

X-RAY DATA BOOKLET

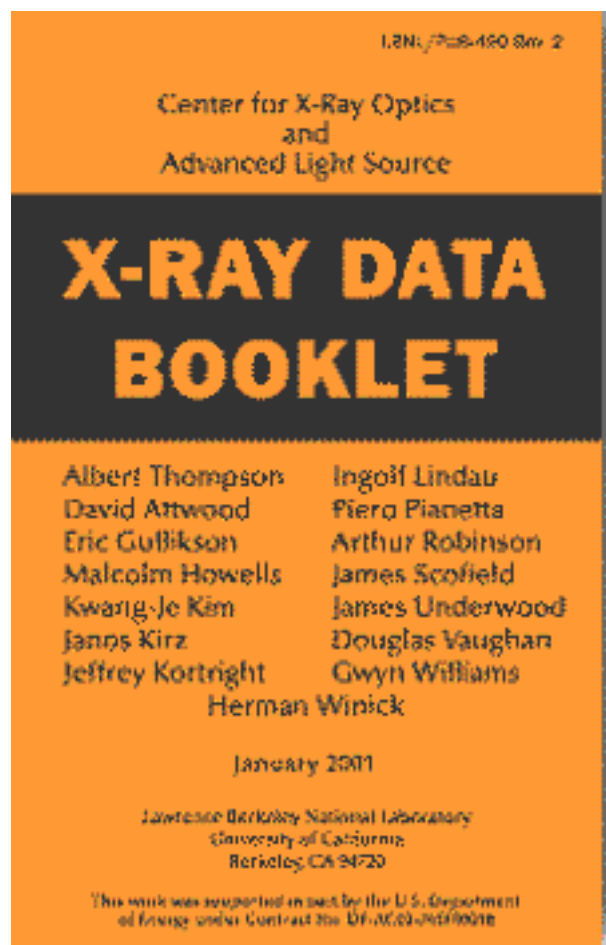
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- [Introduction](#)
- [X-Ray Properties of Elements](#)
- [Electron Binding Energies](#)
- [X-Ray Energy Emission Energies](#)
- [Fluorescence Yields for K and L Shells](#)
- [Principal Auger Electron Energies](#)
- [Subshell Photoionization Cross-Sections](#)
- [Mass Absorption Coefficients](#)
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● [Synchrotron Facilities](#)

● [Scattering Processes](#)

● [Scattering of X-rays from Electrons and Atoms](#)

● [Low-Energy Electron Ranges in Matter](#)

● [Optics and Detectors](#)

● [Crystal and Multilayer Elements](#)

● [Specular Reflectivities for Grazing-Incidence Mirrors](#)

● [Gratings and Monochromators](#)

● [Zone Plates](#)

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● [Miscellaneous](#)



● [Physical Constants](#)

● [Physical Properties of the Elements](#)

● [Electromagnetic Relations](#)

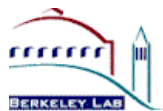
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● [Useful Formulas](#)

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Introduction

- [Preface in PDF format](#)
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CENTER FOR X-RAY OPTICS
ADVANCED LIGHT SOURCE

X-RAY DATA BOOKLET

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CONTENTS

1. X-Ray Properties of the Elements	1-1
1.1 Electron Binding Energies <i>Gwyn P. Williams</i>	1-1
1.2 X-Ray Emission Energies <i>Jeffrey B. Kortright and Albert C. Thompson</i>	1-8
1.3 Fluorescence Yields For <i>K</i> and <i>L</i> Shells <i>Jeffrey B. Kortright</i>	1-28
1.4 Principal Auger Electron Energies	1-30
1.5 Subshell Photoionization Cross Sections <i>Ingolf Lindau</i>	1-32
1.6 Mass Absorption Coefficients <i>Eric M. Gullikson</i>	1-38
1.7 Atomic Scattering Factors <i>Eric M. Gullikson</i>	1-44
1.8 Energy Levels of Few-Electron Ionic Species <i>James H. Scofield</i>	1-53
2. Synchrotron Radiation	2-1
2.1 Characteristics of Synchrotron Radiation <i>Kwang-Je Kim</i>	2-1
2.2 History of Synchrotron Radiation <i>Arthur L. Robinson</i>	2-17
2.3 Operating and Planned Facilities <i>Herman Winick</i>	2-24

3.	Scattering Processes	3-1
3.1	Scattering of X-Rays from Electrons and Atoms <i>Janos Kirz</i>	3-1
3.2	Low-Energy Electron Ranges In Matter <i>Piero Pianetta</i>	3-5
4.	Optics and Detectors	4-1
4.1	Multilayers and Crystals <i>James H. Underwood</i>	4-1
4.2	Specular Reflectivities for Grazing-Incidence Mirrors <i>Eric M. Gullikson</i>	4-14
4.3	Gratings and Monochromators <i>Malcolm R. Howells</i>	4-17
4.4	Zone Plates <i>Janos Kirz and David Attwood</i>	4-28
4.5	X-Ray Detectors <i>Albert C. Thompson</i>	4-33
5.	Miscellaneous	5-1
5.1	Physical Constants	5-1
5.2	Physical Properties of the Elements	5-4
5.3	Electromagnetic Relations	5-11
5.4	Radioactivity and Radiation Protection	5-14
5.5	Useful Equations	5-17

PREFACE

For the first time since its original publication in 1985, the *X-Ray Data Booklet* has undergone a significant revision. Tabulated values and graphical plots have been revised and updated, and the content has been modified to reflect the changing needs of the x-ray community. Further, the *Booklet* is now posted on the web at <http://xdb.lbl.gov>, together with additional detail and further references for many of the sections.

As before, the compilers are grateful to a host of contributors who furnished new material or reviewed and revised their original sections. Also, as in the original edition, many sections draw heavily on work published elsewhere, as indicated in the text and figure captions. Thanks also to Linda Geniesse, Connie Silva, and Jean Wolslegel of the LBNL Technical and Electronic Information Department, whose skills were invaluable and their patience apparently unlimited. Finally, we express continuing thanks to David Attwood for his support of this project, as well as his contributions to the *Booklet*, and to Janos Kirz, who conceived the *Booklet* as a service to the community and who remains an active contributor to the second edition.

As the compilers, we take full responsibility for any errors in this new edition, and we invite readers to bring them to our attention at the Center for X-Ray Optics, 2-400, Lawrence Berkeley National Laboratory, Berkeley, California 94720, or by e-mail at xdb@grace.lbl.gov. Corrections will be posted on the web and incorporated in subsequent printings.

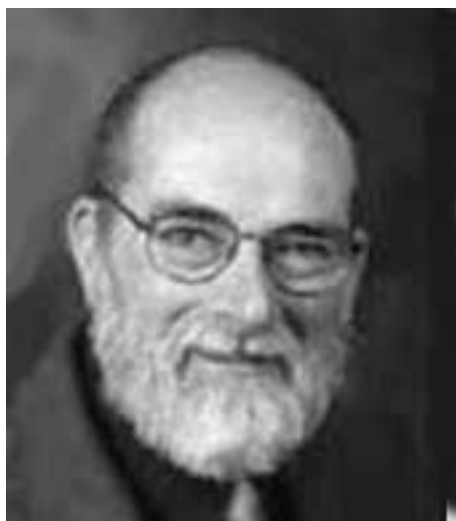
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Gwyn Williams



Herman Winick



1. X-Ray Properties of the Elements

Contents

- [Electron Binding Energies](#) - *Gwyn P. Williams*
- [X-Ray Energy Emission Energies](#)
- *Jeffrey B. Kortright and Albert C. Thompson*
- [Fluorescence Yields for K and L Shells](#) - *Jeffrey B. Kortright*
- [Principal Auger Electron Energies](#)
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X-Ray Data Booklet

Section 1.1 ELECTRON BINDING ENERGIES

Gwyn P. Williams

Table 1-1 gives the electron binding energies for the elements in their natural forms. A [PDF version of this table](#) is also available. The energies are given in electron volts relative to the vacuum level for the rare gases and for H₂, N₂, O₂, F₂, and Cl₂; relative to the Fermi level for the metals; and relative to the top of the valence bands for semiconductors. Values have been taken from Ref. 1 except as follows:

*Values taken from Ref. 2, with additional corrections

†Values taken from Ref. 3.

^aOne-particle approximation not valid owing to short core-hole lifetime.

^bValue derived from Ref. 1.

Thanks also to R. Johnson, G. Ice, M. Olmstead, P. Dowben, M. Seah, E. Gullikson, F. Boscherini, W. O'Brien, R. Alkire, and others.

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1. J. A. Bearden and A. F. Burr, "Reevaluation of X-Ray Atomic Energy Levels," *Rev. Mod. Phys.* **39**, 125 (1967).
2. M. Cardona and L. Ley, Eds., *Photoemission in Solids I: General Principles* (Springer-Verlag, Berlin, 1978).
3. J. C. Fuggle and N. Mårtensson, "Core-Level Binding Energies in Metals," *J. Electron Spectrosc. Relat. Phenom.* **21**, 275 (1980).

[Elements Hydrogen \(1\) to Ag \(47\)](#)

[Elements Cadmium \(48\) to Ytterbium\(70\)](#)

[Elements Lutetium \(71\) to Uranium \(92\)](#)

Table 1-1. *Electron binding energies, in electron volts, for the elements in their natural forms.*

Element	K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}	N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
1 H	13.6											
2 He	24.6*											
3 Li	54.7*											
4 Be	111.5*											
5 B	188*											
6 C	284.2*											
7 N	409.9*	37.3*										
8 O	543.1*	41.6*										
9 F	696.7*											
10 Ne	870.2*	48.5*	21.7*	21.6*								
11 Na	1070.8†	63.5†	30.65	30.81								
12 Mg	1303.0†	88.7	49.78	49.50								
13 Al	1559.6	117.8	72.95	72.55								
14 Si	1839	149.7*b	99.82	99.42								
15 P	2145.5	189*	136*	135*								
16 S	2472	230.9	163.6*	162.5*								
17 Cl	2822.4	270*	202*	200*								
18 Ar	3205.9*	326.3*	250.6†	248.4*	29.3*	15.9*	15.7*					
19 K	3608.4*	378.6*	297.3*	294.6*	34.8*	18.3*	18.3*					
20 Ca	4038.5*	438.4†	349.7†	346.2†	44.3 †	25.4†	25.4†					
21 Sc	4492	498.0*	403.6*	398.7*	51.1*	28.3*	28.3*					
22 Ti	4966	560.9†	460.2†	453.8†	58.7†	32.6†	32.6†					

23	V	5465	626.7†	519.8†	512.1†	66.3†	37.2†	37.2†					
24	Cr	5989	696.0†	583.8†	574.1†	74.1†	42.2†	42.2†					
25	Mn	6539	769.1†	649.9†	638.7†	82.3†	47.2†	47.2†					
26	Fe	7112	844.6†	719.9†	706.8†	91.3†	52.7†	52.7†					
27	Co	7709	925.1†	793.2†	778.1†	101.0†	58.9†	59.9†					
28	Ni	8333	1008.6†	870.0†	852.7†	110.8†	68.0†	66.2†					
29	Cu	8979	1096.7†	952.3†	932.7	122.5†	77.3†	75.1†					
30	Zn	9659	1196.2*	1044.9*	1021.8*	139.8*	91.4*	88.6*	10.2*	10.1*			
31	Ga	10367	1299.0*b	1143.2†	1116.4†	159.5†	103.5†	100.0†	18.7†	18.7†			
32	Ge	11103	1414.6*b	1248.1*b	1217.0*b	180.1*	124.9*	120.8*	29.8	29.2			
33	As	11867	1527.0*b	1359.1*b	1323.6*b	204.7*	146.2*	141.2*	41.7*	41.7*			
34	Se	12658	1652.0*b	1474.3*b	1433.9*b	229.6*	166.5*	160.7*	55.5*	54.6*			
35	Br	13474	1782*	1596*	1550*	257*	189*	182*	70*	69*			
36	Kr	14326	1921	1730.9*	1678.4*	292.8*	222.2*	214.4	95.0*	93.8*	27.5*	14.1*	14.1*
37	Rb	15200	2065	1864	1804	326.7*	248.7*	239.1*	113.0*	112*	30.5*	16.3*	15.3 *
38	Sr	16105	2216	2007	1940	358.7†	280.3†	270.0†	136.0†	134.2†	38.9†	21.3	20.1†
39	Y	17038	2373	2156	2080	392.0*b	310.6*	298.8*	157.7†	155.8†	43.8*	24.4*	23.1*
40	Zr	17998	2532	2307	2223	430.3†	343.5†	329.8†	181.1†	178.8†	50.6†	28.5†	27.1†
41	Nb	18986	2698	2465	2371	466.6†	376.1†	360.6†	205.0†	202.3†	56.4†	32.6†	30.8†
42	Mo	20000	2866	2625	2520	506.3†	411.6†	394.0†	231.1†	227.9†	63.2†	37.6†	35.5†
43	Tc	21044	3043	2793	2677	544*	447.6	417.7	257.6	253.9*	69.5*	42.3*	39.9*
44	Ru	22117	3224	2967	2838	586.1*	483.5†	461.4†	284.2†	280.0†	75.0†	46.3†	43.2†
45	Rh	23220	3412	3146	3004	628.1†	521.3†	496.5†	311.9†	307.2†	81.4*b	50.5†	47.3†
46	Pd	24350	3604	3330	3173	671.6†	559.9†	532.3†	340.5†	335.2†	87.1*b	55.7†a	50.9†
47	Ag	25514	3806	3524	3351	719.0†	603.8†	573.0†	374.0†	368.3	97.0†	63.7†	58.3†

Table 1-1. Electron binding energies (continued).

Element	K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}	N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
48 Cd	26711	4018	3727	3538	772.0†	652.6†	618.4†	411.9†	405.2†	109.8†	63.9†a	63.9†a
49 In	27940	4238	3938	3730	827.2†	703.2†	665.3†	451.4†	443.9†	122.9†	73.5†a	73.5†a
50 Sn	29200	4465	4156	3929	884.7†	756.5†	714.6†	493.2†	484.9†	137.1†	83.6†a	83.6†a
51 Sb	30491	4698	4380	4132	946†	812.7†	766.4†	537.5†	528.2†	153.2†	95.6†a	95.6†a
52 Te	31814	4939	4612	4341	1006†	870.8†	820.0†	583.4†	573.0†	169.4†	103.3†a	103.3†a
53 I	33169	5188	4852	4557	1072*	931*	875*	630.8	619.3	186*	123*	123*
54 Xe	34561	5453	5107	4786	1148.7*	1002.1*	940.6*	689.0*	676.4*	213.2*	146.7	145.5*
55 Cs	35985	5714	5359	5012	1211*b	1071*	1003*	740.5*	726.6*	232.3*	172.4*	161.3*
56 Ba	37441	5989	5624	5247	1293*b	1137*b	1063*b	795.7†	780.5*	253.5†	192	178.6†
57 La	38925	6266	5891	5483	1362*b	1209*b	1128*b	853*	836*	274.7*	205.8	196.0*
58 Ce	40443	6549	6164	5723	1436*b	1274*b	1187*b	902.4*	883.8*	291.0*	223.2	206.5*
59 Pr	41991	6835	6440	5964	1511	1337	1242	948.3*	928.8*	304.5	236.3	217.6
60 Nd	43569	7126	6722	6208	1575	1403	1297	1003.3*	980.4*	319.2*	243.3	224.6
61 Pm	45184	7428	7013	6459	—	1471	1357	1052	1027	—	242	242
62 Sm	46834	7737	7312	6716	1723	1541	1420	1110.9*	1083.4*	347.2*	265.6	247.4
63 Eu	48519	8052	7617	6977	1800	1614	1481	1158.6*	1127.5*	360	284	257
64 Gd	50239	8376	7930	7243	1881	1688	1544	1221.9*	1189.6*	378.6*	286	271
65 Tb	51996	8708	8252	7514	1968	1768	1611	1276.9*	1241.1*	396.0*	322.4*	284.1*
66 Dy	53789	9046	8581	7790	2047	1842	1676	1333	1292.6*	414.2*	333.5*	293.2*
67 Ho	55618	9394	8918	8071	2128	1923	1741	1392	1351	432.4*	343.5	308.2*
68 Er	57486	9751	9264	8358	2207	2006	1812	1453	1409	449.8*	366.2	320.2*
69 Tm	59390	10116	9617	8648	2307	2090	1885	1515	1468	470.9*	385.9*	332.6*
70 Yb	61332	10486	9978	8944	2398	2173	1950	1576	1528	480.5*	388.7*	339.7*

Table 1-1. *Electron binding energies (continued).*

Element	N ₄ 4d _{3/2}	N ₅ 4d _{5/2}	N ₆ 4f _{5/2}	N ₇ 4f _{7/2}	O ₁ 5s	O ₂ 5p _{1/2}	O ₃ 5p _{3/2}	O ₄ 5d _{3/2}	O ₅ 5d _{5/2}	P ₁ 6s	P ₂ 6p _{1/2}	P ₃ 6p _{3/2}
48 Cd	11.7†	10.7†										
49 In	17.7†	16.9†										
50 Sn	24.9†	23.9†										
51 Sb	33.3†	32.1†										
52 Te	41.9†	40.4†										
53 I	50.6	48.9										
54 Xe	69.5*	67.5*	—	—	23.3*	13.4*	12.1*					
55 Cs	79.8*	77.5*	—	—	22.7	14.2*	12.1*					
56 Ba	92.6†	89.9†	—	—	30.3†	17.0†	14.8†					
57 La	105.3*	102.5*	—	—	34.3*	19.3*	16.8*					
58 Ce	109*	—	0.1	0.1	37.8	19.8*	17.0*					
59 Pr	115.1*	115.1*	2.0	2.0	37.4	22.3	22.3					
60 Nd	120.5*	120.5*	1.5	1.5	37.5	21.1	21.1					
61 Pm	120	120	—	—	—	—	—					
62 Sm	129	129	5.2	5.2	37.4	21.3	21.3					
63 Eu	133	127.7*	0	0	32	22	22					
64 Gd	—	142.6*	8.6*	8.6*	36	28	21					
65 Tb	150.5*	150.5*	7.7*	2.4*	45.6*	28.7*	22.6*					
66 Dy	153.6*	153.6*	8.0*	4.3*	49.9*	26.3	26.3					
67 Ho	160*	160*	8.6*	5.2*	49.3*	30.8*	24.1*					
68 Er	167.6*	167.6*	—	4.7*	50.6*	31.4*	24.7*					
69 Tm	175.5*	175.5*	—	4.6	54.7*	31.8*	25.0*					
70 Yb	191.2*	182.4*	2.5*	1.3*	52.0*	30.3*	24.1*					

Table 1-1. Electron binding energies (continued).

Element	K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}	N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
71 Lu	63314	10870	10349	9244	2491	2264	2024	1639	1589	506.8*	412.4*	359.2*
72 Hf	65351	11271	10739	9561	2601	2365	2108	1716	1662	538*	438.2†	380.7†
73 Ta	67416	11682	11136	9881	2708	2469	2194	1793	1735	563.4†	463.4†	400.9†
74 W	69525	12100	11544	10207	2820	2575	2281	1872	1809	594.1†	490.4†	423.6†
75 Re	71676	12527	11959	10535	2932	2682	2367	1949	1883	625.4†	518.7†	446.8†
76 Os	73871	12968	12385	10871	3049	2792	2457	2031	1960	658.2†	549.1†	470.7†
77 Ir	76111	13419	12824	11215	3174	2909	2551	2116	2040	691.1†	577.8†	495.8†
78 Pt	78395	13880	13273	11564	3296	3027	2645	2202	2122	725.4†	609.1†	519.4†
79 Au	80725	14353	13734	11919	3425	3148	2743	2291	2206	762.1†	642.7†	546.3†
80 Hg	83102	14839	14209	12284	3562	3279	2847	2385	2295	802.2†	680.2†	576.6†
81 Tl	85530	15347	14698	12658	3704	3416	2957	2485	2389	846.2†	720.5†	609.5†
82 Pb	88005	15861	15200	13035	3851	3554	3066	2586	2484	891.8†	761.9†	643.5†
83 Bi	90524	16388	15711	13419	3999	3696	3177	2688	2580	939†	805.2†	678.8†
84 Po	93105	16939	16244	13814	4149	3854	3302	2798	2683	995*	851*	705*
85 At	95730	17493	16785	14214	4317	4008	3426	2909	2787	1042*	886*	740*
86 Rn	98404	18049	17337	14619	4482	4159	3538	3022	2892	1097*	929*	768*
87 Fr	101137	18639	17907	15031	4652	4327	3663	3136	3000	1153*	980*	810*
88 Ra	103922	19237	18484	15444	4822	4490	3792	3248	3105	1208*	1058	879*
89 Ac	106755	19840	19083	15871	5002	4656	3909	3370	3219	1269*	1080*	890*
90 Th	109651	20472	19693	16300	5182	4830	4046	3491	3332	1330*	1168*	966.4†
91 Pa	112601	21105	20314	16733	5367	5001	4174	3611	3442	1387*	1224*	1007*
92 U	115606	21757	20948	17166	5548	5182	4303	3728	3552	1439*b	1271*b	1043†

Table 1-1. Electron binding energies (continued).

Element	N ₄ 4d _{3/2}	N ₅ 4d _{5/2}	N ₆ 4f _{5/2}	N ₇ 4f _{7/2}	O ₁ 5s	O ₂ 5p _{1/2}	O ₃ 5p _{3/2}	O ₄ 5d _{3/2}	O ₅ 5d _{5/2}	P ₁ 6s	P ₂ 6p _{1/2}	P ₃ 6p _{3/2}
71 Lu	206.1*	196.3*	8.9*	7.5*	57.3*	33.6*	26.7*					
72 Hf	220.0†	211.5†	15.9†	14.2†	64.2†	38*	29.9†					
73 Ta	237.9†	226.4†	23.5†	21.6†	69.7†	42.2*	32.7†					
74 W	255.9†	243.5†	33.6*	31.4†	75.6†	45.3*b	36.8†					
75 Re	273.9†	260.5†	42.9*	40.5*	83†	45.6*	34.6*b					
76 Os	293.1†	278.5†	53.4†	50.7†	84*	58*	44.5†					
77 Ir	311.9†	296.3†	63.8†	60.8†	95.2*b	63.0*b	48.0†					
78 Pt	331.6†	314.6†	74.5†	71.2†	101.7*b	65.3*b	51.7†					
79 Au	353.2†	335.1†	87.6†	84.0	107.2*b	74.2†	57.2†					
80 Hg	378.2†	358.8†	104.0†	99.9†	127†	83.1†	64.5†	9.6†	7.8†			
81 Tl	405.7†	385.0†	122.2†	117.8†	136.0*b	94.6†	73.5†	14.7†	12.5†			
82 Pb	434.3†	412.2†	141.7†	136.9†	147*b	106.4†	83.3†	20.7†	18.1†			
83 Bi	464.0†	440.1†	162.3†	157.0†	159.3*b	119.0†	92.6†	26.9†	23.8†			
84 Po	500*	473*	184*	184*	177*	132*	104*	31*	31*			
85 At	533*	507	210*	210*	195*	148*	115*	40*	40*			
86 Rn	567*	541*	238*	238*	214*	164*	127*	48*	48*	26		
87 Fr	603*	577*	268*	268*	234*	182*	140*	58*	58*	34	15	15
88 Ra	636*	603*	299*	299*	254*	200*	153*	68*	68*	44	19	19
89 Ac	675*	639*	319*	319*	272*	215*	167*	80*	80*	—	—	—
90 Th	712.1†	675.2†	342.4†	333.1†	290*a	229*a	182*a	92.5†	85.4†	41.4†	24.5†	16.6†
91 Pa	743*	708*	371*	360*	310*	232*	232*	94*	94*	—	—	—
92 U	778.3†	736.2†	388.2*	377.4†	321*ab	257*ab	192*ab	102.8†	94.2†	43.9†	26.8†	16.8†

Table 1-1. *Electron binding energies, in electron volts, for the elements H to Ti in their natural forms.*

Element	K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}
1 H	13.6						
2 He	24.6*						
3 Li	54.7*						
4 Be	111.5*						
5 B	188*						
6 C	284.2*						
7 N	409.9*	37.3*					
8 O	543.1*	41.6*					
9 F	696.7*						
10 Ne	870.2*	48.5*	21.7*	21.6*			
11 Na	1070.8†	63.5†	30.65	30.81			
12 Mg	1303.0†	88.7	49.78	49.50			
13 Al	1559.6	117.8	72.95	72.55			
14 Si	1839	149.7*b	99.82	99.42			

15 P	2145.5	189*	136*	135*			
16 S	2472	230.9	163.6*	162.5*			
17 Cl	2822.4	270*	202*	200*			
18 Ar	3205.9*	326.3*	250.6†	248.4*	29.3*	15.9*	15.7*
19 K	3608.4*	378.6*	297.3*	294.6*	34.8*	18.3*	18.3*
20 Ca	4038.5*	438.4†	349.7†	346.2†	44.3 †	25.4†	25.4†
21 Sc	4492	498.0*	403.6*	398.7*	51.1*	28.3*	28.3*
22 Ti	4966	560.9†	460.2†	453.8†	58.7†	32.6†	32.6†

Table 1-1. *Electron binding energies, in electron volts, for the elements V to Ag in their natural forms.*

Element	K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}	N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
23 V	5465	626.7†	519.8†	512.1†	66.3†	37.2†	37.2†					
24 Cr	5989	696.0†	583.8†	574.1†	74.1†	42.2†	42.2†					
25 Mn	6539	769.1†	649.9†	638.7†	82.3†	47.2†	47.2†					
26 Fe	7112	844.6†	719.9†	706.8†	91.3†	52.7†	52.7†					
27 Co	7709	925.1†	793.2†	778.1†	101.0†	58.9†	59.9†					

Table 1-2

28	Ni	8333	1008.6†	870.0†	852.7†	110.8†	68.0†	66.2†					
29	Cu	8979	1096.7†	952.3†	932.7	122.5†	77.3†	75.1†					
30	Zn	9659	1196.2*	1044.9*	1021.8*	139.8*	91.4*	88.6*	10.2*	10.1*			
31	Ga	10367	1299.0*b	1143.2†	1116.4†	159.5†	103.5†	100.0†	18.7†	18.7†			
32	Ge	11103	1414.6*b	1248.1*b	1217.0*b	180.1*	124.9*	120.8*	29.8	29.2			
33	As	11867	1527.0*b	1359.1*b	1323.6*b	204.7*	146.2*	141.2*	41.7*	41.7*			
34	Se	12658	1652.0*b	1474.3*b	1433.9*b	229.6*	166.5*	160.7*	55.5*	54.6*			
35	Br	13474	1782*	1596*	1550*	257*	189*	182*	70*	69*			
36	Kr	14326	1921	1730.9*	1678.4*	292.8*	222.2*	214.4	95.0*	93.8*	27.5*	14.1*	14.1*
37	Rb	15200	2065	1864	1804	326.7*	248.7*	239.1*	113.0*	112*	30.5*	16.3*	15.3*
38	Sr	16105	2216	2007	1940	358.7†	280.3†	270.0†	136.0†	134.2†	38.9†	21.3	20.1†
39	Y	17038	2373	2156	2080	392.0*b	310.6*	298.8*	157.7†	155.8†	43.8*	24.4*	23.1*
40	Zr	17998	2532	2307	2223	430.3†	343.5†	329.8†	181.1†	178.8†	50.6†	28.5†	27.1†
41	Nb	18986	2698	2465	2371	466.6†	376.1†	360.6†	205.0†	202.3†	56.4†	32.6†	30.8†
42	Mo	20000	2866	2625	2520	506.3†	411.6†	394.0†	231.1†	227.9†	63.2†	37.6†	35.5†
43	Tc	21044	3043	2793	2677	544*	447.6	417.7	257.6	253.9*	69.5*	42.3*	39.9*

Table 1-2

44	Ru	22117	3224	2967	2838	586.1*	483.5†	461.4†	284.2†	280.0†	75.0†	46.3†	43.2†
45	Rh	23220	3412	3146	3004	628.1†	521.3†	496.5†	311.9†	307.2†	81.4*b	50.5†	47.3†
46	Pd	24350	3604	3330	3173	671.6†	559.9†	532.3†	340.5†	335.2†	87.1*b	55.7†a	50.9†
47	Ag	25514	3806	3524	3351	719.0†	603.8†	573.0†	374.0†	368.3	97.0†	63.7†	58.3†

Table 1-1. *Electron binding energies, in electron volts, for the elements Cd (48) to Yb (70) in their natural forms.*

Element	K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}	N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
48 Cd	26711	4018	3727	3538	772.0†	652.6†	618.4†	411.9†	405.2†	109.8†	63.9†a	63.9†a
49 In	27940	4238	3938	3730	827.2†	703.2†	665.3†	451.4†	443.9†	122.9†	73.5†a	73.5†a
50 Sn	29200	4465	4156	3929	884.7†	756.5†	714.6†	493.2†	484.9†	137.1†	83.6†a	83.6†a
51 Sb	30491	4698	4380	4132	946†	812.7†	766.4†	537.5†	528.2†	153.2†	95.6†a	95.6†a
52 Te	31814	4939	4612	4341	1006†	870.8†	820.0†	583.4†	573.0†	169.4†	103.3†a	103.3†a
53 I	33169	5188	4852	4557	1072*	931*	875*	630.8	619.3	186*	123*	123*
54 Xe	34561	5453	5107	4786	1148.7*	1002.1*	940.6*	689.0*	676.4*	213.2*	146.7	145.5*
55 Cs	35985	5714	5359	5012	1211*b	1071*	1003*	740.5*	726.6*	232.3*	172.4*	161.3*
56 Ba	37441	5989	5624	5247	1293*b	1137*b	1063*b	795.7†	780.5*	253.5†	192	178.6†
57 La	38925	6266	5891	5483	1362*b	1209*b	1128*b	853*	836*	274.7*	205.8	196.0*
58 Ce	40443	6549	6164	5723	1436*b	1274*b	1187*b	902.4*	883.8*	291.0*	223.2	206.5*
59 Pr	41991	6835	6440	5964	1511	1337	1242	948.3*	928.8*	304.5	236.3	217.6
60 Nd	43569	7126	6722	6208	1575	1403	1297	1003.3*	980.4*	319.2*	243.3	224.6

61	Pm	45184	7428	7013	6459	—	1471	1357	1052	1027	—	242	242
62	Sm	46834	7737	7312	6716	1723	1541	1420	1110.9*	1083.4*	347.2*	265.6	247.4
63	Eu	48519	8052	7617	6977	1800	1614	1481	1158.6*	1127.5*	360	284	257
64	Gd	50239	8376	7930	7243	1881	1688	1544	1221.9*	1189.6*	378.6*	286	271
65	Tb	51996	8708	8252	7514	1968	1768	1611	1276.9*	1241.1*	396.0*	322.4*	284.1*
66	Dy	53789	9046	8581	7790	2047	1842	1676	1333	1292.6*	414.2*	333.5*	293.2*
67	Ho	55618	9394	8918	8071	2128	1923	1741	1392	1351	432.4*	343.5	308.2*
68	Er	57486	9751	9264	8358	2207	2006	1812	1453	1409	449.8*	366.2	320.2*
69	Tm	59390	10116	9617	8648	2307	2090	1885	1515	1468	470.9*	385.9*	332.6*
70	Yb	61332	10486	9978	8944	2398	2173	1950	1576	1528	480.5*	388.7*	339.7*

Table 1-1. Electron binding energies (continued).

Element	N ₄ 4d _{3/2}	N ₅ 4d _{5/2}	N ₆ 4f _{5/2}	N ₇ 4f _{7/2}	O ₁ 5s	O ₂ 5p _{1/2}	O ₃ 5p _{3/2}
48 Cd	11.7†	10.7†					
49 In	17.7†	16.9†					
50 Sn	24.9†	23.9†					

Table 1-2

51	Sb	33.3†	32.1†					
52	Te	41.9†	40.4†					
53	I	50.6	48.9					
54	Xe	69.5*	67.5*	—	—	23.3*	13.4*	12.1*
55	Cs	79.8*	77.5*	—	—	22.7	14.2*	12.1*
56	Ba	92.6†	89.9†	—	—	30.3†	17.0†	14.8†
57	La	105.3*	102.5*	—	—	34.3*	19.3*	16.8*
58	Ce	109*	—	0.1	0.1	37.8	19.8*	17.0*
59	Pr	115.1*	115.1*	2.0	2.0	37.4	22.3	22.3
60	Nd	120.5*	120.5*	1.5	1.5	37.5	21.1	21.1
61	Pm	120	120	—	—	—	—	—
62	Sm	129	129	5.2	5.2	37.4	21.3	21.3
63	Eu	133	127.7*	0	0	32	22	22
64	Gd	—	142.6*	8.6*	8.6*	36	28	21
65	Tb	150.5*	150.5*	7.7*	2.4*	45.6*	28.7*	22.6*
66	Dy	153.6*	153.6*	8.0*	4.3*	49.9*	26.3	26.3

Table 1-2

67	Ho	160*	160*	8.6*	5.2*	49.3*	30.8*	24.1*
68	Er	167.6*	167.6*	—	4.7*	50.6*	31.4*	24.7*
69	Tm	175.5*	175.5*	—	4.6	54.7*	31.8*	25.0*
70	Yb	191.2*	182.4*	2.5*	1.3*	52.0*	30.3*	24.1*

Table 1-1. *Electron binding energies, in electron volts, for the elements in their natural forms.*

Element	K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}	N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
71 Lu	63314	10870	10349	9244	2491	2264	2024	1639	1589	506.8*	412.4*	359.2*
72 Hf	65351	11271	10739	9561	2601	2365	2108	1716	1662	538*	438.2†	380.7†
73 Ta	67416	11682	11136	9881	2708	2469	2194	1793	1735	563.4†	463.4†	400.9†
74 W	69525	12100	11544	10207	2820	2575	2281	1872	1809	594.1†	490.4†	423.6†
75 Re	71676	12527	11959	10535	2932	2682	2367	1949	1883	625.4†	518.7†	446.8†
76 Os	73871	12968	12385	10871	3049	2792	2457	2031	1960	658.2†	549.1†	470.7†
77 Ir	76111	13419	12824	11215	3174	2909	2551	2116	2040	691.1†	577.8†	495.8†
78 Pt	78395	13880	13273	11564	3296	3027	2645	2202	2122	725.4†	609.1†	519.4†
79 Au	80725	14353	13734	11919	3425	3148	2743	2291	2206	762.1†	642.7†	546.3†
80 Hg	83102	14839	14209	12284	3562	3279	2847	2385	2295	802.2†	680.2†	576.6†
81 Tl	85530	15347	14698	12658	3704	3416	2957	2485	2389	846.2†	720.5†	609.5†
82 Pb	88005	15861	15200	13035	3851	3554	3066	2586	2484	891.8†	761.9†	643.5†
83 Bi	90524	16388	15711	13419	3999	3696	3177	2688	2580	939†	805.2†	678.8†

84	Po	93105	16939	16244	13814	4149	3854	3302	2798	2683	995*	851*	705*
85	At	95730	17493	16785	14214	4317	4008	3426	2909	2787	1042*	886*	740*
86	Rn	98404	18049	17337	14619	4482	4159	3538	3022	2892	1097*	929*	768*
87	Fr	101137	18639	17907	15031	4652	4327	3663	3136	3000	1153*	980*	810*
88	Ra	103922	19237	18484	15444	4822	4490	3792	3248	3105	1208*	1058	879*
89	Ac	106755	19840	19083	15871	5002	4656	3909	3370	3219	1269*	1080*	890*
90	Th	109651	20472	19693	16300	5182	4830	4046	3491	3332	1330*	1168*	966.4†
91	Pa	112601	21105	20314	16733	5367	5001	4174	3611	3442	1387*	1224*	1007*
92	U	115606	21757	20948	17166	5548	5182	4303	3728	3552	1439*b	1271*b	1043†

Table 1-1. Electron binding energies (continued).

Element	N ₄ 4d _{3/2}	N ₅ 4d _{5/2}	N ₆ 4f _{5/2}	N ₇ 4f _{7/2}	O ₁ 5s	O ₂ 5p _{1/2}	O ₃ 5p _{3/2}	O ₄ 5d _{3/2}	O ₅ 5d _{5/2}	P ₁ 6s	P ₂ 6p _{1/2}	P ₃ 6p _{3/2}
71 Lu	206.1*	196.3*	8.9*	7.5*	57.3*	33.6*	26.7*					
72 Hf	220.0†	211.5†	15.9†	14.2†	64.2†	38*	29.9†					
73 Ta	237.9†	226.4†	23.5†	21.6†	69.7†	42.2*	32.7†					
74 W	255.9†	243.5†	33.6*	31.4†	75.6†	45.3*b	36.8†					

Table 1-2

75	Re	273.9†	260.5†	42.9*	40.5*	83†	45.6*	34.6*b				
76	Os	293.1†	278.5†	53.4†	50.7†	84*	58*	44.5†				
77	Ir	311.9†	296.3†	63.8†	60.8†	95.2*b	63.0*b	48.0†				
78	Pt	331.6†	314.6†	74.5†	71.2†	101.7*b	65.3*b	51.7†				
79	Au	353.2†	335.1†	87.6†	84.0	107.2*b	74.2†	57.2†				
80	Hg	378.2†	358.8†	104.0†	99.9†	127†	83.1†	64.5†	9.6†	7.8†		
81	Tl	405.7†	385.0†	122.2†	117.8†	136.0*b	94.6†	73.5†	14.7†	12.5†		
82	Pb	434.3†	412.2†	141.7†	136.9†	147*b	106.4†	83.3†	20.7†	18.1†		
83	Bi	464.0†	440.1†	162.3†	157.0†	159.3*b	119.0†	92.6†	26.9†	23.8†		
84	Po	500*	473*	184*	184*	177*	132*	104*	31*	31*		
85	At	533*	507	210*	210*	195*	148*	115*	40*	40*		
86	Rn	567*	541*	238*	238*	214*	164*	127*	48*	48*	26	
87	Fr	603*	577*	268*	268*	234*	182*	140*	58*	58*	34	15
88	Ra	636*	603*	299*	299*	254*	200*	153*	68*	68*	44	19
89	Ac	675*	639*	319*	319*	272*	215*	167*	80*	80*	—	—
90	Th	712.1†	675.2†	342.4†	333.1†	290*a	229*a	182*a	92.5†	85.4†	41.4†	24.5†
											16.6†	

91	Pa	743*	708*	371*	360*	310*	232*	232*	94*	94*	—	—	—
92	U	778.3†	736.2†	388.2*	377.4†	321*ab	257*ab	192*ab	102.8†	94.2†	43.9†	26.8†	16.8†

X-Ray Data Booklet

Section 1.2 X-RAY EMISSION ENERGIES

Jeffrey B. Kortright and Albert C. Thompson

In [Table 1-2 \(pdf format\)](#), characteristic K , L , and M x-ray line energies are given for elements with $3 \leq Z \leq 95$. Only the strongest lines are included: $K\alpha_1$, $K\alpha_2$, $K\beta_1$, $L\alpha_1$, $L\alpha_2$, $L\beta_1$, $L\beta_2$, $L\gamma_1$, and $M\alpha_1$. Wavelengths, in angstroms, can be obtained from the relation $\lambda = 12,3984/E$, where E is in eV. The data in the table were based on Ref. 1, which should be consulted for a more complete listing. Widths of the $K\alpha$ lines can be found in Ref. 2.

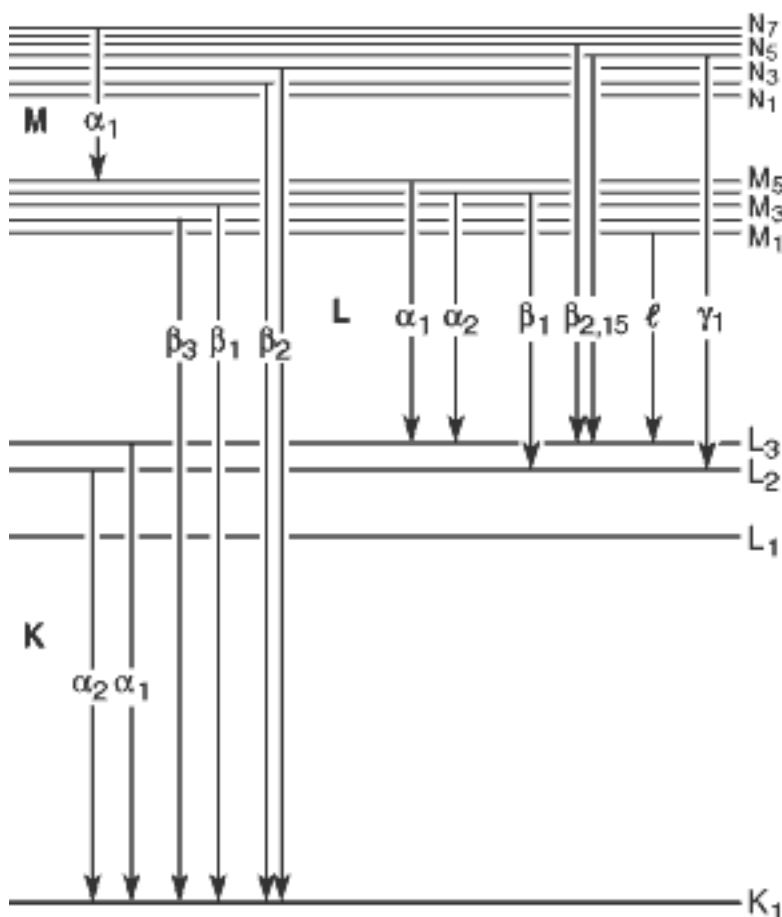


Fig 1-1. Transitions that give rise to the various emission lines.

[Table 1-3 \(pdf format\)](#) provides a listing of these, and additional, lines (arranged by increasing energy), together with relative intensities. An intensity of 100 is assigned to the strongest line in each shell for each element. Figure 1-1 illustrates the transitions that give rise to the lines in Table 1-3.

REFERENCES

1. J. A. Bearden, "X-Ray Wavelengths," *Rev. Mod. Phys.* **39**, 78 (1967).
2. M. O. Krause and J. H. Oliver, "Natural Widths of Atomic *K* and *L* Levels, *K* α X-Ray Lines and Several *KLL* Auger Lines," *J. Phys. Chem. Ref. Data* **8**, 329 (1979).

X-Ray Data Booklet Table 1-2. Photon energies, in electron volts, of principal K-, L-, and M-shell emission lines.

Element	$K\alpha_1$	$K\alpha_2$	$K\beta_1$	$L\alpha_1$	$L\alpha_2$	$L\beta_1$	$L\beta_2$	$L\gamma_1$	$M\alpha_1$
3 Li	54.3								
4 Be	108.5								
5 B	183.3								
6 C	277								
7 N	392.4								
8 O	524.9								
9 F	676.8								
10 Ne	848.6	848.6							
11 Na	1,040.98	1,040.98	1,071.1						
12 Mg	1,253.60	1,253.60	1,302.2						
13 Al	1,486.70	1,486.27	1,557.45						
14 Si	1,739.98	1,739.38	1,835.94						
15 P	2,013.7	2,012.7	2,139.1						
16 S	2,307.84	2,306.64	2,464.04						
17 Cl	2,622.39	2,620.78	2,815.6						
18 Ar	2,957.70	2,955.63	3,190.5						
19 K	3,313.8	3,311.1	3,589.6						
20 Ca	3,691.68	3,688.09	4,012.7	341.3	341.3	344.9			
21 Sc	4,090.6	4,086.1	4,460.5	395.4	395.4	399.6			

Table 1-2. Energies of x-ray emission lines (continued).

