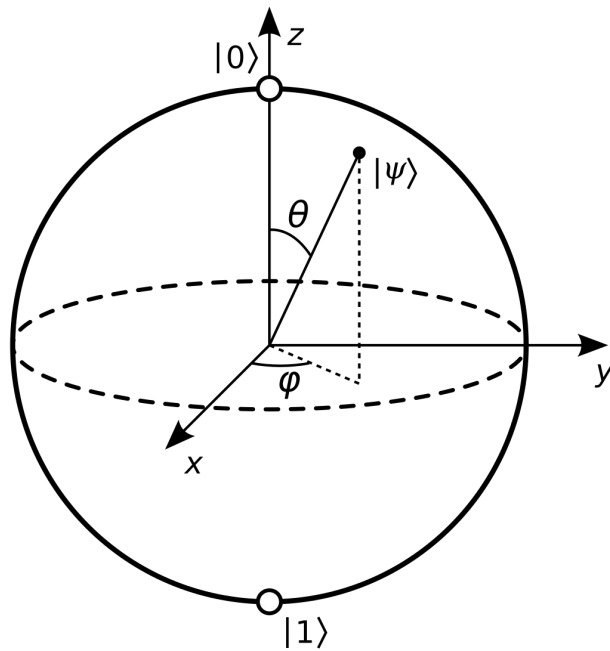


HW2 Questions?

- Phase troubles?
- Flipping a sign vs. flipping a qubit

Pure States

- So far, we've been talking in terms of *statevectors*.
- A statevector represents what is known as a *pure state*
- Each element tells of how much of each basis state we have in the superposition.
- **What happens if qubits aren't perfect?**



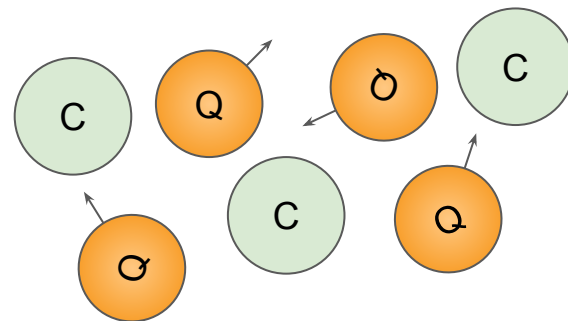
The Bloch sphere represents a pure state

$$|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad |1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

$$|\psi\rangle = \alpha|0\rangle + \beta|1\rangle \quad |\alpha|^2 + |\beta|^2 = 1$$

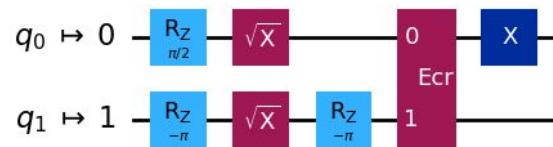
Density Matrices

- Consider the case where we have a bunch of qubits in the same system...
 - Clouds of atoms
 - Liquid NMR
 - NV Centers
- ...or a single state prepared and measured over and over.
 - Quantum circuits and algorithms
 - Quantum communication protocols
- The density matrix gives us a unified representation of quantum coherence and classical mixtures.



A mixture of quantum and classical states

Global Phase: $3\pi/4$



Running a circuit over and over

Density Matrices

- Last time (and on HW3) we showed that the outer product of two statevectors gives us a matrix.
 - Projectors, if we're talking about the basis states.
- In general, the outer product of two quantum statevectors gives us a *density matrix*.
- Density matrices represent a statistical mixture of quantum *and* classical states.

$$P_0 = |0\rangle\langle 0| \quad P_0 = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

$$P_1 = |1\rangle\langle 1| \quad P_1 = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$$

$$\rho = \begin{pmatrix} \rho_{00} & \rho_{01} \\ \rho_{11} & \rho_{11} \end{pmatrix}$$

Density Matrices & Mixed States

- Consider the Hadamard state.
- The statevector shows we have a 50-50 superposition.
- The density matrix shows we also have a 50-50 classical mixture!
 - The off-diagonal elements are the quantum, or “coherent” parts.
 - The diagonal elements are the classical parts.
- How do we know if we have a quantum state?

