RCSJ model using python3

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This script is based on Tinkham's "Introduction to superconductivity"[1] and Gross' "Applied superconductivity", chapter 3[2]. All simulations were done using `python3` and the `rcsj` module written by Felix[3].

CONTENTS

I. CURRENT BIASING

A. Junction model and ODE

We model the Josephson junction in the simplest way possible: A series network of a resistor R, a capacitor C and a nonlinear inductor. The total current running through the network is

$$I(t) = I_s(t) + I_R(t) + I_C(t) \tag{1}$$

$$I(t) = I_c \sin(\phi) + \tfrac{V}{R} + C\frac{\mathrm{d}V}{\mathrm{d}t} \tag{2}$$

$$V(t) = \tfrac{\hbar}{2e}\frac{\mathrm{d}\phi}{\mathrm{d}t},\ \frac{\mathrm{d}V}{\mathrm{d}t} = \frac{\hbar}{2e}\frac{\mathrm{d}^2\phi}{\mathrm{d}t^2} \tag{3}$$

$$\rightarrow I(t) = I_c \sin(\phi) + \tfrac{\hbar}{2eR}\frac{\mathrm{d}\phi}{\mathrm{d}t} + \frac{\hbar C}{2e}\frac{\mathrm{d}^2\phi}{\mathrm{d}t^2} \tag{4}$$

There are two approaches to simplifying this equation in terms of normalization. We here call them the **Tinkham approach**[1] and the **Gross approach**[2].

1. Tinkham approach

The equation is simplified with normalized time via the plasma frequency, $\tau = \omega_p t$:

$$\omega_p = \frac{1}{\tau_p} = \frac{1}{\sqrt{L_c C}} = \sqrt{\frac{2eI_c}{\hbar C}} \tag{5}$$

$$\mathrm{d}t = \frac{1}{\omega_p}\mathrm{d}\tau \rightarrow \frac{\mathrm{d}^n}{\mathrm{d}t^n} = \omega_p^n \frac{\mathrm{d}^n}{\mathrm{d}\tau^n} \tag{6}$$

$$\frac{I}{I_c} - \sin(\phi) = \underbrace{\frac{\hbar}{2eI_cR}\sqrt{\frac{2eI_c}{\hbar C}}}_{\sqrt{\frac{\hbar C}{2eI_c}}\frac{1}{RC} \equiv Q^{-1}} \frac{\mathrm{d}\phi}{\mathrm{d}\tau} + \underbrace{\frac{\hbar C}{2eI_c}\frac{2eI_c}{\hbar C}}_{1} \frac{\mathrm{d}^2\phi}{\mathrm{d}\tau^2} \tag{7}$$

Hence the final ODE is

$$\frac{\mathrm{d}^2\phi}{\mathrm{d}\tau^2} = \frac{I}{I_c} - \sin(\phi) - \frac{1}{Q}\frac{\mathrm{d}\phi}{\mathrm{d}\tau} \tag{8}$$

2. Gross approach

This approach is different in that it seems more intuitive, but the time scale and damping are different: We normalize the time not by the plasma frequency, but the L_c/R_n time constant that yields the characteristic junction frequency:

$$\omega_c = \frac{1}{\tau_c} = \frac{R_n}{L_c} = \frac{2e}{\hbar} I_c R_n = \frac{\Phi_0}{2\pi} V_c \tag{9}$$

$$\tau = \frac{t}{\tau_c},\ \tau_c = \frac{\hbar}{2eI_cR} \tag{10}$$

$$\mathrm{d}t = \tau_c \mathrm{d}\tau \to \frac{\mathrm{d}^n}{\mathrm{d}t^n} = \frac{1}{\tau_c^n}\frac{\mathrm{d}^n}{\mathrm{d}\tau^n} \tag{11}$$

The above ODE then can be resorted into

$$I - I_c\sin(\phi) = \frac{\hbar\tau_c}{2eR}\frac{\mathrm{d}\phi}{\mathrm{d}\tau} + \frac{\hbar C\tau^2}{2e}\frac{\mathrm{d}^2\phi}{\mathrm{d}\tau^2} \tag{12}$$

$$i - \sin(\phi) = \underbrace{\frac{\hbar}{2eRI_c}\left(\frac{2eI_cR}{\hbar}\right)}_{1}\frac{\mathrm{d}\phi}{\mathrm{d}\tau} + \underbrace{\frac{\hbar C}{2e}\left(\frac{2eI_cR}{\hbar}\right)^2}_{\frac{2eI_cR^2C}{\hbar}=\beta_c}\frac{\mathrm{d}^2\phi}{\mathrm{d}\tau^2} \tag{13}$$

$$\frac{\mathrm{d}^2\phi}{\mathrm{d}\tau^2} = \frac{1}{\beta_c}\left(i - \sin(\phi) - \frac{\mathrm{d}\phi}{\mathrm{d}\tau}\right) \tag{14}$$