Element	$K\alpha_1$	$K\alpha_2$	$K\beta_1$	$L\alpha_1$	$L\alpha_2$	$L\beta_1$	$L\beta_2$	$L\gamma_1$	$M\alpha_1$
22 Ti	4,510.84	4,504.86	4,931.81	452.2	452.2	458.4			
23 V	4,952.20	4,944.64	5,427.29	511.3	511.3	519.2			
24 Cr	5,414.72	5,405.509	5,946.71	572.8	572.8	582.8			
25 Mn	5,898.75	5,887.65	6,490.45	637.4	637.4	648.8			
26 Fe	6,403.84	6,390.84	7,057.98	705.0	705.0	718.5			
27 Co	6,930.32	6,915.30	7,649.43	776.2	776.2	791.4			
28 Ni	7,478.15	7,460.89	8,264.66	851.5	851.5	868.8			
29 Cu	8,047.78	8,027.83	8,905.29	929.7	929.7	949.8			
30 Zn	8,638.86	8,615.78	9,572.0	1,011.7	1,011.7	1,034.7			
31 Ga	9,251.74	9,224.82	10,264.2	1,097.92	1,097.92	1,124.8			
32 Ge	9,886.42	9,855.32	10,982.1	1,188.00	1,188.00	1,218.5			
33 As	10,543.72	10,507.99	11,726.2	1,282.0	1,282.0	1,317.0			
34 Se	11,222.4	11,181.4	12,495.9	1,379.10	1,379.10	1,419.23			
35 Br	11,924.2	11,877.6	13,291.4	1,480.43	1,480.43	1,525.90			
36 Kr	12,649	12,598	14,112	1,586.0	1,586.0	1,636.6			
37 Rb	13,395.3	13,335.8	14,961.3	1,694.13	1,692.56	1,752.17			
38 Sr	14,165	14,097.9	15,835.7	1,806.56	1,804.74	1,871.72			
39 Y	14,958.4	14,882.9	16,737.8	1,922.56	1,920.47	1,995.84			
40 Zr	15,775.1	15,690.9	17,667.8	2,042.36	2,039.9	2,124.4	2,219.4	2,302.7	

41 Nb	16,615.1	16,521.0	18,622.5	2,165.89	2,163.0	2,257.4	2,367.0	2,461.8	
42 Mo	17,479.34	17,374.3	19,608.3	2,293.16	2,289.85	2,394.81	2,518.3	2,623.5	
43 Tc	18,367.1	18,250.8	20,619	2,424	2,420	2,538	2,674	2,792	
44 Ru	19,279.2	19,150.4	21,656.8	2,558.55	2,554.31	2,683.23	2,836.0	2,964.5	
45 Rh	20,216.1	20,073.7	22,723.6	2,696.74	2,692.05	2,834.41	3,001.3	3,143.8	
46 Pd	21,177.1	21,020.1	23,818.7	2,838.61	2,833.29	2,990.22	3,171.79	3,328.7	
47 Ag	22,162.92	21,990.3	24,942.4	2,984.31	2,978.21	3,150.94	3,347.81	3,519.59	
48 Cd	23,173.6	22,984.1	26,095.5	3,133.73	3,126.91	3,316.57	3,528.12	3,716.86	
49 In	24,209.7	24,002.0	27,275.9	3,286.94	3,279.29	3,487.21	3,713.81	3,920.81	
50 Sn	25,271.3	25,044.0	28,486.0	3,443.98	3,435.42	3,662.80	3,904.86	4,131.12	
51 Sb	26,359.1	26,110.8	29,725.6	3,604.72	3,595.32	3,843.57	4,100.78	4,347.79	
52 Te	27,472.3	27,201.7	30,995.7	3,769.33	3,758.8	4,029.58	4,301.7	4,570.9	
53 I	28,612.0	28,317.2	32,294.7	3,937.65	3,926.04	4,220.72	4,507.5	4,800.9	
54 Xe	29,779	29,458	33,624	4,109.9	—	—	—	—	
55 Cs	30,972.8	30,625.1	34,986.9	4,286.5	4,272.2	4,619.8	4,935.9	5,280.4	
56 Ba	32,193.6	31,817.1	36,378.2	4,466.26	4,450.90	4,827.53	5,156.5	5,531.1	
57 La	33,441.8	33,034.1	37,801.0	4,650.97	4,634.23	5,042.1	5,383.5	5,788.5	833
58 Ce	34,719.7	34,278.9	39,257.3	4,840.2	4,823.0	5,262.2	5,613.4	6,052	883
59 Pr	36,026.3	35,550.2	40,748.2	5,033.7	5,013.5	5,488.9	5,850	6,322.1	929
60 Nd	37,361.0	36,847.4	42,271.3	5,230.4	5,207.7	5,721.6	6,089.4	6,602.1	978
61 Pm	38,724.7	38,171.2	43,826	5,432.5	5,407.8	5,961	6,339	6,892	—
62 Sm	40,118.1	39,522.4	45,413	5,636.1	5,609.0	6,205.1	6,586	7,178	1,081

Table 1-2. Energies of x-ray emission lines (continued).

Element	$K\alpha_1$	$K\alpha_2$	$K\beta_1$	$L\alpha_1$	$L\alpha_2$	$L\beta_1$	$L\beta_2$	$L\gamma_1$	$M\alpha_1$
63 Eu	41,542.2	40,901.9	47,037.9	5,845.7	5,816.6	6,456.4	6,843.2	7,480.3	1,131
64 Gd	42,996.2	42,308.9	48,697	6,057.2	6,025.0	6,713.2	7,102.8	7,785.8	1,185
65 Tb	44,481.6	43,744.1	50,382	6,272.8	6,238.0	6,978	7,366.7	8,102	1,240
66 Dy	45,998.4	45,207.8	52,119	6,495.2	6,457.7	7,247.7	7,635.7	8,418.8	1,293
67 Ho	47,546.7	46,699.7	53,877	6,719.8	6,679.5	7,525.3	7,911	8,747	1,348
68 Er	49,127.7	48,221.1	55,681	6,948.7	6,905.0	7,810.9	8,189.0	9,089	1,406
69 Tm	50,741.6	49,772.6	57,517	7,179.9	7,133.1	8,101	8,468	9,426	1,462
70 Yb	52,388.9	51,354.0	59,370	7,415.6	7,367.3	8,401.8	8,758.8	9,780.1	1,521.4
71 Lu	54,069.8	52,965.0	61,283	7,655.5	7,604.9	8,709.0	9,048.9	10,143.4	1,581.3
72 Hf	55,790.2	54,611.4	63,234	7,899.0	7,844.6	9,022.7	9,347.3	10,515.8	1,644.6
73 Ta	57,532	56,277	65,223	8,146.1	8,087.9	9,343.1	9,651.8	10,895.2	1,710
74 W	59,318.24	57,981.7	67,244.3	8,397.6	8,335.2	9,672.35	9,961.5	11,285.9	1,775.4
75 Re	61,140.3	59,717.9	69,310	8,652.5	8,586.2	10,010.0	10,275.2	11,685.4	1,842.5
76 Os	63,000.5	61,486.7	71,413	8,911.7	8,841.0	10,355.3	10,598.5	12,095.3	1,910.2
77 Ir	64,895.6	63,286.7	73,560.8	9,175.1	9,099.5	10,708.3	10,920.3	12,512.6	1,979.9
78 Pt	66,832	65,112	75,748	9,442.3	9,361.8	11,070.7	11,250.5	12,942.0	2,050.5
79 Au	68,803.7	66,989.5	77,984	9,713.3	9,628.0	11,442.3	11,584.7	13,381.7	2,122.9
80 Hg	70,819	68,895	80,253	9,988.8	9,897.6	11,822.6	11,924.1	13,830.1	2,195.3
81 Tl	72,871.5	70,831.9	82,576	10,268.5	10,172.8	12,213.3	12,271.5	14,291.5	2,270.6

82 Pb	74,969.4	72,804.2	84,936	10,551.5	10,449.5	12,613.7	12,622.6	14,764.4	2,345.5
83 Bi	77,107.9	74,814.8	87,343	10,838.8	10,730.91	13,023.5	12,979.9	15,247.7	2,422.6
84 Po	79,290	76,862	89,800	11,130.8	11,015.8	13,447	13,340.4	15,744	—
85 At	81,520	78,950	92,300	11,426.8	11,304.8	13,876	—	16,251	—
86 Rn	83,780	81,070	94,870	11,727.0	11,597.9	14,316	—	16,770	—
87 Fr	86,100	83,230	97,470	12,031.3	11,895.0	14,770	14,450	17,303	—
88 Ra	88,470	85,430	100,130	12,339.7	12,196.2	15,235.8	14,841.4	17,849	—
89 Ac	90,884	87,670	102,850	12,652.0	12,500.8	15,713	—	18,408	—
90 Th	93,350	89,953	105,609	12,968.7	12,809.6	16,202.2	15,623.7	18,982.5	2,996.1
91 Pa	95,868	92,287	108,427	13,290.7	13,122.2	16,702	16,024	19,568	3,082.3
92 U	98,439	94,665	111,300	13,614.7	13,438.8	17,220.0	16,428.3	20,167.1	3,170.8
93 Np	—	—	—	13,944.1	13,759.7	17,750.2	16,840.0	20,784.8	—
94 Pu	—	—	—	14,278.6	14,084.2	18,293.7	17,255.3	21,417.3	—
95 Am	—	—	—	14,617.2	14,411.9	18,852.0	17,676.5	22,065.2	—

X-Ray Data Booklet Table 1-3. Photon energies and relative intensities of K-, L-, and M-shell lines shown in Fig. 1-1, arranged by increasing energy. An intensity of 100 is assigned to the strongest line in each shell for each element.

Energy (eV)	Element	Line	Relative intensity								
54.3	3 Li	K $\alpha_{1,2}$	150	524.9	8 O	K $\alpha_{1,2}$	151	851.5	28 Ni	L $\alpha_{1,2}$	111
108.5	4 Be	K $\alpha_{1,2}$	150	556.3	25 Mn	Ll	15	868.8	28 Ni	L β_1	68
183.3	5 B	K $\alpha_{1,2}$	151	572.8	24 Cr	L $\alpha_{1,2}$	111	883	58 Ce	M α_1	100
277	6 C	K $\alpha_{1,2}$	147	582.8	24 Cr	L β_1	79	884	30 Zn	Ll	7
348.3	21 Sc	Ll	21	615.2	26 Fe	Ll	10	929.2	59 Pr	M α_1	100
392.4	7 N	K $\alpha_{1,2}$	150	637.4	25 Mn	L $\alpha_{1,2}$	111	929.7	29 Cu	L $\alpha_{1,2}$	111
395.3	22 Ti	Ll	46	648.8	25 Mn	L β_1	77	949.8	29 Cu	L β_1	65
395.4	21 Sc	L $\alpha_{1,2}$	111	676.8	9 F	K $\alpha_{1,2}$	148	957.2	31 Ga	Ll	7
399.6	21 Sc	L β_1	77	677.8	27 Co	Ll	10	978	60 Nd	M α_1	100
446.5	23 V	Ll	28	705.0	26 Fe	L $\alpha_{1,2}$	111	1,011.7	30 Zn	L $\alpha_{1,2}$	111
452.2	22 Ti	L $\alpha_{1,2}$	111	718.5	26 Fe	L β_1	66	1,034.7	30 Zn	L β_1	65
458.4	22 Ti	L β_1	79	742.7	28 Ni	Ll	9	1,036.2	32 Ge	Ll	6
500.3	24 Cr	Ll	17	776.2	27 Co	L $\alpha_{1,2}$	111	1,041.0	11 Na	K $\alpha_{1,2}$	150
511.3	23 V	L $\alpha_{1,2}$	111	791.4	27 Co	L β_1	76	1,081	62 Sm	M α_1	100
519.2	23 V	L β_1	80	811.1	29 Cu	Ll	8	1,097.9	31 Ga	L $\alpha_{1,2}$	111
				833	57 La	M α_1	100	1,120	33 As	Ll	6
				848.6	10 Ne	K $\alpha_{1,2}$	150	1,124.8	31 Ga	L β_1	66

Table 1-3. *Energies and intensities of x-ray emission lines (continued).*

Energy (eV)	Element	Line	Relative intensity								
1,131	63 Eu	M α_1	100	1,462	69 Tm	M α_1	100	1,740.0	14 Si	K α_1	100
1,185	64 Gd	M α_1	100	1,480.4	35 Br	L $\alpha_{1,2}$	111	1,752.2	37 Rb	L β_1	58
1,188.0	32 Ge	L $\alpha_{1,2}$	111	1,482.4	37 Rb	Ll	5	1,775.4	74 W	M α_1	100
1,204.4	34 Se	Ll	6	1,486.3	13 Al	K α_2	50	1,792.0	40 Zr	Ll	5
1,218.5	32 Ge	L β_1	60	1,486.7	13 Al	K α_1	100	1,804.7	38 Sr	L α_2	11
1,240	65 Tb	M α_1	100	1,521.4	70 Yb	M α_1	100	1,806.6	38 Sr	L α_1	100
1,253.6	12 Mg	K $\alpha_{1,2}$	150	1,525.9	35 Br	L β_1	59	1,835.9	14 Si	K β_1	2
1,282.0	33 As	L $\alpha_{1,2}$	111	1,557.4	13 Al	K β_1	1	1,842.5	75 Re	M α_1	100
1,293	66 Dy	M α_1	100	1,581.3	71 Lu	M α_1	100	1,871.7	38 Sr	L β_1	58
1,293.5	35 Br	Ll	5	1,582.2	38 Sr	Ll	5	1,902.2	41 Nb	Ll	5
1,317.0	33 As	L β_1	60	1,586.0	36 Kr	L $\alpha_{1,2}$	111	1,910.2	76 Os	M α_1	100
1,348	67 Ho	M α_1	100	1,636.6	36 Kr	L β_1	57	1,920.5	39 Y	L α_2	11
1,379.1	34 Se	L $\alpha_{1,2}$	111	1,644.6	72 Hf	M α_1	100	1,922.6	39 Y	L α_1	100
1,386	36 Kr	Ll	5	1,685.4	39 Y	Ll	5	1,979.9	77 Ir	M α_1	100
1,406	68 Er	M α_1	100	1,692.6	37 Rb	L α_2	11	1,995.8	39 Y	L β_1	57
1,419.2	34 Se	L β_1	59	1,694.1	37 Rb	L α_1	100	2,012.7	15 P	K α_2	50
				1,709.6	73 Ta	M α_1	100	2,013.7	15 P	K α_1	100
				1,739.4	14 Si	K α_2	50	2,015.7	42 Mo	Ll	5

2,039.9	40 Zr	L α_2	11
2,042.4	40 Zr	L α_1	100
2,050.5	78 Pt	M α_1	100
2,122	43 Tc	Ll	5
2,122.9	79 Au	M α_1	100
2,124.4	40 Zr	L β_1	54
2,139.1	15 P	K β_1	3
2,163.0	41 Nb	L α_2	11
2,165.9	41 Nb	L α_1	100
2,195.3	80 Hg	M α_1	100
2,219.4	40 Zr	L $\beta_{2,15}$	1
2,252.8	44 Ru	Ll	4
2,257.4	41 Nb	L β_1	52
2,270.6	81 Tl	M α_1	100
2,289.8	42 Mo	L α_2	11
2,293.2	42 Mo	L α_1	100
2,302.7	40 Zr	L γ_1	2
2,306.6	16 S	K α_2	50
2,307.8	16 S	K α_1	100
2,345.5	82 Pb	M α_1	100

2,367.0	41 Nb	L $\beta_{2,15}$	3
2,376.5	45 Rh	Ll	4
2,394.8	42 Mo	L β_1	53
2,420	43 Tc	L α_2	11
2,422.6	83 Bi	M α_1	100
2,424	43 Tc	L α_1	100
2,461.8	41 Nb	L γ_1	2
2,464.0	16 S	K β_1	5
2,503.4	46 Pd	Ll	4
2,518.3	42 Mo	L $\beta_{2,15}$	5
2,538	43 Tc	L β_1	54
2,554.3	44 Ru	L α_2	11
2,558.6	44 Ru	L α_1	100
2,620.8	17 Cl	K α_2	50
2,622.4	17 Cl	K α_1	100
2,623.5	42 Mo	L γ_1	3
2,633.7	47 Ag	Ll	4
2,674	43 Tc	L $\beta_{2,15}$	7
2,683.2	44 Ru	L β_1	54
2,692.0	45 Rh	L α_2	11

2,696.7	45 Rh	L α_1	100
2,767.4	48 Cd	Ll	4
2,792	43 Tc	L γ_1	3
2,815.6	17 Cl	K β_1	6
2,833.3	46 Pd	L α_2	11
2,834.4	45 Rh	L β_1	52
2,836.0	44 Ru	L $\beta_{2,15}$	10
2,838.6	46 Pd	L α_1	100
2,904.4	49 In	Ll	4
2,955.6	18 Ar	K α_2	50
2,957.7	18 Ar	K α_1	100
2,964.5	44 Ru	L γ_1	4
2,978.2	47 Ag	L α_2	11
2,984.3	47 Ag	L α_1	100
2,990.2	46 Pd	L β_1	53
2,996.1	90 Th	M α_1	100
3,001.3	45 Rh	L $\beta_{2,15}$	10
3,045.0	50 Sn	Ll	4
3,126.9	48 Cd	L α_2	11
3,133.7	48 Cd	L α_1	100

Table 1-3. Energies and intensities of x-ray emission lines (continued).

Energy (eV)	Element	Line	Relative intensity								
3,143.8	45 Rh	L γ_1	5	3,487.2	49 In	L β_1	58	3,937.6	53 I	L α_1	100
3,150.9	47 Ag	L β_1	56	3,519.6	47 Ag	L γ_1	6	3,954.1	56 Ba	Ll	4
3,170.8	92 U	M α_1	100	3,528.1	48 Cd	L $\beta_{2,15}$	15	4,012.7	20 Ca	K $\beta_{1,3}$	13
3,171.8	46 Pd	L $\beta_{2,15}$	12	3,589.6	19 K	K $\beta_{1,3}$	11	4,029.6	52 Te	L β_1	61
3,188.6	51 Sb	Ll	4	3,595.3	51 Sb	L α_2	11	4,086.1	21 Sc	K α_2	50
3,190.5	18 Ar	K $\beta_{1,3}$	10	3,604.7	51 Sb	L α_1	100	4,090.6	21 Sc	K α_1	100
3,279.3	49 In	L α_2	11	3,636	54 Xe	Ll	4	4,093	54 Xe	L α_2	11
3,286.9	49 In	L α_1	100	3,662.8	50 Sn	L β_1	60	4,100.8	51 Sb	L $\beta_{2,15}$	17
3,311.1	19 K	K α_2	50	3,688.1	20 Ca	K α_2	50	4,109.9	54 Xe	L α_1	100
3,313.8	19 K	K α_1	100	3,691.7	20 Ca	K α_1	100	4,124	57 La	Ll	4
3,316.6	48 Cd	L β_1	58	3,713.8	49 In	L $\beta_{2,15}$	15	4,131.1	50 Sn	L γ_1	7
3,328.7	46 Pd	L γ_1	6	3,716.9	48 Cd	L γ_1	6	4,220.7	53 I	L β_1	61
3,335.6	52 Te	Ll	4	3,758.8	52 Te	L α_2	11	4,272.2	55 Cs	L α_2	11
3,347.8	47 Ag	L $\beta_{2,15}$	13	3,769.3	52 Te	L α_1	100	4,286.5	55 Cs	L α_1	100
3,435.4	50 Sn	L α_2	11	3,795.0	55 Cs	Ll	4	4,287.5	58 Ce	Ll	4
3,444.0	50 Sn	L α_1	100	3,843.6	51 Sb	L β_1	61	4,301.7	52 Te	L $\beta_{2,15}$	18
3,485.0	53 I	Ll	4	3,904.9	50 Sn	L $\beta_{2,15}$	16	4,347.8	51 Sb	L γ_1	8
				3,920.8	49 In	L γ_1	6	4,414	54 Xe	L β_1	60
				3,926.0	53 I	L α_2	11	4,450.9	56 Ba	L α_2	11

4,453.2	59 Pr	Ll	4
4,460.5	21 Sc	K $\beta_{1,3}$	15
4,466.3	56 Ba	L α_1	100
4,504.9	22 Ti	K α_2	50
4,507.5	53 I	L $\beta_{2,15}$	19
4,510.8	22 Ti	K α_1	100
4,570.9	52 Te	L γ_1	8
4,619.8	55 Cs	L β_1	61
4,633.0	60 Nd	Ll	4
4,634.2	57 La	L α_2	11
4,651.0	57 La	L α_1	100
4,714	54 Xe	L $\beta_{2,15}$	20
4,800.9	53 I	L γ_1	8
4,809	61 Pm	Ll	4
4,823.0	58 Ce	L α_2	11
4,827.5	56 Ba	L β_1	60
4,840.2	58 Ce	L α_1	100
4,931.8	22 Ti	K $\beta_{1,3}$	15
4,935.9	55 Cs	L $\beta_{2,15}$	20
4,944.6	23 V	K α_2	50

4,952.2	23 V	K α_1	100
4,994.5	62 Sm	Ll	4
5,013.5	59 Pr	L α_2	11
5,033.7	59 Pr	L α_1	100
5,034	54 Xe	L γ_1	8
5,042.1	57 La	L β_1	60
5,156.5	56 Ba	L $\beta_{2,15}$	20
5,177.2	63 Eu	Ll	4
5,207.7	60 Nd	L α_2	11
5,230.4	60 Nd	L α_1	100
5,262.2	58 Ce	L β_1	61
5,280.4	55 Cs	L γ_1	8
5,362.1	64 Gd	Ll	4
5,383.5	57 La	L $\beta_{2,15}$	21
5,405.5	24 Cr	K α_2	50
5,408	61 Pm	L α_2	11
5,414.7	24 Cr	K α_1	100
5,427.3	23 V	K $\beta_{1,3}$	15
5,432	61 Pm	L α_1	100
5,488.9	59 Pr	L β_1	61

5,531.1	56 Ba	L γ_1	9
5,546.7	65 Tb	Ll	4
5,609.0	62 Sm	L α_2	11
5,613.4	58 Ce	L $\beta_{2,15}$	21
5,636.1	62 Sm	L α_1	100
5,721.6	60 Nd	L β_1	60
5,743.1	66 Dy	Ll	4
5,788.5	57 La	L γ_1	9
5,816.6	63 Eu	L α_2	11
5,845.7	63 Eu	L α_1	100
5,850	59 Pr	L $\beta_{2,15}$	21
5,887.6	25 Mn	K α_2	50
5,898.8	25 Mn	K α_1	100
5,943.4	67 Ho	Ll	4
5,946.7	24 Cr	K $\beta_{1,3}$	15
5,961	61 Pm	L β_1	61
6,025.0	64 Gd	L α_2	11
6,052	58 Ce	L γ_1	9
6,057.2	64 Gd	L α_1	100
6,089.4	60 Nd	L $\beta_{2,15}$	21

Table 1-3. Energies and intensities of x-ray emission lines (continued).

Energy (eV)	Element	Line	Relative intensity								
6,152	68 Er	Ll	4	6,713.2	64 Gd	L β_1	62	7,367.3	70 Yb	L α_2	11
6,205.1	62 Sm	L β_1	61	6,719.8	67 Ho	L α_1	100	7,387.8	74 W	Ll	5
6,238.0	65 Tb	L α_2	11	6,752.8	71 Lu	Ll	4	7,415.6	70 Yb	L α_1	100
6,272.8	65 Tb	L α_1	100	6,843.2	63 Eu	L $\beta_{2,15}$	21	7,460.9	28 Ni	K α_2	51
6,322.1	59 Pr	L γ_1	9	6,892	61 Pm	L γ_1	10	7,478.2	28 Ni	K α_1	100
6,339	61 Pm	L β_2	21	6,905.0	68 Er	L α_2	11	7,480.3	63 Eu	L γ_1	10
6,341.9	69 Tm	Ll	4	6,915.3	27 Co	K α_2	51	7,525.3	67 Ho	L β_1	64
6,390.8	26 Fe	K α_2	50	6,930.3	27 Co	K α_1	100	7,603.6	75 Re	Ll	5
6,403.8	26 Fe	K α_1	100	6,948.7	68 Er	L α_1	100	7,604.9	71 Lu	L α_2	11
6,456.4	63 Eu	L β_1	62	6,959.6	72 Hf	Ll	5	7,635.7	66 Dy	L β_2	20
6,457.7	66 Dy	L α_2	11	6,978	65 Tb	L β_1	61	7,649.4	27 Co	K $\beta_{1,3}$	17
6,490.4	25 Mn	K $\beta_{1,3}$	17	7,058.0	26 Fe	K $\beta_{1,3}$	17	7,655.5	71 Lu	L α_1	100
6,495.2	66 Dy	L α_1	100	7,102.8	64 Gd	L $\beta_{2,15}$	21	7,785.8	64 Gd	L γ_1	11
6,545.5	70 Yb	Ll	4	7,133.1	69 Tm	L α_2	11	7,810.9	68 Er	L β_1	64
6,587.0	62 Sm	L $\beta_{2,15}$	21	7,173.1	73 Ta	Ll	5	7,822.2	76 Os	Ll	5
6,602.1	60 Nd	L γ_1	10	7,178.0	62 Sm	L γ_1	10	7,844.6	72 Hf	L α_2	11
6,679.5	67 Ho	L α_2	11	7,179.9	69 Tm	L α_1	100	7,899.0	72 Hf	L α_1	100
				7,247.7	66 Dy	L β_1	62	7,911	67 Ho	L $\beta_{2,15}$	20
				7,366.7	65 Tb	L $\beta_{2,15}$	21	8,027.8	29 Cu	K α_2	51

8,045.8	77 Ir	Ll	5	8,721.0	80 Hg	Ll	5	9,442.3	78 Pt	L α_1	100
8,047.8	29 Cu	K α_1	100	8,747	67 Ho	L γ_1	11	9,572.0	30 Zn	K $\beta_{1,3}$	17
8,087.9	73 Ta	L α_2	11	8,758.8	70 Yb	L $\beta_{2,15}$	20	9,628.0	79 Au	L α_2	11
8,101	69 Tm	L β_1	64	8,841.0	76 Os	L α_2	11	9,651.8	73 Ta	L β_2	20
8,102	65 Tb	L γ_1	11	8,905.3	29 Cu	K $\beta_{1,3}$	17	9,672.4	74 W	L β_1	67
8,146.1	73 Ta	L α_1	100	8,911.7	76 Os	L α_1	100	9,713.3	79 Au	L α_1	100
8,189.0	68 Er	L $\beta_{2,15}$	20	8,953.2	81 Tl	Ll	6	9,780.1	70 Yb	L γ_1	12
8,264.7	28 Ni	K $\beta_{1,3}$	17	9,022.7	72 Hf	L β_1	67	9,855.3	32 Ge	K α_2	51
8,268	78 Pt	Ll	5	9,048.9	71 Lu	L β_2	19	9,886.4	32 Ge	K α_1	100
8,335.2	74 W	L α_2	11	9,089	68 Er	L γ_1	11	9,897.6	80 Hg	L α_2	11
8,397.6	74 W	L α_1	100	9,099.5	77 Ir	L α_2	11	9,961.5	74 W	L β_2	21
8,401.8	70 Yb	L β_1	65	9,175.1	77 Ir	L α_1	100	9,988.8	80 Hg	L α_1	100
8,418.8	66 Dy	L γ_1	11	9,184.5	82 Pb	Ll	6	10,010.0	75 Re	L β_1	66
8,468	69 Tm	L $\beta_{2,15}$	20	9,224.8	31 Ga	K α_2	51	10,143.4	71 Lu	L γ_1	12
8,493.9	79 Au	Ll	5	9,251.7	31 Ga	K α_1	100	10,172.8	81 Tl	L α_2	11
8,586.2	75 Re	L α_2	11	9,343.1	73 Ta	L β_1	67	10,260.3	31 Ga	K β_3	5
8,615.8	30 Zn	K α_2	51	9,347.3	72 Hf	L β_2	20	10,264.2	31 Ga	K β_1	66
8,638.9	30 Zn	K α_1	100	9,361.8	78 Pt	L α_2	11	10,268.5	81 Tl	L α_1	100
8,652.5	75 Re	L α_1	100	9,420.4	83 Bi	Ll	6	10,275.2	75 Re	L β_2	22
8,709.0	71 Lu	L β_1	66	9,426	69 Tm	L γ_1	12	10,355.3	76 Os	L β_1	67

Table 1-3. Energies and intensities of x-ray emission lines (continued).

Energy (eV)	Element	Line	Relative intensity								
10,449.5	82 Pb	L α_2	11	11,250.5	78 Pt	L β_2	23	12,598	36 Kr	K α_2	52
10,508.0	33 As	K α_2	51	11,285.9	74 W	L γ_1	13	12,613.7	82 Pb	L β_1	66
10,515.8	72 Hf	L γ_1	12	11,442.3	79 Au	L β_1	67	12,622.6	82 Pb	L β_2	25
10,543.7	33 As	K α_1	100	11,584.7	79 Au	L β_2	23	12,649	36 Kr	K α_1	100
10,551.5	82 Pb	L α_1	100	11,618.3	92 U	Ll	7	12,652	34 Se	K β_2	1
10,598.5	76 Os	L β_2	22	11,685.4	75 Re	L γ_1	13	12,809.6	90 Th	L α_2	11
10,708.3	77 Ir	L β_1	66	11,720.3	33 As	K β_3	6	12,942.0	78 Pt	L γ_1	13
10,730.9	83 Bi	L α_2	11	11,726.2	33 As	K β_1	13	12,968.7	90 Th	L α_1	100
10,838.8	83 Bi	L α_1	100	11,822.6	80 Hg	L β_1	67	12,979.9	83 Bi	L β_2	25
10,895.2	73 Ta	L γ_1	12	11,864	33 As	K β_2	1	13,023.5	83 Bi	L β_1	67
10,920.3	77 Ir	L β_2	22	11,877.6	35 Br	K α_2	52	13,284.5	35 Br	K β_3	7
10,978.0	32 Ge	K β_3	6	11,924.1	80 Hg	L β_2	24	13,291.4	35 Br	K β_1	14
10,982.1	32 Ge	K β_1	60	11,924.2	35 Br	K α_1	100	13,335.8	37 Rb	K α_2	52
11,070.7	78 Pt	L β_1	67	12,095.3	76 Os	L γ_1	13	13,381.7	79 Au	L γ_1	13
11,118.6	90 Th	Ll	6	12,213.3	81 Tl	L β_1	67	13,395.3	37 Rb	K α_1	100
11,181.4	34 Se	K α_2	52	12,271.5	81 Tl	L β_2	25	13,438.8	92 U	L α_2	11
11,222.4	34 Se	K α_1	100	12,489.6	34 Se	K β_3	6	13,469.5	35 Br	K β_2	1
				12,495.9	34 Se	K β_1	13	13,614.7	92 U	L α_1	100
				12,512.6	77 Ir	L γ_1	13	13,830.1	80 Hg	L γ_1	14

14,097.9	38 Sr	K α_2	52
14,104	36 Kr	K β_3	7
14,112	36 Kr	K β_1	14
14,165.0	38 Sr	K α_1	100
14,291.5	81 Tl	L γ_1	14
14,315	36 Kr	K β_2	2
14,764.4	82 Pb	L γ_1	14
14,882.9	39 Y	K α_2	52
14,951.7	37 Rb	K β_3	7
14,958.4	39 Y	K α_1	100
14,961.3	37 Rb	K β_1	14
15,185	37 Rb	K β_2	2
15,247.7	83 Bi	L γ_1	14
15,623.7	90 Th	L β_2	26
15,690.9	40 Zr	K α_2	52
15,775.1	40 Zr	K α_1	100
15,824.9	38 Sr	K β_3	7
15,835.7	38 Sr	K β_1	14
16,084.6	38 Sr	K β_2	3

16,202.2	90 Th	L β_1	69
16,428.3	92 U	L β_2	26
16,521.0	41 Nb	K α_2	52
16,615.1	41 Nb	K α_1	100
16,725.8	39 Y	K β_3	8
16,737.8	39 Y	K β_1	15
17,015.4	39 Y	K β_2	3
17,220.0	92 U	L β_1	61
17,374.3	42 Mo	K α_2	52
17,479.3	42 Mo	K α_1	100
17,654	40 Zr	K β_3	8
17,667.8	40 Zr	K β_1	15
17,970	40 Zr	K β_2	3
18,250.8	43 Tc	K α_2	53
18,367.1	43 Tc	K α_1	100
18,606.3	41 Nb	K β_3	8
18,622.5	41 Nb	K β_1	15
18,953	41 Nb	K β_2	3
18,982.5	90 Th	L γ_1	16

19,150.4	44 Ru	K α_2	53
19,279.2	44 Ru	K α_1	100
19,590.3	42 Mo	K β_3	8
19,608.3	42 Mo	K β_1	15
19,965.2	42 Mo	K β_2	3
20,073.7	45 Rh	K α_2	53
20,167.1	92 U	L γ_1	15
20,216.1	45 Rh	K α_1	100
20,599	43 Tc	K β_3	8
20,619	43 Tc	K β_1	16
21,005	43 Tc	K β_2	4
21,020.1	46 Pd	K α_2	53
21,177.1	46 Pd	K α_1	100
21,634.6	44 Ru	K β_3	8
21,656.8	44 Ru	K β_1	16
21,990.3	47 Ag	K α_2	53
22,074	44 Ru	K β_2	4
22,162.9	47 Ag	K α_1	100
22,698.9	45 Rh	K β_3	8

Table 1-3. Energies and intensities of x-ray emission lines (continued).

Energy (eV)	Element	Line	Relative intensity	Energy	Element	Line	Relative intensity	Energy	Element	Line	Relative intensity
22,723.6	45 Rh	K β_1	16	26,359.1	51 Sb	K α_1	100	30,972.8	55 Cs	K α_1	100
22,984.1	48 Cd	K α_2	53	26,643.8	48 Cd	K β_2	4	30,995.7	52 Te	K β_1	18
23,172.8	45 Rh	K β_2	4	27,201.7	52 Te	K α_2	54	31,700.4	52 Te	K β_2	5
23,173.6	48 Cd	K α_1	100	27,237.7	49 In	K β_3	9	31,817.1	56 Ba	K α_2	54
23,791.1	46 Pd	K β_3	8	27,275.9	49 In	K β_1	17	32,193.6	56 Ba	K α_1	100
23,818.7	46 Pd	K β_1	16	27,472.3	52 Te	K α_1	100	32,239.4	53 I	K β_3	9
24,002.0	49 In	K α_2	53	27,860.8	49 In	K β_2	5	32,294.7	53 I	K β_1	18
24,209.7	49 In	K α_1	100	28,317.2	53 I	K α_2	54	33,034.1	57 La	K α_2	54
24,299.1	46 Pd	K β_2	4	28,444.0	50 Sn	K β_3	9	33,042	53 I	K β_2	5
24,911.5	47 Ag	K β_3	9	28,486.0	50 Sn	K β_1	17	33,441.8	57 La	K α_1	100
24,942.4	47 Ag	K β_1	16	28,612.0	53 I	K α_1	100	33,562	54 Xe	K β_3	9
25,044.0	50 Sn	K α_2	53	29,109.3	50 Sn	K β_2	5	33,624	54 Xe	K β_1	18
25,271.3	50 Sn	K α_1	100	29,458	54 Xe	K α_2	54	34,278.9	58 Ce	K α_2	55
25,456.4	47 Ag	K β_2	4	29,679.2	51 Sb	K β_3	9	34,415	54 Xe	K β_2	5
26,061.2	48 Cd	K β_3	9	29,725.6	51 Sb	K β_1	18	34,719.7	58 Ce	K α_1	100
26,095.5	48 Cd	K β_1	17	29,779	54 Xe	K α_1	100	34,919.4	55 Cs	K β_3	9
26,110.8	51 Sb	K α_2	54	30,389.5	51 Sb	K β_2	5	34,986.9	55 Cs	K β_1	18
				30,625.1	55 Cs	K α_2	54	35,550.2	59 Pr	K α_2	55
				30,944.3	52 Te	K β_3	9	35,822	55 Cs	K β_2	6

36,026.3	59 Pr	K α_1	100
36,304.0	56 Ba	K β_3	10
36,378.2	56 Ba	K β_1	18
36,847.4	60 Nd	K α_2	55
37,257	56 Ba	K β_2	6
37,361.0	60 Nd	K α_1	100
37,720.2	57 La	K β_3	10
37,801.0	57 La	K β_1	19
38,171.2	61 Pm	K α_2	55
38,724.7	61 Pm	K α_1	100
38,729.9	57 La	K β_2	6
39,170.1	58 Ce	K β_3	10
39,257.3	58 Ce	K β_1	19
39,522.4	62 Sm	K α_2	55
40,118.1	62 Sm	K α_1	100
40,233	58 Ce	K β_2	6
40,652.9	59 Pr	K β_3	10
40,748.2	59 Pr	K β_1	19
40,901.9	63 Eu	K α_2	56

41,542.2	63 Eu	K α_1	100
41,773	59 Pr	K β_2	6
42,166.5	60 Nd	K β_3	10
42,271.3	60 Nd	K β_1	19
42,308.9	64 Gd	K α_2	56
42,996.2	64 Gd	K α_1	100
43,335	60 Nd	K β_2	6
43,713	61 Pm	K β_3	10
43,744.1	65 Tb	K α_2	56
43,826	61 Pm	K β_1	19
44,481.6	65 Tb	K α_1	100
44,942	61 Pm	K β_2	6
45,207.8	66 Dy	K α_2	56
45,289	62 Sm	K β_3	10
45,413	62 Sm	K β_1	19
45,998.4	66 Dy	K α_1	100
46,578	62 Sm	K β_2	6
46,699.7	67 Ho	K α_2	56
46,903.6	63 Eu	K β_3	10

47,037.9	63 Eu	K β_1	19
47,546.7	67 Ho	K α_1	100
48,221.1	68 Er	K α_2	56
48,256	63 Eu	K β_2	6
48,555	64 Gd	K β_3	10
48,697	64 Gd	K β_1	20
49,127.7	68 Er	K α_1	100
49,772.6	69 Tm	K α_2	57
49,959	64 Gd	K β_2	7
50,229	65 Tb	K β_3	10
50,382	65 Tb	K β_1	20
50,741.6	69 Tm	K α_1	100
51,354.0	70 Yb	K α_2	57
51,698	65 Tb	K β_2	7
51,957	66 Dy	K β_3	10
52,119	66 Dy	K β_1	20
52,388.9	70 Yb	K α_1	100
52,965.0	71 Lu	K α_2	57
53,476	66 Dy	K β_2	7

Table 1-3. Energies and intensities of x-ray emission lines (continued).

Energy (eV)	Element	Line	Relative intensity	59,370	70 Yb	K β_1	21	66,989.5	79 Au	K α_2	59
				59,717.9	75 Re	K α_2	58	66,990	73 Ta	K β_2	7
53,711	67 Ho	K β_3	11	60,980	70 Yb	K β_2	7	67,244.3	74 W	K β_1	22
53,877	67 Ho	K β_1	20	61,050	71 Lu	K β_3	11	68,803.7	79 Au	K α_1	100
54,069.8	71 Lu	K α_1	100	61,140.3	75 Re	K α_1	100	68,895	80 Hg	K α_2	59
54,611.4	72 Hf	K α_2	57	61,283	71 Lu	K β_1	21	68,994	75 Re	K β_3	12
55,293	67 Ho	K β_2	7	61,486.7	76 Os	K α_2	58	69,067	74 W	K β_2	8
55,494	68 Er	K β_3	11	62,970	71 Lu	K β_2	7	69,310	75 Re	K β_1	22
55,681	68 Er	K β_1	21	62,980	72 Hf	K β_3	11	70,819	80 Hg	K α_1	100
55,790.2	72 Hf	K α_1	100	63,000.5	76 Os	K α_1	100	70,831.9	81 Tl	K α_2	60
56,277	73 Ta	K α_2	57	63,234	72 Hf	K β_1	22	71,077	76 Os	K β_3	12
57,210	68 Er	K β_2	7	63,286.7	77 Ir	K α_2	58	71,232	75 Re	K β_2	8
57,304	69 Tm	K β_3	11	64,895.6	77 Ir	K α_1	100	71,413	76 Os	K β_1	23
57,517	69 Tm	K β_1	21	64,948.8	73 Ta	K β_3	11	72,804.2	82 Pb	K α_2	60
57,532	73 Ta	K α_1	100	64,980	72 Hf	K β_2	7	72,871.5	81 Tl	K α_1	100
57,981.7	74 W	K α_2	58	65,112	78 Pt	K α_2	58	73,202.7	77 Ir	K β_3	12
59,090	69 Tm	K β_2	7	65,223	73 Ta	K β_1	22	73,363	76 Os	K β_2	8
59,140	70 Yb	K β_3	11	66,832	78 Pt	K α_1	100	73,560.8	77 Ir	K β_1	23
59,318.2	74 W	K α_1	100	66,951.4	74 W	K β_3	11	74,814.8	83 Bi	K α_2	60

74,969.4	82 Pb	$K\alpha_1$	100
75,368	78 Pt	$K\beta_3$	12
75,575	77 Ir	$K\beta_2$	8
75,748	78 Pt	$K\beta_1$	23
77,107.9	83 Bi	$K\alpha_1$	100
77,580	79 Au	$K\beta_3$	12
77,850	78 Pt	$K\beta_2$	8
77,984	79 Au	$K\beta_1$	23
79,822	80 Hg	$K\beta_3$	12
80,150	79 Au	$K\beta_2$	8
80,253	80 Hg	$K\beta_1$	23

82,118	81 Tl	$K\beta_3$	12
82,515	80 Hg	$K\beta_2$	8
82,576	81 Tl	$K\beta_1$	23
84,450	82 Pb	$K\beta_3$	12
84,910	81 Tl	$K\beta_2$	8
84,936	82 Pb	$K\beta_1$	23
86,834	83 Bi	$K\beta_3$	12
87,320	82 Pb	$K\beta_2$	8
87,343	83 Bi	$K\beta_1$	23
89,830	83 Bi	$K\beta_2$	9

89,953	90 Th	$K\alpha_2$	62
93,350	90 Th	$K\alpha_1$	100
94,665	92 U	$K\alpha_2$	62
98,439	92 U	$K\alpha_1$	100
104,831	90 Th	$K\beta_3$	12
105,609	90 Th	$K\beta_1$	24
108,640	90 Th	$K\beta_2$	9
110,406	92 U	$K\beta_3$	13
111,300	92 U	$K\beta_1$	24
114,530	92 U	$K\beta_2$	9

X-Ray Data Booklet

Section 1.3 FLUORESCENCE YIELDS FOR *K* and *L* SHELLS

Jeffrey B. Kortright

Fluorescence yields for the *K* and *L* shells for the elements $5 \leq Z \leq 110$ are plotted in Fig. 1-2; the data are based on Ref. 1. These yields represent the probability of a core hole in the *K* or *L* shells being filled by a radiative process, in competition with nonradiative processes. Auger processes are the only nonradiative processes competing with fluorescence for the *K* shell and *L*₃ subshell holes. Auger and Coster-Kronig nonradiative processes compete with fluorescence to fill *L*₁ and *L*₂ subshell holes. Only one curve is presented for the three *L* subshells, representing the average of the *L*₁, *L*₂, and *L*₃ effective fluorescence yields in Ref. 1, which differ by less than about 10% over most of the periodic table. See Ref. 1 for more detail on the *L* subshell rates and the nonradiative rates, and for an appendix containing citations to the theoretical and experimental work upon which Fig. 1-2 is based. Widths of *K* and *L* fluorescence lines can be found in Ref. 2.

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1. M. O. Krause, "Atomic Radiative and Radiationless Yields for *K* and *L* Shells," *J. Phys. Chem. Ref. Data* **8**, 307 (1979).
2. M. O. Krause and J. H. Oliver, "Natural Widths of Atomic *K* and *L* Levels, *K*α X-Ray Lines and Several *KLL* Auger Lines," *J. Phys. Chem. Ref. Data* **8**, 329 (1979).