$$|+\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

$$\rho_+ = |+\rangle \langle +|$$

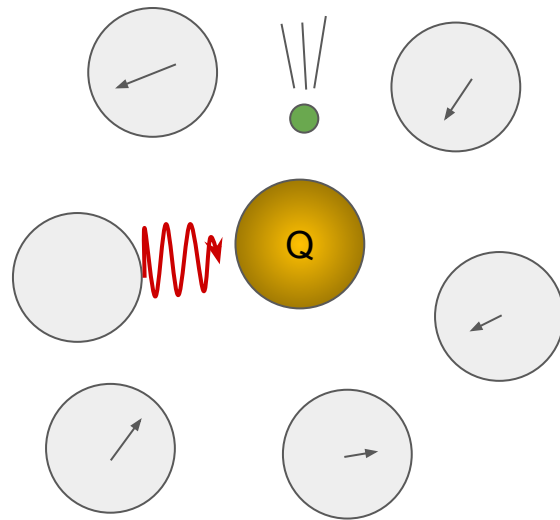
$$\rho_+ = \begin{pmatrix} \boxed{\frac{1}{2}} & \boxed{\frac{1}{2}} \\ \boxed{\frac{1}{2}} & \boxed{\frac{1}{2}} \end{pmatrix}$$

Mixed States

- The quantum, or coherent, part we can detect and estimate by measuring in different bases.
- Mixed states can only be described with density matrices.
- Effects in the environment transfer information from the coherent elements into the classical elements.
- This leakage of quantum information is termed Decoherence.
 - A measurement is intentional decoherence.
 - Unintentional decoherence occurs when:
 - Our qubits get entangled with something else in the environment.
 - Our qubits get measured by something else in the environment.

Physical Decoherence

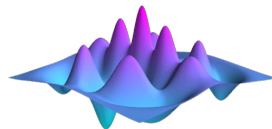
- Schrodinger's equation assumes a quantum system that is perfectly isolated.
 - Real systems lose energy to the environment
 - Formally, the environment is called a “bath”
- Open quantum systems are governed by the *Lindblad Master Equation*.
 - The real world is non-unitary.
 - Quantum term plus a dissipative term.
 - In the limit of zero dissipation, the Schrodinger equation is recovered.
- Lindblad simulations are used in qubit and quantum device design.



$$\frac{d\rho}{dt} = \boxed{-i[H, \rho]} + \boxed{\sum_k \left(L_k \rho L_k^\dagger - \frac{1}{2} \{ L_k^\dagger L_k, \rho \} \right)}$$

Coherent Evolution

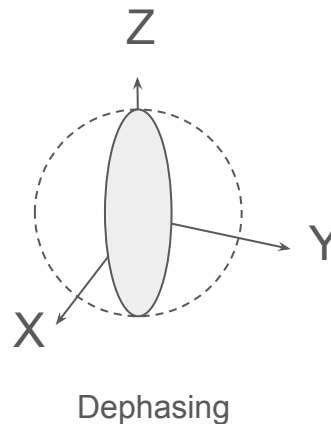
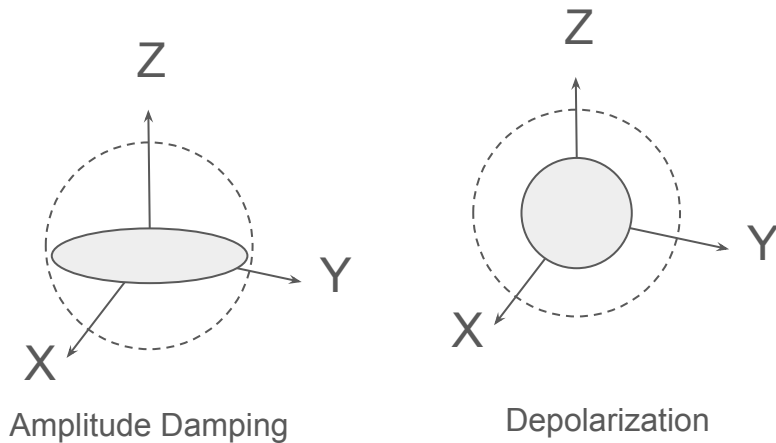
Decoherence