Note that $\beta_c \equiv Q^2 \equiv \tau_c RC$. For this ODE it is much easier to distinguish between $\beta_c \gg 1$ and $\beta_c \ll 1$, simply because of its form. However, it is important to keep in mind that the timescales are different!

B. Unit discussion

We consider first τ_c. The units of $\hbar/(2e) = \Phi_0/(2\pi)$ are $\mathrm{Wb}=\mathrm{kg\,m^2\,s^{-2}\,A^{-1}}$. Critical current is in A, resistance in $\Omega=\mathrm{kg\,m^2\,s^{-3}\,A^{-2}}$. Therefore

$$[\tau_c] = \left[\frac{\hbar}{2e}\frac{1}{I_c}\frac{1}{R}\right] = \frac{\mathrm{kg\,m^2}}{\mathrm{s^2\,A}}\frac{1}{\mathrm{A}}\frac{\mathrm{s^3\,A^2}}{\mathrm{kg\,m^2}} = \mathrm{s} \tag{15}$$

$$\to [\tau] = \left[\frac{t}{\tau_c}\right] = \frac{\mathrm{s}}{\mathrm{s}} = 1 \tag{16}$$

Now the same for the plasma frequency: Capacitance is given in units of $\mathrm{F=s^4\,A^2\,m^{-2}\,kg^{-1}}$

$$[\omega_p] = \left[\sqrt{\frac{2e}{\hbar}\frac{I_c}{C}}\right] = \sqrt{\frac{\mathrm{s^2\,A}}{\mathrm{kg\,m^2}}\mathrm{A}\frac{\mathrm{m^2\,kg}}{\mathrm{s^4\,A^2}}} = \sqrt{\frac{1}{\mathrm{s^2}}} = \frac{1}{\mathrm{s}} \tag{17}$$

$$\to [\tau] = [\omega_p t] = \frac{\mathrm{s}}{\mathrm{s}} = 1 \tag{18}$$

C. Damping

We consider eq. 14 for simplicity. For **very weak damping**, i.e. $\beta_c \gg 1$, the left hand side of the ODE can be set to zero. The result is a linear IVC on the retrapping branch, with $I_r \leq I_c$:

$$\frac{\mathrm{d}^2\phi}{\mathrm{d}\tau^2} = 0 \to \langle V(t)\rangle = IR_n \tag{19}$$

$$\frac{I_r}{I_c} \approx \frac{4}{\pi\sqrt{\beta_c}} \equiv \frac{4}{\pi Q} \tag{20}$$

Figure 1. Current-voltage curves for various damping cases. As predicted by theory, high damping ($Q \ll 1$) corresponds to a square-root like IVC with no hysteresis, while low damping ($Q \gg 1$) results in strong hystersis with a almost linear retrapping branch.

An open question is why Gross (and others) report significant hysteresis $I_r/I_c \approx 0.85$ already for $\beta_c = 1$, while we achieve at least $I_r/I_c = 0.95$ for the implemented python simulations.

For **very strong damping**, i.e. $\beta_c \ll 1$, the right hand side of the ODE is zero:

$$\frac{\mathrm{d}\phi}{\mathrm{d}\tau} = i - \sin(\phi) \rightarrow \langle V(t) \rangle = I_c R \sqrt{(I/I_c)^2 - 1} \tag{21}$$

without hysteresis.

D. Frequency analysis

II. VOLTAGE BIASING

The system equation only slightly changes. What remains to be seen is the normalization and the exact values needed to achieve voltage biasing.

REFERENCES

[1] M. Tinkham, *Introduction to Superconductivity: Second Edition (Dover Books on Physics) (Vol i)*, 2nd ed. (Dover Publications, 2004).
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[3] F. E. Schmidt, "RCSJ model," https://github.com/feschmidt/rcsj (2017).