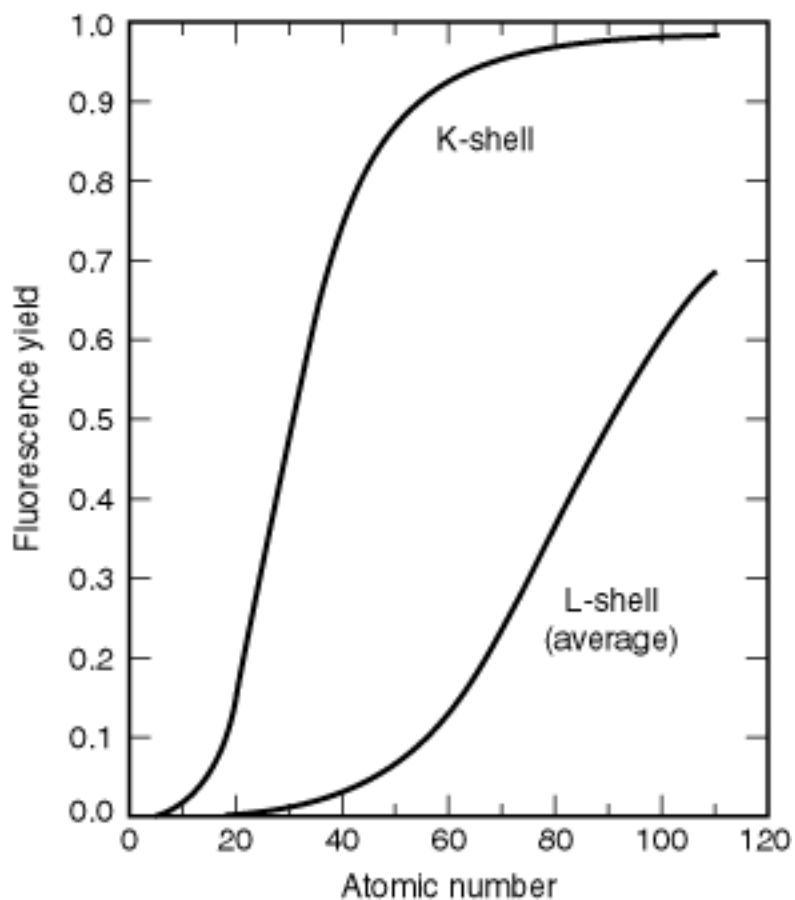


Fig. 1-2. Fluorescence yields for K and L shells for $5 \leq Z \leq 110$. The plotted curve for the L shell represents an average of L_1 , L_2 , and L_3 effective yields.

X-Ray Data Booklet

Section 1.4 PRINCIPAL AUGER ELECTRON ENERGIES

Figure 1-3 has been reproduced by permission of Physical Electronics, Inc., and is taken from Ref. 1. For each element, dots indicate the energies of principal Auger peaks, the pre-dominant ones represented by the heavier dots. The families of Auger transitions are denoted by labels of the form WXY , where W is the shell in which the original vacancy occurs, X is the shell from which the W vacancy is filled, and Y is the shell from which the Auger electron is ejected. The listed references should be consulted for detailed tabulations and for shifted values in several common compounds.

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1. K. D. Childs, B. A. Carlson, L. A. Vanier, J. F. Moulder, D. F. Paul, W. F. Stickle, and D. G. Watson, in C. L. Hedberg, Ed., *Handbook of Auger Electron Spectroscopy* (Physical Electronics, Eden Prairie, MN, 1995).
2. J. F. Moulder, W. F. Stickle, P. E. Sobol, and K. D. Bomben, *Handbook of X-Ray Photoelectron Spectroscopy* (Physical Electronics, Eden Prairie, MN, 1995).
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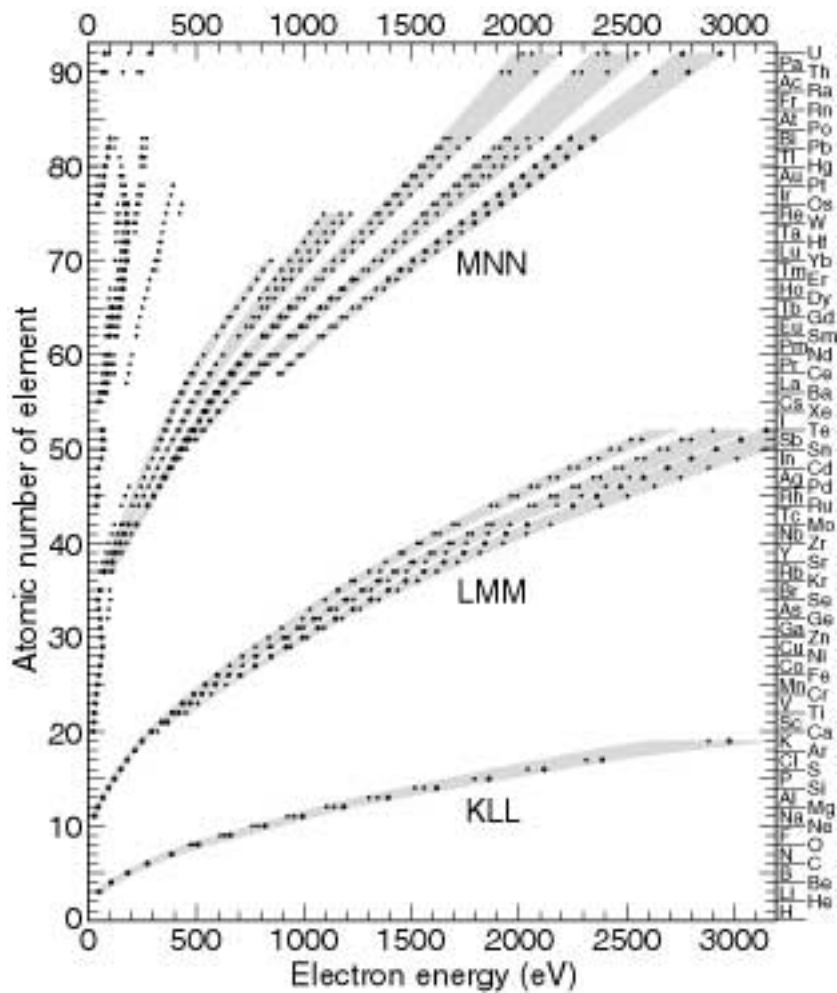


Fig. 1-3. Auger electron energies for the elements. Points indicate the electron energies of the principal Auger peaks for each element. The larger points represent the most intense peaks. (Reproduced by permission of Physical Electronics, Inc.)

X-Ray Data Booklet

Section 1.5 SUBSHELL PHOTOIONIZATION CROSS SECTIONS

Ingolf Lindau

The atomic subshell photoemission cross sections plotted in Fig. 1-4 have been calculated for isolated atoms by Yeh and Lindau [1,2]. The calculations were done with a one-electron central-field frozen-core model using first-order perturbation theory. No single model accurately predicts the photoionization process of all orbitals for all elements from the VUV to 1.5 keV. The complexity of the physics of different atomic orbitals makes it impossible for any single rule to describe all of them. The accuracy of the model used has been discussed in detail by Cooper and Manson [3–5]. A [PDF version of this section](#) is also available.

REFERENCES

1. J.-J. Yeh and I. Lindau, “Atomic Subshell Photoionization Cross Sections and Asymmetry Parameters: $1 < Z < 103$,” *At. Data Nucl. Data Tables* **32**, 1 (1985).
2. J.-J. Yeh, *Atomic Calculations of Photoionization Cross Sections and Asymmetry Parameters* (Gordon and Breach, Langhorne, PA, 1993).
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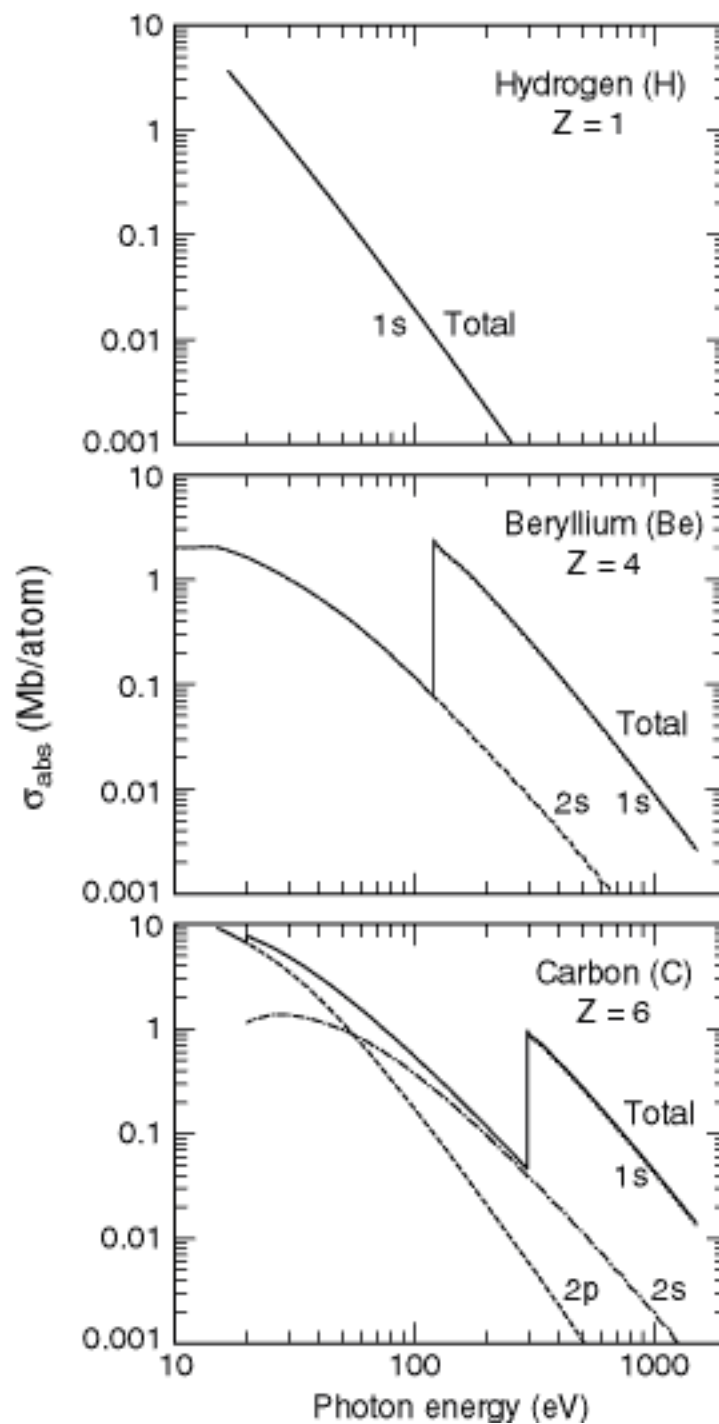


Fig. 1-4. Plots of atomic subshell photoemission cross sections for H, Be and C.

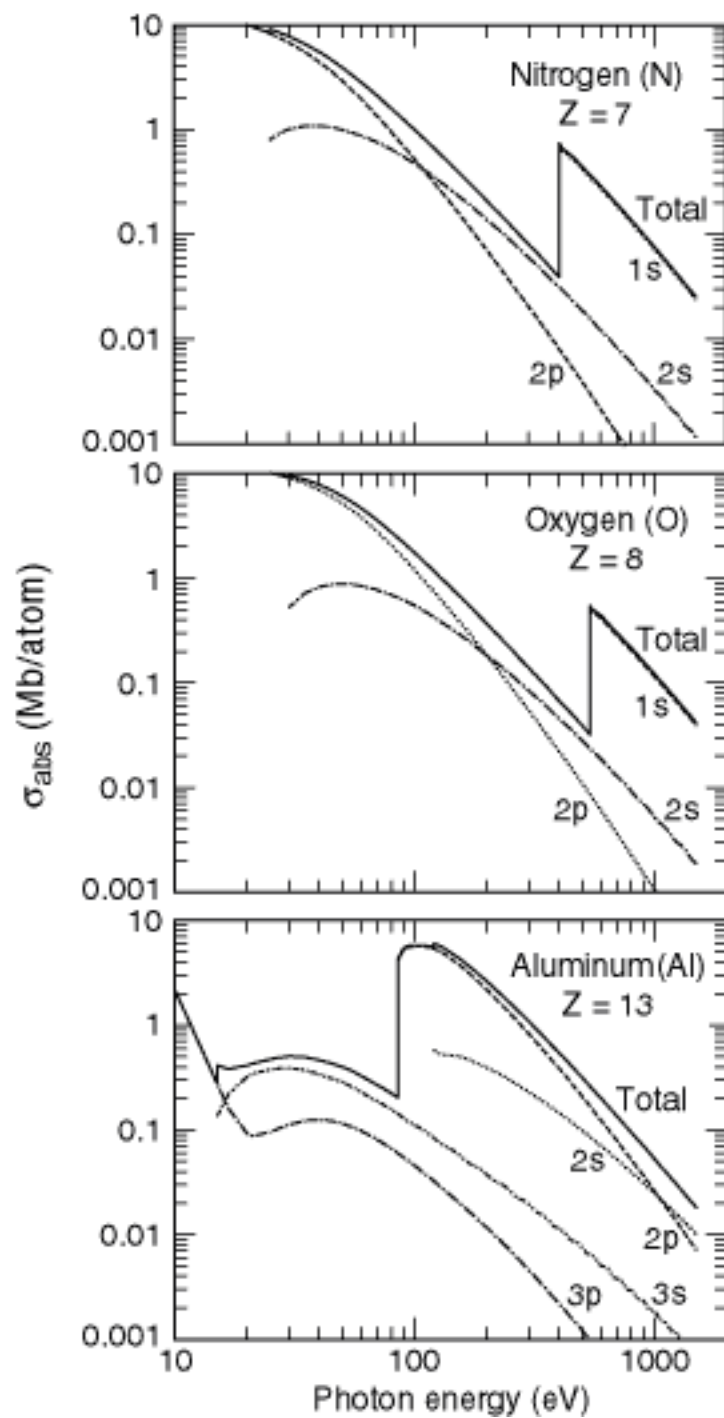


Fig. 1-4. Subshell photoemission cross sections N, O and Al.

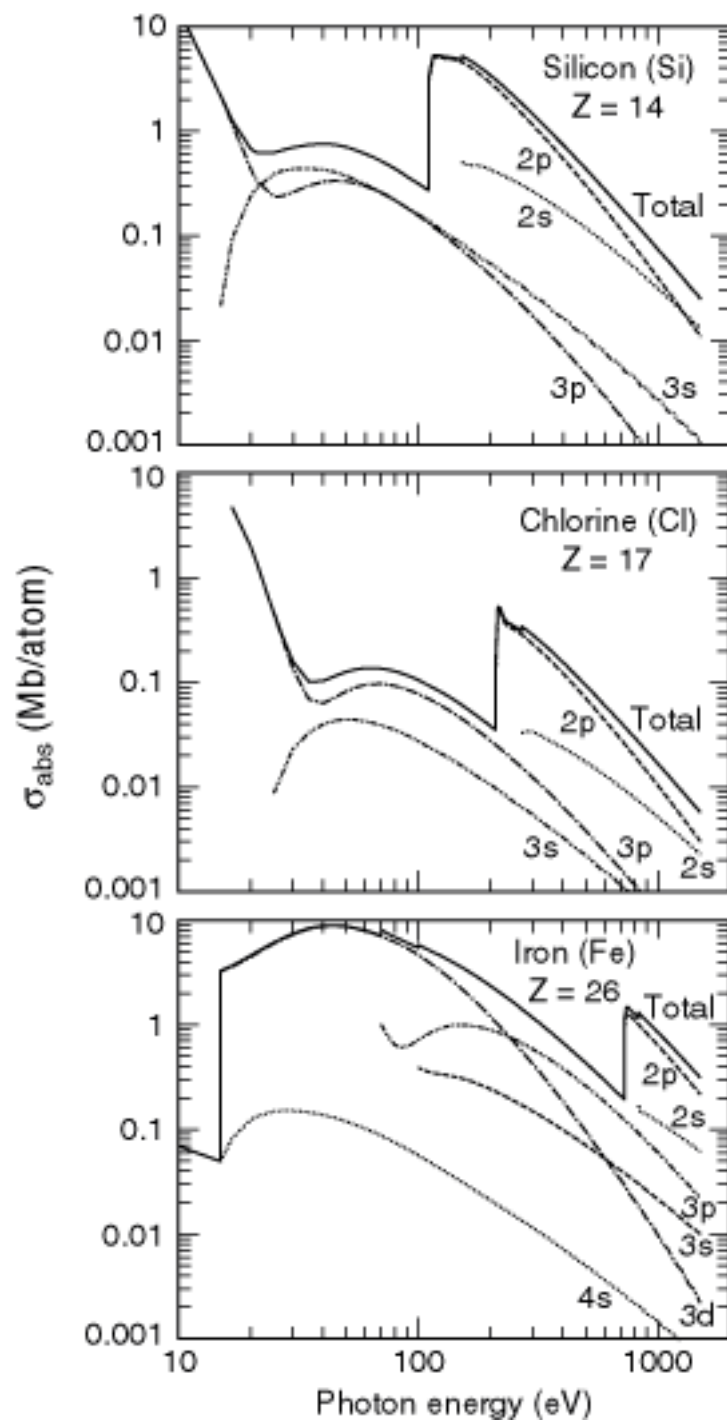


Fig. 1-4. Subshell photoemission cross sections for Si, Cl and Fe.

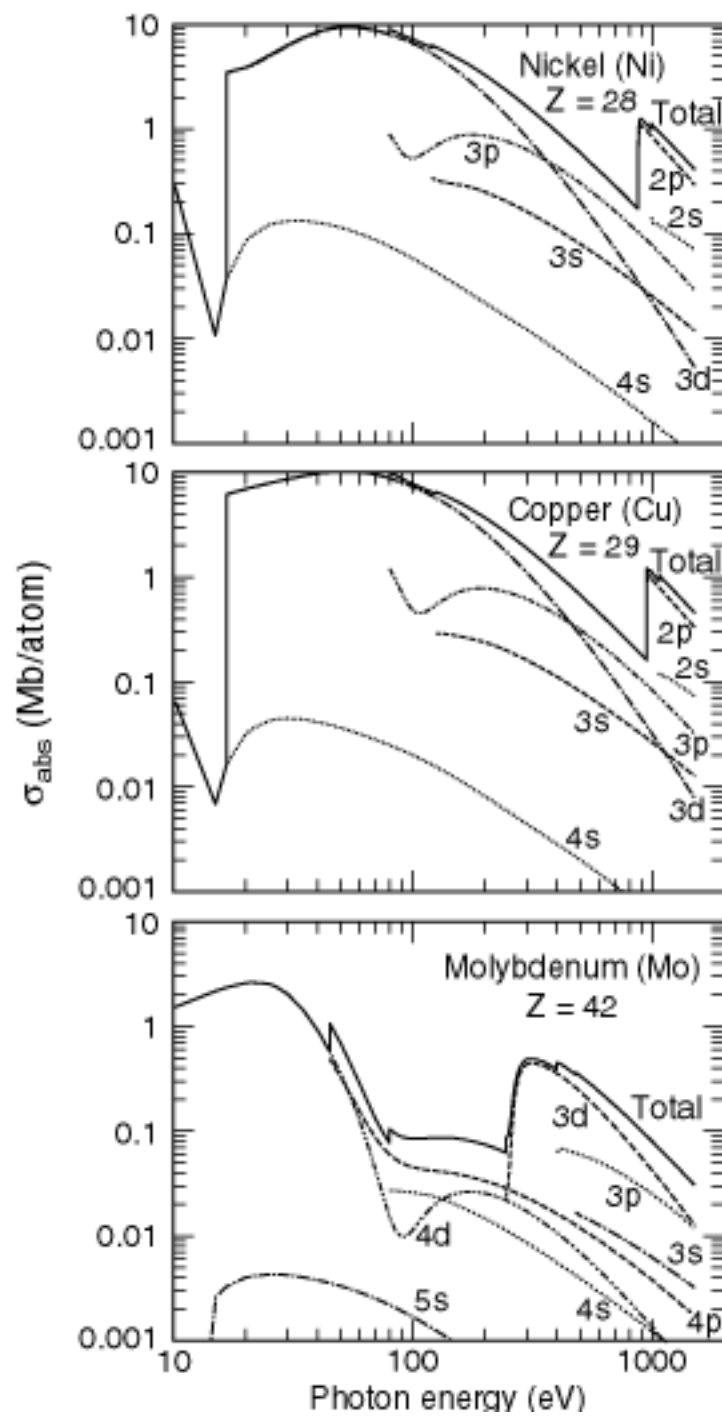


Fig. 1-4. Subshell photoemission cross sections for Ni, Cu and Mo.

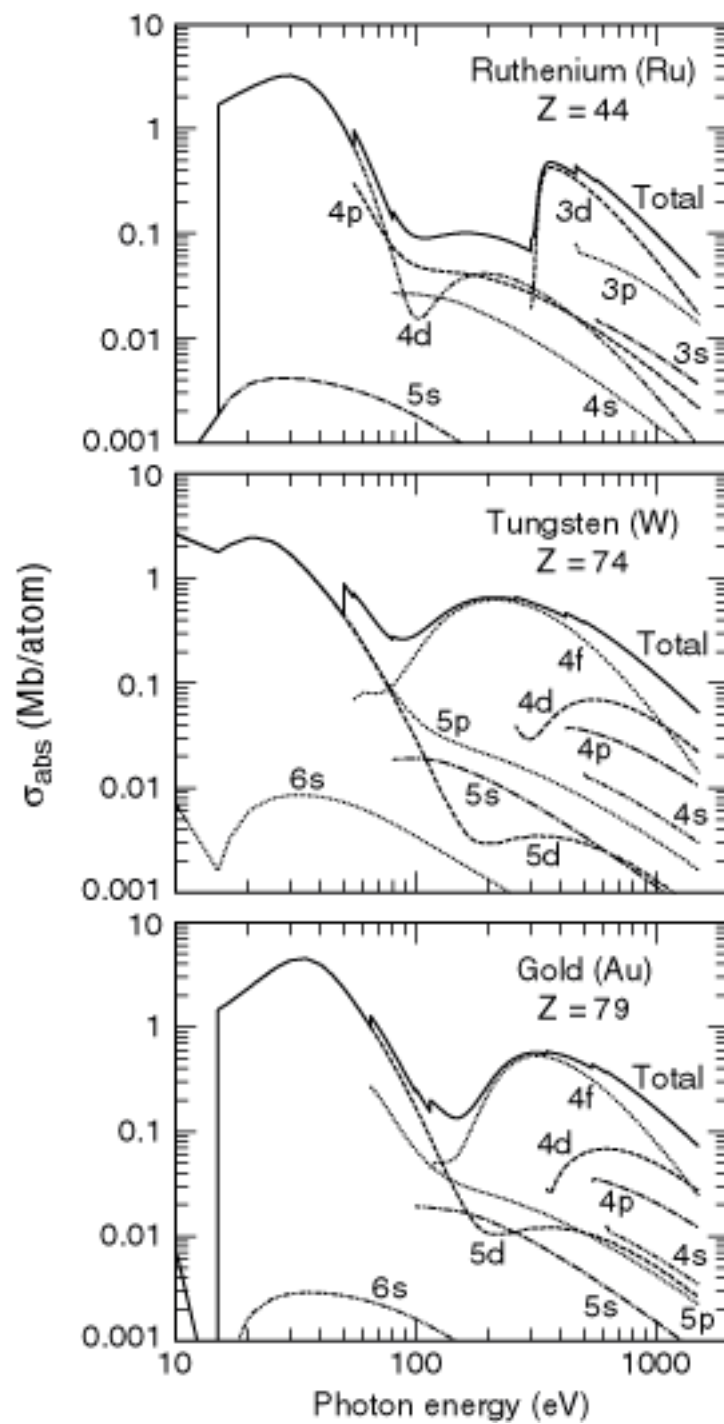


Fig. 1-4. Subshell photoemission cross sections for Ru, W and Au.

X-Ray Data Booklet

Section 1.5 SUBSHELL PHOTOIONIZATION CROSS SECTIONS

Ingolf Lindau

The atomic subshell photoemission cross sections plotted in Fig. 1-4 have been calculated for isolated atoms by Yeh and Lindau [1,2]. The calculations were done with a one-electron central-field frozen-core model using first-order perturbation theory. No single model accurately predicts the photoionization process of all orbitals for all elements from the VUV to 1.5 keV. The complexity of the physics of different atomic orbitals makes it impossible for any single rule to describe all of them. The accuracy of the model used has been discussed in detail by Cooper and Manson [3–5].

REFERENCES

1. J.-J. Yeh and I. Lindau, “Atomic Subshell Photoionization Cross Sections and Asymmetry Parameters: $1 < Z < 103$,” *At. Data Nucl. Data Tables* **32**, 1 (1985).
2. J.-J. Yeh, *Atomic Calculations of Photoionization Cross Sections and Asymmetry Parameters* (Gordon and Breach, Langhorne, PA, 1993).
3. J. W. Cooper, *Phys. Rev.* **128**, 681 (1962).
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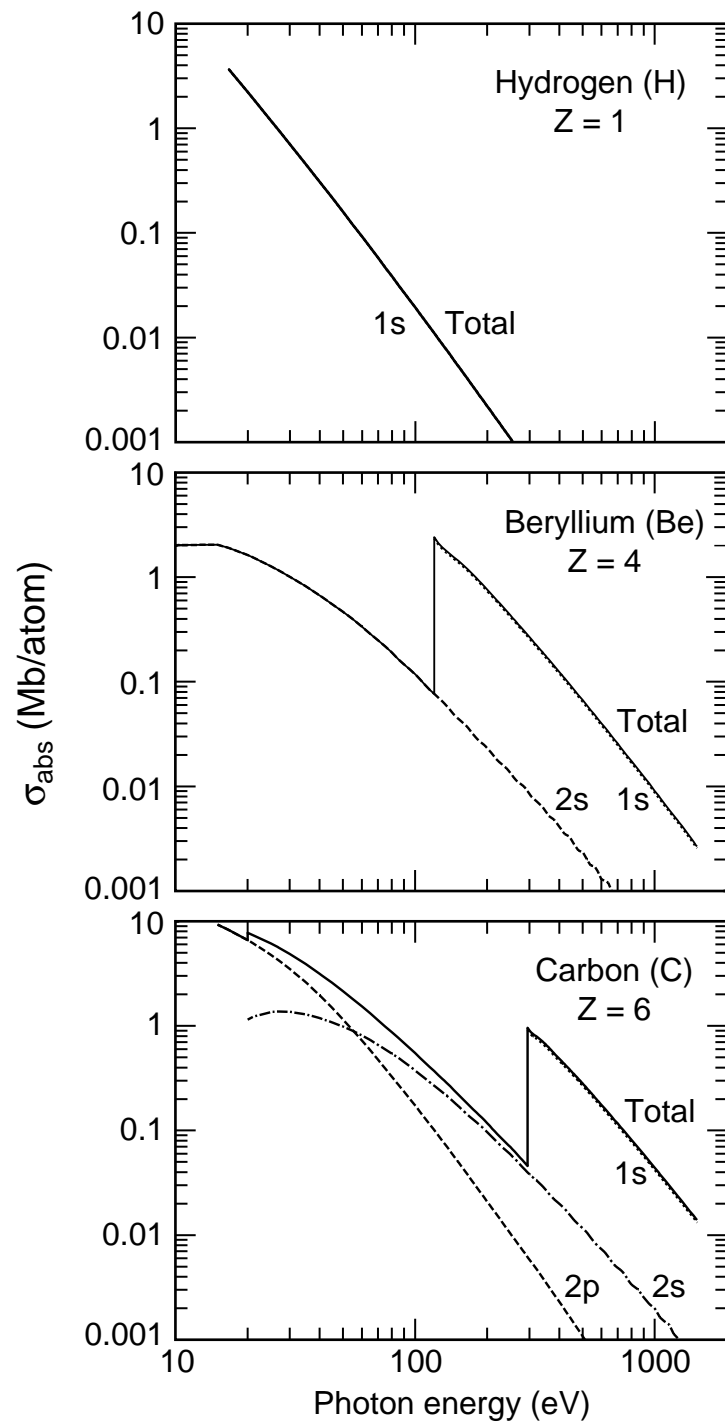


Fig. 1-4. Plots of atomic subshell photoemission cross sections, calculated for isolated atoms.

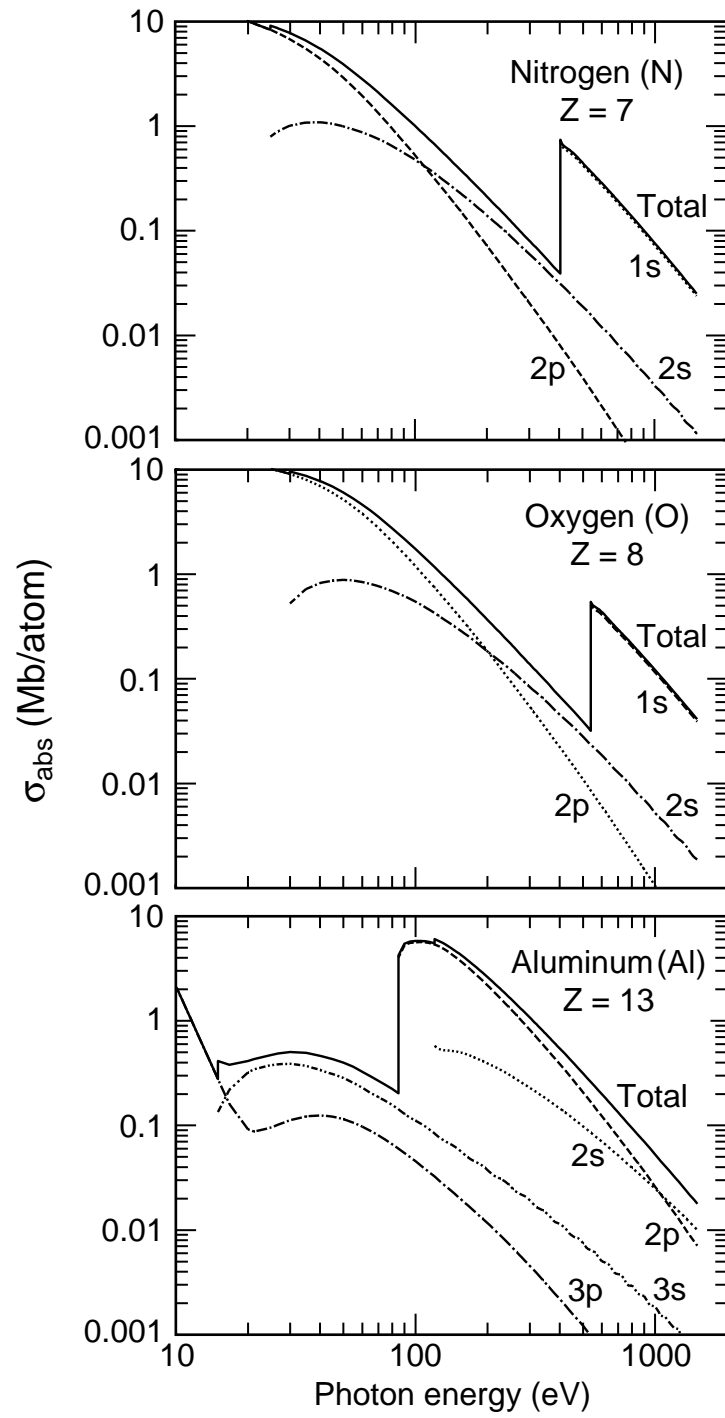


Fig. 1-4. Subshell photoemission cross sections (continued).

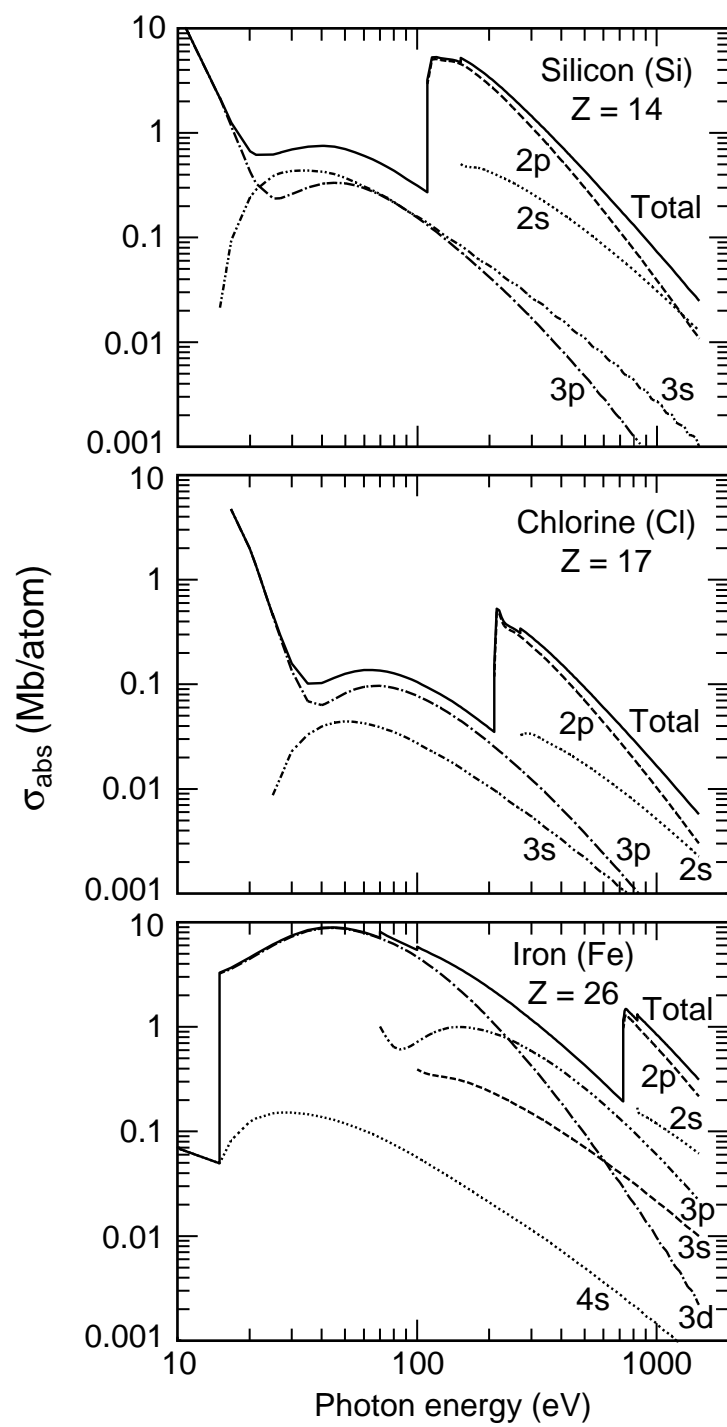


Fig. 1-4. Subshell photoemission cross sections (continued).

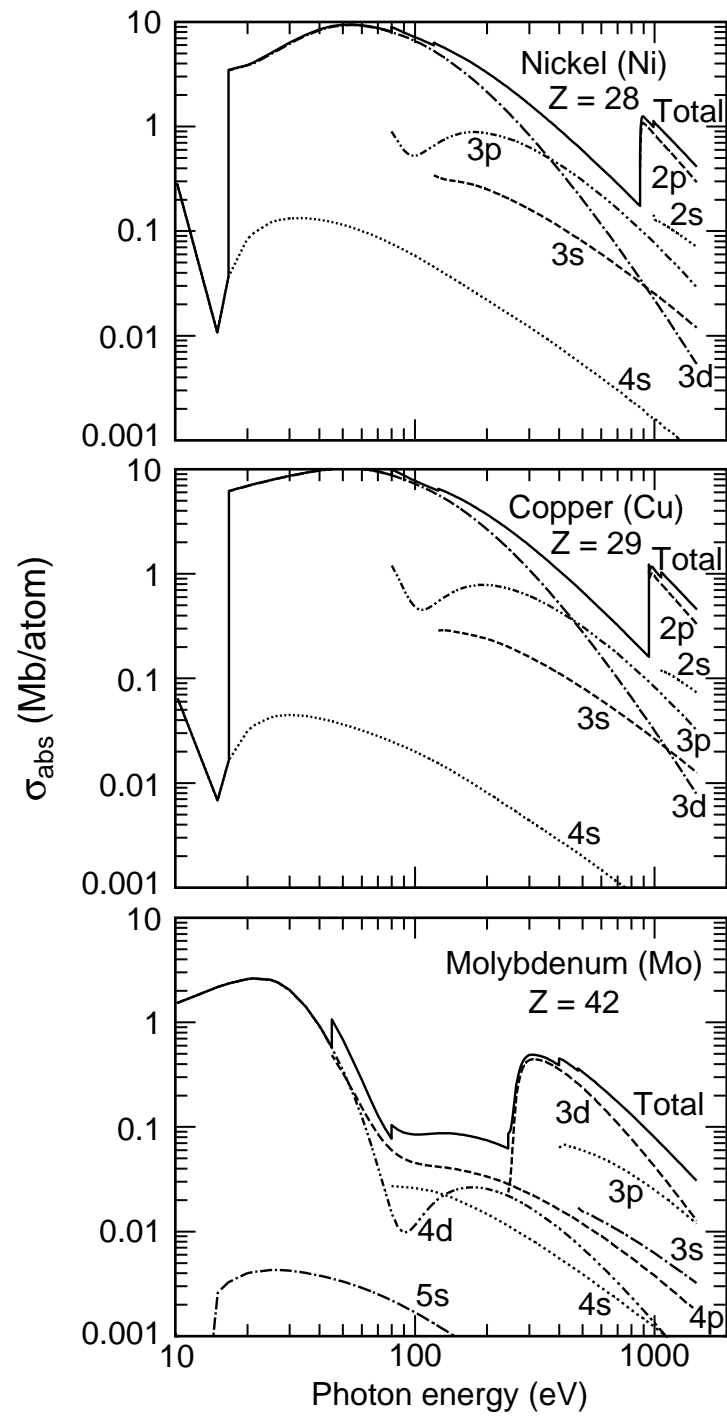


Fig. 1-4. Subshell photoemission cross sections (continued).

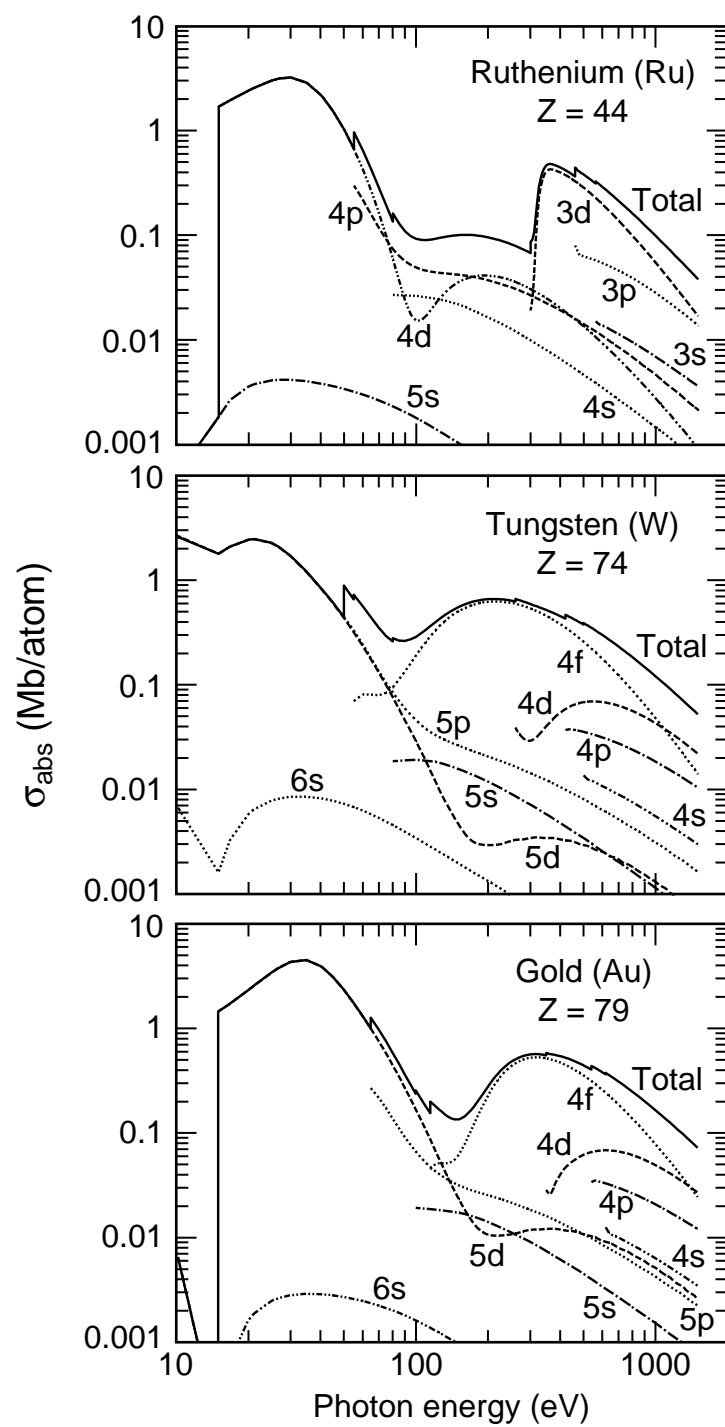


Fig. 1-4. Subshell photoemission cross sections (continued).

X-Ray Data Booklet

Section 1.6 MASS ABSORPTION COEFFICIENTS

Eric M. Gullikson

Mass absorption coefficients have been tabulated for elements $Z = 1-92$, based on both measured values and theoretical calculations [see B. L. Henke, E. M. Gullikson, and J. C. Davis, "X-Ray Interactions: Photoabsorption, Scattering, Transmission, and Reflection at $E = 50-30,000$ eV, $Z = 1-92$," *At. Data Nucl. Data Tables* **54**, 181 (1993); for updated values, see http://www-cxro.lbl.gov/optical_constants/]. The mass absorption coefficient μ (cm^2/g) is related to the transmitted intensity through a material of density ρ (g/cm^3) and thickness d by

$$I = I_0 e^{-\mu d} \quad (1)$$

Thus, the linear absorption coefficient is $\mu_{\text{lin}} (\text{cm}^{-1}) = \mu \rho$. For a pure material, the mass absorption coefficient is directly related to the total atomic absorption cross section σ_a (cm^2/atom) by

$$\mu = \frac{N_A}{A} \sigma_a \quad (2)$$

where N_A is Avogadro's number and A is the atomic weight. For a compound material, the mass absorption coefficient is obtained from the sum of the absorption cross sections of the constituent atoms by

$$\mu = \frac{N_A}{MW} \sum_i x_i \sigma_{ai} \quad (3)$$

where the molecular weight of a compound containing x_i atoms of type i is $MW = \sum_i x_i A_i$. This approximation, which neglects interactions among the atoms in the material, is generally applicable for photon energies above about 30 eV and sufficiently far from absorption edges.

In Fig. 1-5, the mass absorption coefficient is plotted for 15 elements over the photon energy range 10–30,000 eV. In much of this range, the absorption coefficient is dominated by photoabsorption. However, for H, Be, C, N, and O, Compton (in-elastic) scattering is significant at the higher energies. In these cases, the total cross section is shown as a solid curve and the photoabsorption cross section as a separate dashed curve.

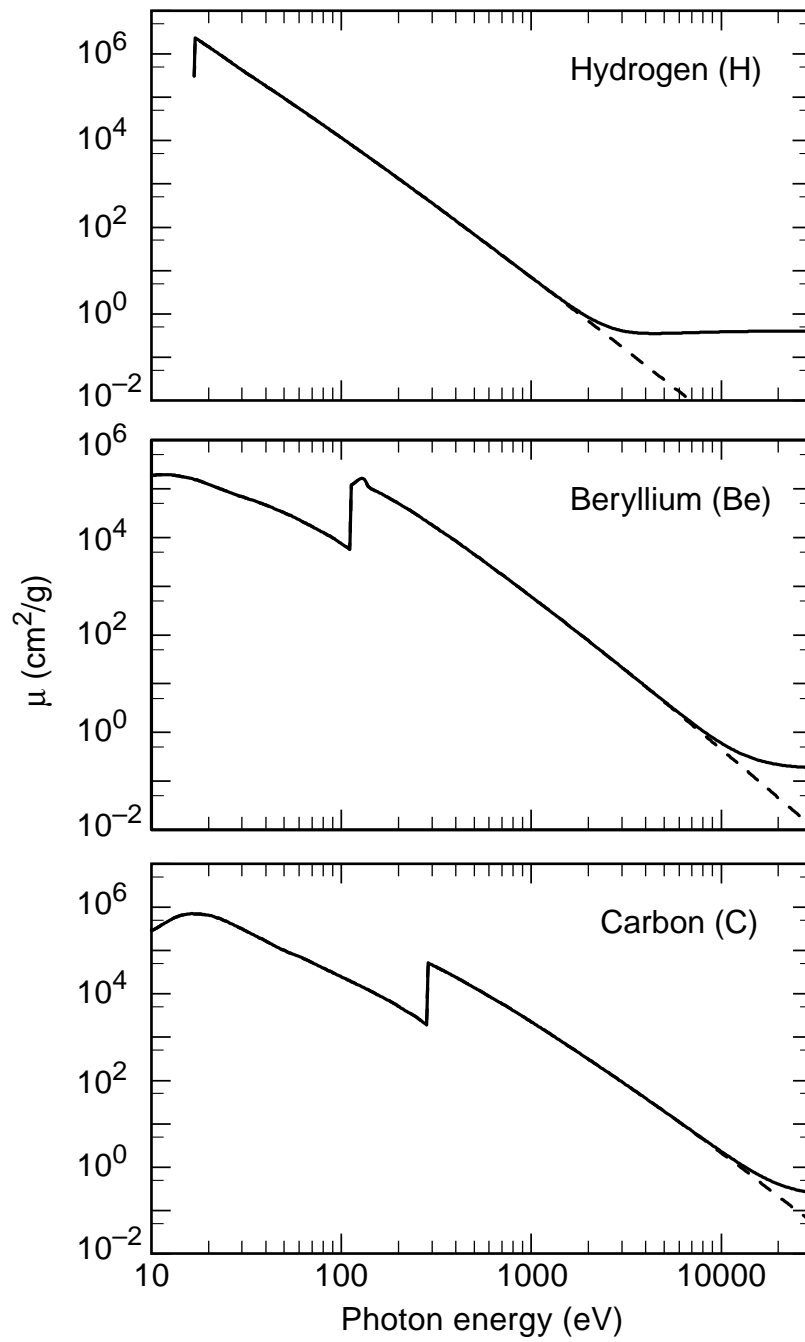


Fig. 1-5. Plots of mass absorption coefficients for several elements in their natural forms. For H, Be, C, N, and O, the photoabsorption cross section is shown as a dashed curve.