QuTiP - Quantum
Toolbox in Python

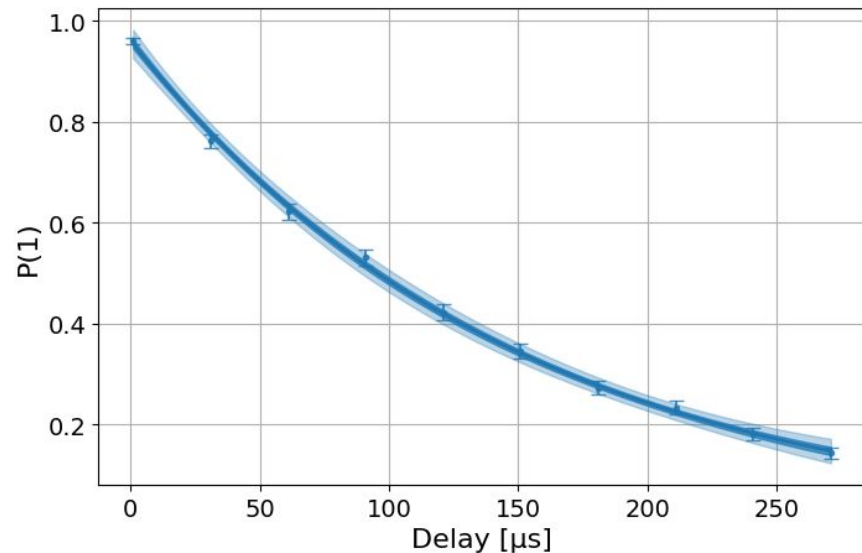
Decoherence of Qubits

- The quantum state in a qubit (or set of qubits) is very fragile
 - Thermal vibration
 - Stray E/M fields
 - Radiation
- Decoherence is the process that returns quantum states back to classical ones.
 - Measurement is a “controlled” decoherence process.
 - Noise is measurement done by the environment.
- There are several “channels” for decoherence.
 - Amplitude damping
 - Phase damping
 - Depolarization (combination of the first 2)
- Decoherence channels “squish” the reachable parts of the Bloch sphere



Amplitude Damping

- A qubit placed in an excited state only stays there temporarily before losing energy to the environment.
 - De-excitation
 - Spin relaxation
- This decoherence process is called *amplitude damping*.
- The “decay constant” that represents its de-excitation back to the 0 state is denoted T_1 .



$T_1 = 148 \pm 9.86 \mu\text{s}$
reduced- $\chi^2 = 0.5777$

For superconducting qubits, T_1 values of $\sim 150 \mu\text{s}$ are typical

Amplitude Damping

- We can measure T_1 using individual qubits, or many in parallel.
- Experiment Algorithm:
 - Start with a qubit in the 0 state.
 - Apply an X gate to put it in the 1 state.
 - Wait for a predefined delay time, t_d .
 - Measure the qubit.
 - Tally the expectation value with the $(I-Z)/2$ operator.
 - Repeat for N_d delay times.
- Fitting an exponential function to the data points gives us T_1 .

Phase Damping

- When we place a qubit in superposition, it has a well-defined phase.
 - Consider the + and - Hadamard states.
- As it undergoes interactions with the environment, the phase drifts.
- Coherent phase information leaks into the classical elements of the density matrix.
- The time constant for this phase coherence loss is called T_2 or T_2^* , depending on how it is measured.

Phase Damping

- T_2 is measured using a *Hahn Echo* sequence
 - Apply a Hadamard gate
 - Wait for a delay time
 - Apply an X gate
 - Wait another delay time
 - Apply a Hadamard
 - Measure
- During the delay the qubit precesses around the Z-axis due to its natural frequency.
 - Applying an X gate flips it 180 degrees.
 - Applying the last Hadamard should bring us back to 0, if the coherent phase is well defined.