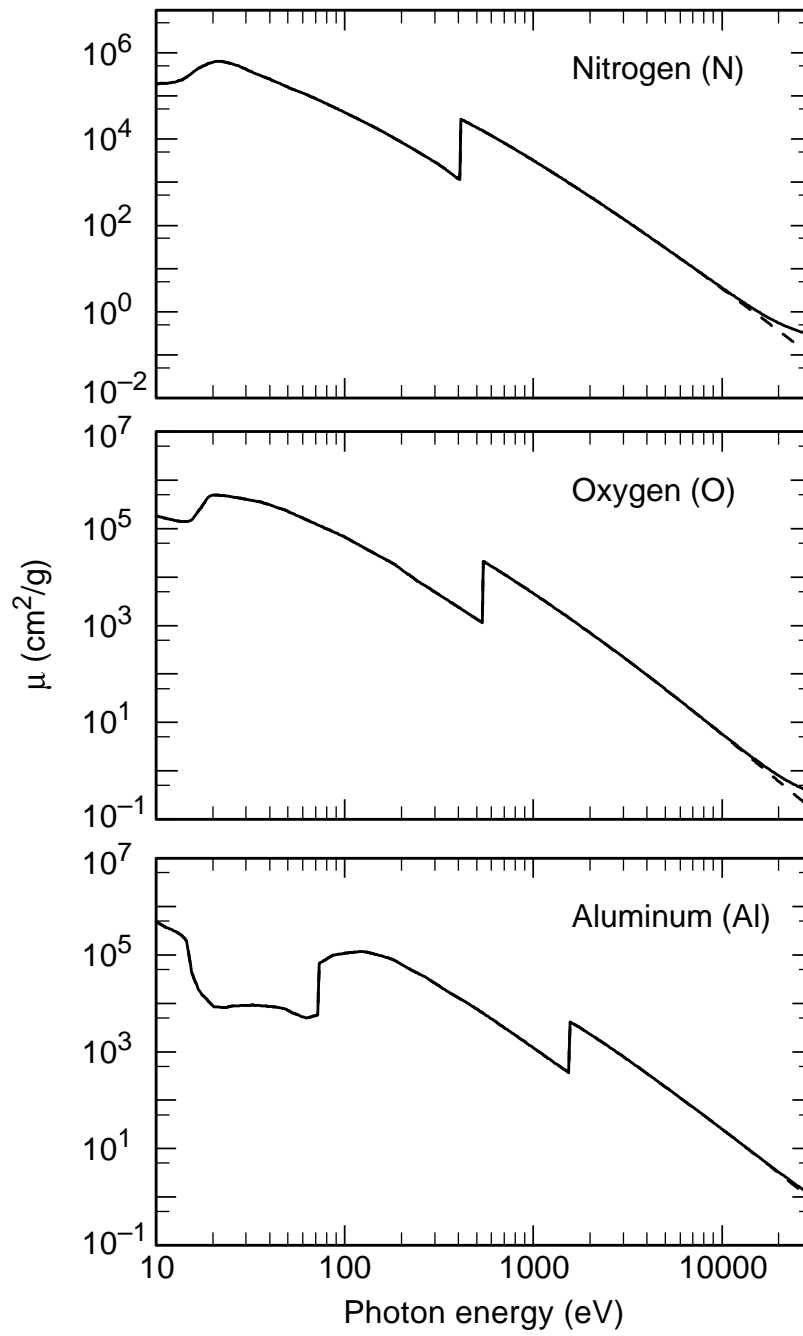


Fig. 1-5. Nitrogen, oxygen and aluminium mass absorption coefficients

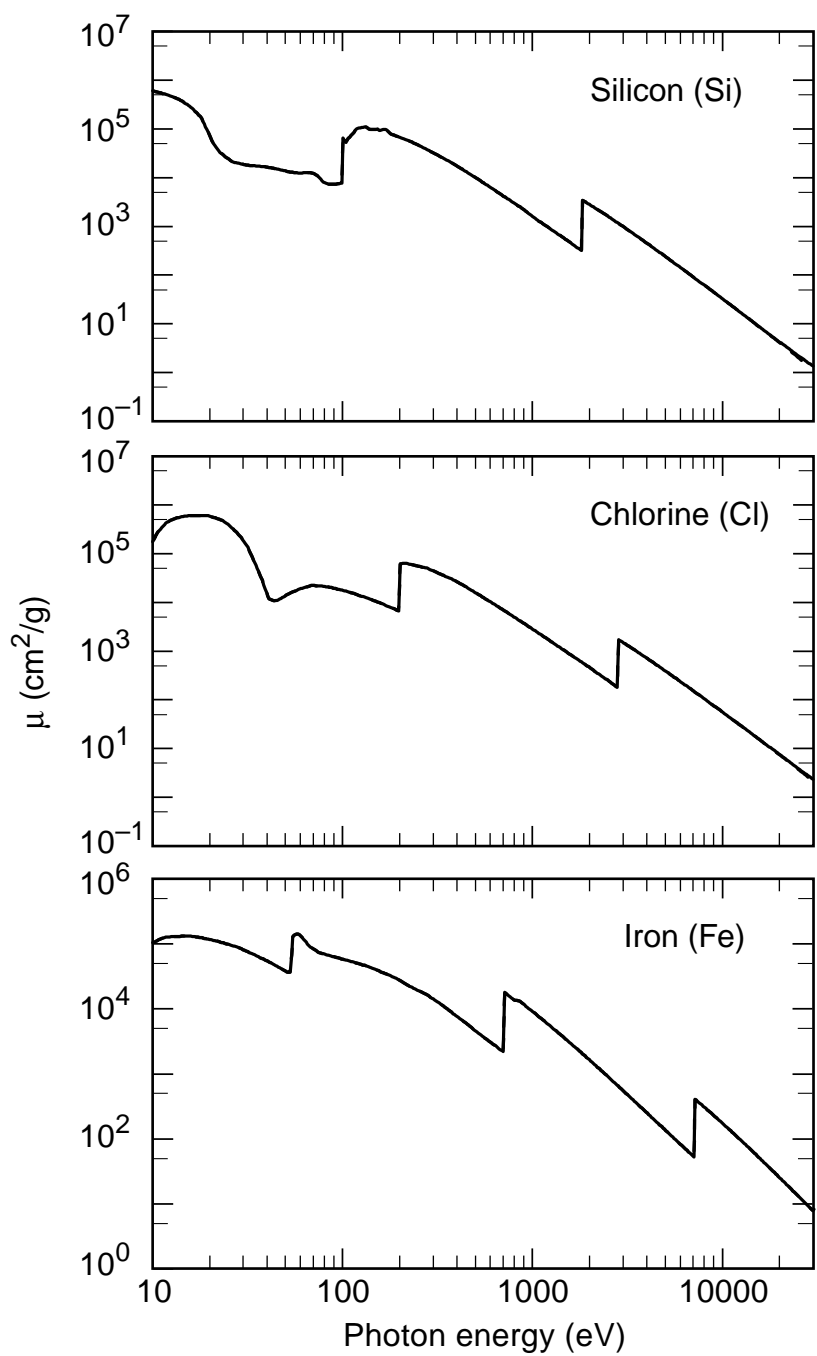


Fig. 1-5. Silicon, chlorine and iron mass absorption coefficients.

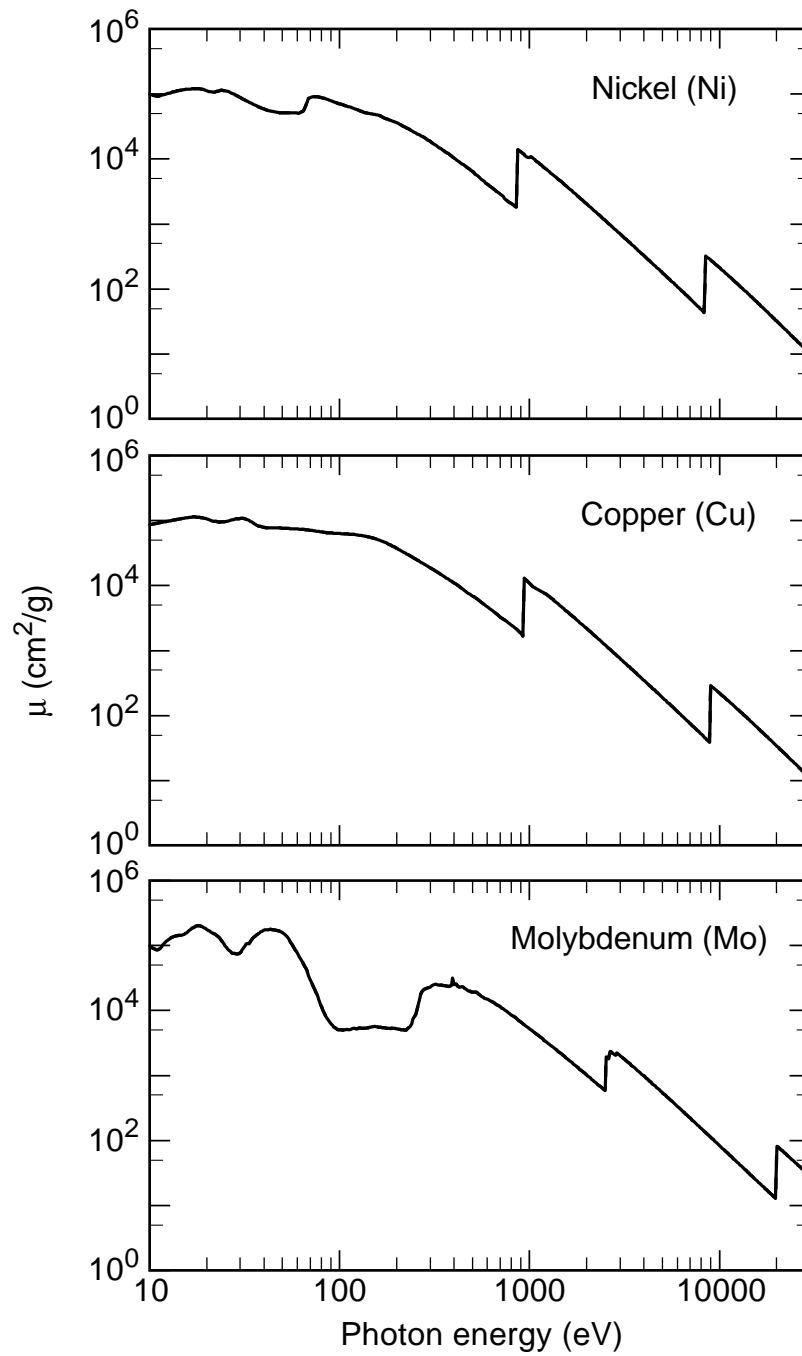


Fig. 1-5. Nickel, copper and molybdenum mass absorption coefficients.

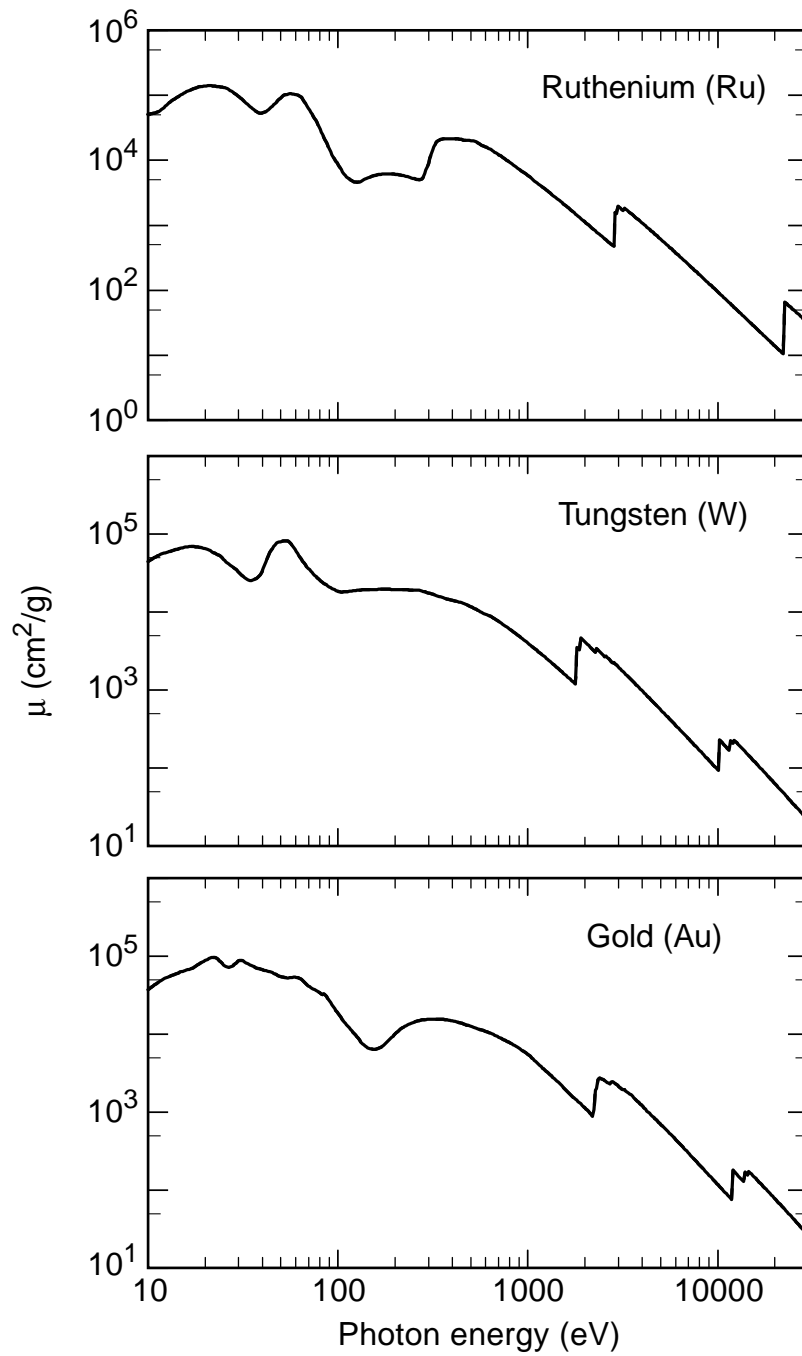


Fig. 1-5. Ruthenium, tungsten and gold mass absorption coefficients.

X-Ray Data Booklet

Section 1.7 ATOMIC SCATTERING FACTORS

Eric M. Gullikson

The optical properties of materials in the photon energy range above about 30 eV can be described by the atomic scattering factors. The index of refraction of a material is related to the scattering factors of the individual atoms by

$$n = 1 - \delta - i\beta = 1 - \frac{r_e}{2\pi} \lambda^2 \sum_i n_i f_i(0) \quad , \quad (1)$$

where r_e is the classical electron radius, λ is the wavelength, and n_i is the number of atoms of type i per unit volume. The parameters δ and β are called the refractive index decrement and the absorption index, respectively. The complex atomic scattering factor for the forward scattering direction is

$$f(0) = f_1 + if_2 \quad . \quad (2)$$

The imaginary part is derived from the atomic photoabsorption cross section:

$$f_2 = \frac{\sigma_a}{2r_e\lambda} \quad . \quad (3)$$

The real part of the atomic scattering factor is related to the imaginary part by the Kramers-Kronig dispersion relation:

$$f_1 = Z^* + \frac{1}{\pi r_e h c} \int_0^\infty \frac{\epsilon^2 \sigma_a(\epsilon)}{E^2 - \epsilon^2} d\epsilon \quad . \quad (4)$$

In the high-photon-energy limit, f_1 approaches Z^* , which differs from the atomic number Z by a small relativistic correction:

$$Z^* \approx Z - (Z/82.5)^{2.37} \quad . \quad (5)$$

On the following pages, Fig. 1-6 presents the scattering factors for 15 elements in their natural forms. Complete tables are given in B. L. Henke, E. M. Gullikson, and J. C. Davis, "X-Ray Interactions: Photoabsorption, Scattering, Transmission, and Reflection at $E = 50\text{--}30,000$ eV, $Z = 1\text{--}92$," *At. Data Nucl. Data Tables* **54**, 181 (1993).

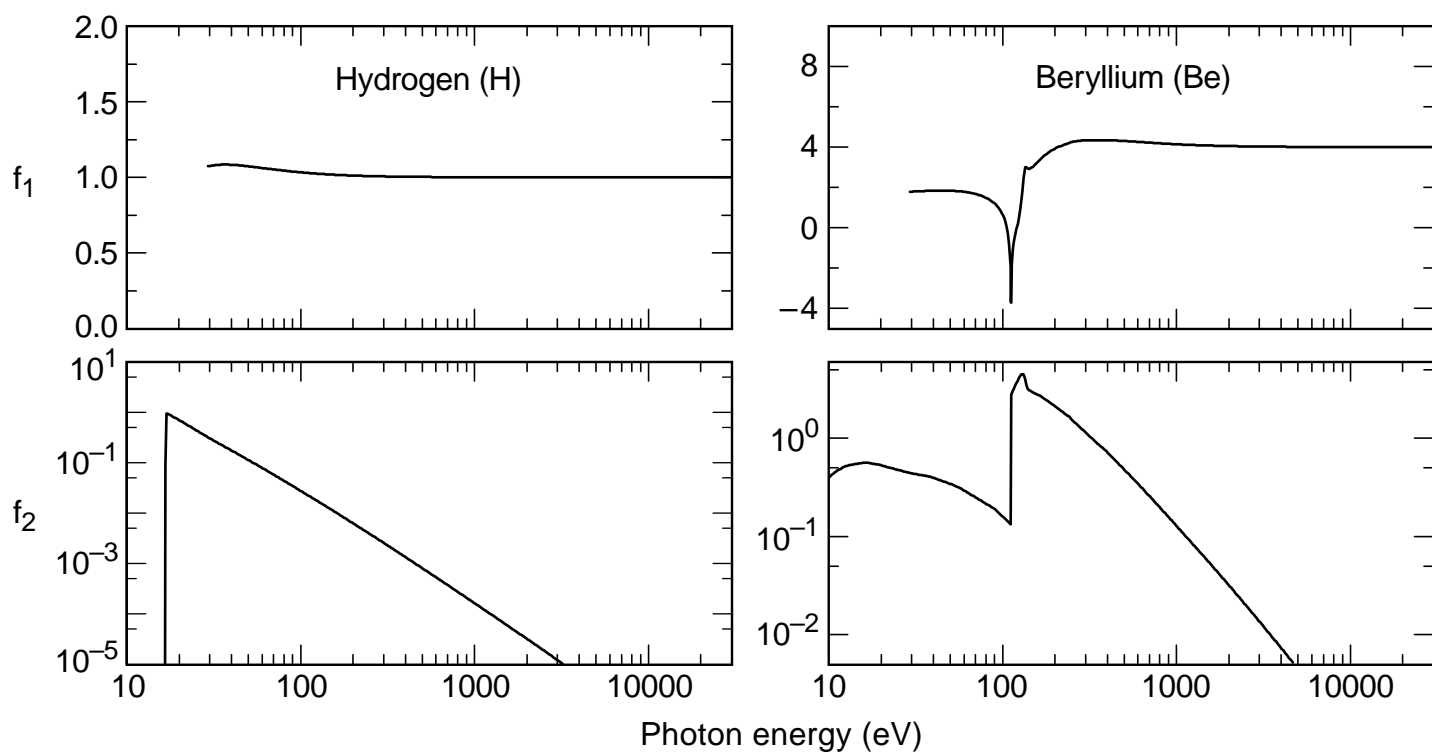


Fig. 1-6. Plots of scattering factors for several elements in their natural forms.

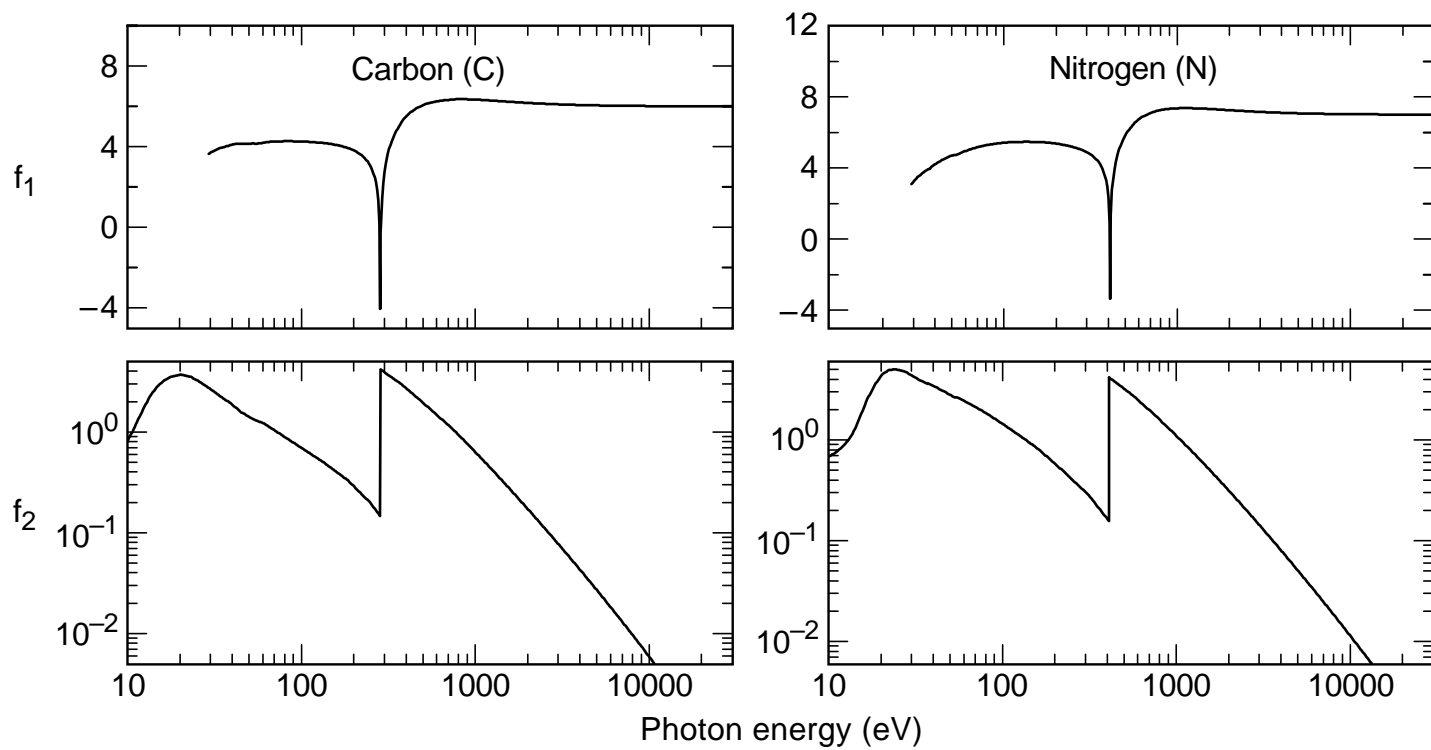


Fig. 1-6. Carbon and Nitrogen scattering factors.

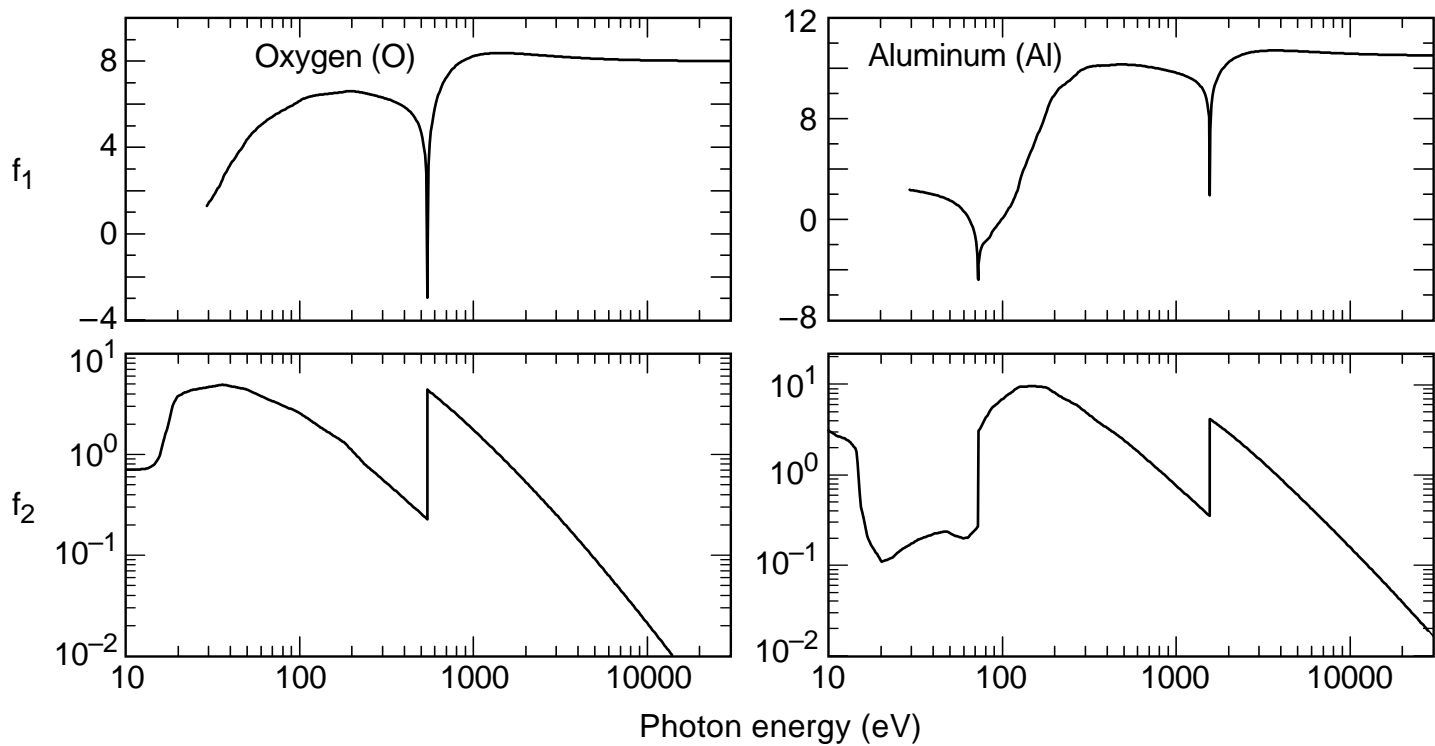


Fig. 1-6. Oxygen and Aluminium scattering factors.

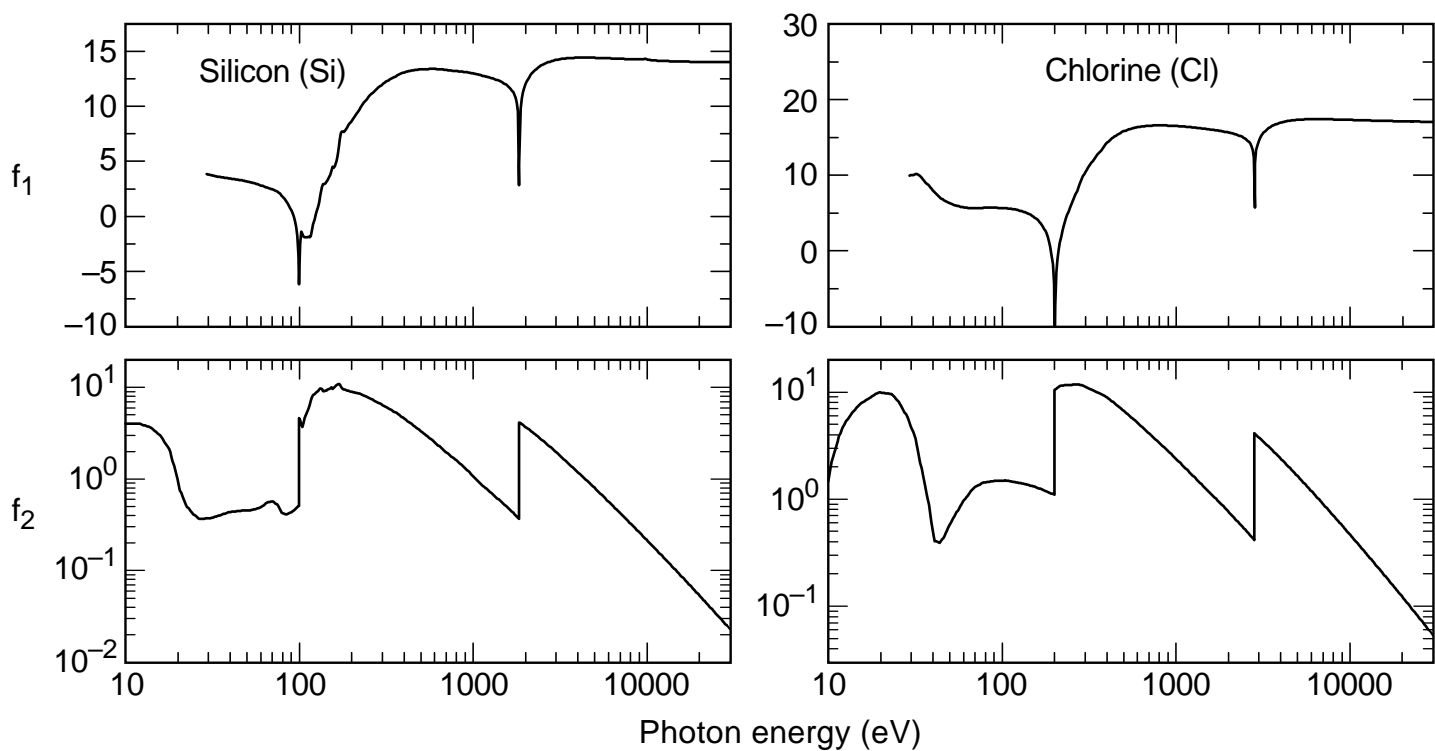


Fig. 1-6. Silicon and Chlorine scattering factors.

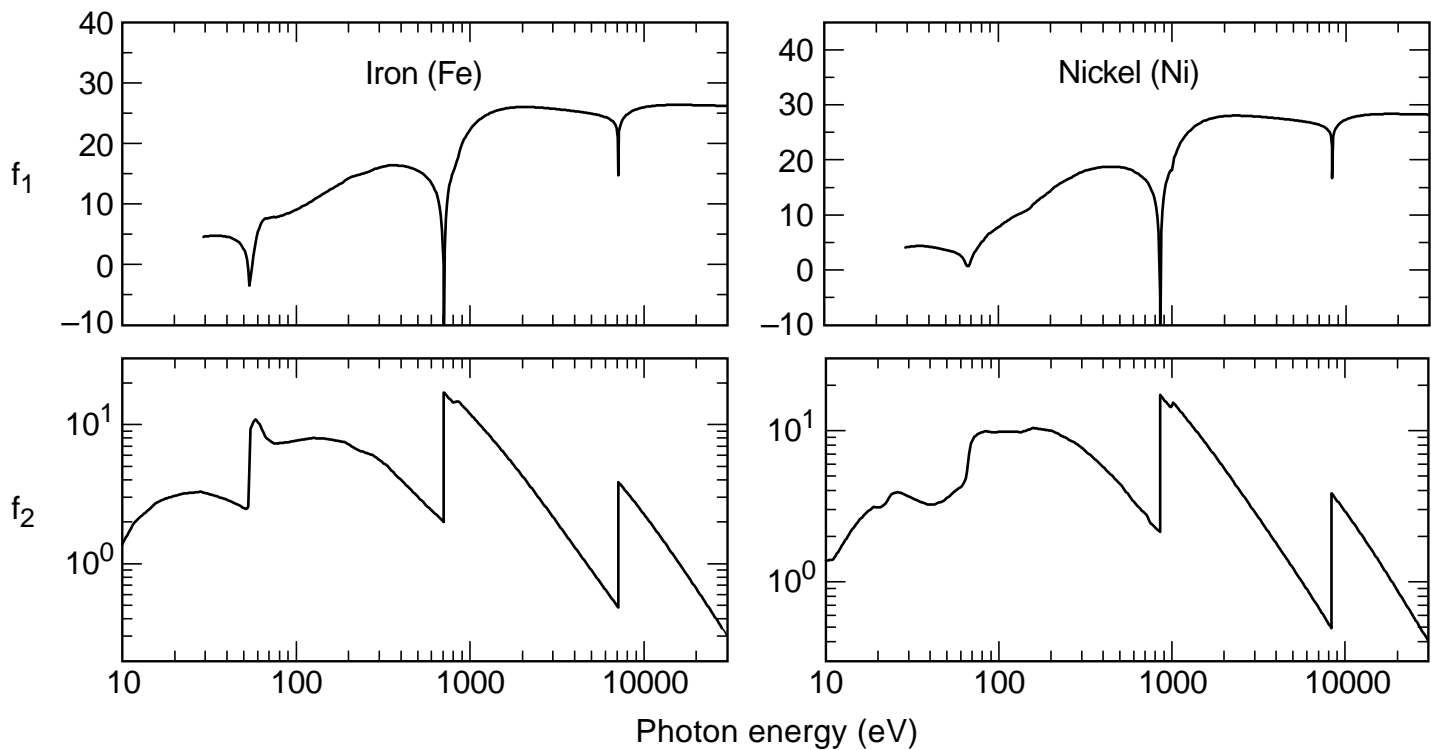


Fig. 1-6. Fe and Ni Scattering factors.

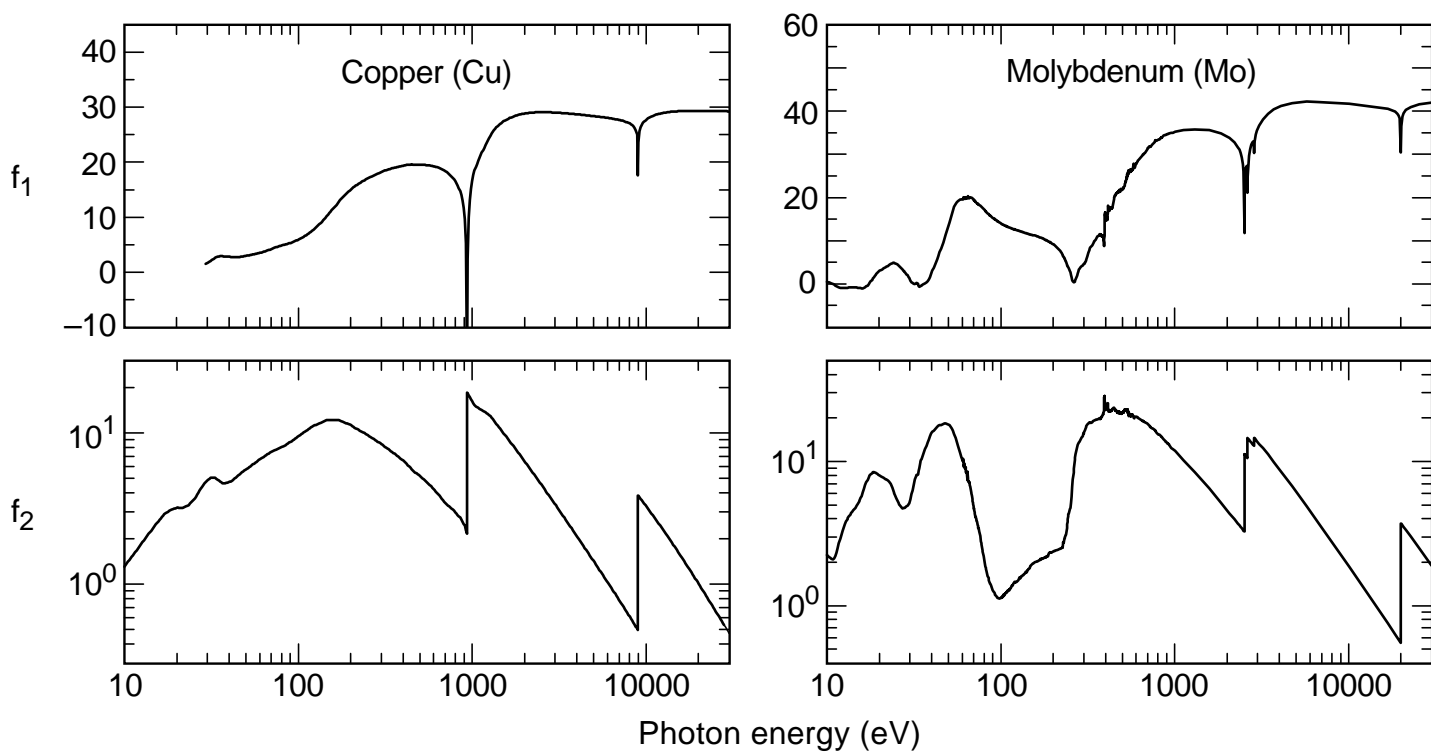


Fig. 1-6. Copper and Molybdenum scattering factors.

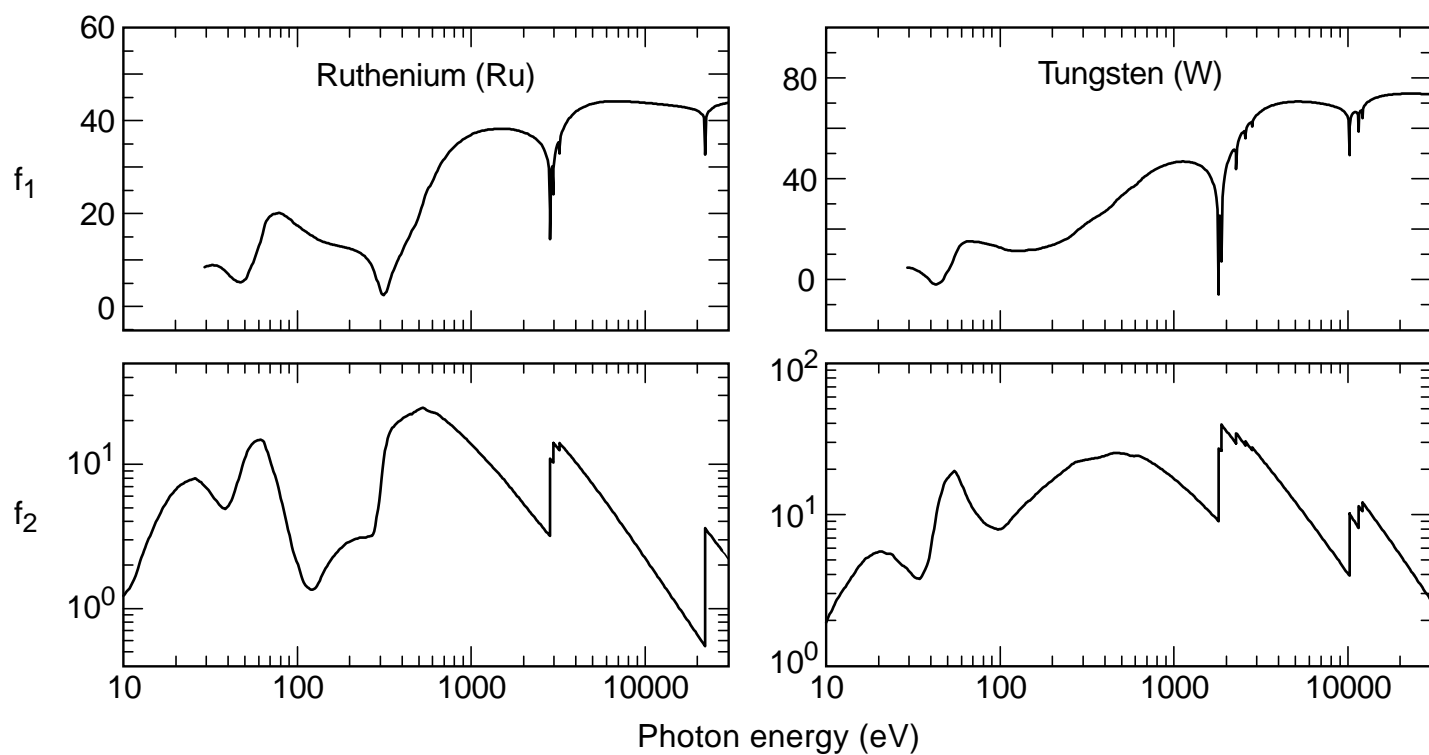


Fig. 1-6. Ruthenium and Tungsten scattering factors.

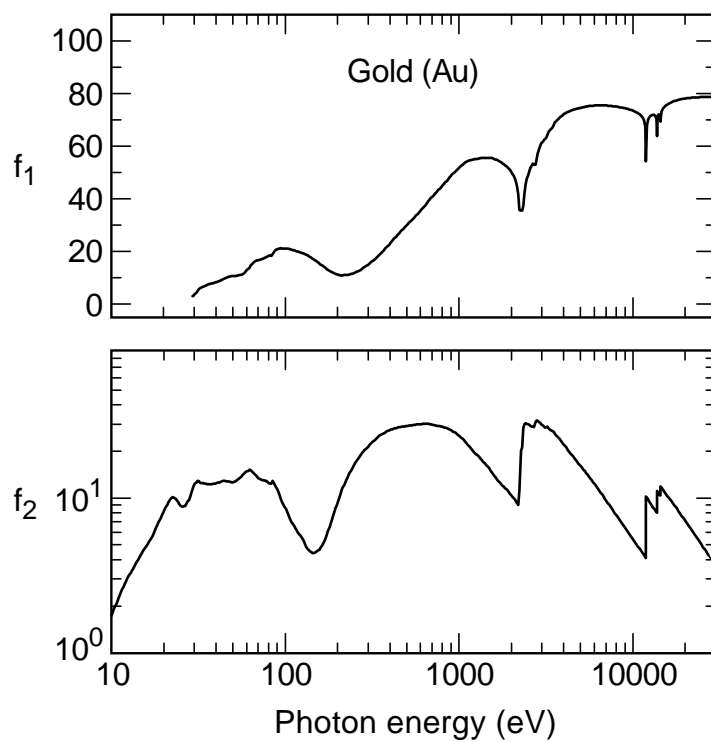


Fig. 1-6. Gold scattering factors.

X-Ray Data Booklet

Section 1.8 ENERGY LEVELS OF FEW-ELECTRON IONIC SPECIES

James H. Scofield

Table 1-4 presents ionization energies for selected few-electron ions with $6 \leq Z \leq 54$. Table 1-5 gives the energies of the resonant $2p$ transitions in hydrogen- and heliumlike ions. The energy values in this section have been generated using the relativistic Hartree-Fock code of I. P. Grant and collaborators [1] with a correction term of the form $A + B/(Z - Q)$ added to bring about agreement with the experimental values known for low atomic numbers. Nuclear size effects, radiative corrections, and the Breit interaction accounting for retardation and the magnetic electron-electron interaction are included in the calculations. The hydrogenic values are uncorrected as they come from the code, but to the accuracy given here, they agree with more detailed calculations. The values in Table 1-4 for Co-, Ni-, and Cu-like ions are based on data from C. E. Moore [2], J. Sugar and A. Musgrove [3], and others referenced therein. A [PDF version of these tables](#) is also available.

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2. C. E. Moore, *Ionization Potentials and Ionization Limits Derived from the Analysis of Optical Spectra*, NBS Pub. NSRDS-NBS 34 (1970).
3. J. Sugar and A. Musgrove, "Energy Levels of Zinc, Zn I through Zn XXX,"

J. Phys. Chem. Ref. Data **24** , 1803 (1995).


Table 1-4. *Ionization energies, in electron volts, for selected few-electron ionic species. Each column is labeled with the number of electrons in the ion before ionization and with the symbol for the neutral atom with the same number of electrons.*

Element	1 (H)	2 (He)	3 (Li)	4 (Be)	10 (Ne)	11 (Na)	12 (Mg)	27 (Co)	28 (Ni)	29 (Cu)
6 C	490.0	392.1	64.49	47.89						
7 N	667.1	552.1	97.89	77.48						
8 O	871.4	739.3	138.11	113.90						
9 F	1103.1	953.9	185.18	157.15						
10 Ne	1362.2	1195.8	239.09	207.26	21.564					
11 Na	1648.7	1465.1	299.86	264.21	47.286	5.139				
12 Mg	1962.7	1761.8	367.5	328.0	80.143	15.035	7.646			
13 Al	2304.2	2086.0	442.0	398.7	119.99	28.447	18.828			
14 Si	2673.2	2437.7	523.4	476.3	166.42	45.12	33.64			
15 P	3070	2816.9	611.7	560.8	220.31	65.02	51.50			
16 S	3494	3224	707.0	652.1	281.00	88.05	72.59			
17 Cl	3946	3658	809.2	750.5	348.5	114.20	96.84			
18 Ar	4426	4121	918.4	855.8	422.8	143.46	124.24			

19 K	4934	4611	1034.6	968.0	503.9	175.82	154.75				
20 Ca	5470	5129	1157.7	1087.3	591.9	211.28	188.38				
21 Sc	6034	5675	1288.0	1213.6	686.6	249.84	225.13				
22 Ti	6626	6249	1425.3	1346.9	788.2	291.50	264.98				
23 V	7246	6851	1569.7	1487.3	896.6	336.3	307.9				
24 Cr	7895	7482	1721.2	1634.8	1011.8	384.2	354.0				
25 Mn	8572	8141	1879.9	1789.5	1133.8	435.2	403.2				
26 Fe	9278	8828	2045.8	1951.3	1262.7	489.3	455.6				
27 Co	10012	9544	2218.9	2120.4	1398.3	546.6	511.0	7.86			
28 Ni	10775	10289	2399.3	2296.7	1540.8	607.0	569.7	18.17	7.63		
29 Cu	11568	11063	2587.0	2480.2	1690.2	670.6	631.4	36.83	20.29	7.73	
30 Zn	12389	11865	2782.0	2671.1	1846.4	737.3	696.4	59.57	39.72	17.96	
31 Ga	13239	12696	2984.4	2869.4	2009.4	807.3	764.5	86.0	63.4	30.7	
32 Ge	14119	13557	3194	3075	2179.3	880.4	835.8	115.9	90.5	45.72	
33 As	15029	14448	3412	3288	2356.0	956.8	910.3	149.2	121.2	62.3	
34 Se	15968	15367	3637	3509	2539.6	1036.3	988.1	185.5	155.4	81.7	
35 Br	16937	16317	3869	3737	2730.1	1119.1	1069.1	225.4	192.8	103.0	
36 Kr	17936	17296	4109	3973	2927.4	1205.2	1153.3	268.2	233.4	125.9	
37 Rb	18965	18306	4357	4216	3132	1294.5	1240.8	314.2	277.1	150.7	

38 Sr	20025	19345	4612	4467	3343	1387.2	1331.5	363.3	324.1	177.3
39 Y	21115	20415	4876	4726	3561	1483.1	1425.6	413.6	374.0	205.9
40 Zr	22237	21516	5147	4993	3786	1582.4	1523.0	471	427.4	236.2
41 Nb	23389	22648	5426	5268	4017	1684.9	1623.7	530	483.8	268.5
42 Mo	24572	23810	5713	5550	4256	1790.9	1727.8	592	541.7	302.6
43 Tc	25787	25004	6008	5841	4502	1900.3	1835.2	656	605.8	338.5
44 Ru	27033	26230	6312	6140	4754	2013.0	1946.1	724	671.4	376.3
45 Rh	28312	27487	6623	6447	5014	2129.2	2060.3	795	740.1	416.0
46 Pd	29623	28776	6943	6762	5280	2248.9	2178.0	869	811.8	457.5
47 Ag	30966	30097	7271	7086	5553	2372.0	2299.2	946	886.6	500.9
48 Cd	32341	31451	7608	7418	5834	2498.6	2423.9	1026	964.5	546.2
49 In	33750	32837	7953	7758	6121	2628.8	2552.1	1109	1045.4	593.3
50 Sn	35192	34257	8307	8107	6415	2762.5	2683.9	1196	1129.1	642.3
51 Sb	36668	35710	8670	8465	6717	2899.8	2819.2	1285	1215.3	693.2
52 Te	38177	37196	9041	8832	7025	3041	2958.1	1377	1306.3	746.1
53 I	39721	38716	9421	9207	7340	3185	3101	1472	1399.3	800.8
54 Xe	41300	40271	9810	9591	7663	3334	3247	1571	1495.4	857.4

Table 1-5. *Transition energies, in electron volts, for transitions from the $n = 2$ states to the $n = 1$ ground state of H- and He-like ions.*

Element	Hydrogenlike		Heliumlike	
	$2p_{1/2}$	$2p_{3/2}$	$2p\ ^3P_1$	$2p\ ^1P_1$
5 B	255.17	255.20	202.78	205.37
6 C	367.5	367.5	304.3	307.8
7 N	500.3	500.4	426.3	430.7
8 O	653.5	653.7	568.7	574.0
9 F	827.3	827.6	731.5	737.8
10 Ne	1021.5	1022.0	914.9	922.1
11 Na	1236.3	1237.0	1118.8	1126.9
12 Mg	1471.7	1472.7	1343.2	1352.3
13 Al	1727.7	1729.0	1588.3	1598.4 
14 Si	2004.3	2006.1	1853.9	1865.1
15 P	2301.7	2304.0	2140.3	2152.6
16 S	2619.7	2622.7	2447.3	2460.8
17 Cl	2958.5	2962.4	2775.1	2789.8
18 Ar	3318	3323	3124	3140
19 K	3699	3705	3493	3511
20 Ca	4100	4108	3883	3903
21 Sc	4523	4532	4295	4316

22 Ti	4966	4977	4727	4750
23 V	5431	5444	5180	5205
24 Cr	5917	5932	5655	5682
25 Mn	6424	6442	6151	6181
26 Fe	6952	6973	6668	6701
27 Co	7502	7526	7206	7242
28 Ni	8073	8102	7766	7806
29 Cu	8666	8699	8347	8392
30 Zn	9281	9318	8950	8999
31 Ga	9917	9960	9575	9628
32 Ge	10575	10624	10221	10280
33 As	11255	11311	10889	10955
34 Se	11958	12021	11579	11652
35 Br	12682	12753	12292	12372
36 Kr	13429	13509	13026	13114
37 Rb	14199	14288	13783	13880
38 Sr	14990	15090	14562	14669
39 Y	15805	15916	15364	15482
40 Zr	16643	16765	16189	16318

41 Nb	17503	17639	17036	17178
42 Mo	18387	18537	17907	18062
43 Tc	19294	19459	18800	18971
44 Ru	20224	20406	19717	19904
45 Rh	21178	21377	20658	20861
46 Pd	22156	22374	21622	21843
47 Ag	23157	23396	22609	22851
48 Cd	24183	24444	23621	23884
49 In	25233	25518	24657	24942
50 Sn	26308	26617	25717	26027

X-Ray Data Booklet

Section 1.8 ENERGY LEVELS OF FEW-ELECTRON IONIC SPECIES

James H. Scofield

Table 1-4 presents ionization energies for selected few-electron ions with $6 \leq Z \leq 54$. Table 1-5 gives the energies of the resonant $2p$ transitions in hydrogen- and heliumlike ions. The energy values in this section have been generated using the relativistic Hartree-Fock code of I. P. Grant and collaborators [1] with a correction term of the form $A + B/(Z - Q)$ added to bring about agreement with the experimental values known for low atomic numbers. Nuclear size effects, radiative corrections, and the Breit interaction accounting for retardation and the magnetic electron-electron interaction are included in the calculations. The hydrogenic values are uncorrected as they come from the code, but to the accuracy given here, they agree with more detailed calculations. The values in Table 1-4 for Co-, Ni-, and Cu-like ions are based on data from C. E. Moore [2], J. Sugar and A. Musgrove [3], and others referenced therein.

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1. I. P. Grant, B. J. McKenzie, P. H. Norrington, D. F. Mayers, and N. C. Pyper, "An Atomic Multiconfigurational Dirac-Fock Package," *Comput. Phys. Commun.* **21**, 207 (1980).
2. C. E. Moore, *Ionization Potentials and Ionization Limits Derived from the Analysis of Optical Spectra*, NBS Pub. NSRDS-NBS 34 (1970).
3. J. Sugar and A. Musgrove, "Energy Levels of Zinc, Zn I through Zn XXX," *J. Phys. Chem. Ref. Data* **24**, 1803 (1995).

X-Ray Data Booklet Table 1-4. *Ionization energies, in electron volts, for selected few-electron ionic species.*
Each column is labeled with the number of electrons in the ion before ionization and with the symbol for the neutral atom with the same number of electrons.

Element	1 (H)	2 (He)	3 (Li)	4 (Be)	10 (Ne)	11 (Na)	12 (Mg)	27 (Co)	28 (Ni)	29 (Cu)
6 C	490.0	392.1	64.49	47.89						
7 N	667.1	552.1	97.89	77.48						
8 O	871.4	739.3	138.11	113.90						
9 F	1103.1	953.9	185.18	157.15						
10 Ne	1362.2	1195.8	239.09	207.26	21.564					
11 Na	1648.7	1465.1	299.86	264.21	47.286	5.139				
12 Mg	1962.7	1761.8	367.5	328.0	80.143	15.035	7.646			
13 Al	2304.2	2086.0	442.0	398.7	119.99	28.447	18.828			
14 Si	2673.2	2437.7	523.4	476.3	166.42	45.12	33.64			
15 P	3070	2816.9	611.7	560.8	220.31	65.02	51.50			
16 S	3494	3224	707.0	652.1	281.00	88.05	72.59			
17 Cl	3946	3658	809.2	750.5	348.5	114.20	96.84			
18 Ar	4426	4121	918.4	855.8	422.8	143.46	124.24			
19 K	4934	4611	1034.6	968.0	503.9	175.82	154.75			
20 Ca	5470	5129	1157.7	1087.3	591.9	211.28	188.38			
21 Sc	6034	5675	1288.0	1213.6	686.6	249.84	225.13			
22 Ti	6626	6249	1425.3	1346.9	788.2	291.50	264.98			
23 V	7246	6851	1569.7	1487.3	896.6	336.3	307.9			
24 Cr	7895	7482	1721.2	1634.8	1011.8	384.2	354.0			
25 Mn	8572	8141	1879.9	1789.5	1133.8	435.2	403.2			
26 Fe	9278	8828	2045.8	1951.3	1262.7	489.3	455.6			
27 Co	10012	9544	2218.9	2120.4	1398.3	546.6	511.0	7.86		
28 Ni	10775	10289	2399.3	2296.7	1540.8	607.0	569.7	18.17	7.63	
29 Cu	11568	11063	2587.0	2480.2	1690.2	670.6	631.4	36.83	20.29	7.73
30 Zn	12389	11865	2782.0	2671.1	1846.4	737.3	696.4	59.57	39.72	17.96
31 Ga	13239	12696	2984.4	2869.4	2009.4	807.3	764.5	86.0	63.4	30.7
32 Ge	14119	13557	3194	3075	2179.3	880.4	835.8	115.9	90.5	45.72
33 As	15029	14448	3412	3288	2356.0	956.8	910.3	149.2	121.2	62.3
34 Se	15968	15367	3637	3509	2539.6	1036.3	988.1	185.5	155.4	81.7
35 Br	16937	16317	3869	3737	2730.1	1119.1	1069.1	225.4	192.8	103.0

Table 1-4. *Ionization energies, in electron volts, for selected few-electron ionic species. Each column is labeled with the number of electrons in the ion before ionization and with the symbol for the neutral atom with the same number of electrons.*

Element	1 (H)	2 (He)	3 (Li)	4 (Be)	10 (Ne)	11 (Na)	12 (Mg)	27 (Co)	28 (Ni)	29 (Cu)
36 Kr	17936	17296	4109	3973	2927.4	1205.2	1153.3	268.2	233.4	125.9
37 Rb	18965	18306	4357	4216	3132	1294.5	1240.8	314.2	277.1	150.7
38 Sr	20025	19345	4612	4467	3343	1387.2	1331.5	363.3	324.1	177.3
39 Y	21115	20415	4876	4726	3561	1483.1	1425.6	413.6	374.0	205.9
40 Zr	22237	21516	5147	4993	3786	1582.4	1523.0	471	427.4	236.2
41 Nb	23389	22648	5426	5268	4017	1684.9	1623.7	530	483.8	268.5
42 Mo	24572	23810	5713	5550	4256	1790.9	1727.8	592	541.7	302.6
43 Tc	25787	25004	6008	5841	4502	1900.3	1835.2	656	605.8	338.5
44 Ru	27033	26230	6312	6140	4754	2013.0	1946.1	724	671.4	376.3
45 Rh	28312	27487	6623	6447	5014	2129.2	2060.3	795	740.1	416.0
46 Pd	29623	28776	6943	6762	5280	2248.9	2178.0	869	811.8	457.5
47 Ag	30966	30097	7271	7086	5553	2372.0	2299.2	946	886.6	500.9
48 Cd	32341	31451	7608	7418	5834	2498.6	2423.9	1026	964.5	546.2
49 In	33750	32837	7953	7758	6121	2628.8	2552.1	1109	1045.4	593.3
50 Sn	35192	34257	8307	8107	6415	2762.5	2683.9	1196	1129.1	642.3
51 Sb	36668	35710	8670	8465	6717	2899.8	2819.2	1285	1215.3	693.2
52 Te	38177	37196	9041	8832	7025	3041	2958.1	1377	1306.3	746.1
53 I	39721	38716	9421	9207	7340	3185	3101	1472	1399.3	800.8
54 Xe	41300	40271	9810	9591	7663	3334	3247	1571	1495.4	857.4

X-Ray Data Booklet Table 1-5. Transition energies, in electron volts, for transitions from the $n = 2$ states to the $n = 1$ ground state of H- and He- like ions.

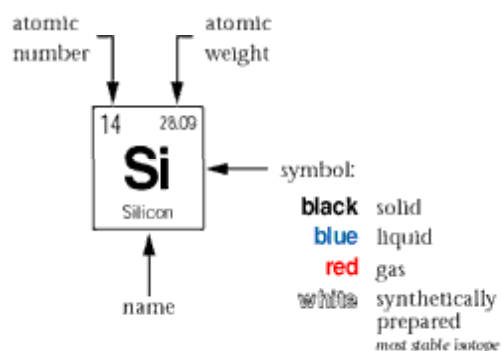
Element	Hydrogenlike		Heliumlike	
	$2p_{1/2}$	$2p_{3/2}$	$2p\ ^3P_1$	$2p\ ^1P_1$
5 B	255.17	255.20	202.78	205.37
6 C	367.5	367.5	304.3	307.8
7 N	500.3	500.4	426.3	430.7
8 O	653.5	653.7	568.7	574.0
9 F	827.3	827.6	731.5	737.8
10 Ne	1021.5	1022.0	914.9	922.1
11 Na	1236.3	1237.0	1118.8	1126.9
12 Mg	1471.7	1472.7	1343.2	1352.3
13 Al	1727.7	1729.0	1588.3	1598.4
14 Si	2004.3	2006.1	1853.9	1865.1
15 P	2301.7	2304.0	2140.3	2152.6
16 S	2619.7	2622.7	2447.3	2460.8
17 Cl	2958.5	2962.4	2775.1	2789.8
18 Ar	3318	3323	3124	3140
19 K	3699	3705	3493	3511
20 Ca	4100	4108	3883	3903
21 Sc	4523	4532	4295	4316
22 Ti	4966	4977	4727	4750
23 V	5431	5444	5180	5205
24 Cr	5917	5932	5655	5682
25 Mn	6424	6442	6151	6181
26 Fe	6952	6973	6668	6701
27 Co	7502	7526	7206	7242
28 Ni	8073	8102	7766	7806
29 Cu	8666	8699	8347	8392
30 Zn	9281	9318	8950	8999
31 Ga	9917	9960	9575	9628
32 Ge	10575	10624	10221	10280
33 As	11255	11311	10889	10955
34 Se	11958	12021	11579	11652
35 Br	12682	12753	12292	12372
36 Kr	13429	13509	13026	13114
37 Rb	14199	14288	13783	13880
38 Sr	14990	15090	14562	14669
39 Y	15805	15916	15364	15482
40 Zr	16643	16765	16189	16318
41 Nb	17503	17639	17036	17178
42 Mo	18387	18537	17907	18062
43 Tc	19294	19459	18800	18971
44 Ru	20224	20406	19717	19904
45 Rh	21178	21377	20658	20861
46 Pd	22156	22374	21622	21843
47 Ag	23157	23396	22609	22851
48 Cd	24183	24444	23621	23884
49 In	25233	25518	24657	24942
50 Sn	26308	26617	25717	26027

X-ray Properties of the Elements

Click on an element to see its properties



1 1.01 H Hydrogen	2 4.003 He Helium
3 6.94 Li Lithium	4 9.01 Be Beryllium
11 22.99 Na Sodium	12 24.31 Mg Magnesium
19 39.10 K Potassium	20 40.08 Ca Calcium
27 58.93 Co Cobalt	28 58.70 Ni Nickel
35 79.90 Br Bromine	36 83.80 Kr Krypton
43 98 Tc Technetium	44 101.07 Ru Ruthenium
51 121.75 Sb Antimony	52 127.60 Te Tellurium
59 138.91 La Lanthanum	60 144.24 Nd Neodymium
67 164.93 Ho Holmium	68 167.26 Er Erbium
75 186.21 Re Rhenium	76 190.23 Os Osmium
83 208.98 Bi Bismuth	84 (209) Po Polonium
91 231.04 Pa Protactinium	92 238.03 U Uranium
99 (252) Es Einsteinium	100 (257) Fm Fermium
107 (262) Bh Bohrium	108 (265) Hs Hassium
115 (288) Mc Moscovium	116 (289) Lv Livermorium
117 (293) Ts Tennessine	118 (293) Og Oganesson



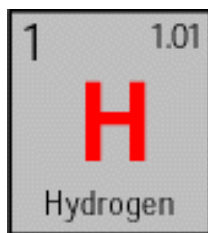
- alkali metals
- alkaline earth metals
- transitional metals
- other metals
- nonmetals
- noble gases

Lanthanide series

Actinide series

58 140.12 Ce Cerium	59 140.91 Pr Praseodymium	60 144.24 Nd Neodymium	61 (145) Pm Promethium	62 150.36 Sm Samarium	63 151.96 Eu Europium	64 157.25 Gd Gadolinium	65 158.93 Tb Terbium	66 162.50 Dy Dysprosium	67 164.93 Ho Holmium	68 167.26 Er Erbium	69 168.93 Tm Thulium	70 173.04 Yb Ytterbium	71 174.97 Lu Lutetium
90 232.04 Th Thorium	91 231.04 Pa Protactinium	92 238.03 U Uranium	93 237.05 Np Neptunium	94 (244) Pu Plutonium	95 (243) Am Americium	96 (247) Cm Curium	97 (247) Bk Berkelium	98 (251) Cf Californium	99 (252) Es Einsteinium	100 (257) Fm Fermium	101 (258) Md Mendelevium	102 (259) No Nobelium	103 (262) Lr Lawrencium

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Electron Binding Energies in electron volts

K 1s

13.6

Electron Configuration = $1s^1$

Oxidation State = 1

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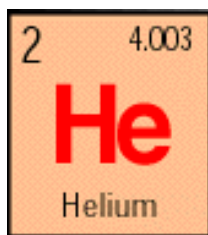
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Electron Binding Energies in electron volts

K 1s

24.6

Electron Configuration = $1s^2$

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Properties
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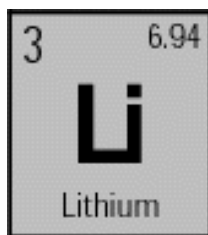
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}
54.7			

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1
54.3		

Electron Configuration = 1s²2s¹

Oxidation State = 1

Density = 0.53 g/cm³ Concentration = 4.63 x 10²² atoms/cm³

**Physical
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}
111.5			

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1
108.5		

Electron Configuration = 1s²2s²

Oxidation State = **2** Nearest Neighbor = 0.223 nm

Density = 1.85 g/cm³ Concentration = 12.3 x 10²² atoms/cm³

**Physical
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}
188.			

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1
183.3		

Electron Configuration = $1s^2 2s^2 2p^1$

Oxidation State = **3** Nearest Neighbor = 0.178 nm

Density = 2.47 g/cm³ Concentration = 13.7×10^{22} atoms/cm³

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6	12.01
C	
Carbon	

Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}
284.2			

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1
277.		

Electron Configuration = 1s²2s²2p²

Oxidation State = $\pm 4, 2$ Nearest Neighbor = 0.142 nm

Density = 2.27 g/cm³ Concentration = 11.4 x 10²² atoms/cm³

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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}
409.9	37.3		

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1
392.4		

Electron Configuration = 1s²2s²2p³

Oxidation State = $\pm 3, 5, 4, 2$

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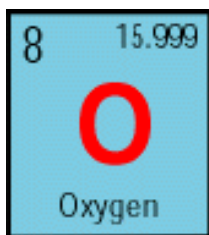
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}
543.1	41.6		

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1
524.9		

Electron Configuration = $1s^2 2s^2 2p^4$

Oxidation State = **-2**

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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}
696.7			

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1
676.8		

Electron Configuration = 1s²2s²2p⁵

Oxidation State = -1

Physical
Properties
from NIST

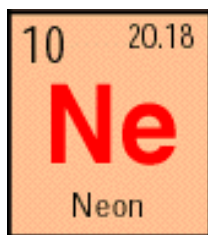
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}
870.2	48.5	21.7	21.6

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1
848.6	848.6	

Electron Configuration = 1s²2s²2p⁶

Physical
Properties
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VISUAL
ELEMENTS
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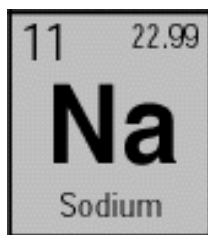
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}
1070.8	63.5	30.65	30.81

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1
1,040.98	1,040.98	1,071.1

Electron Configuration = [Ne]3s¹

Oxidation State = 1

Density = 0.97 g/cm³ Concentration = 2.53 x 10²² atoms/cm³

**Physical
Properties
from NIST**

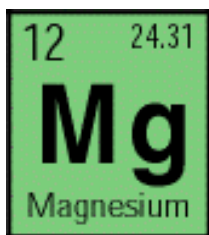
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}
1303.0	88.7	49.78	49.50

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1
1,253.60	1,253.60	1,302.2

Electron Configuration = [Ne]3s²

Oxidation State = **2** Nearest Neighbor = 0.320 nm

Density = 1.74 g/cm³ Concentration = 4.30 x 10²² atoms/cm³

**Physical
Properties
from NIST**

**VISUAL
ELEMENTS
from CHEMSOC**

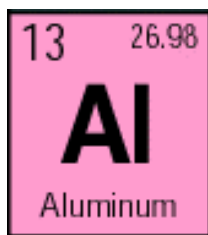
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}
1559.6	117.8	72.95	72.55

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1
1,486.70	1,486.27	1,557.45

Electron Configuration = [Ne]3s²3p¹

Oxidation State = **3** Nearest Neighbor = 0.286 nm

Density = 2.70 g/cm³ Concentration = 6.02 x 10²² atoms/cm³

**Physical
Properties
from NIST**

**VISUAL
ELEMENTS
from CHEMSOC**

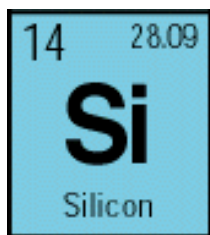
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}
1839.	149.7	99.82	99.42

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1
1,739.98	1,739.38	1,835.94

Electron Configuration = [Ne]3s²3p²

Oxidation State = 4 Nearest Neighbor = 0.235 nm

Density = 2.33 g/cm³ Concentration = 4.99 x 10²² atoms/cm³

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Properties
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ELEMENTS
from CHEMSOC**

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15	30.97
P	
Phosphorus	

Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}
2145.5	189.	136.	135.

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1
2,013.7	2,012.7	2,139.1

Electron Configuration = [Ne]3s²3p³

Oxidation State = $\pm 3, 5, 4$

Density = 1.82 g/cm³ Concentration = 3.54 x 10²² atoms/cm³

Physical
Properties
from NIST

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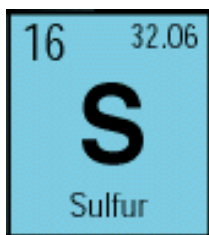
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}
2,472.	230.9	163.6	162.5

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1
2,307.84	2,306.64	2,464.04

Electron Configuration = [Ne]3s²3p⁴

Oxidation State = $\pm 2, 4, 6$

Density = 2.09 g/cm³ Concentration = 3.92 x 10²² atoms/cm³

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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}
2,822.4	270.	202.	200.

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1
2,622.39	2,620.78	2,815.6

Electron Configuration = [Ne]3s²3p⁵

Oxidation State = $\pm 1, 3, 5, 7$

Physical
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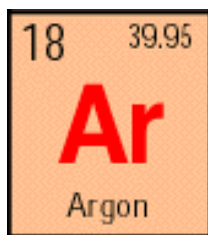
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}
3,205.9	326.3	250.6	248.4	29.3	15.9	15.7

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1
2,957.70	2,955.63	3,190.5

Electron Configuration = [Ne]3s²3p⁶

Oxidation State = $\pm 1, 3, 5, 7$

Physical
Properties
from NIST

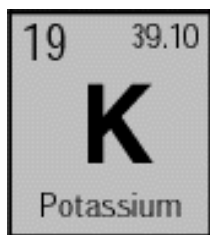
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}
3,608.4	378.6	297.3	294.6	34.8	18.3	18.3

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1
3,313.8	3,311.1	3,589.6

Electron Configuration = [Ar]4s¹

Oxidation State = 1

Density = 0.86 g/cm³ Concentration = 1.33 x 10²² atoms/cm³

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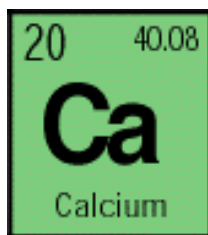
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}
4,038.5	438.4	349.7	346.2	44.3	25.4	25.4

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1
3,691.68	3,688.09	4,012.7	341.3	341.3	344.9

Electron Configuration = [Ar]4s²

Oxidation State = **2** Nearest Neighbor = 0.395 nm

Density = 1.53 g/cm³ Concentration = 2.30 x 10²² atoms/cm³

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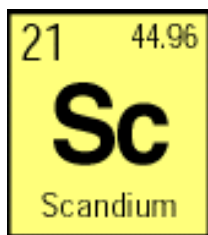
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}
4,492.	498.0	403.6	398.7	51.1	28.3	28.3

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1
4,090.6	4,086.1	4,460.5	395.4	395.4	399.6

Electron Configuration = [Ar]3d¹4s²

Oxidation State = **3** Nearest Neighbor = 0.325 nm

Density = 2.99 g/cm³ Concentration = 4.01 x 10²² atoms/cm³

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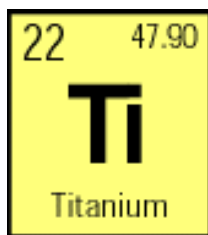
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}
4,966.	560.9	460.2	453.8	58.7	32.6	32.6

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1
4,510.84	4,504.86	4,931.81	452.2	452.2	458.4

Electron Configuration = [Ar]3d²4s²

Oxidation State = 4,3 Nearest Neighbor = 0.289 nm

Density = 4.51 g/cm³ Concentration = 5.67 x 10²² atoms/cm³

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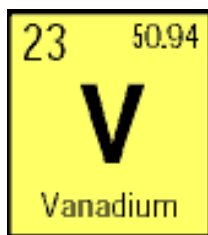
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}
5,465.	626.7	519.8	512.1	66.3	37.2	37.2

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1
4,952.20	4,944.64	5,427.29	511.3	511.3	519.2

Electron Configuration = [Ar]3d³4s²

Oxidation State = 5,4,3,2 Nearest Neighbor = 0.262 nm

Density = 6.09 g/cm³ Concentration = 7.20 x 10²² atoms/cm³

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ELEMENTS
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}
5,989.	696.0	583.8	574.1	74.1	42.2	42.2

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1
5,414.72	5,405.509	5,946.71	572.8	572.8	582.8

Electron Configuration = [Ar]3d⁵4s¹

Oxidation State = 6,3,2 Nearest Neighbor = 0.250 nm

Density = 7.19 g/cm³ Concentration = 8.33 x 10²² atoms/cm³

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**VISUAL
ELEMENTS
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}
6539.	769.1	649.9	638.7	82.3	47.2	47.2

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1
5,898.75	5,887.65	6,490.45	637.4	637.4	648.8

Electron Configuration = [Ar]3d⁵4s²

Oxidation State = 7,6,4,3,2

Density = 7.47 g/cm³ Concentration = 8.19 x 10²² atoms/cm³

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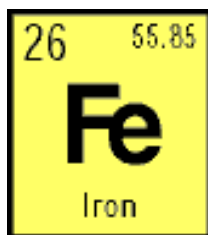
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}
7,112.	844.6	719.9	706.8	91.3	52.7	52.7

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1
6,403.84	6,390.84	7,057.98	705.0	705.0	718.5

Electron Configuration = [Ar]3d⁶4s²

Oxidation State = 2,3 Nearest Neighbor = 0.248 nm

Density = 7.87 g/cm³ Concentration = 8.49 x 10²² atoms/cm³

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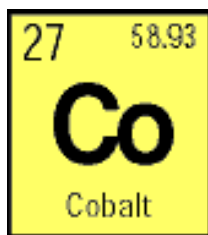
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}
7,709.	925.1	793.2	778.1	101.0	58.9	59.9

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1
6,930.32	6,915.30	7,649.43	776.2	776.2	791.4

Electron Configuration = [Ar]3d⁷4s²

Oxidation State = 2,3 Nearest Neighbor = 0.250 nm

Density = 8.82 g/cm³ Concentration = 9.01 x 10²² atoms/cm³

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Properties
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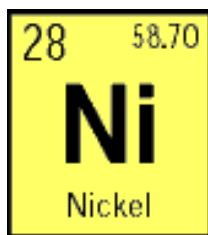
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}
8,333.	1,008.6	870.0	852.7	110.8	68.0	66.2

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1
7,478.15	7,460.89	8,264.66	851.5	851.5	868.8

Electron Configuration = [Ar]3d⁸4s²

Oxidation State = 2,3 Nearest Neighbor = 0.249 nm

Density = 8.91 g/cm³ Concentration = 9.14 x 10²² atoms/cm³

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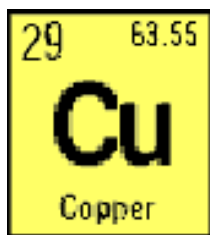
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}
8,979.0	1,096.7	952.3	932.7	122.5	77.3	75.1

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1
8,047.78	8,027.83	8,905.29	929.7	929.7	949.8

Electron Configuration = [Ar]3d¹⁰4s¹

Oxidation States = 2,1 Nearest Neighbor = 0.256 nm

Density = 8.93 g/cm³ Concentration = 8.47 x 10²² atoms/cm³

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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
9,659.	1,196.2	1,044.9	1,021.8	139.8	91.4	88.6	10.2	10.1

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1
8,638.86	8,615.78	9,572.0	1,011.7	1,011.7	1,034.7

Electron Configuration = [Ar]3d¹⁰4s²

Oxidation States = **2** Nearest Neighbor = 0.267 nm

Density = 7.13 g/cm³ Concentration = 6.57 x 10²² atoms/cm³

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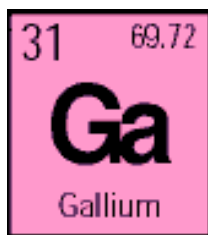
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
10,367.	1,299.0	1,143.2	1,116.4	159.5	103.5	100.0	18.7	18.7

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1
9,251.74	9,224.82	10,264.2	1,097.92	1,097.92	1,124.8

Electron Configuration = [Ar]3d¹⁰4s² 4p¹

Oxidation State = **3** Nearest Neighbor = 0.244 nm

Density = 5.91 g/cm³ Concentration = 5.10 x 10²² atoms/cm³

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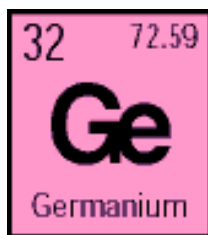
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
11,103.	1,414.6	1,248.1	1,217.0	180.1	124.9	120.8	29.8	29.2

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1
9,886.42	9,855.32	10,982.1	1,188.00	1,188.00	1,218.5

Electron Configuration = [Ar]3d¹⁰4s² 4p²

Oxidation State = 4 Nearest Neighbor = 0.245 nm

Density = 5.32 g/cm³ Concentration = 4.42 x 10²² atoms/cm³

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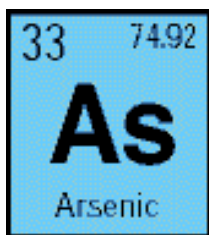
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
11,867.	1,527.0	1,359.1	1,323.6	204.7	146.2	141.2	41.7	41.7

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1
10,543.72	10,507.99	11,726.2	1,282.0	1,282.0	1,317.0

Electron Configuration = [Ar]3d¹⁰4s² 4p³

Oxidation State = $\pm 3, 5$ Nearest Neighbor = 0.251 nm

Density = 5.78 g/cm³ Concentration = 4.64 x 10²² atoms/cm³

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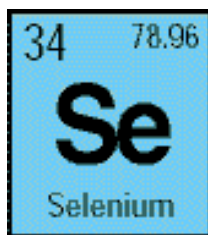
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
12,658.	1,652.0	1,474.3	1,433.9	229.6	166.5	160.7	55.5	54.6

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1
11,222.4	11,181.4	12,495.9	1,379.10	1,379.10	1,419.23

Electron Configuration = [Ar]3d¹⁰4s² 4p⁴

Oxidation State = -2,4,6 Nearest Neighbor = 0.232 nm

Density = 4.81 g/cm³ Concentration = 3.67 x 10²² atoms/cm³

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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
13,474.	1,782.	1,596.	1,550.	257.	189.	182.	70.	69.

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1
11,924.2	11,877.6	13,291.4	1,480.43	1,480.43	1,525.90

Electron Configuration = [Ar]3d¹⁰4s² 4p⁵

Oxidation State = $\pm 1, 5$

Density = 3.12 g/cm³ Concentration = 2.35 x 10²² atoms/cm³

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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
14,326.	1,921.	1,730.9	1,678.4	292.8	222.2	214.4	95.0	93.8

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
27.5	14.1	14.1

K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁
12,649.	12,598.	14,112.	1,586.0	1,586.0	1,636.6

Electron Configuration = [Ar]3d¹⁰4s² 4p⁶

Physical
Properties
from NIST

VISUAL
ELEMENTS
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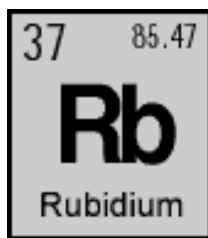
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
15,200.	2,065.	1,864.	1,804.	326.7	248.7	239.1	113.0	112.

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
30.5	16.3	15.3

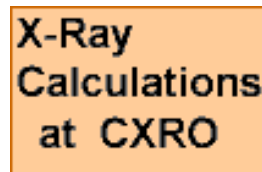
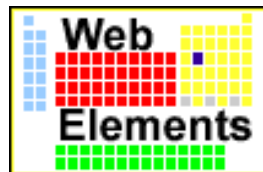
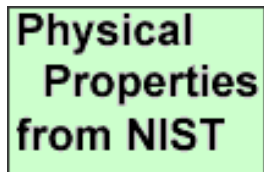
K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1
13,395.3	13,335.8	14,961.3	1,694.13	1,692.56	1,752.17

Electron Configuration = [Kr]5s¹

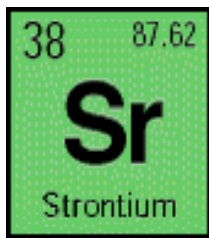
Oxidation State = 1

Density = 1.53 g/cm³ Concentration = 1.08 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
16,105.	2,216.	2,007.	1,940.	358.7	280.3	270.0	136.0	134.2

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
38.9	21.3	20.1

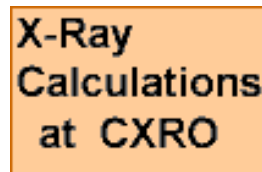
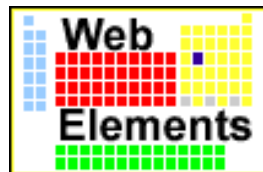
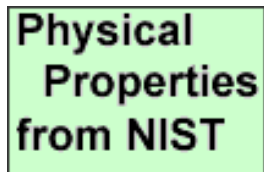
K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁
14,165.	14,097.9	15,835.7.	1,806.56	1,804.74	1,871.72

Electron Configuration = [Kr]5s²

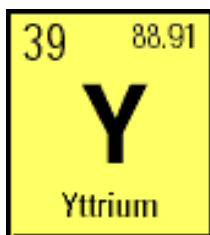
Oxidation State = 2 Nearest Neighbor = 0.430 nm

Density = 2.58 g/cm³ Concentration = 1.78 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

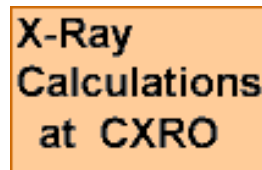
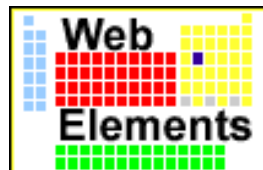
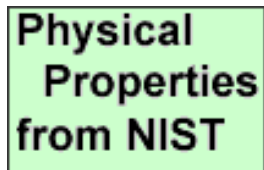
K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
17,038.	2,373.	2,156.	2,080.	392.0	310.6	298.8	157.7	155.8

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
43.8	24.4	23.1

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1
14,958.4	14,882.9	16,737.8	1,922.56	1,920.47	1,995.84

Electron Configuration = [Kr]4d¹5s²Oxidation State = **3** Nearest Neighbor = 0.355 nmDensity = 4.48 g/cm³ Concentration = 3.03 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
17,998.	2,532.	2,307	2,223	430.3	343.5	329.8	181.1	178.8

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
50.6	28.5	27.1

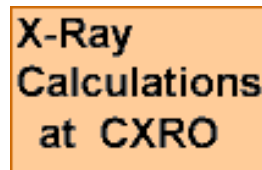
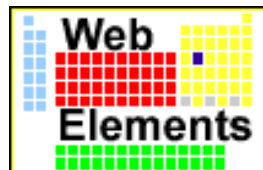
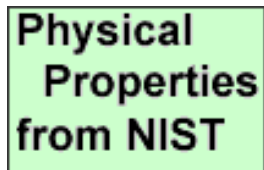
K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁
15,775.1	15,690.9	17,667.8	2,042.36	2,039.9	2,124.4	2,219.4	2,302.7

Electron Configuration = [Kr]4d²5s²

Oxidation State = 4 Nearest Neighbor = 0.318 nm

Density = 6.51 g/cm³ Concentration = 4.30 x 10²² atoms/cm³



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41	92.91
Nb	
Niobium	

Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
18,986.	2,698.	2,465.	2,371.	466.6	376.1	360.6	205.0	202.3

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
56.4	32.6	30.8

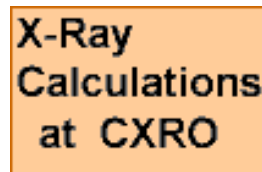
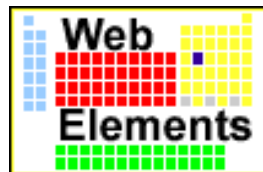
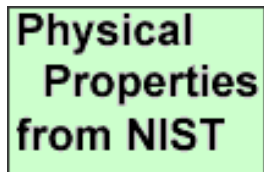
K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁
16,615.1	16,521.0	18,622.5	2,165.89	2,163.0	2,257.4	2,367.0	2,461.8

Electron Configuration = [Kr]4d⁴5s¹

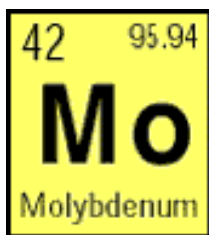
Oxidation State = 5,3 Nearest Neighbor = 0.286 nm

Density = 8.58 g/cm³ Concentration = 5.56 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
20,000.	2,866.	2,625.	2,520.	506.3	411.6	394.0	231.1	227.9

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
63.2	37.6	35.5

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1
17,479.34	17,374.3	19,608.3	2,293.16	2,289.85	2,394.81	2,518.3	2,623.5

Electron Configuration = [Kr]4d⁵5s¹

Oxidation State = 6,3,2 Nearest Neighbor = 0.273 nm

Density = 10.2 g/cm³ Concentration = 6.42 x 10²² atoms/cm³

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from CHEMSOC**

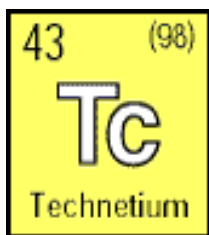
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Electron Binding Energies in electron volts

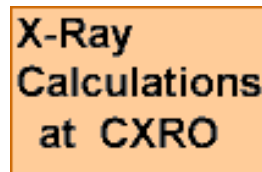
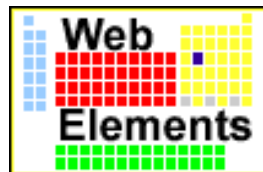
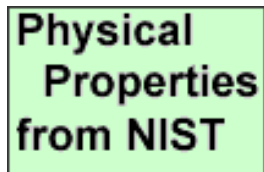
K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
21,044.	3,043.	2,793.	3,677.	544.	447.6	417.7	257.6	253.9

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
69.5	42.3	39.9

K and L shell emission lines in electron volts

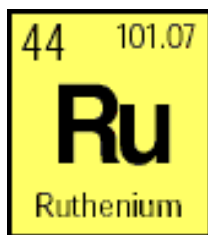
K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1
18,367.1	18,250.8	20,619.	2,424.	2,420.	2,538.	2,674.	2,792.