Phase Damping

- T_2^* is an alternate measurement that can also provide an estimate of the qubit's intrinsic frequency.
- Measurement protocol:
 - Apply a Hadamard gate
 - Wait a delay time
 - Apply an R_z rotation around the Z-axis.
 - Apply another Hadamard gate
 - Measure
- If the R_z gate and the natural precession offset each other perfectly, then the qubit will return to zero.
- Offsets from zero occur as the phase coherence decays.

Qubit Noise

- The combination of the T_1 and T_2 processes lead to single-qubit errors.
- Errors occur when the statevector is not aligned with where you'd expect it to be, based on the operations you've performed.
 - Bit flip - T_1 errors
 - Phase flip - T_2 errors
- If your circuit takes substantially longer than T_1 or T_2 to run, you have a higher probability of errors.
- Many other noise sources too.

Gate Noise

- Applying gates isn't always a perfect process.
 - Frequency drift in lasers or RF
 - Qubit manufacturing variations
- Two-qubit gates are more challenging than single qubit gates.
- We can measure gate fidelity to estimate how many gates we can run before the noise adds up.

Noise Models

- Quantum simulators can model the noise behavior of quantum computers.
 - Useful when you want to check if something is even worth running.
 - Essential when you don't have hardware access.
- Most quantum computing research is done with noise models.
- “All models are wrong, some are useful!”

Noise Models

- A simulator with noise models the evolution of the density matrix with each gate application.
- The effects of decoherence are applied using *Kraus operators*.
- Qiskit Aer lets you define your own Kraus operators, or pick from a library of common error channels.
- Once we have T_1 and T_2 , we can define Kraus operators for these processes.

$$\rho' = \sum_k E_k \rho E_k^\dagger$$

The new density matrix after the application of k different noise channels with Kraus operators

$$\begin{aligned}\gamma &= 1 - e^{t/T_1} \\ E_0 &= \begin{pmatrix} 1 & 0 \\ 0 & \sqrt{1-\gamma} \end{pmatrix} \\ E_1 &= \begin{pmatrix} 0 & \sqrt{\gamma} \\ 0 & 0 \end{pmatrix}\end{aligned}$$

Amplitude Damping

$$\begin{aligned}\lambda &= 1 - e^{t/T_2} \\ E_0 &= \begin{pmatrix} 1 & 0 \\ 0 & \sqrt{1-\lambda} \end{pmatrix} \\ E_1 &= \begin{pmatrix} 0 & 0 \\ 0 & \sqrt{\lambda} \end{pmatrix}\end{aligned}$$

Phase Damping

Using T_1 and T_2 , we define *damping parameters* that we use to make the Kraus operators

Kraus Operators - Example

- Let's look at how T_1 and T_2 processes affect a Hadamard gate.
 - Let $T_1 = 150 \mu\text{s}$
 - Let $T_2 = 50 \mu\text{s}$
- Assume the gate time T_g is 60 ns
- Demo Notebook

Noise & Performance Benchmarks

- There are so many noise sources, a zoo of benchmarking techniques has been developed to try to characterize them.
- There's a hierarchy of benchmarks:
 - Qubit and gate-level
 - Circuit-level
 - Application-level
- Not all benchmarks are created equal.
 - Physical timings and fidelity of operations are relatively objective.
 - Integrated measures at the circuit and gate level can vary widely.
 - Application-oriented benchmarks
- Pay attention to which benchmarks are preferred by different hardware vendors!

Break

What to do with noise?

- The maximum circuit length is constrained by the coherence time and the gate error (especially the two-qubit gate error); circuit depth
- On superconducting systems, circuit depths greater than ~ 100 start to become very challenging to run.
- On trapped ion systems, circuit depths of 500-2000 are possible.
- No qubit modality has currently gotten better than a 10^{-4} total error rate (all error sources combined)
- For comparison, a typical CPU has error rates below 10^{-14}
- We need something to make use of the qubits we have.

The Fault Tolerance Theorem

- In the 1990's researchers demonstrated that extra qubits can be used to catch and correct errors. Quantum Error Correction
- This means that reliable quantum computations are possible even with imperfect qubits.
- There is a threshold at which error-corrected qubits can be stacked up to recover all lost information.
- That threshold is not known exactly, but is around 10^{-2} - 10^{-4} .
- If we can make qubits better than this (we can), then we can achieve fault tolerance.