Electron Configuration = [Kr]4d⁵5s²Oxidation State = **7** Nearest Neighbor = 0.270 nmDensity = 11.5 g/cm³ Concentration = 7.07 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
22,117.	3,224.	2,967.	2,838.	586.1	483.5	461.4	284.2	280.0

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
75.0	46.3	43.2

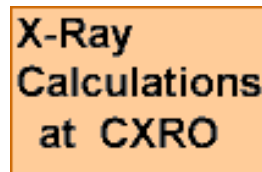
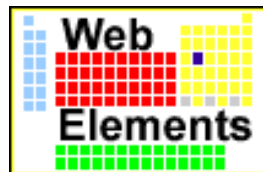
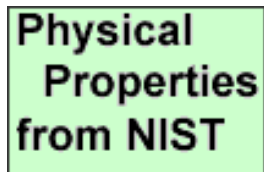
K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1
19,279.2	19,150.4	21,656.8	2,558.55	2,544.31	2,683.23	2,836.0	2,964.5

Electron Configuration = [Kr]4d⁷5s¹

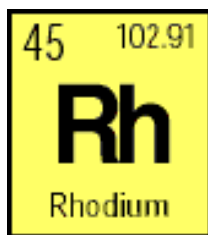
Oxidation State = 2,3,4,6,8 Nearest Neighbor = 0.265 nm

Density = 12.4 g/cm³ Concentration = 7.36 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
23,220.	3,412.	3,146	3,004	628.1	521.3	496.5	311.9	307.2

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
81.4	50.5	47.3

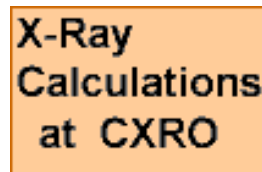
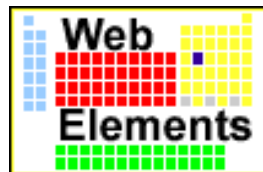
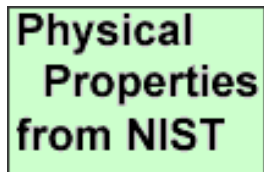
K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1
20,216.1	20,073.7	22,723.6	2,696.74	2,692.05	2,834.41	3,001.3	3,143.8

Electron Configuration = [Kr]4d⁶5s¹

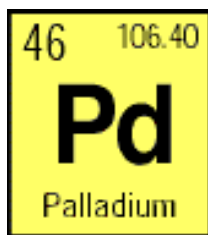
Oxidation State = 2,3,4 Nearest Neighbor = 0.269 nm

Density = 12.4 g/cm³ Concentration = 7.27 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
24,350.	3,604.	3,330.	3,173.	671.6	559.9	532.3	340.5	335.2

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
87.1	55.7	50.9

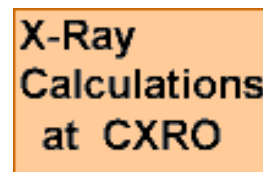
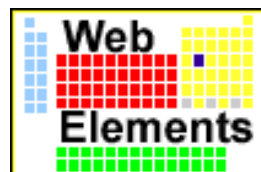
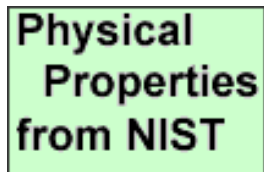
K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁
21,177.1	21,020.1	23,818.7	2,838.61	2,833.29	2,990.22	3,171.79	3,328.7

Electron Configuration = [Kr]4d¹⁰

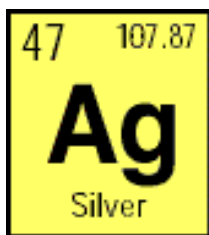
Oxidation State = 2,4 Nearest Neighbor = 0.275 nm

Density = 12.0 g/cm³ Concentration = 6.79 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
25,514.	3,806.	3,524.	3,351.	719.0	603.8	573.0	374.0	368.3

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
97.0	63.7	58.3

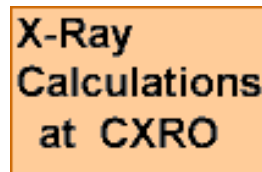
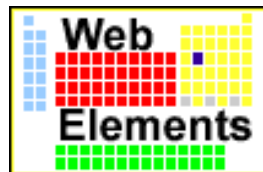
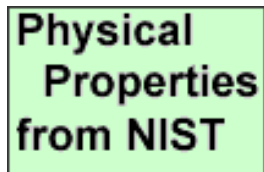
K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁
22,162.92	21,990.3	24,942.4	2,984.31	2,978.21	3,150.94	3,347.81	3,519.59

Electron Configuration = [Kr]4d¹⁰5s¹

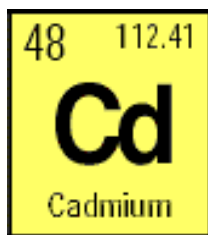
Oxidation State = 1 Nearest Neighbor = 0.289 nm

Density = 10.5 g/cm³ Concentration = 5.86 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
26,711.	4,018.	3,727.	3,538.	772.0	652.6	618.4	411.9	405.2

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
109.8	63.9	63.9

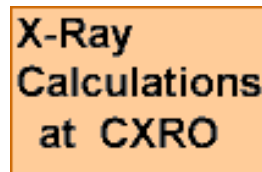
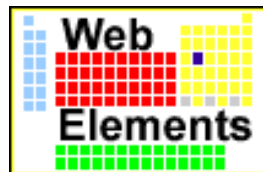
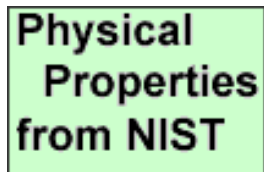
K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁
23,173.6	22,984.1	26,095.5	3,133.73	3,126.91	3,316.57	3,528.12	3,716.86

Electron Configuration = [Kr]4d¹⁰5s²

Oxidation State = 2 Nearest Neighbor = 0.298 nm

Density = 8.65 g/cm³ Concentration = 4.63 x 10²² atoms/cm³



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49	114.82
In	
Indium	

Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
27,940.	4,238.	3,938.	3,730.	827.2	703.2	665.3	451.4	443.9

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
122.9	73.5	73.5

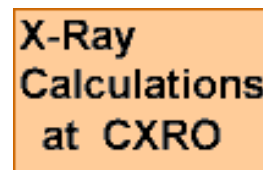
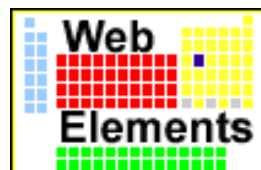
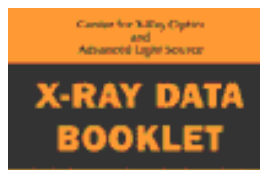
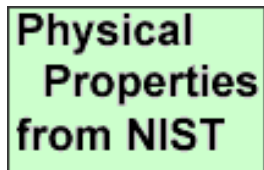
K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁
24,209.7	24,002.0	27,275.9	3,286.94	3,279.29	3,487.21	3,713.81	3,920.81

Electron Configuration = [Kr]4d¹⁰5s²5p¹

Oxidation State = 3 Nearest Neighbor = 0.325 nm

Density = 7.29 g/cm³ Concentration = 3.82 x 10²² atoms/cm³



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50	118.69
Sn	
Tin	

Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
29,200.	4,465.	4,156.	3,929.	884.7	756.5	714.6	493.2	484.9

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
137.1	83.6	83.6

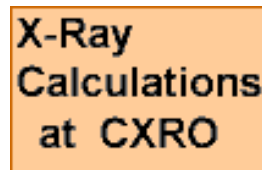
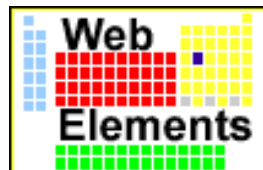
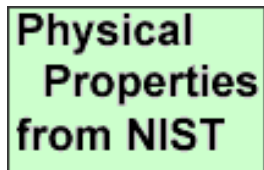
K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁
25,271.3	25,044.0	28,486.0	3,443.98	3,435.42	3,662.80	3,904.86	4,131.12

Electron Configuration = [Kr]4d¹⁰5s²5p²

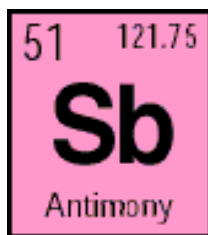
Oxidation State = 4,2 Nearest Neighbor = 0.302 nm

Density = 7.29 g/cm³ Concentration = 3.70 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
30,491.	4,698.	4,380.	4,132.	946.	812.7	766.4	537.5	528.2

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
153.2	95.6	95.6

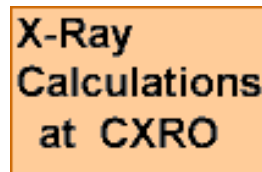
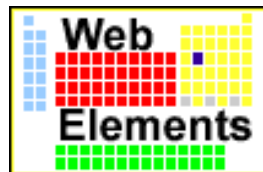
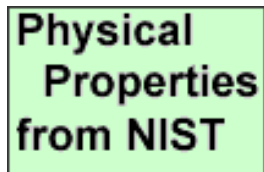
K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁
26,359.1	26,110.8	29,725.6	3,604.72	3,595.32	3,843.57	4,100.78	4,347.79

Electron Configuration = [Kr]4d¹⁰5s² 5p³

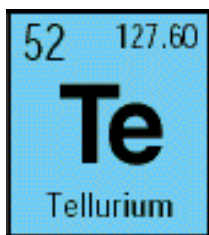
Oxidation State = ±3,5 Nearest Neighbor = 0.290 nm

Density = 6.69 g/cm³ Concentration = 3.31 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
31,814.	4,939.	4,612.	4,341.	1,006.	870.8	820.0	583.4	573.0

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
169.4	103.3	103.3

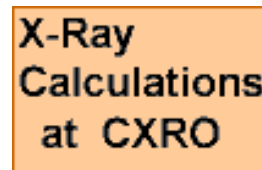
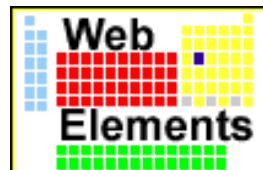
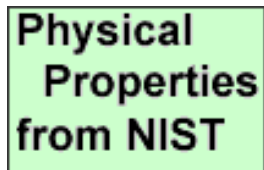
K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1
27,472.3	27,201.7	30,995.7	3,769.33	3,758.8	4,029.58	4,301.7	4,570.9

Electron Configuration = [Kr]4d¹⁰5s² 5p⁴

Oxidation State = -2,4,6 Nearest Neighbor = 0.286 nm

Density = 6.25 g/cm³ Concentration = 2.95 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
33,169.	5,188.	4,852.	4,557.	1,072.	931.	875.	630.8	619.3

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
186.	123.	123.

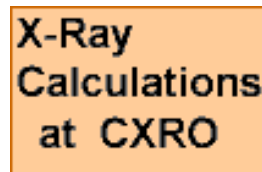
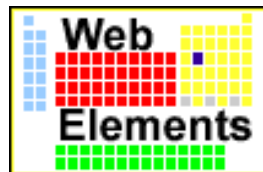
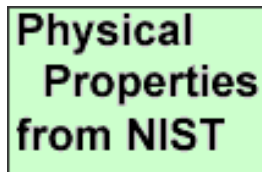
K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1
28,612.0	28,317.2	32,294.7	3,937.65	3,926.04	4,220.72	4,507.5	4,800.9

Electron Configuration = [Kr]4d¹⁰5s² 5p⁵

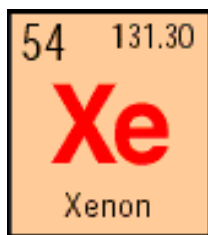
Oxidation State = $\pm 1, 5, 7$

Density = 4.95 g/cm³ Concentration = 2.35 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
34,561.	5,453.	5,107.	4,786.	1,148.7	1,002.1	940.6	689.0	676.4

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
213.2	146.7	145.5

K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁
29,779.	29,458.	33,624.	4,109.9	—	—	—	—

Electron Configuration = [Kr]4d¹⁰5s² 5p⁶

Physical
Properties
from NIST

VISUAL
ELEMENTS
from CHEMSOC

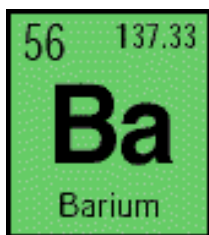
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
37,441.	5,989.	5,624.	5,247.	1,293.	1,137.	1,063.	795.7	780.5

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
253.5	192.	178.6

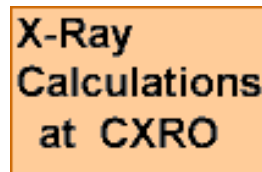
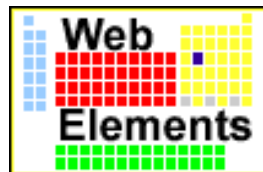
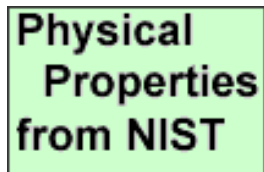
K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1
32,193.6	31,817.1	36,378.2	4,466.26	4,450.90	4,827.53	5,156.5	5,531.1

Electron Configuration = [Xe]6s²

Oxidation State = 2 Nearest Neighbor = 0.435 nm

Density = 3.59 g/cm³ Concentration = 1.58 x 10²² atoms/cm³



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57	138.91
La	
Lanthanum	

Electron Binding Energies in electron volts

K 1s	L₁2s	L₂2p_{1/2}	L₃2p_{3/2}	M₁3s	M₂3p_{1/2}	M₃3p_{3/2}	M₄3d_{3/2}	M₅3d_{5/2}
38,925.	6,266.	5,891.	5,483.	1,362.	1,209.	1,128.	853.	836.

N₁4s	N₂4p_{1/2}	N₃4p_{3/2}
274.7	205.8	196.0

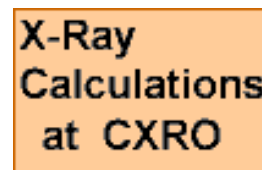
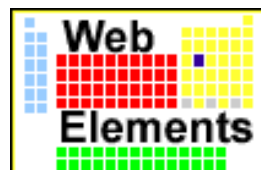
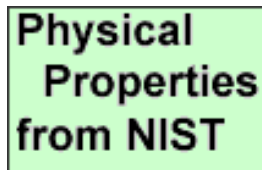
K and L shell emission lines in electron volts

Kα_1	Kα_2	Kβ_1	Lα_1	Lα_2	Lβ_1	Lβ_2	Lγ_1	Mα_1
33,441.8	33,034.1	37,801.0	4,650.97	4,634.23	5,042.1	5,383.5	5,788.5	833.

Electron Configuration = [Xe]5d¹6s²

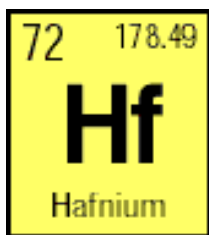
Oxidation State = **3** Nearest Neighbor = 0.373 nm

Density = 6.15 g/cm³ Concentration = 2.68 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
65,351.	11,271.	10,739.	9,561.	2,601.	2,365.	2,108.	1,716.	1,662.

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
538.	438.2	380.7

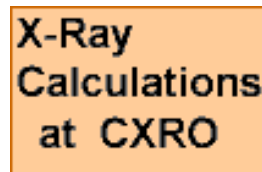
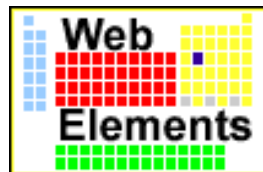
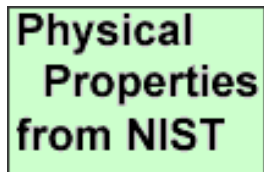
K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁	Mα ₁
55,790.2	54,611.4	63,234.	7,899.0	7,844.6	9,022.7	9,347.3	10,515.8	1,644.6

Electron Configuration = [Xe]4f¹⁴5d²6s²

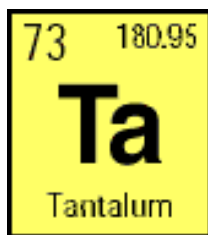
Oxidation State = 4 Nearest Neighbor = 0.313 nm

Density = 13.3 g/cm³ Concentration = 4.48 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
67,416.	11,682.	11,136.	9,881.	2,708.	2,469.	2,194.	1,793.	1,735.

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
563.4	463.4	400.9

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1	M α_1
57,532.	56,277.	65,223.	8,146.1	8,087.9	9,343.1	9,651.8	10,895.2	1,710.

Electron Configuration = [Xe]4f¹⁴5d³6s²Oxidation State = **5** Nearest Neighbor = 0.286 nmDensity = 16.7 g/cm³ Concentration = 5.55 x 10²² atoms/cm³

**Physical
Properties
from NIST**

**VISUAL
ELEMENTS
from CHEMSOC**

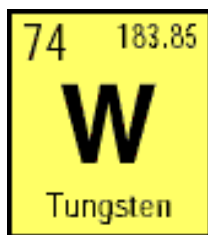
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
69,525.	12,100.	11,544.	10,207.	2,820.	2,575.	2,281.	1,872.	1,809.

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
594.1	490.4	423.6

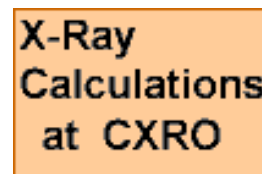
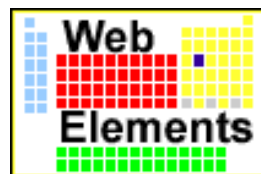
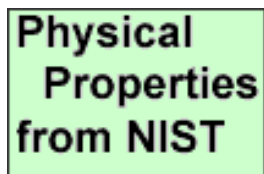
K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1	M α_1
59,318.24	57,981.7	67,244.3	8,397.6	8,335.2	9,672.35	9,961.5	11,285.9	1,775.4

Electron Configuration = [Xe]4f¹⁴5d⁴6s²

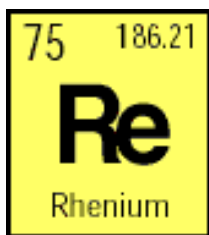
Oxidation State = 6,5,4,3,2 Nearest Neighbor = 0.274 nm

Density = 19.3 g/cm³ Concentration = 6.31 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
71,676.	12,527.	11,959.	10,535.	2,932.	2,682.	2,367.	1,949.	1,883.

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
625.4	518.7	446.8

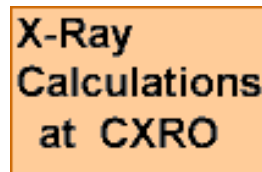
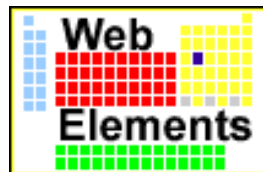
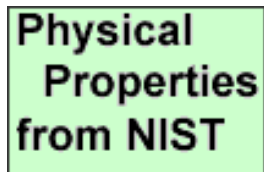
K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁	Mα ₁
61,140.3	59,717.9	69,310.	8,652.5	8,586.2	10,010.0	10,275.2	11,685.4	1,842.5

Electron Configuration = [Xe]4f¹⁴5d⁵6s²

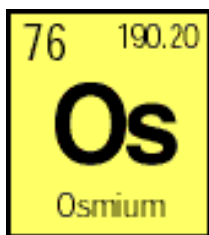
Oxidation State = 7,6,4,2,-1 Nearest Neighbor = 0.274 nm

Density = 21.0 g/cm³ Concentration = 6.80 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
73,871.	12,968.	12,385.	10,871.	3,049.	2,792.	2,457.	2,031.	1,960.

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
658.2	549.1	470.7

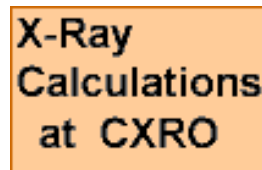
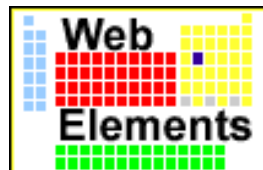
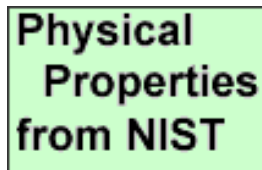
K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁	Mα ₁
63,000.5	61,486.7	71,413.	8,911.7	8,841.0	10,355.3	10,598.5	12,095.3	1,910.2

Electron Configuration = [Xe]4f¹⁴5d⁶6s²

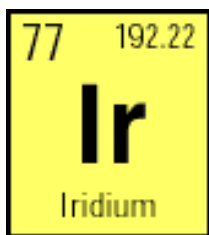
Oxidation State = 2,3,**4**,6,8 Nearest Neighbor = 0.268 nm

Density = 22.6 g/cm³ Concentration = 7.15 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
67,111.	13,419.	12,824.	11,215.	3,174.	2,909.	2,551.	2,116.	2,040.

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
691.1	577.8	495.8

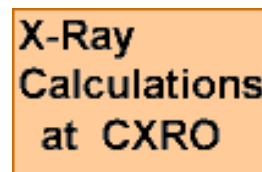
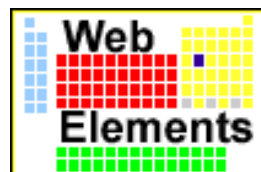
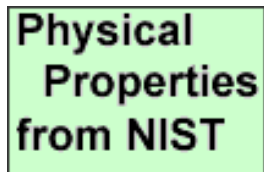
K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1	M α_1
64,895.6	63,286.7	73,560.8	9,175.1	9,099.5	10,708.3	10,920.3	12,512.6	1,979.9

Electron Configuration = [Xe]4f¹⁴5d⁷6s²

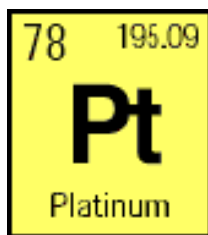
Oxidation State = 2,3,4,6 Nearest Neighbor = 0.271 nm

Density = 22.6 g/cm³ Concentration = 7.07 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
78,395.	13,880.	13,273.	11,564.	3,296.	3,027.	2,645.	2,202.	2,122.

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
725.4	609.1	519.4

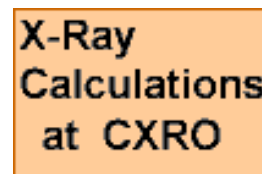
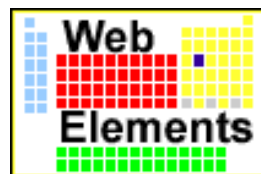
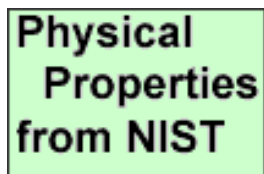
K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁	Mα ₁
66,832.	65,112.	75,748.	9,442.3	9,361.8	11,070.7	11,250.5	12,942.0	2,050.5

Electron Configuration = [Xe]4f¹⁴5d⁹6s¹

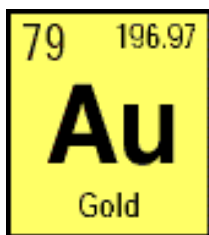
Oxidation State = 2,4 Nearest Neighbor = 0.278 nm

Density = 21.4 g/cm³ Concentration = 6.62 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
80,725.	14,353.	13,734.	11,919.	3,425.	3,148.	2,743.	2,291.	2,206.

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
762.1	642.7	546.3

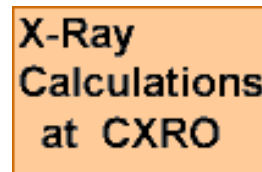
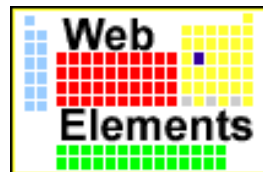
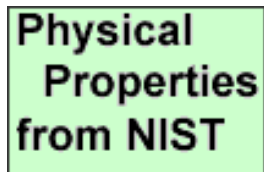
K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁	Mα ₁
68,803.7	66,989.5	77,984.	9,713.3	9,628.0	11,442.3	11,584.7	13,381.7	2,122.9

Electron Configuration = [Xe]4f¹⁴5d¹⁰6s¹

Oxidation State = 3,1 Nearest Neighbor = 0.288 nm

Density = 19.3 g/cm³ Concentration = 5.89 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
83,102.	14,839.	14,209.	12,284.	3,562.	3,279.	2,847.	2,385.	2,295.

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
802.2	680.2	576.6

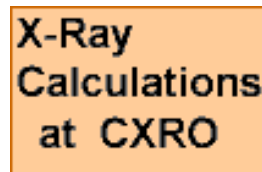
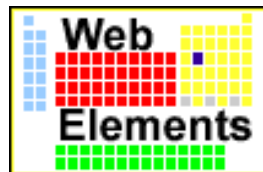
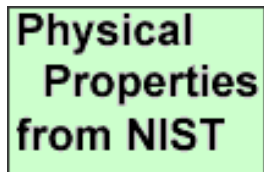
K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁	Mα ₁
70,819.	68,895.	80,253.	9,988.8	9,897.6	11,822.6	11,924.1	13,830.1	2,195.3

Electron Configuration = [Xe]4f¹⁴5d¹⁰6s²

Oxidation State = 2, 1

Density = 13.59 g/cm³



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81	204.37
Tl	
Thallium	

Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
85,530.	15,347.	14,698.	12,658.	3,704.	3,416.	2,957.	2,485.	2,389.

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
846.2	720.5	609.5

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1	M α_1
72,871.5	70,831.9	82,576.	10,268.5	10,172.8	12,213.3	12,271.5	14,291.5	2,270.6

Electron Configuration = [Xe]4f¹⁴5d¹⁰6s²6p¹

Oxidation State = 3,1 Nearest Neighbor = 0.341 nm

Density = 11.9 g/cm³ Concentration = 3.50 x 10²² atoms/cm³

**Physical
Properties
from NIST**

**VISUAL
ELEMENTS
from CHEMSOC**

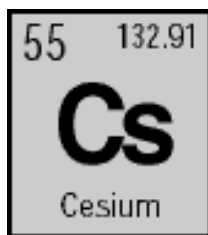
Center for X-Ray Optics
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Advanced Light Source
**X-RAY DATA
BOOKLET**

**Web
Elements**

**X-Ray
Calculations
at CXRO**

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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
35,985.	5,714.	5,359.	5,012.	1,211.	1,071.	1,003.	740.5	726.6

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
232.3	172.4	161.3

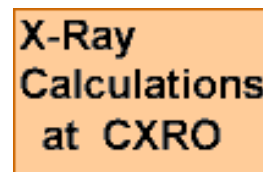
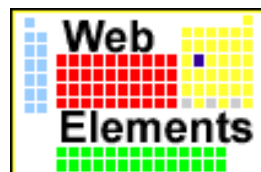
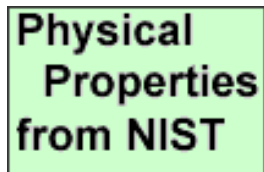
K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁
30,972.8	30,625.1	34,986.9	4,286.5	4,272.2	4,619.8	4,935.9	5,280.4

Electron Configuration = [Xe]6s¹

Oxidation State = 1

Density = 1.90 g/cm³ Concentration = 0.86 x 10²² atoms/cm³



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82	207.19
Pb	
Lead	

Electron Binding Energies in electron volts

K 1s	L₁2s	L₂2p_{1/2}	L₃2p_{3/2}	M₁3s	M₂3p_{1/2}	M₃3p_{3/2}	M₄3d_{3/2}	M₅3d_{5/2}
88,005.	15,861.	15,200.	13,035.	3,851.	3,554.	3,066.	2,586.	2,484.

N₁4s	N₂4p_{1/2}	N₃4p_{3/2}
891.8	761.9	643.5

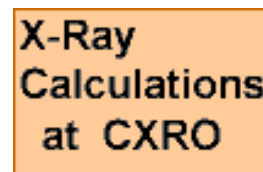
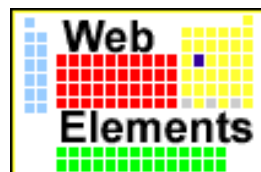
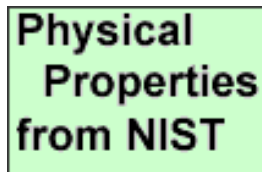
K and L shell emission lines in electron volts

Kα_1	Kα_2	Kβ_1	Lα_1	Lα_2	Lβ_1	Lβ_2	Lγ_1	Mα_1
74,969.4	72,804.2	84,936.	10,551.5	10,449.5	12,613.7	12,622.6	14,764.4	2,345.5

Electron Configuration = [Xe]4f¹⁴5d¹⁰6s² 6p²

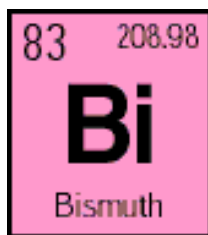
Oxidation State = 4,2 Nearest Neighbor = 0.350 nm

Density = 11.3 g/cm³ Concentration = 3.30 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
90,524.	16,388.	15,711.	13,419.	3,999.	3,696.	3,177.	2,688.	2,580.

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
939.	805.2	678.8

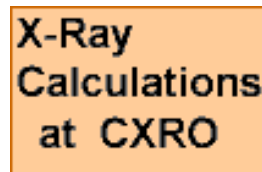
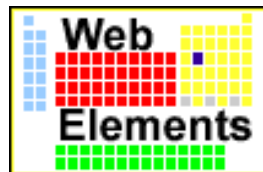
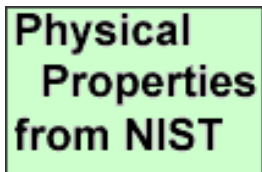
K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1	M α_1
77,107.9	74,814.8	87,343.	10,838.8	10,730.91	13,023.5	12,979.9	15,247.7	2,422.6

Electron Configuration = [Xe]4f¹⁴5d¹⁰6s² 6p³

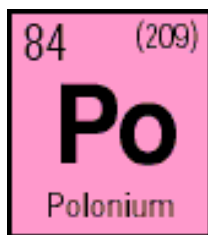
Oxidation State = 3,5 Nearest Neighbor = 0.311 nm

Density = 9.80 g/cm³ Concentration = 2.82 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
93,105.	16,939.	16,244.	13,814.	4,149.	3,854.	3,302.	2,798.	2,683.

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
995.	851.	705.

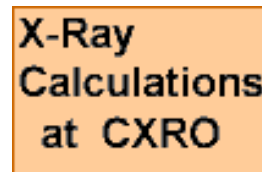
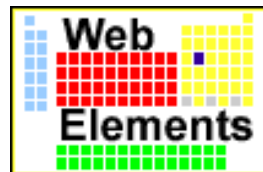
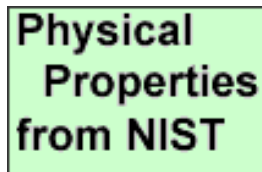
K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1	M α_1
79,290.	76,862.	89,800.	11,130.8	11,015.8	13,447.	13,340.4	15,744.	—

Electron Configuration = [Xe]4f¹⁴5d¹⁰6s²6p⁴

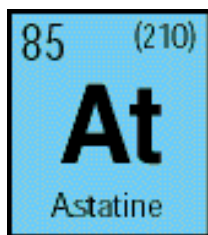
Oxidation State = 4,2 Nearest Neighbor = 0.335 nm

Density = 9.27 g/cm³ Concentration = 2.67 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
95,730.	17,493.	16,785.	14,214.	4,317.	4,008.	3,426.	2,909.	2,787.

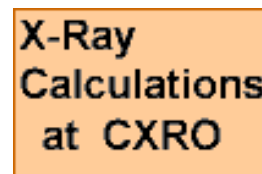
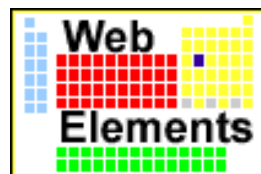
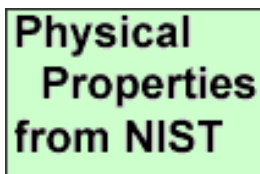
N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
1,042.	886.	740.

K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁	Mα ₁
81,520.	78,950.	92,300.	11,426.8	11,304.8	13,876.	—	16,251.	—

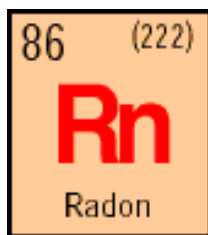
Electron Configuration = [Xe]4f¹⁴5d¹⁰6s²6p⁵

Oxidation State = ±1,3,5,7



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
98,404.	18,049.	17,337.	14,619.	4,482.	4,159.	3,538.	3,022.	2,892.

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
1,097.	929.	768.

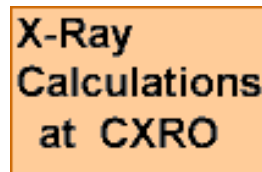
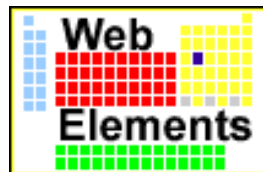
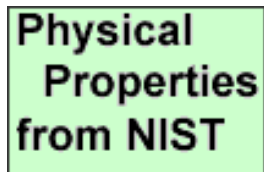
K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1	M α_1
83,780.	81,070.	94,870.	11,727.0	11,597.9	14,316.	—	16,770.	—

Electron Configuration = [Xe]4f¹⁴5d¹⁰6s²6p⁶

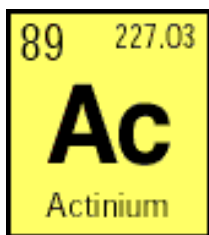
Oxidation State = 4 Nearest Neighbor = 0.313 nm

Density = 13.3 g/cm³ Concentration = 4.48 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

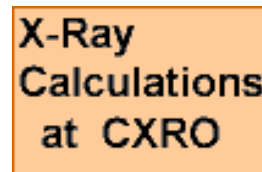
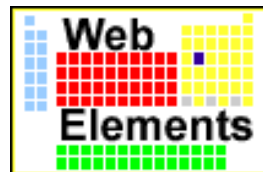
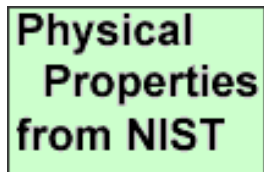
K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
106,755.	19,840.	19,083.	15,871.	5,002.	4,656.	3,909.	3,370.	3,219.

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
1,269.	1,080.	890.

K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁	Mα ₁
90,884.	87,670.	102,850.	12,652.0	12,500.8	15,713.	—	18,408.	—

Electron Configuration = [Rn]5f¹⁴6d²7s²Oxidation State = **3** Nearest Neighbor = 0.376 nmDensity = 10.07 g/cm³ Concentration = 2.67 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
109,651.	20,472.	19,693.	16,300.	5,182.	4,830.	4,046.	3,491.	3,332.

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1	M α_1
93,350.	89,953.	105,609.	12,968.7	12,809.6	16,202.2	15,623.7	18,982.5	2,996.1

Electron Configuration = [Rn]6d²7s²

**Physical
Properties
from NIST**

**VISUAL
ELEMENTS**
from CHEMSOC

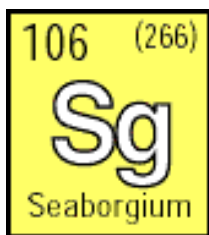
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
109,651.	20,472.	19,693.	16,300.	5,182.	4,830.	4,046.	3,491.	3,332.

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1	M α_1
93,350.	89,953.	105,609.	12,968.7	12,809.6	16,202.2	15,623.7	18,982.5	2,996.1

Electron Configuration = [Rn]5f¹⁴6d³7s²

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from NIST

VISUAL
ELEMENTS
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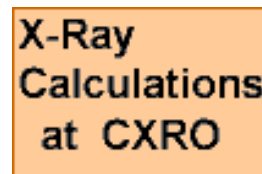
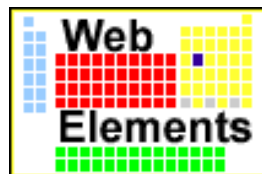
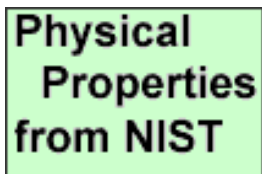
Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
109,651.	20,472.	19,693.	16,300.	5,182.	4,830.	4,046.	3,491.	3,332.

K and L shell emission lines in electron volts

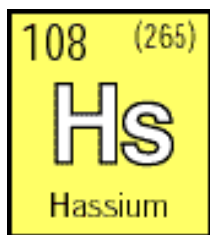
K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1	M α_1
93,350.	89,953.	105,609.	12,968.7	12,809.6	16,202.2	15,623.7	18,982.5	2,996.1

Electron Configuration = [Rn]5f¹⁴6d⁵7s²



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
109,651.	20,472.	19,693.	16,300.	5,182.	4,830.	4,046.	3,491.	3,332.

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1	M α_1
93,350.	89,953.	105,609.	12,968.7	12,809.6	16,202.2	15,623.7	18,982.5	2,996.1

Electron Configuration = [Rn]5f¹⁴6d⁶7s²

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Properties
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
109,651.	20,472.	19,693.	16,300.	5,182.	4,830.	4,046.	3,491.	3,332.

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1	M α_1
93,350.	89,953.	105,609.	12,968.7	12,809.6	16,202.2	15,623.7	18,982.5	2,996.1

Electron Configuration = [Rn]5f¹⁴6d⁷7s²

Physical
Properties
from NIST

VISUAL
ELEMENTS
from CHEMSOC

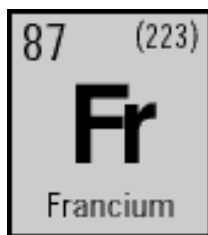
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
101,137.	18,639.	17,907.	15,031.	4,652.	4,327.	3,663.	3,136.	3,000.

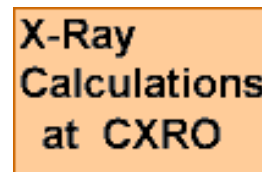
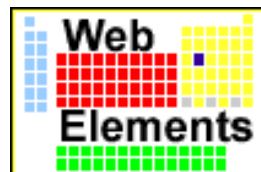
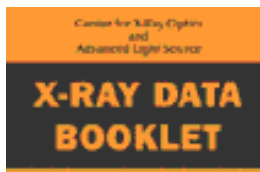
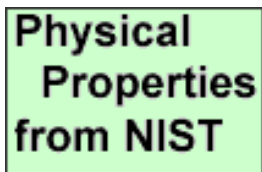
N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
1,153.	980.	810.

K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁	Mα ₁
86,100.	83,230.	97,470.	12,031.3	11,895.0	14,770.	14,450.	17,303.	—

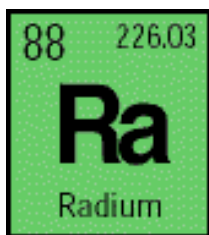
Electron Configuration = [Rn]7s¹

Oxidation State = 1



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
103,922.	19,237.	18,484.	15,444.	4,822.	4,490.	3,792.	3,248.	3,105.

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
1,208.	1,058.	879.

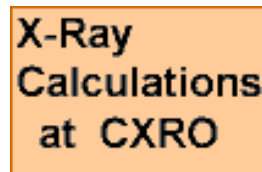
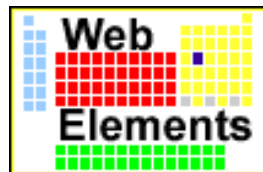
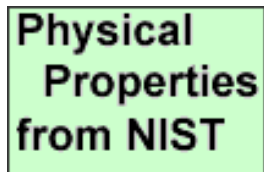
K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁	Mα ₁
88,470.	85,430.	100,130.	12,339.7	12,196.2	15,235.8	14,841.4	17,849.	—

Electron Configuration = [Rn]7s²

Oxidation State = 2

Density = 5 g/cm³



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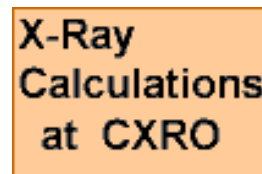
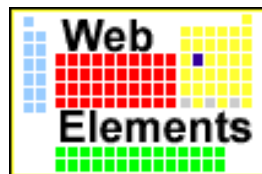
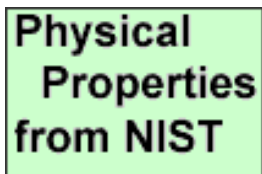
Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
109,651.	20,472.	19,693.	16,300.	5,182.	4,830.	4,046.	3,491.	3,332.

K and L shell emission lines in electron volts

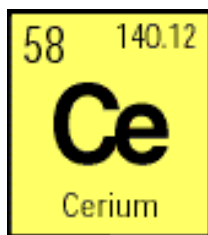
K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1	M α_1
93,350.	89,953.	105,609.	12,968.7	12,809.6	16,202.2	15,623.7	18,982.5	2,996.1

Electron Configuration = [Rn]5f¹⁴6d³7s²



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
40,443.	6,549.	6,164.	5,723.	1,436.	1,274.	1,187.	902.4.	883.8

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
291.0	223.2	206.5

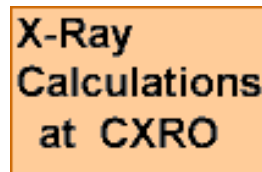
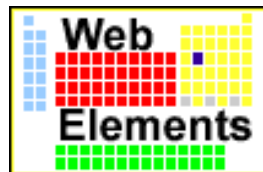
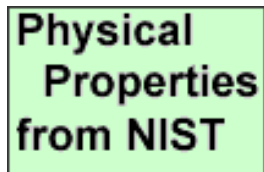
K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁	Mα ₁
34,719.7	34,278.9	39,257.3	4,840.2	4,823.0	5,262.2	5,613.4	6,052.	883

Electron Configuration = [Xe]4f¹⁵5d¹6s²

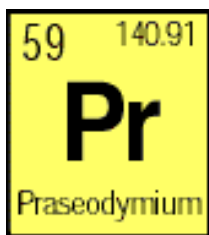
Oxidation State = 3,4 Nearest Neighbor = 0.365 nm

Density = 8.16 g/cm³ Concentration = 2.91 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
41,991.	6,835.	6,440.	5,964.	1,511.	1,337.	1,242.	948.3	928.8

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
304.5	236.3	217.6

K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁	Mα ₁
36,026.3	35,550.2	40,748.2	5,033.7	5,013.5	5,488.9	5,850.	6,322.1	929.

Electron Configuration = [Xe]4f³6s²Oxidation State = **3,4** Nearest Neighbor = 0.364 nmDensity = 6.77 g/cm³ Concentration = 2.90 x 10²² atoms/cm³

**Physical
Properties
from NIST**

**VISUAL
ELEMENTS
from CHEMSOC**

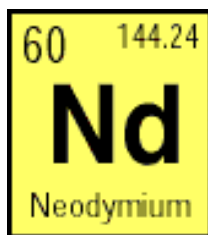
Center for X-Ray Optics
and
Advanced Light Source
**X-RAY DATA
BOOKLET**

**Web
Elements**

**X-Ray
Calculations
at CXRO**

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Electron Binding Energies in electron volts

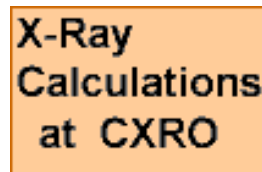
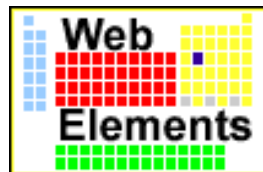
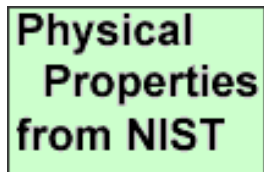
K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
43,569.	7,126.	6,722.	6,208.	1,575.	1,403.	1,297.	1,003.3	980.4

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
319.2	243.3	224.6

K and L shell emission lines in electron volts

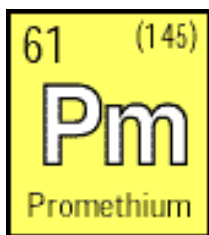
Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁	Mα ₁
37,361.0	36,847.4	42,271.3	5,230.4	5,207.7	5,721.6	6,089.4	6,602.1	978.

Electron Configuration = [Xe]4f⁴6s²Oxidation State = **3** Nearest Neighbor = 0.363 nmDensity = 7.01g/cm³ Concentration = 2.92 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
45,148.	7,428.	7,013.	6,459.	—	1,471.	1,357.	1,052.	1,027.

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
—	242.	242.

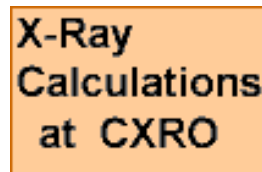
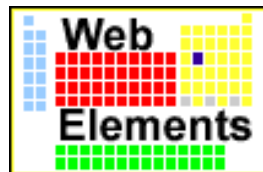
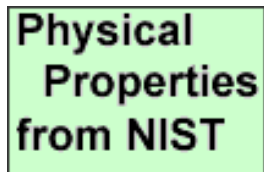
K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁	Mα ₁
38,724.7	38,171.2	43,826.	5,432.5	5,407.8	5,961.	6,339.	6,892.	—

Electron Configuration = [Xe]4f⁵6s²

Oxidation State = 3

Density = 7.26 g/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
46,834.	7,737.	7,312.	6,716.	1,723.	1,541.	1,420.	1,110.9	1083.4

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
347.2	265.6	247.4

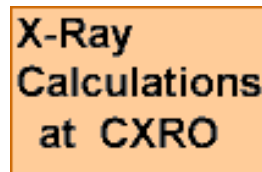
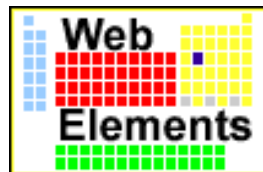
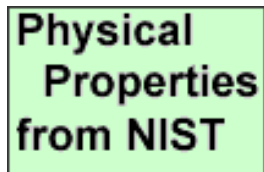
K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1	M α_1
40,118.1	39,522.4	45,413.	5,636.1	5,609.0	6,205.1	6,586.	7,178.	1,081.

Electron Configuration = [Xe]4f⁶6s²

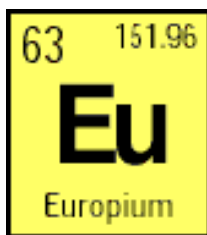
Oxidation State = 3,2 Nearest Neighbor = 0.359 nm

Density = 7.52 g/cm³ Concentration = 3.02 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
48,519.	8,052.	7,617.	6,977.	1,800.	1,614.	1,481.	1,158.6	1,127.5

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
360.	284.	257.

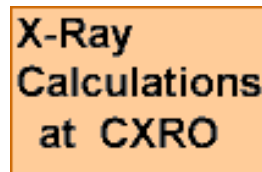
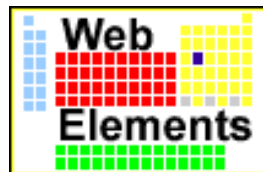
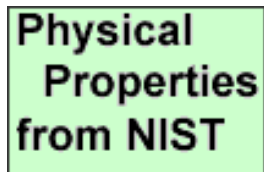
K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁	Mα ₁
42,542.2	40,901.9	47,037.9	5,845.7	5,816.6	6,456.4	6,843.2	7,480.3	1,131.

Electron Configuration = [Xe]4f⁷6s²

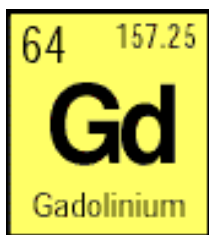
Oxidation State = 3,2 Nearest Neighbor = 0.397 nm

Density = 5.24 g/cm³ Concentration = 2.08 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

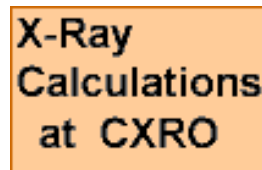
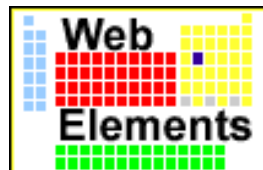
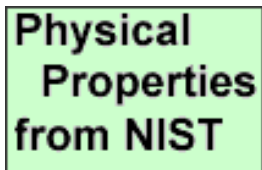
K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
50,239.	8,376.	7,930.	7,243.	1,881.	1,688.	1,544.	1,221.9	1,189.6

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
378.6	286.	271.

K and L shell emission lines in electron volts

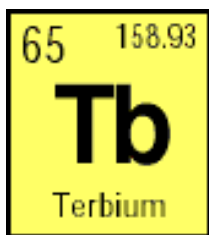
Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁	Mα ₁
42,996.2	42,308.9	48,697.	6,057.2	6,025.0	6,713.2	7,102.8	7,785.8	1,185

Electron Configuration = [Xe]4f⁷5d¹6s²Oxidation State = **3** Nearest Neighbor = 0.358 nmDensity = 7.90 g/cm³ Concentration = 3.01 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
51,996.	8,708.	8,252.	7,514.	1,968.	1,768	1,611	1,276.9	1,241.1

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
396.0	322.4	384.1

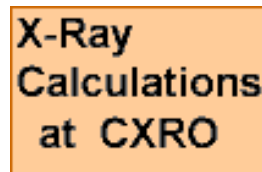
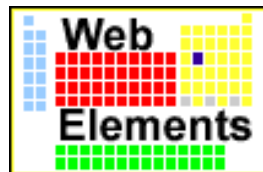
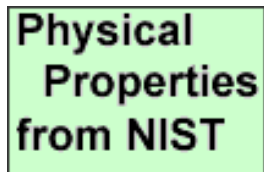
K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1	M α_1
44,481.6	43,744.1	50,382.	6,272.8	6,238.0	6,978.	7,366.7	8,102.	1,240.

Electron Configuration = [Xe]4f⁹6s²

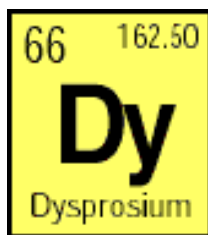
Oxidation State = 3,4 Nearest Neighbor = 0.353 nm

Density = 8.23 g/cm³ Concentration = 3.13 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

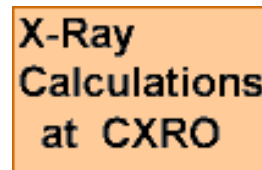
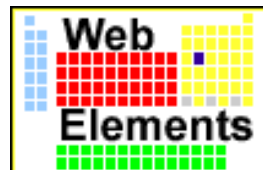
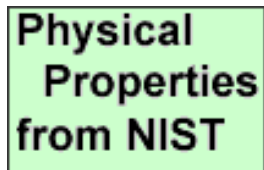
K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
53,789.	9,046.	8,581.	7,790.	2,047.	1,842.	1,676.	1,333.	1,292.6

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
414.2	333.5	293.2

K and L shell emission lines in electron volts

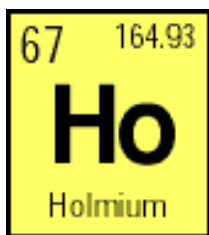
Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁	Mα ₁
45,998.4	45,207.8	52,119.	6,495.2	6,457.7	7,247.7	7,635.7	8,418.8	1,293.

Electron Configuration = [Xe]4f¹⁰6s²Oxidation State = **3** Nearest Neighbor = 0.351 nmDensity = 8.55 g/cm³ Concentration = 3.16 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

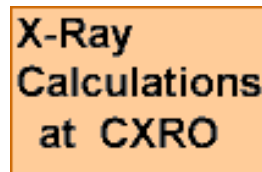
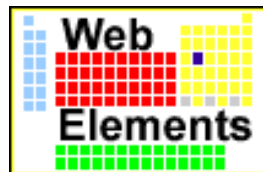
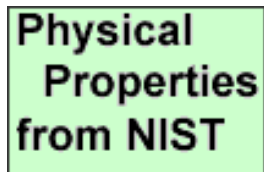
K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
55,618.	9,394.	8,918	8,071.	2,128.	1,923.	1,741.	1,392.	1,351.

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
432.4	434.5	308.2

K and L shell emission lines in electron volts

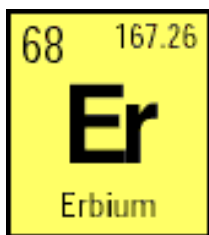
K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1	M α_1
47,546.7	46,699.7	53,877.	6,719.8	6,679.5	7,525.3	7,911.	8,747.	1,348.

Electron Configuration = [Xe]4f¹¹6s²Oxidation State = **3** Nearest Neighbor = 0.349 nmDensity = 8.80 g/cm³ Concentration = 3.21 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

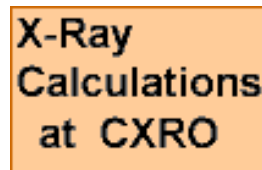
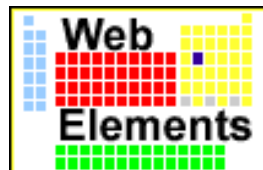
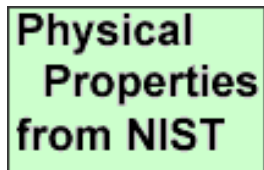
K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
57,486.	9,751.	9,264.	8,358.	2,207.	2,006.	1,812.	1,453.	1,409.

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
449.8	366.2	320.2

K and L shell emission lines in electron volts

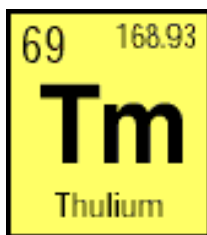
Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁	Mα ₁
49,127.7	48,221.1	55,681.	6,948.7	6,905.0	7,810.9	8,189.0	9,089.	1,406.

Electron Configuration = [Xe]4f¹²6s²Oxidation State = **3** Nearest Neighbor = 0.347 nmDensity = 9.07 g/cm³ Concentration = 3.26 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
59,390.	10,116.	9,617.	8,648.	2,307.	2,090.	1,885.	1,515.	1,468.

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
470.9	385.9	332.6

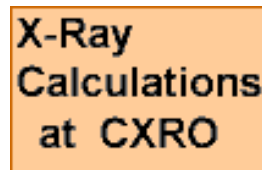
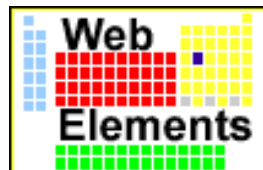
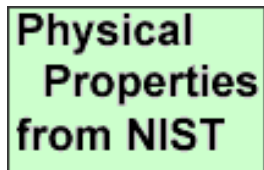
K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁	Mα ₁
50,741.6	49,772.6	57,517.	7,179.9	7,133.1	8,101.	8,468.	9,426.	1,462.

Electron Configuration = [Xe]4f¹³6s²

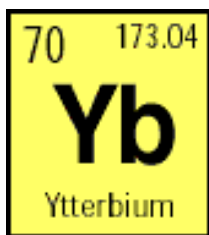
Oxidation State = 3,2 Nearest Neighbor = 0.345 nm

Density = 9.32 g/cm³ Concentration = 3.32 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
61,332.	10,486.	9,978.	8,944.	2,398.	2,173.	1,950.	1,576.	1,528.

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
480.5	388.7	339.7

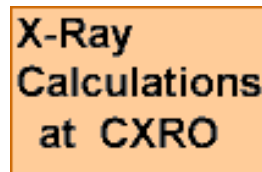
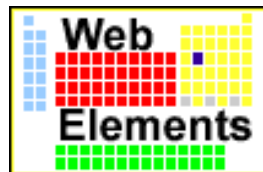
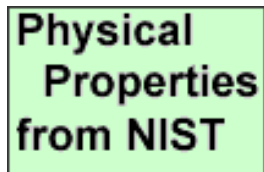
K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁	Mα ₁
52,388.9	51,354.0	59,370.	7,415.6	7,367.3	8,401.8	8,758.8	9,780.1	1,521.4

Electron Configuration = [Xe]4f¹⁴6s²

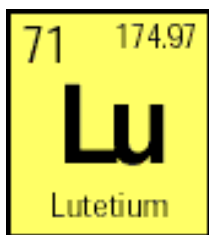
Oxidation State = 3,2 Nearest Neighbor = 0.388 nm

Density = 6.90 g/cm³ Concentration = 2.42 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
63,314.	10,870.	10,349.	9,244.	2,491.	2,264.	2,024.	1,639.	1,589.

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
506.8	412.4	359.2

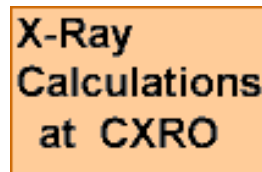
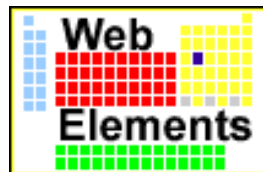
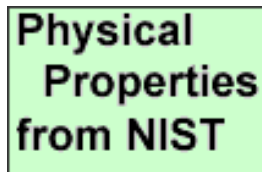
K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁	Mα ₁
54,069.8	52,965.0	61,283.	7,655.5	7,604.9	8,709.0	9,048.9	10,143.4	1,581.3

Electron Configuration = [Xe]4f¹⁴5d¹6s²

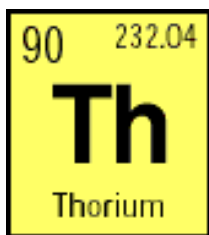
Oxidation State = **3** Nearest Neighbor = 0.343 nm

Density = 9.84 g/cm³ Concentration = 3.39 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
109,651.	20,472.	19,693.	16,300.	5,182.	4,830.	4,046.	3,491.	3,332.

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
1,330.	1,168.	966.4

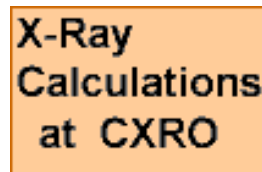
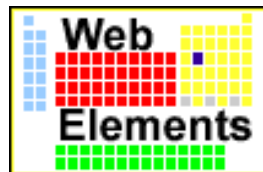
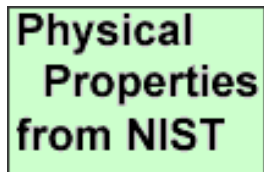
K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁	Mα ₁
93,350.	89,953.	105,609.	12,968.7	12,809.6	16,202.2	15,623.7	18,982.5	2,996.1

Electron Configuration = [Rn]6d²7s²

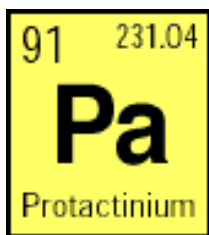
Oxidation State = 4 Nearest Neighbor = 0.360 nm

Density = 11.72 g/cm³ Concentration = 3.04 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
112,601.	21,105.	20,314.	16,733.	5,367.	5,001.	4,174.	3,611.	3,442.

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
1,387.	1,224.	1,007.

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1	M α_1
95,868.	92,287.	108,427.	13,290.7	13,122.2	16,702.	16,024.	19,568.	3,082.3

Electron Configuration = [Rn]5f²6d¹7s²

Oxidation State = 5,4 Nearest Neighbor = 0.321 nm

Density = 15.37 g/cm³ Concentration = 4.01 x 10²² atoms/cm³

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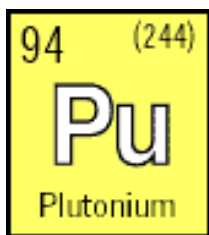
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1	M α_1
103,734.	99,525.	117,228.	14,278.6	14,084.2	18,293.7	17,255.3	21,417.3	—

Electron Configuration = [Rn]5f⁶7s²

Oxidation State = 6,5,4,3

Density = 19.84 g/cm³ Concentration = 4.89 x 10²² atoms/cm³

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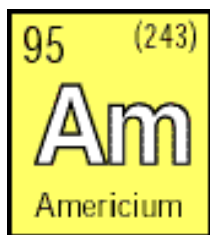
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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1	M α_1
106,472.	102,030.	120,284.	14,617.2	14,411.9	18,852.0	17,676.5	22,065.2	—

Electron Configuration = [Rn]5f⁷7s²

Oxidation State = 6,5,4,**3** Nearest Neighbor = 0.361 nm

Density = 13.69 g/cm³ Concentration = 2.94 x 10²² atoms/cm³

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ELEMENTS
from CHEMSOC**

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Electron Binding Energies in electron volts

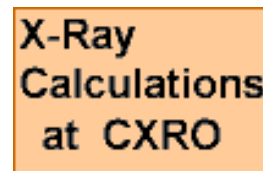
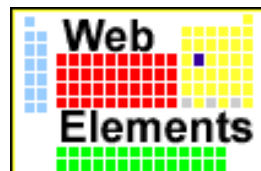
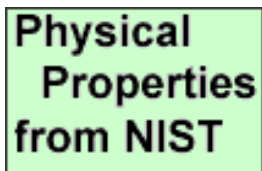
K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
-	-	-	-	-	-	-	-	-

K and L shell emission lines in electron volts

Kα ₁	Kα ₂	Kβ ₁	Lα ₁	Lα ₂	Lβ ₁	Lβ ₂	Lγ ₁	Mα ₁
109,271.	104,590.	126,403.	14,961.	14,746.	19,427.	18,081.	22,735.	-

Electron Configuration = [Rn]5f⁷6d¹7s²

Oxidation State = **3** Density = 13.51 g/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
-	-	-	-	-	-	-	-	-

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1	M α_1
112,121.	107,185.	126,580.	15,308.	15,082.	20,018.	18,509.	23,416.	-

Electron Configuration = [Rn]5f⁹7s²

Oxidation State = 4,3 Density = 14 g/cm³

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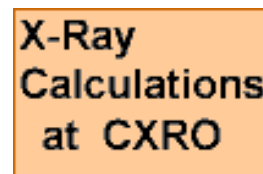
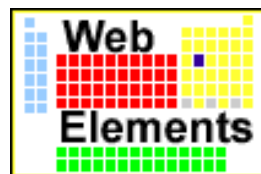
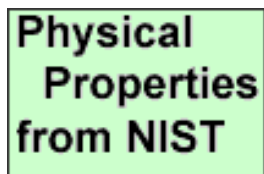
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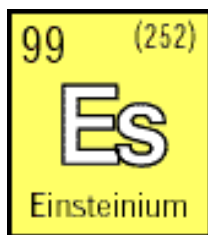
Electron Configuration = $[\text{Rn}]5f^{10}7s^2$

Oxidation State = 3



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Electron Configuration = [Rn]5f¹¹7s²

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ELEMENTS**
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Electron Configuration = [Rn]5f¹²7s²

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Properties
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ELEMENTS
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Electron Configuration = [Rn]5f¹³7s²

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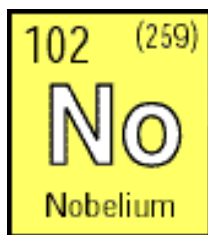
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Electron Configuration = [Rn]5f¹⁴7s²

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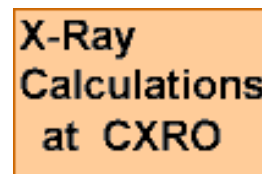
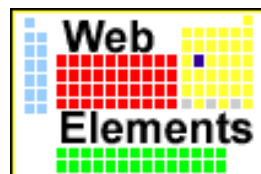
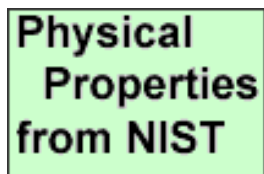
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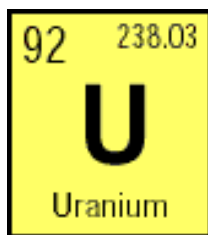


Electron Configuration = $[\text{Rn}]5f^{14}6d^17s^2$



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}
115,606.	21,757.	20,948.	17,166.	5,548.	5,182.	4,303.	3,728.	3,552.

N ₁ 4s	N ₂ 4p _{1/2}	N ₃ 4p _{3/2}
1,439.	1,271.	1,043.

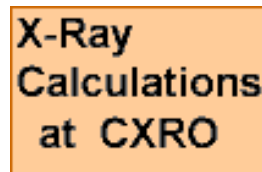
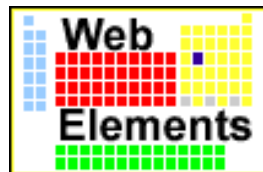
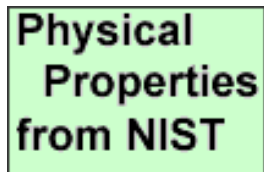
K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1	M α_1
98,439.	94,665.	111,300.	13,614.7	13,438.8	17,220.0	16,428.3	20,167.1	3,170.8

Electron Configuration = [Rn]5f³6d¹7s²

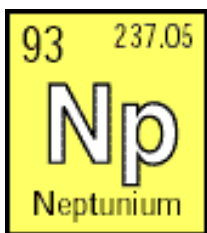
Oxidation State = 6,5,4,3 Nearest Neighbor = 0.275 nm

Density = 18.95 g/cm³ Concentration = 4.82 x 10²² atoms/cm³



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Electron Binding Energies in electron volts

K 1s	L ₁ 2s	L ₂ 2p _{1/2}	L ₃ 2p _{3/2}	M ₁ 3s	M ₂ 3p _{1/2}	M ₃ 3p _{3/2}	M ₄ 3d _{3/2}	M ₅ 3d _{5/2}

K and L shell emission lines in electron volts

K α_1	K α_2	K β_1	L α_1	L α_2	L β_1	L β_2	L γ_1	M α_1
101,059.	97,069.	114,234.	13,944.1	13,759.7	17,750.2	16,840.0	20,784.8	—

Electron Configuration = [Rn]5f⁴6d¹7s²

Oxidation State = 6,5,4,3

Density = 20.25 g/cm³ Concentration = 5.20 x 10²² atoms/cm³

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Mailing Address:

Institution:

Address line 1:

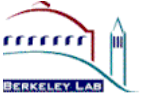
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(There is no charge for the booklet)



2. Synchrotron Radiation

Contents

- [Characteristics of Synchrotron Radiation](#) - *Kwang-Je Kim*
- [History of X-rays and Synchrotron Radiation](#) - *Arthur L. Robinson*
- [Synchrotron Facilities](#) - *Herman Winick*

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Section 2.1 Characteristics of Synchrotron Radiation

Kwang-Je Kim

Synchrotron radiation occurs when a charge moving at relativistic speeds follows a curved trajectory. In this section, formulas and supporting graphs are used to quantitatively describe characteristics of this radiation for the cases of circular motion (bending magnets) and sinusoidal motion (periodic magnetic structures). We will first discuss the ideal case, where the effects due to the angular divergence and the finite size of the electron beam—the emittance effects—can be neglected.

A. BENDING MAGNETS

The angular distribution of radiation emitted by electrons moving through a bending magnet with a circular trajectory in the horizontal plane is given by

$$\frac{d^2 \check{S}_B(\omega)}{d\theta d\psi} = \frac{3\alpha}{4\pi^2} \gamma^2 \frac{\Delta\omega}{\omega} \frac{I}{e} y^2 (1 + X^2)^2 \times \left[K_{2/3}^2(\xi) + \frac{X^2}{1 + X^2} K_{1/3}^2(\xi) \right] \quad , \quad (1)$$

where

\check{S}_B = photon flux (number of photons per second)

θ = observation angle in the horizontal plane

ψ = observation angle in the vertical plane

α = fine-structure constant

γ = electron energy/ $m_e c^2$ (m_e = electron mass, c = velocity of light)

ω = angular frequency of photon ($\varepsilon = \hbar\omega$ = energy of photon)

I = beam current

e = electron charge = 1.602×10^{-19} coulomb

y = $\omega/\omega_c = \varepsilon/\varepsilon_c$

ω_c = critical frequency, defined as the frequency that divides the emitted power into equal halves,

$$= 3\gamma^3 c / 2\rho$$

ρ = radius of instantaneous curvature of the electron trajectory (in practical units,

$$\rho[\text{m}] = 3.3 E[\text{GeV}] / B[\text{T}])$$

E = electron beam energy

B = magnetic field strength

ε_c = $\hbar\omega_c$ (in practical units,

$$\varepsilon_c [\text{keV}] = 0.665 E^2 [\text{GeV}] B[\text{T}])$$

X = $\gamma\psi$

$$\xi = y(1 + X^2)^{3/2} / 2$$

The subscripted K 's are modified Bessel functions of the second kind. In the horizontal direction ($\psi = 0$), Eq. (1) becomes

$$\left. \frac{d^2 \check{S}_B}{d\theta d\psi} \right|_{\psi=0} = \frac{3\alpha}{4\pi^2} \gamma^2 \frac{\hbar\omega}{\omega} \frac{I}{e} H_2(y) \quad , \quad (2)$$

where

$$H_2(y) = y^2 K_{2/3}^2(y/2) \quad . \quad (3)$$

In practical units [$\text{photons} \cdot \text{s}^{-1} \cdot \text{mr}^{-2} \cdot (0.1\% \text{ bandwidth})^{-1}$],

$$\left. \frac{d^2 \check{S}_B}{d\theta d\psi} \right|_{\psi=0} = 1.327 \times 10^{13} E^2 [\text{GeV}] I [\text{A}] H_2(y) \quad .$$

The function $H_2(y)$ is shown in Fig. 2-1.

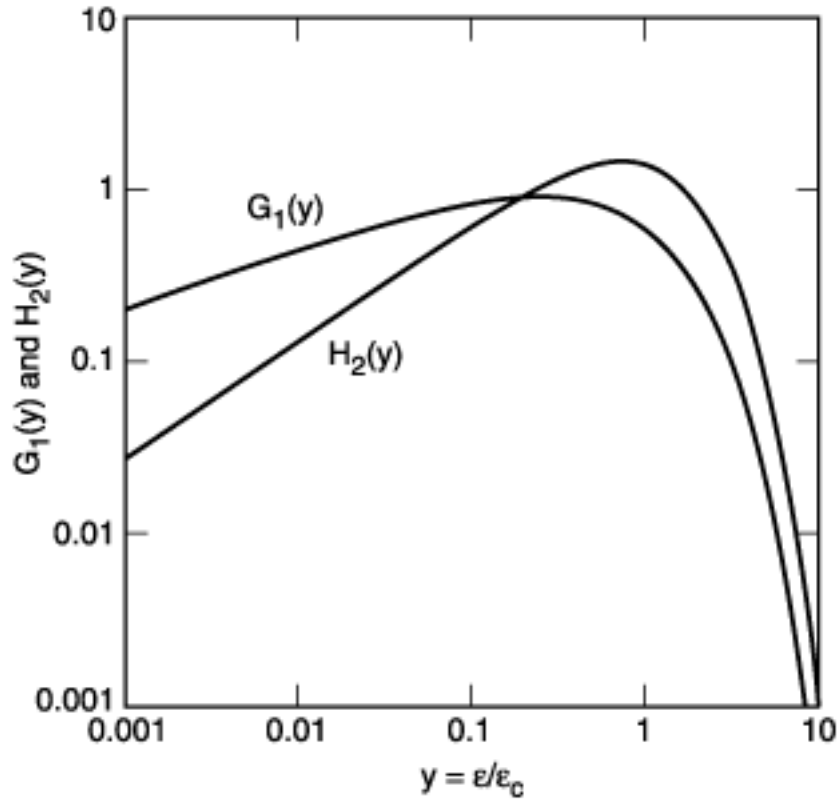


Fig. 2-1. The functions $G_1(y)$ and $H_2(y)$, where y is the ratio of photon energy to critical photon energy.

The distribution integrated over ψ is given by

$$\frac{d\check{S}_B}{d\theta} = \frac{\sqrt{3}}{2\pi} \alpha \gamma \frac{\Delta \omega}{\omega} \frac{I}{e} G_1(y) \quad , \quad (4)$$

where

$$G_1(y) = y \int_y^\infty K_{5/3}(y') dy' \quad .$$

(5)

In practical units [photons·s⁻¹·mr⁻¹·(0.1% bandwidth)⁻¹],

$$\frac{dS_B}{d\theta} = 2.457 \times 10^{13} E[\text{GeV}] I[\text{A}] G_1(y) \quad .$$

The function $G_1(y)$ is also plotted in Fig. 2-1.

Radiation from a bending magnet is linearly polarized when observed in the bending plane. Out of this plane, the polarization is elliptical and can be decomposed into its horizontal and vertical components. The first and second terms in the last bracket of Eq. (1) correspond, respectively, to the intensity of the horizontally and vertically polarized radiation. Figure 2-2 gives the normalized intensities of these two components, as functions of emission angle, for different energies. The square root of the ratio of these intensities is the ratio of the major and minor axes of the polarization ellipse. The sense of the electric field rotation reverses as the vertical observation angle changes from positive to negative.

Synchrotron radiation occurs in a narrow cone of nominal angular width $\sim 1/\gamma$. To provide a more specific measure of this angular width, in terms of electron and photon energies, it is convenient to introduce the effective rms half-angle σ_ψ as follows:

$$\left. \frac{d\check{S}_B}{d\theta} \frac{d^2\check{S}_B}{d\theta d\psi} \right|_{\psi=0} = \sqrt{2\pi} \sigma_\psi \quad ,$$

(6)

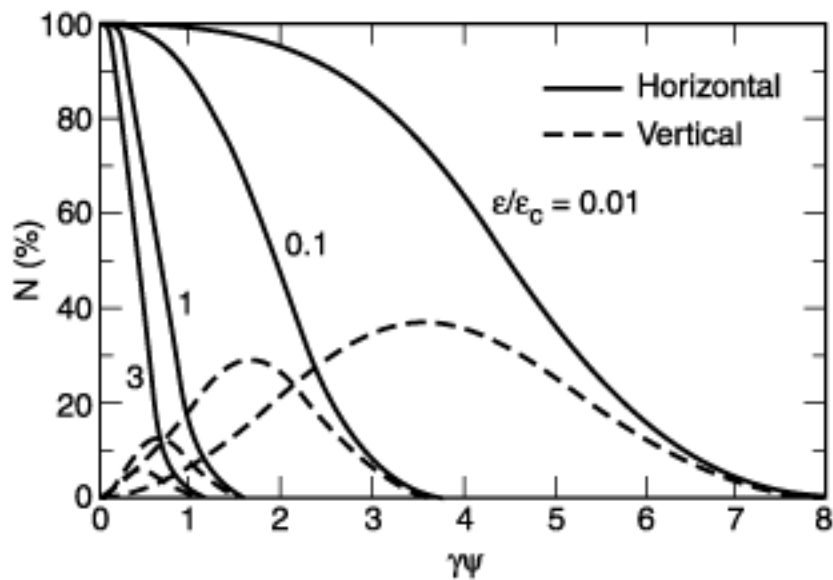


Fig. 2-2. Normalized intensities of horizontal and vertical polarization components, as functions of the vertical observation angle ψ , for different photon energies. (Adapted from Ref. 1.)

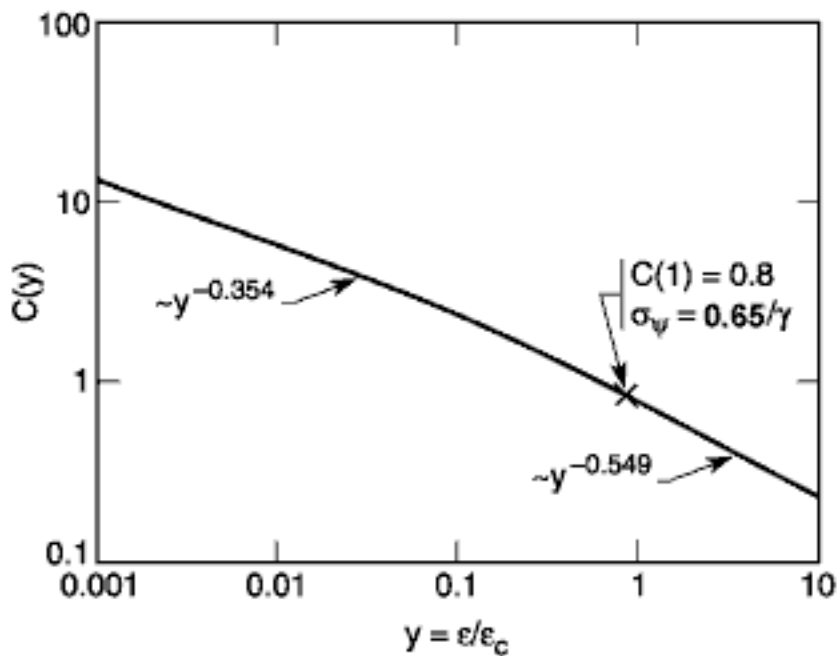


Fig. 2-3. The function $C(y)$. The limiting slopes, for $\epsilon/\epsilon_c \ll 1$ and $\epsilon/\epsilon_c \gg 1$, are indicated.

where σ_ψ is given by

$$\sigma_\psi = \frac{2}{\gamma \sqrt{2\pi}} C(y) = 0.408 \frac{C(y)[\text{mr}]}{E[\text{GeV}]} \quad (7)$$

The function $C(y)$ is plotted in Fig. 2-3. In terms of σ_ψ , Eq. (2) may now be rewritten as

$$\left. \frac{d^2 \check{S}_B}{d\theta d\psi} \right|_{\psi=0} = \frac{d\check{S}_B}{d\theta} \frac{1}{\sigma_\psi \sqrt{2\pi}} \quad (2a)$$

B. PERIODIC MAGNETIC STRUCTURES

In a wiggler or an undulator, electrons travel through a periodic magnetic structure. We consider the case where the magnetic field B varies sinusoidally and is in the vertical direction:

$$B(z) = B_0 \cos(2\pi z/\lambda_u) \quad , \quad (8)$$

where z is the distance along the wiggler axis, B_0 the peak magnetic field, and λ_u the magnetic period. Electron motion is also sinusoidal and lies in the horizontal plane. An important parameter characterizing the electron motion is the deflection parameter K given by

$$K = eB_0\lambda_u / 2\pi mc = 0.934\lambda_u[\text{cm}] B_0[\text{T}] \quad . \quad (9)$$

In terms of K , the maximum angular deflection of the orbit is $\delta = K/\gamma$. For $K \lesssim 1$, radiation from the various periods can exhibit strong interference phenomena, because the angular excursions of the electrons are within the nominal $1/\gamma$ radiation cone; in this case, the structure is referred to as an undulator. In the case $K \gg 1$, interference effects are less important, and the structure is referred to as a

wiggler.

B.1 Wiggler radiation

In a wiggler, K is large (typically $\gtrsim 10$) and radiation from different parts of the electron trajectory adds incoherently. The flux distribution is then given by $2N$ (where N is the number of magnet periods) times the appropriate formula for bending magnets, either Eq. (1) or Eq. (2). However, ρ or B must be taken at the point of the electron's trajectory tangent to the direction of observation. Thus, for a horizontal angle θ ,

$$\varepsilon_c(\theta) = \varepsilon_{c\max} \sqrt{1 - (\theta / \delta)^2} \quad , \quad (10)$$

where

$$\varepsilon_{c\max} = 0.665 E^2[\text{GeV}] B_0[\text{T}] \quad .$$

When $\psi = 0$, the radiation is linearly polarized in the horizontal plane, as in the case of the bending magnet. As ψ increases, the direction of the polarization changes, but because the elliptical polarization from one half-period of the motion combines with the elliptical polarization (of opposite sense of rotation) from the next, the polarization remains linear.

B.2 Undulator radiation

In an undulator, K is moderate ($\lesssim 1$) and radiation from different periods interferes coherently, thus producing sharp peaks at harmonics of the fundamental ($n = 1$). The wavelength of the fundamental on axis ($\theta = \psi = 0$) is given by

$$\lambda_1 = \frac{(1 + K^2/2)}{2\gamma^2} \lambda_u \quad (11)$$

or

$$\lambda_1[\text{\AA}] = \frac{13.056 \lambda_u[\text{cm}]}{E^2[\text{GeV}]} (1 + K^2/2) \quad .$$

The corresponding energy, in practical units, is

$$\varepsilon_1[\text{keV}] = 0.950 \frac{E^2[\text{GeV}]}{(1 + K^2/2)\lambda_u[\text{cm}]} .$$

The relative bandwidth at the n th harmonic is

$$\frac{\Delta\lambda}{\lambda} \cong \frac{\Delta\omega}{\omega} \cong \frac{1}{nN} \quad (n = 1, 2, 3, K) . \quad (12)$$

On axis the peak intensity of the n th harmonic is given by

$$\left. \frac{d^2 \check{S}_n}{d\theta d\psi} \right|_0 = \alpha N^2 \gamma^2 \frac{\Delta\omega}{\omega} \frac{I}{e} F_n(K) \quad (n = 1, 3, 5, K) \\ = 0 \quad (n = 2, 4, 6, K) , \quad (13)$$

where

$$F_n(K) = \frac{K^2 n^2}{(1 + K^2/2)^2} \left\{ J_{\frac{n-1}{2}} \left[\frac{nK^2}{4(1 + K^2/2)} \right] \right. \\ \left. - J_{\frac{n+1}{2}} \left[\frac{nK^2}{4(1 + K^2/2)} \right] \right\}^2 . \quad (14)$$

Here, the J 's are Bessel functions. The function $F_n(K)$ is plotted in Fig. 2-4. In practical units [photons·s⁻¹·mr⁻²·(0.1% band-width)⁻¹], Eq. (13) becomes

$$\left. \frac{d^2 \check{S}_n}{d\theta d\psi} \right|_0 = 1.744 \times 10^{14} N^2 E^2 [\text{GeV}] I[A] F_n(K) \quad .$$

The angular distribution of the n th harmonic is concentrated in a narrow cone whose half-width is given by

$$\sigma_{\gamma'} \cong \sqrt{\frac{\lambda_n}{L}} = \frac{1}{\gamma} \sqrt{\frac{(1 + K^2/2)}{2Nn}} \quad . \quad (15)$$

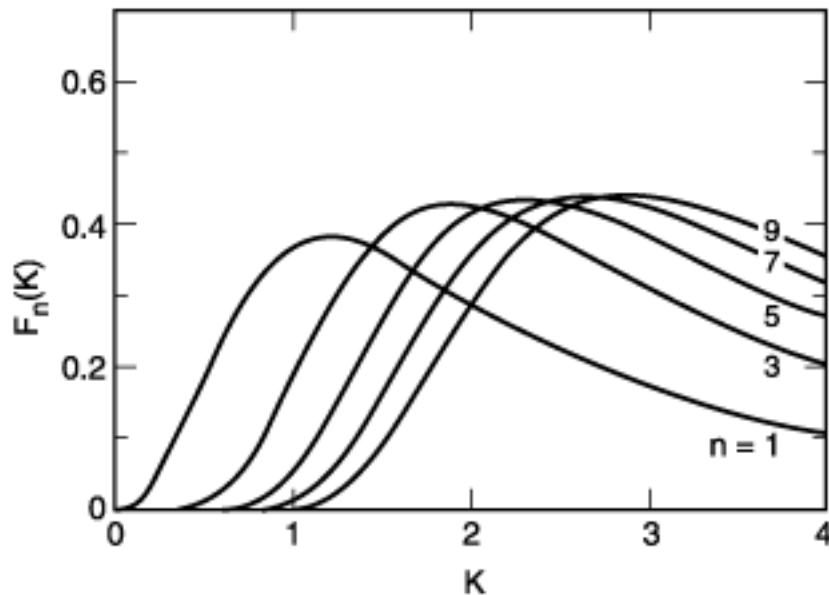


Fig. 2-4. The function $F_n(K)$ for different values of n , where K is the deflection parameter.

Here L is the length of the undulator ($L = N\lambda_u$). Additional rings of radiation of the same frequency also appear at angular distances

$$\theta_{n,1} = \frac{1}{\gamma} \sqrt{\frac{1}{n} (1 + K^2/2)} \quad (n = 1, 2, 3, \dots) \quad (16)$$

The angular structure of undulator radiation is illustrated in Fig. 2-5 for the limiting case of zero beam emittance.

We are usually interested in the central cone. An approximate formula for the flux integrated over the central cone is

$$\check{S}_n = \pi \alpha N \frac{\Delta\omega}{\omega} \frac{I}{e} Q_n(K) \quad (17)$$

or, in units of photons \cdot s $^{-1}$.(0.1% bandwidth) $^{-1}$,

$$\check{S}_n = 1.431 \times 10^{14} N Q_n [A] \quad .$$

The function $Q_n(K) = (1 + K^2/2)F_n/n$ is plotted in Fig. 2-6. Equation (13) can also be written as

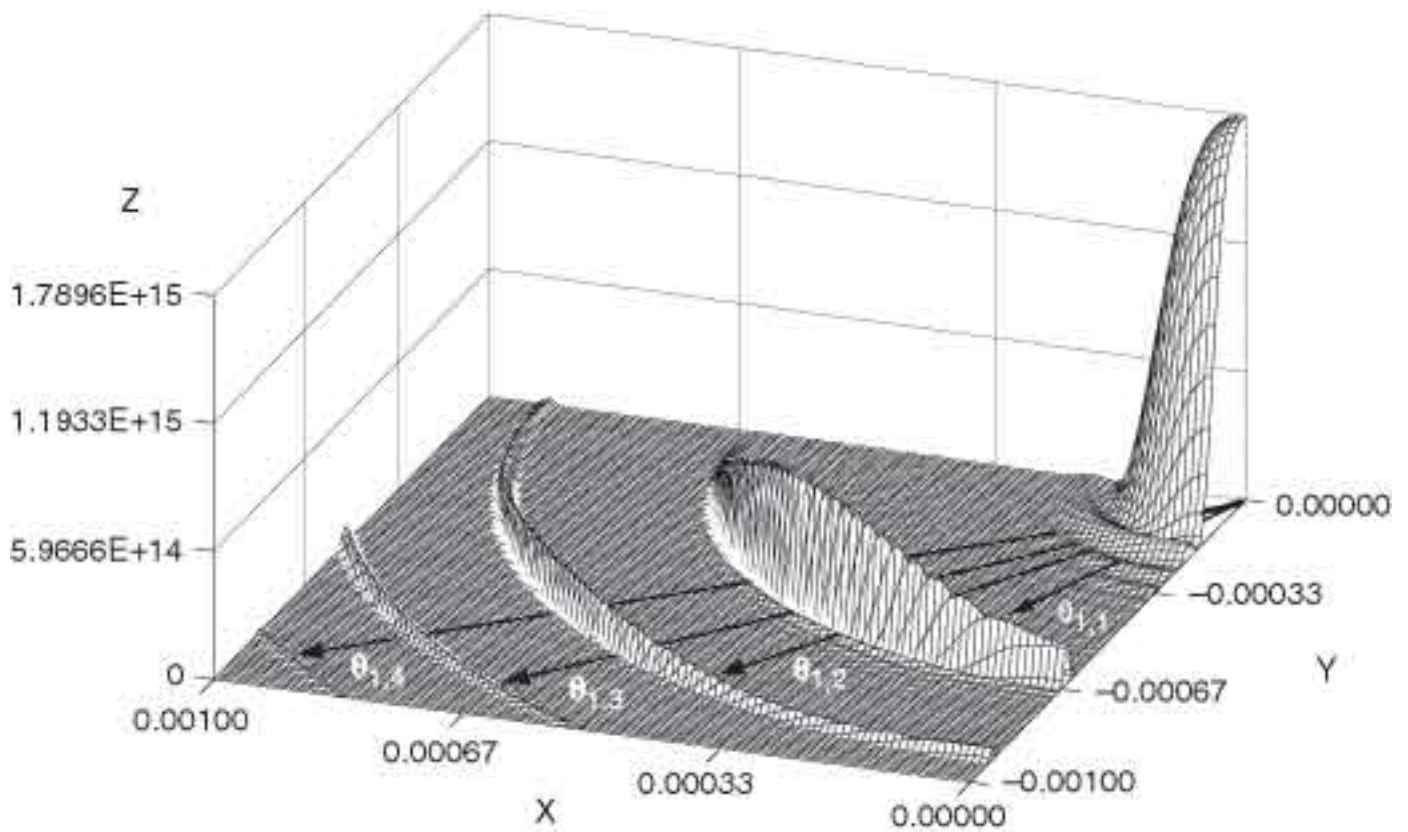


Fig. 2-5. The angular distribution of fundamental ($n = 1$) undulator radiation for the limiting case of zero beam emittance. The x and y axes correspond to the observation angles θ and ψ (in radians), respectively, and the z axis is the intensity in $\text{photons} \cdot \text{s}^{-1} \cdot \text{A}^{-1} \cdot (0.1 \text{ mr})^{-2} \cdot (1\% \text{ bandwidth})^{-1}$. The undulator parameters for this theoretical calculation were $N = 14$, $K = 1.87$, $\lambda_u = 3.5 \text{ cm}$, and $E = 1.3 \text{ GeV}$. (Figure courtesy of R. Tatchyn, Stanford University.)

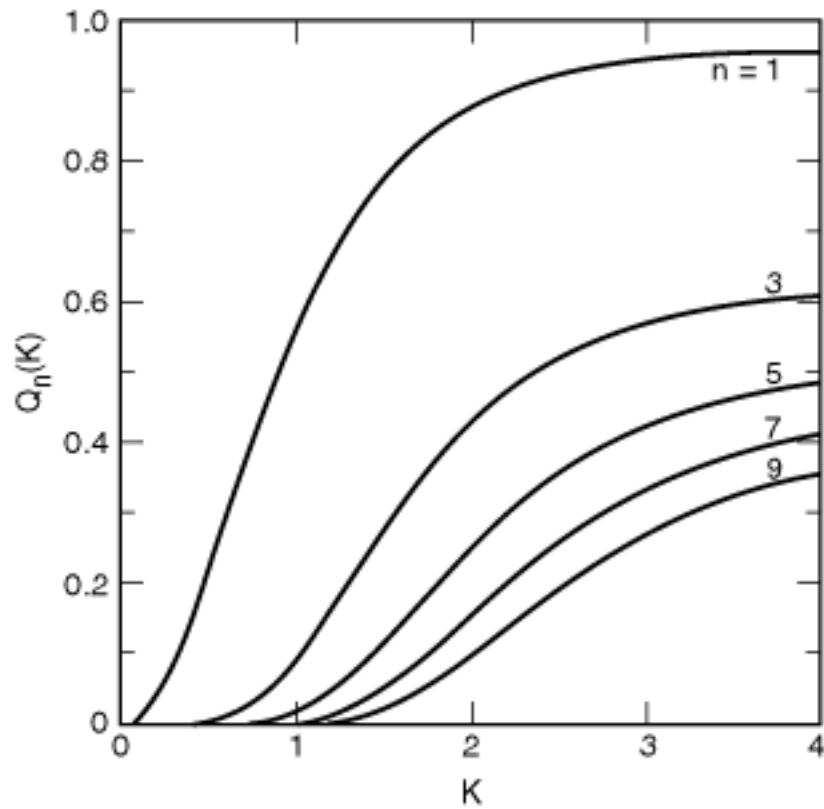


Fig. 2-6. The function $Q_n(K)$ for different values of n .

$$\left. \frac{d^2 \check{S}_n}{d\theta d\psi} \right|_0 = \frac{\check{S}_n}{2\pi\sigma_\gamma^2}, \quad (13a)$$

Away from the axis, there is also a change in wavelength: The factor $(1 + K^2/2)$ in Eq. (11) must be replaced by $[1 + K^2/2 + \gamma^2 (\theta^2 + \psi^2)]$. Because of this wavelength shift with emission angle, the angle-integrated spectrum consists of peaks at λ_n superposed on a continuum. The peak-to-continuum ratio is large for $K \ll 1$, but the continuum increases with K , as one shifts from undulator to wiggler conditions.

B.3 Power

The total power radiated by an undulator or wiggler is

$$P_T = \frac{N}{6} Z_0 I e \frac{2\pi c}{\lambda_u} \gamma^2 K^2 \quad , \quad (18)$$

where $Z_0 = 377$ ohms, or, in practical units,

$$P_T[\text{kW}] = 0.633 E^2[\text{GeV}] B_0^2[\text{T}] L[\text{m}] I[\text{A}] \quad .$$

The angular distribution of the radiated power is

$$\frac{d^2 P}{d\theta d\psi} = P_T \frac{21\gamma^2}{16\pi K} G(K) f_K(\gamma\theta, \gamma\psi) \quad , \quad (19)$$

or, in units of $\text{W}\cdot\text{mr}^{-2}$,

$$\frac{d^2 P}{d\theta d\psi} = 10.84 B_0[\text{T}] E^4[\text{GeV}] I[\text{A}] N G(K) f_K(\gamma\theta, \gamma\psi) \quad .$$

The behavior of the angular function $f_K(\gamma\theta, \gamma\psi)$, which is normalized as $f_K(0,0) = 1$, is shown in Fig. 2-7. The function $G(K)$, shown in Fig. 2-8, quickly approaches unity as K increases from zero.

C. EMITTANCE EFFECTS

Electrons in storage rings are distributed in a finite area of transverse phase space—position \times angle. We introduce the rms beam sizes σ_x (horizontal) and σ_y (vertical), and beam divergences $\sigma_{x'}$ (horizontal) and $\sigma_{y'}$ (vertical). The quantities $\epsilon_x = \sigma_x \sigma_{x'}$ and $\epsilon_y = \sigma_y \sigma_{y'}$ are known as the horizontal and vertical emittances, respectively. In general, owing to the finite emittances of real electron beams, the intensity of the radiation observed in the forward direction is less than that given

by Eqs. (2a) and (13a). Finite emittances can be taken into account approximately by replacing these equations by

$$\left. \frac{d^2 \check{S}_B}{d\theta d\psi} \right|_{\psi=0} = \frac{d\check{S}_B}{d\theta} \frac{1}{\sqrt{2\pi(\sigma_{\psi'}^2 + \sigma_{y'}^2)}} \quad (20)$$

and

$$\left. \frac{d^2 \check{S}_n}{d\theta d\psi} \right|_0 = \frac{\check{S}_n}{2\pi \sqrt{(\sigma_{r'}^2 + \sigma_{x'}^2)(\sigma_{r'}^2 + \sigma_{y'}^2)}} \quad (21)$$

for bends and undulators, respectively. For bending magnets, the electron beam divergence effect is usually negligible in the horizontal plane.

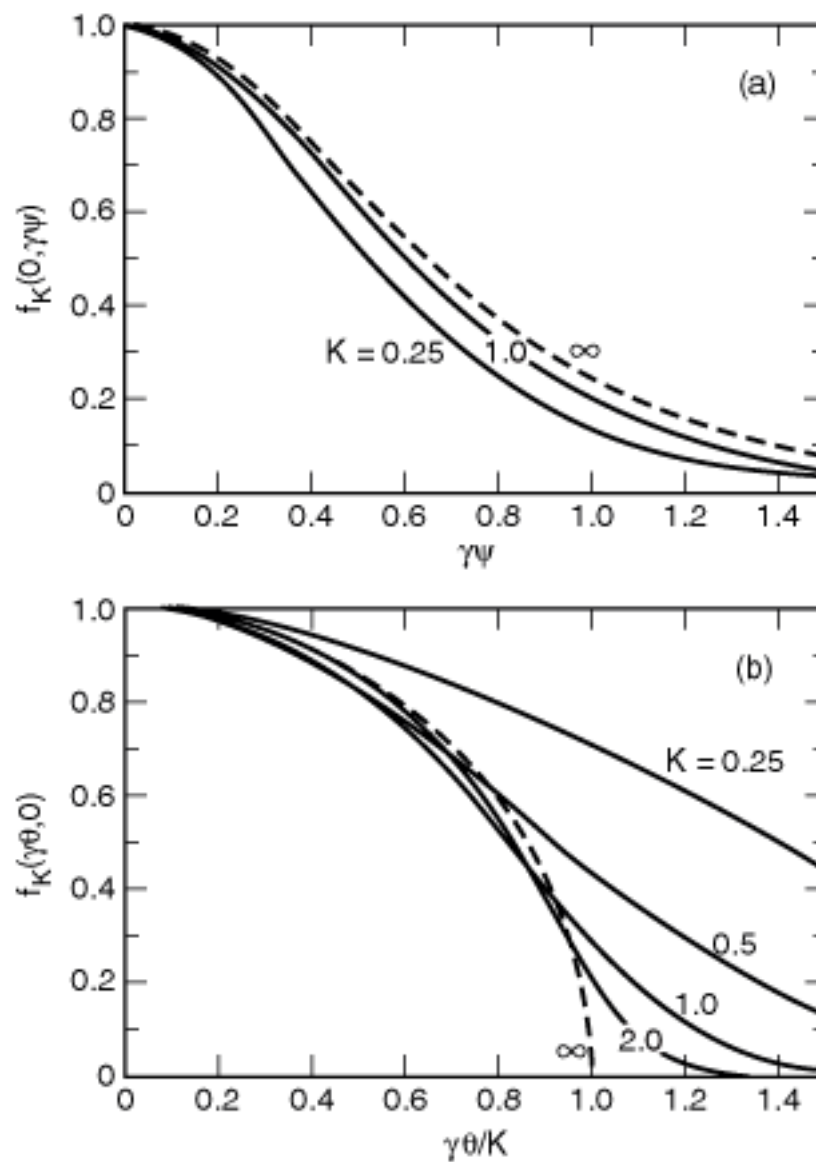


Fig. 2-7. The angular function f_K , for different values of the deflection parameter K , (a) as a function of the vertical observation angle ψ when the horizontal observation angle $\theta = 0$ and (b) as a function of θ when $\psi = 0$.

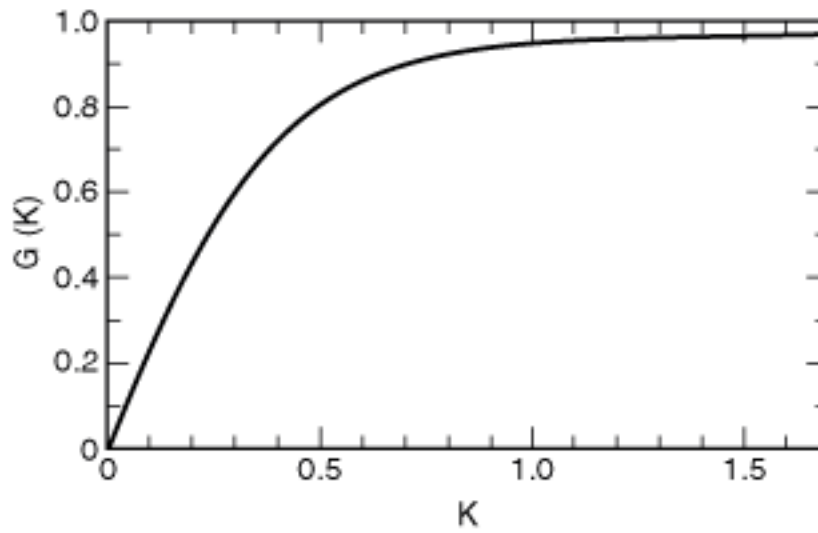


Fig. 2-8. The function $G(K)$.

D. SPECTRAL BRIGHTNESS AND TRANSVERSE COHERENCE

For experiments that require a small angular divergence and a small irradiated area, the relevant figure of merit is the beam brightness B , which is the photon flux per unit phase space volume, often given in units of $\text{photons} \cdot \text{s}^{-1} \cdot \text{mr}^{-2} \cdot \text{mm}^{-2} \cdot (0.1\% \text{ bandwidth})^{-1}$. For an undulator, an approximate formula for the peak brightness is

$$B_n(0, 0) = \frac{S_n}{(2\pi)^2 \sigma_{Tx} \sigma_{Ty} \sigma_{Tx'} \sigma_{Ty'}} \quad , \quad (22)$$

where, for example,

$$\begin{aligned} \sigma_{Tx} &= \sqrt{\sigma_x^2 + \sigma_y^2} \quad , \\ \sigma_{Tx'} &= \sqrt{\sigma_{x'}^2 + \sigma_{y'}^2} \quad , \end{aligned} \quad (23)$$

and where the single-electron radiation from an axially extended source of finite wavelength is described by

$$\begin{aligned}\sigma_r &= \frac{1}{4\pi} \sqrt{\lambda L} \quad , \\ \sigma_{r'} &= \sqrt{\lambda / L} \quad .\end{aligned}\tag{24}$$

Brightness is shown in Fig. 2-9 for several sources of synchrotron radiation, as well as some conventional x-ray sources.

That portion of the flux that is transversely coherent is given by

$$\check{S}_c = B_n \left(\frac{\lambda}{2} \right)^2 = \frac{\check{S}_n \lambda^2}{(4\pi)^2 \sigma_{Tx} \sigma_{Ty} \sigma_{Tx'} \sigma_{Ty'}} \quad .\tag{25}$$

A substantial fraction of undulator flux is thus transversely coherent for a low-emittance beam satisfying $\epsilon_x \epsilon_y \lesssim (\lambda/4\pi)^2$.

E. LONGITUDINAL COHERENCE

Longitudinal coherence is described in terms of a coherence length

$$l_c = \lambda^2 / \Delta\lambda \quad .$$

For an undulator, the various harmonics have a natural spectral purity of $\Delta\lambda/\lambda = 1/nN$ [see Eq. (12)]; thus, the coherence length is given by

$$l_c = nN\lambda \quad ,\tag{27}$$

which corresponds to the relativistically contracted length of the undulator. Thus, undulator radiation from low-emittance electron beams [$\epsilon_x \epsilon_y \lesssim (\lambda/4\pi)^2$] is

transversely coherent and is longitudinally coherent within a distance described by Eq. (27). In the case of finite beam emittance or finite angular acceptance, the longitudinal coherence is reduced because of the change in wavelength with emission angle. In this sense, undulator radiation is partially coherent. Transverse and longitudinal coherence can be enhanced when necessary by the use of spatial and spectral filtering (i.e., by use of apertures and monochromators, respectively).

The references listed below provide more detail on the characteristics of synchrotron radiation.

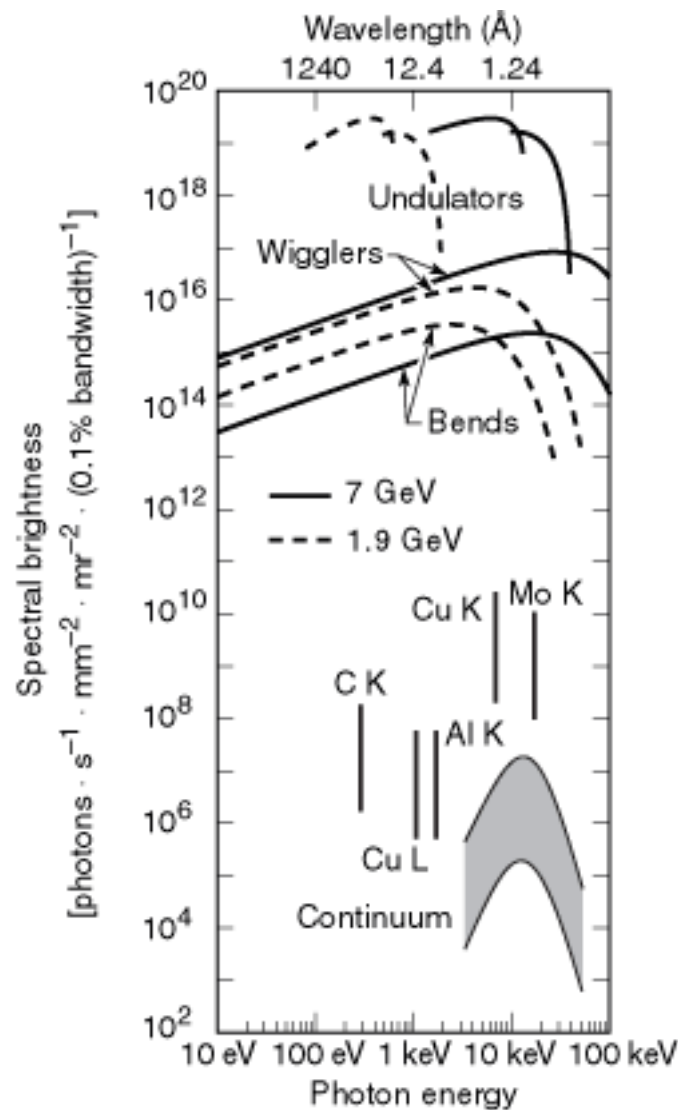


Fig. 2-9. Spectral brightness for several synchrotron radiation sources and conventional x-ray sources. The data for conventional x-ray tubes should be taken as rough estimates only, since brightness depends strongly on such

parameters as operating voltage and take-off angle. The indicated two-order-of-magnitude ranges show the approximate variation that can be expected among stationary-anode tubes (lower end of range), rotating-anode tubes (middle), and rotating-anode tubes with microfocusing (upper end of range).

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X-Ray Data Booklet

Section 2.2 HISTORY of SYNCHROTRON RADIATION

Arthur L. Robinson

(This is an expanded version of the section published in the booklet)

Although natural synchrotron radiation from charged particles spiraling around magnetic-field lines in space is as old as the stars—for example the light we see from the Crab Nebula—short-wavelength synchrotron radiation generated by relativistic electrons in circular accelerators is only a half-century old. The first observation, literally since it was visible light that was seen, came at the General Electric Research Laboratory in Schenectady, New York, on April 24, 1947. In the 50 years since, synchrotron radiation has become a premier research tool for the study of matter in all its varied manifestations, as facilities around the world constantly evolved to provide this light in ever more useful forms.

A. X-RAY BACKGROUND

From the time of their discovery in 1895, both scientists and society have recognized the exceptional importance of x rays, beginning with the awarding of the very first Nobel Prize in Physics in 1901 to Röntgen "in recognition of the extraordinary services he has rendered by the discovery of the remarkable rays subsequently named after him." By the time synchrotron radiation was observed almost a half-century later, the scientific use of x rays was well established. Some highlights include

- 1909: Barkla and Sadler discover characteristic x-ray radiation (1917 Nobel Prize to Barkla)
- 1912: von Laue, Friedrich, and Knipping observe x-ray diffraction (1914 Nobel Prize to von Laue)
- 1913: Bragg, father and son, build an x-ray spectrometer (1915 Nobel Prize)

- 1913: Moseley develops quantitative x-ray spectroscopy and Moseley's Law
- 1916: Siegbahn and Stenstrom observe emission satellites (1924 Nobel Prize to Siegbahn)
- 1921: Wentzel observes two-electron excitations
- 1922: Meitner discovers Auger electrons
- 1924: Lindh and Lundquist resolve chemical shifts
- 1927: Coster and Druyvesteyn observe valence-core multiplets
- 1931: Johann develops bent-crystal spectroscopy

B. EARLY HISTORY

The theoretical basis for synchrotron radiation traces back to the time of Thomson's discovery of the electron. In 1897, Larmor derived an expression from classical electrodynamics for the instantaneous total power radiated by an accelerated charged particle. The following year, Liénard extended this result to the case of a relativistic particle undergoing centripetal acceleration in a circular trajectory. Liénard's formula showed the radiated power to be proportional to $(E/mc^2)^4/R^2$, where E is particle energy, m is the rest mass, and R is the radius of the trajectory. A decade later in 1907, Schott reported his attempt to explain the discrete nature of atomic spectra by treating the motion of a relativistic electron in a homogeneous magnetic field. In so doing, he obtained expressions for the angular distribution of the radiation as a function of the harmonic of the orbital frequency. Schott wrote a book-length essay on the subject in 1912.

When attempts to understand atomic structure took a different turn with the work of Bohr, attention to radiation from circulating electrons waned. Interest in the radiation as an energy-loss mechanism was reawakened in the 1920s after

physicists began contemplating magnetic-induction electron accelerators (betatrons) as machines to produce intense beams of x rays by directing the accelerated beam to a suitable target. (Commercial betatrons are still used as x-ray sources.) The first betatron to operate successfully was a 2.3-MeV device built in 1940 by Kerst at the University of Illinois, followed at GE by a 20-MeV and then a 100-MeV machine to produce high-energy x rays for nuclear research. Meanwhile, in the Soviet Union, Ivanenko and Pomeranchuk published their 1944 calculations showing that energy losses due to radiating electrons would set a limit on the energy obtainable in a betatron, which they estimated to be around 0.5 GeV.

Subsequent theoretical work proceeded independently by Pomeranchuk and others in the Soviet Union and in the U. S., where by 1945 Schwinger had worked out in considerable detail the classical (i.e., non-quantum) theory of radiation from accelerated relativistic electrons. Major features demonstrated for the case of circular trajectories included the warping of the globular non-relativistic dipole radiation pattern into the strongly forward peaked distribution that gives synchrotron radiation its highly collimated property and the shift of the spectrum of the radiation to higher photon energies (higher harmonics of the orbital frequency) as the electron energy increased, with the photon energy at the peak of the distribution varying as E^3/R . Schwinger did not publish his complete findings until 1949, but he made them available to interested parties. Quantum calculations, such as those by Sokolov and Tersov in the Soviet Union, later confirmed the classical results for electron energies below about 10 TeV.

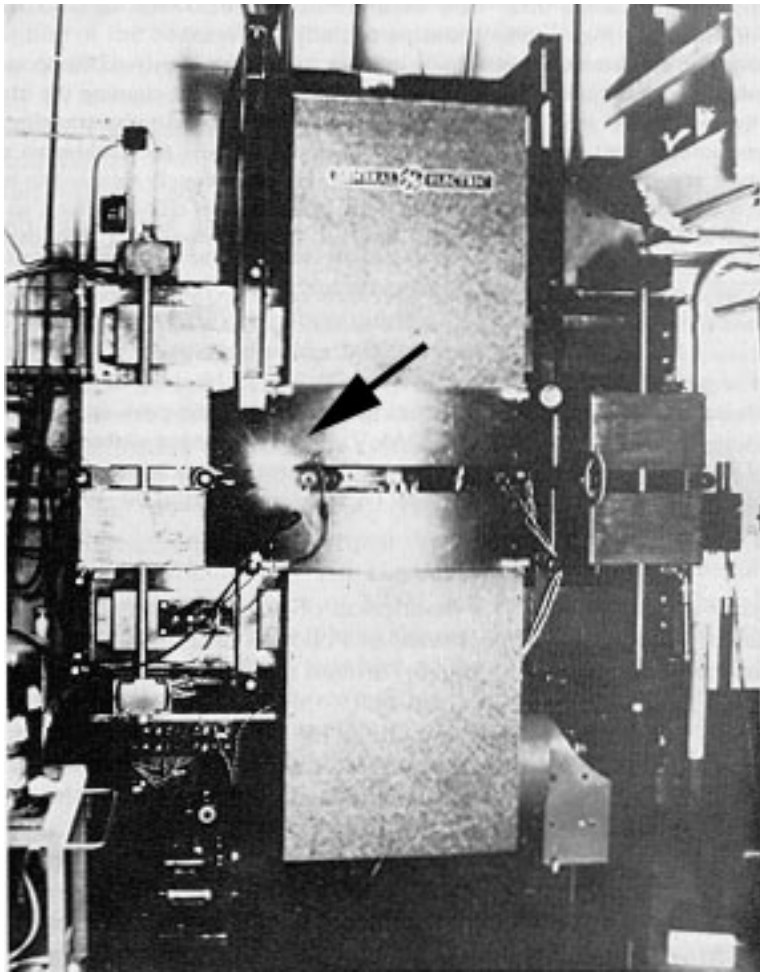
C. DISCOVERY OF SYNCHROTRON RADIATION

After testing of GE's 100-MeV betatron commenced in 1944, Blewett suggested a search for the radiation losses, which he expected from the work of Ivanenko and Pomeranchuk to be significant at this energy. However, two factors prevented success: whereas, according to Schwinger's calculations, the radiation spectrum for the 100-MeV betatron should peak in the near- infrared/visible range, the search took place in the radio and microwave regions at the orbital frequency (and

low harmonics) and the tube in which the electrons circulated was opaque. Although quantitative measurements reported in 1946 of the electron-orbit radius as it shrunk with energy were in accord with predicted losses, there was also another proposed explanation with the result that, while Blewett remained convinced the losses were due to synchrotron radiation, his colleagues were not.

Advances on another accelerator front led to the 1947 visual observation of synchrotron radiation at GE. The mass of particles in a cyclotron grows as the energy increases into the relativistic range. The heavier particles then arrive too late at the electrodes for a radio-frequency (RF) voltage of fixed frequency to accelerate them, thereby limiting the maximum particle energy. To deal with this problem, in 1945 McMillan in the U. S. and Veksler in the Soviet Union independently proposed decreasing the frequency of the RF voltage as the energy increases to keep the voltage and the particle in synch. This was a specific application of their phase-stability principle for RF accelerators, which explains how particles that are too fast get less acceleration and slow down relative to their companions while particles that are too slow get more and speed up, thereby resulting in a stable bunch of particles that are accelerated together.

At GE, Pollack got permission to assemble a team to build a 70-MeV electron synchrotron to test the idea. Fortunately for the future of synchrotron radiation, the machine was not fully shielded and the coating on the doughnut-shaped electron tube was transparent, which allowed a technician to look around the shielding with a large mirror to check for sparking in the tube. Instead, he saw a bright arc of light, which the GE group quickly realized was actually coming from the electron beam. Langmuir is credited as recognizing it as synchrotron radiation or, as he called it, "Schwinger radiation." Subsequent measurements by the GE group began the experimental establishment of its spectral and polarization properties. Characterization measurements were also carried out in the 1950s at a 250-MeV synchrotron at the Lebedev Institute in Moscow.



Synchrotron light from the 70-MeV electron synchrotron at GE.

The next step came with the 1956 experiments of Tombouliau and Hartman, who were granted a two-week run at the 320-MeV electron synchrotron at Cornell. Despite the limited time, they were able to confirm the spectral and angular distribution of the radiation with a grazing-incidence spectrograph in the ultraviolet from 80 Å to 300 Å. They also reported the first soft x-ray spectroscopy experiments with synchrotron radiation, measuring the transmission of beryllium and aluminum foils near the K and L edges. However, despite the advantages of synchrotron radiation that were detailed by the Cornell scientists and the interest their work stimulated, it wasn't until 1961 that an experimental program using synchrotron radiation got under way when the National Bureau of Standards (now National Institute of Standards and Technology) modified its 180-MeV electron synchrotron to allow access to the radiation via a tangent section into the machine's vacuum system.

D. THE FIRST GENERATION: PARASITIC OPERATION

Under Madden and Codling, measurements began at the new NBS facility (Synchrotron Ultraviolet Radiation Facility or SURF) to determine the potential of synchrotron radiation for standards and as a source for spectroscopy in the ultraviolet (the wavelength for peak radiated power per unit wavelength was 335 Å). Absorption spectra of noble gases revealed a large number of previously unobserved resonances due to inner-shell and two-electron excitations, including doubly excited helium, which remains today a prime test bed for studying electron-electron correlations. These findings further stimulated the growing interest in synchrotron radiation. Establishment of SURF began the first generation of synchrotron-radiation facilities, sometimes also called parasitic facilities because the accelerators were built and usually operated primarily for high-energy or nuclear physics. However, the NBS synchrotron had outlived its usefulness for nuclear physics and was no longer used for this purpose.

If SURF headed the first generation, it was not by much, as activity was also blossoming in both Europe and Asia. At the Frascati laboratory near Rome, researchers began measuring absorption in thin metal films using a 1.15-GeV synchrotron. In 1962, scientists in Tokyo formed the INS- SOR (Institute for Nuclear Studies-Synchrotron Orbital Radiation) group and by 1965 were making measurements of soft x-ray absorption spectra of solids using light from a 750-MeV synchrotron. The trend toward higher energy and shorter wavelengths took a big leap with the use of the 6-GeV Deutsches Elektronen-Synchrotron (DESY) in Hamburg, which began operating for both high-energy physics and synchrotron radiation in 1964. With synchrotron radiation available at wavelengths in the x-ray region down to 0.1 Å, experimenters at DESY were able to carefully check the spectral distribution against Schwinger's theory, as well as begin absorption measurements of metals and alkali halides and of photoemission in aluminum.

E. THE FIRST GENERATION: STORAGE RINGS

While the number of synchrotrons with budding synchrotron-radiation facilities

was growing, the next major advance was the development of electron storage rings, the basis for all of today's synchrotron sources. In the 1950s, the Midwest Universities Research Association (MURA) was formed to develop a proposal for a high-current accelerator for particle physics. As part of the project, Mills and Rowe designed a 240-MeV storage ring, then a new idea, as a test bed for advanced accelerator concepts. Politics intervened, however, and the decision was made in 1963 to build a new high-energy accelerator in Illinois, a facility that became the Fermi National Accelerator Laboratory. With this decision, MURA eventually dissolved, but in the meantime construction of the storage ring proceeded.

Thanks to the rapidly swelling interest in synchrotron radiation for solid-state research that stimulated a 1965 study by the U. S. National Research Council documenting this promise, MURA agreed to alterations in the storage-ring vacuum chamber that would provide access to synchrotron radiation without interfering with the accelerator studies. With MURA's demise in 1967, funding for the original purpose of the storage ring also disappeared, but supported by the U. S. Air Force Office of Scientific Research, the University of Wisconsin took on the responsibility of completing the storage ring, known as Tantalus I, and operating it for synchrotron-radiation research. The first spectrum was measured in 1968. In subsequent years, improvements enabled Tantalus I to reach its peak performance, add a full complement of ten beamlines with monochromators, and become in many respects a model for today's multi-user synchrotron-radiation facilities.

With Tantalus I, the superiority of the electron storage ring as a source of synchrotron radiation became evident. In a storage ring, the beam continuously circulates current at a fixed energy for periods up to many hours, whereas the synchrotron beam undergoes a repeated sequence of injection, acceleration, and extraction at rates up to 50 Hz. Among the advantages stemming from this feature are a much higher "duty cycle" when the beam is available, higher beam currents and hence higher fluxes of radiation, a synchrotron-radiation spectrum that does not change with time, greater beam stability, and a reduced radiation hazard.

A surge of interest in storage rings soon followed. In 1971, synchrotron-radiation work began on the 540-MeV ACO storage ring at the Orsay laboratory in France. With the help of the Wisconsin group, the NBS converted its synchrotron into a 250-MeV storage ring (SURF II) in 1974. The same year the INS-SOR group (now part of the Institute for Solid State Physics) in Tokyo began commissioning a 300-MeV storage ring, generally considered the first machine designed from the start specifically for the production of synchrotron radiation. The first storage ring in the multi- GeV class to provide x rays to a large community of synchrotron-radiation users was the 2.5-GeV SPEAR ring at the Stanford Linear Accelerator Center (SLAC), where a beamline with five experimental stations was added in 1974 under the auspices of the Stanford Synchrotron Radiation Project. Other large storage rings to which synchrotron-radiation capabilities were added early on include DORIS at the DESY laboratory, VEPP-3 at the Institute for Nuclear Physics in Novosibirsk, DCI at Orsay, and CESR at Cornell (the CHESS facility).

F. THE SECOND GENERATION: DEDICATED SOURCES

The larger storage rings just cited were electron-positron colliding-beam machines that were operated to provide the highest possible collision rates without blowing up the beams, a condition that generally meant low beam currents. Moreover, while studying the then-fashionable J/ψ particle and its relatives, they often ran at low beam energies. Under these conditions, parasitic operation meant a severely limited output of synchrotron radiation, thereby motivating a clamor for storage rings designed for and dedicated to the production of synchrotron radiation. The Synchrotron Radiation Source (SRS) at the Daresbury Laboratory in the UK was the first fruit of this movement. Synchrotron-radiation research had begun at Daresbury around 1970 with the addition of a beamline to the 5-GeV NINA electron synchrotron. When NINA shut down in 1977, a plan was already approved to build a 2-GeV electron storage ring at the same site expressly for synchrotron radiation. Experiments began at the new facility in 1981.

In the U. S., after a 1976 National Research Council study documented an increasing imbalance between demand for synchrotron radiation and its availability, construction of the National Synchrotron Light Source (NSLS) at the Brookhaven National Laboratory was approved. With construction completed in 1981, the NSALS complex included separate 700-MeV and 2.5-GeV storage rings for production of UV and x rays, respectively. During this same period, the University of Wisconsin Synchrotron Radiation Center built a new 1-GeV storage ring named Aladdin, which replaced the old Tantalus I (part of which was later sent to the Smithsonian for eventual exhibit). In Japan, the Photon Factory was completed in 1982 at the KEK laboratory in Tsukuba. And in Berlin, the BESSY facility began serving users in 1982 with a 0.8-GeV storage ring. And at Orsay, LURE (Laboratoire pour l'Utilisation du Rayonnement Electromagnétique) began operating an 800-MeV storage ring, SuperACO, in 1984.

Elsewhere, some of the first-generation facilities gradually evolved toward second-generation status by means of upgrades and agreements with laboratory managements to dedicate a fraction and sometimes all of the yearly machine operations to synchrotron radiation as the high-energy physics frontier advanced. The Stanford Synchrotron Radiation Laboratory at SLAC and HASYLAB (Hamburger Synchrotronstrahlungslabor) at DESY are prime examples. All of these second-generation facilities provide fine examples of the productivity of a dedicated source of synchrotron radiation. Over the years, for example, the SRS has grown to about 40 experimental stations serving around 4000 users from physics and biology to engineering, and the NSLS has around 80 operating beamlines and more than 2200 users each year. (By this time, the number of synchrotron-radiation facilities has grown too large to mention them all here; the reader should turn to Chapter 8 for a list of current facilities and the spectral ranges they serve.)

Major experimental developments included a major enhancement of photoemission for studying the electronic structure of solids and surfaces (for example, angle-resolved photoemission began at Tantalus), the development of

extended x-ray absorption fine-structure spectroscopy (EXAFS) for the measurement of local atomic structure (which got its start at SSRL), and the extension of high-resolution protein crystallography to small, difficult to grow, or otherwise unstable samples (beginning with the work at SSRL and expanding rapidly to DESY, CHESS, and elsewhere).

G. BRIGHTNESS

As the clamor for facilities dedicated to synchrotron radiation expanded in the 1970's, users increasingly appreciated that spectral brightness or brilliance (the flux per unit area of the radiation source per unit solid angle of the radiation cone per unit spectral bandwidth) was often more important than flux alone for many experiments. For example, since the photon beam is most often ribbon shaped with a larger horizontal than vertical size, typical spectroscopy experiments at synchrotron facilities use monochromators with horizontal slits. Because spectroscopy experiments achieve the highest spectral resolution when the slits are narrowed, obtaining a useful flux through the monochromator exit slits requires that the photon beam have a small vertical size and angular divergence so that most of the flux from the source can pass through the narrowed entrance slits and strike the dispersing element at nearly the same angle (i.e., when the vertical brightness is high). Crystallography experiments, especially those with small crystals and large unit cells, also place a premium on brightness, since it is necessary to match the incident beam to the crystal size while maintaining sufficient angular resolution to resolve closely spaced diffraction spots.

Brightness (flux density in phase space) is an invariant quantity in statistical mechanics, so that no optical technique can improve it. For example, focusing the beam to a smaller size necessarily increases the beam divergence, and vice-versa; apertures can help reduce beam size and divergence but only at the expense of flux. The cure therefore is proper design of the source, the electron beam in the storage ring. The size and divergence of the electron beam are determined by the storage-ring lattice—the arrangement and strengths of the dipole, quadrupole, and sextupole magnets. As planning of the NSLS progressed, Chasman and Green

designed what has become the prototype lattice (a so-called double-bend achromat) for storage rings with a low emittance (product of beam size and divergence) and hence a light source with high brightness. The Chasman-Green lattice and variations are the basis for most of today's synchrotron sources.

H. INSERTION DEVICES

Undulators provide a way to take maximum advantage of the intrinsic brightness of the synchrotron-radiation source. The magnetic structure of today's most common (planar) undulator is an array of closely spaced vertically oriented dipole magnets of alternating polarity. As the electron beam passes longitudinally through the array, its trajectory oscillates in the horizontal plane. Owing to the relatively weak field, the radiation cones emitted at each bend in the trajectory overlap, giving rise to a constructive interference effect that results in one or a few spectrally narrow peaks (a fundamental and harmonics) in a beam that is highly collimated in both the horizontal and vertical directions; that is, the beam has a high spectral brightness (see Chapter 4). Tuning the wavelengths of the harmonics is by means of mechanically adjusting the vertical spacing (gap) between the pole tips.

The undulator concept traces back to the 1947 theoretical work Ginzburg in the Soviet Union. Motz and coworkers experimentally verified the idea in 1953 by building an undulator and using it to produce radiation from the millimeter-wave to the visible range in experiments with a linear accelerator at Stanford University. The next step came in the 1970's with the installation of undulators in storage rings at the Lebedev Institute in Moscow and the Tomsk Polytechnic Institute. Measurements at these laboratories began to provide the information needed for a comprehensive description of undulator radiation. Owing to the large number of closely spaced dipoles, electromagnets or superconducting magnets are not that best choice for undulators, which became practical devices for producing synchrotron radiation in storage rings in 1981 when Halbach at Lawrence Berkeley Laboratory and coworkers constructed a device based on permanent magnets and successfully tested it at SSRL. Parallel work was also under way in

Novosibirsk.



Klaus Halbach shown in 1986 with Kwang-Je Kim discussing a model of an undulator that Halbach designed.

Wigglers are similar to undulators but generally have higher fields and fewer dipoles, with the result that they produce a continuous spectrum with a higher flux and a spectrum that extends to shorter wavelengths than bend magnets. Despite the similarity, wigglers evolved independently from undulators at the start. A decade after the initial suggestion by Robinson, a wiggler was installed in 1966 at the Cambridge Electron Accelerator (a 3-GeV storage ring that was actually the first multi-GeV storage ring to produce x rays before it was shut down in 1972) to enhance beam storage. In 1979, a wiggler comprising just seven electromagnet poles at SSRL was the first to be used for producing synchrotron radiation. Nowadays, wigglers may be permanent-magnet devices following the Halbach design or be based on high-field superconductors that shift the spectrum to the shortest wavelengths. Together, wigglers and undulators are called insertion devices because they are placed in one of the generally empty straight sections that connect the curved arcs of large storage rings, where the magnets that guide and focus the electron beam reside.

The planar insertion devices just described produce radiation that is linearly polarized in the horizontal plane. However, a feature available from bend-magnet sources is the generation of elliptically polarized radiation, with the most obvious applications in the study of magnetic materials. The radiation from a bend magnet is elliptically polarized above and below the horizontal plane of the electron-beam orbit, and this feature has now been exploited at many facilities, including the pioneering work on magnetic materials in the hard x-ray region by Schütz and coworkers at HASYLAB (1987) and in the soft x-ray region by Chen and colleagues at the NSLS (1990). Now, among several designs for both wigglers and undulators that produce elliptically polarized synchrotron radiation, some have been implemented, tested, and are in regular use.

I. THE THIRD GENERATION: OPTIMIZED FOR BRIGHTNESS

Both undulators and wigglers have been retrofitted into older storage rings, and in some cases, the second-generation rings, such as those at the NSLS, were designed with the possibility of incorporating insertion devices. Nonetheless, even before the second-generation facilities were broken in, synchrotron users recognized that a new generation of storage rings with a still lower emittance and long straight sections for undulators would permit achieving even higher brightness and with it, a considerable degree of spatial coherence. Beneficiaries of high brightness would include those who need spatially resolved information, ranging from x-ray microscopy to spectromicroscopy (the combination of spectroscopy and microscopy) and those who need temporal resolution, as well as spectroscopists, crystallographers, and anyone who needs to collect higher resolution data faster.

Construction of third-generation synchrotron-radiation facilities brings us to the present day. Following the NSLS two-ring model, third-generation facilities specialize in either short- wavelength (high-energy or hard) x rays or vacuum-ultraviolet and long-wavelength (low energy or soft) x rays. The range in between (intermediate-energy x rays) is accessible by both. The European Synchrotron Radiation Facility (ESRF) in Grenoble was the first of the third-generation hard x-

ray sources to operate, coming on line for experiments by users with a 6-GeV storage ring and a partial complement of commissioned beamlines in 1994. The ESRF has been followed by the Advanced Photon Source at Argonne National Laboratory (7 GeV) in late 1996, and SPring-8 (8 GeV) in Harima Science Garden City in Japan in late 1997. These machines are physically large (850 to 1440 meters in circumference) with a capability for 30 or more insertion-device, and a comparable number of bend-magnet, beamlines.

Among the long-wavelength sources, the Advanced Light Source at Berkeley (1.9 GeV) began its scientific program in early 1994, as did the Sincrotrone Trieste (2.0 GeV) in Italy, followed by the Synchrotron Radiation Research Center (1.3 GeV) in Hsinchu, Taiwan, and the Pohang Light Source (2.0 GeV) in Pohang, Korea. These physically smaller machines (120 to 280 meters in circumference) have fewer straight sections and therefore can service fewer insertion-device beamlines than the larger machines, but since they are also less expensive, many more of them have been and are being constructed around the world, from Canada in North America; to Brazil in South America; to Japan, China, Thailand, and India in Asia; and to Sweden, Germany, Switzerland, and other European countries, although in some cases, these are not truly third-generation machines in terms of performance. Addition of superconducting bend magnets to the storage-ring lattice in these smaller machines, as some facilities are planning to do, allows them to extend their spectral coverage to higher photon energies without sacrificing their performance at lower photon energies.

J. NEXT: THE FOURTH GENERATION

The race to develop a new generation of synchrotron radiation sources with vastly enhanced performance has already begun, even as the third-generation facilities enter their prime, which takes us past the present into the future; namely, to the fourth generation. The candidate with the best scientific case for a fourth-generation source is the hard x-ray (wavelength less than 1\AA) free-electron laser (FEL) based on a very long undulator in a high-energy electron linear accelerator. Such a device would have a peak brightness many orders of magnitude beyond

that of the third- generation sources, as well as pulse lengths of 100 fs or shorter, and would be fully coherent. Research and development on the many technical challenges that must be overcome are well under way at many laboratories around the world. In the United States, effort is centering around the multi-institutional "Linac Coherent Light Source" proposal to use 15-GeV electrons from the SLAC linac as the source for a 1.5-Å FEL, which if successful would lay the foundation for a later sub-angstrom x-ray FEL. In Europe, HASYLAB at DESY is hosting the two-phase TTF-FEL project culminating in a device operating at 6.4 Å several years from now. The project would pave the way to a still more ambitious 0.1-Å FEL (TESLA-FEL) farther in the future.

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X-Ray Data Booklet

Section 2.3 OPERATING AND PLANNED FACILITIES

Herman Winick

The number of synchrotron radiation facilities is growing rapidly. As a result, the internet is the most reliable source of up-to-date information on facilities around the world. Useful sites include

http://www-ssrl.slac.stanford.edu/sr_sources.html,

<http://www.esrf.fr/navigate/synchrotrons.html> and

http://www.spring8.or.jp/ENGLISH/other_sr/.

Table 2-1 was compiled in October 1999, with input from Masami Ando, Ronald Frahm, and Gwyn Williams. A [PDF version of this table](#) is also available.

Table 2-1. *Storage ring synchrotron radiation sources both planned and operating (October 1999).*

Location	Ring (Institution)	Energy (GeV)	Internet address
<i>Australia</i>	Boomerang	3 —	
<i>Brazil</i>			
Campinas	LNLS-1	1.35 —	
	LNLS-2	2 —	
<i>Canada</i>			
Saskatoon	CLS (Canadian Light Source)	2.5–2.9	www.cls.usask.ca/

China (PRC)

Beijing	BEPC (Inst. High Energy	1.5–2.8	www.friends-
	Phys.)	2.2–2.5	partners.org/~china/ins/IHEP/bsrf/bsrf.html
Hefei	BLS (Inst. High Energy	0.8	—
Shanghai	Phys.)	3.5	—
	NSRL (Univ. Sci. Tech.		—
	China)		
	SSRF (Inst. Nucl. Res.)		

Denmark

Aarhus	ASTRID (ISA)	0.6	www.isa.au.dk/
	ASTRID II (ISA)	1.4	www.isa.au.dk/

England

Daresbury	SRS (Daresbury)	2	srs.dl.ac.uk/index.htm
	DIAMOND	3.0	srs.dl.ac.uk/top/diamond.html
	(Daresbury/Appleton)	0.6	—
	SINBAD (Daresbury)		

France

Grenoble	ESRF	6	www.esrf.fr/
Orsay	DCI (LURE)	1.8	www.lure.u-psud.fr/
	SuperACO (LURE)	0.8	—
	SOLEIL	2.5–2.75	www.SOLEIL.u-psud.fr/

Germany

Berlin	BESSY I	0.8	www.bessy.de/
	BESSY II	1.7–1.9	www.bessy.de/BII/
Bonn	ELSA (Bonn Univ.)	1.5–3.5	www-elsa.physik.uni-bonn.de/elsahome.html
Dortmund	DELTA (Dortmund Univ.)	1.5	www.delta.uni-dortmund.de/home_e.html
Hamburg	DORIS III	4.5–5.3	www-hasylab.desy.de/
	(HASYLAB/DESY)	7–14	www-hasylab.desy.de/
Karlsruhe	PETRA II	2.5	www.fzk.de/anka/english/welcome.html
	(HASYLAB/DESY)		
	ANKA (FZK)		

Location	Ring (Institution)	Energy (GeV)	Internet address
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India

Indore	INDUS-I (Ctr. Adv. Tech.)	0.45	www.ee.ualberta.ca/~naik/cataccel.html
	INDUS-II (Ctr. Adv. Tech.)	2.5	www.ee.ualberta.ca/~naik/cataccel.html

Italy

Frascati	DAFNE	0.51	www.lnf.infn.it/acceleratori/
Trieste	ELETTRA (Synch. Trieste)	1.5–2	waxa.elettra.trieste.it/ELETTRA.html

Japan

Hiroshima	HISOR (Hiroshima Univ.)	0.7	www.hiroshima-u.ac.jp/Organization/src.html
Ichihara	Nano-hana (Japan SOR Inc.)	1.5–2	www.ijnet.or.jp/NANO-HANA/en.html
Kashiwa	VSX (Univ. of Tokyo–ISSP)	2–2.5	www.issp.u-tokyo.ac.jp/labs/sor/vsxs/index.html
Kusatsu	AURORA (Ritsumaiken Univ.)	0.6	www.ritsumeit.ac.jp/se/d11/index-e.html
Kyoto		0.3	www.wal.kuicr.kyoto-u.ac.jp/www/Buttons-e.imagemap?112,64
Nishi Harima	KSR (Kyoto University)	8	
	SPring-8 (JASRI)	1–1.5	www.spring8.or.jp/
	Subaru (Himeji Inst. Tech.)		www.lasti.himeji-tech.ac.jp/NS/Index.html
Okasaki	UVSOR (Inst. Mol. Science)	0.75	www.ims.ac.jp/about_ims/facilities.html
	UVSOR-II (Inst. Mol. Science)	1.0	—
Sendai		1.5	—
Tsukuba	TLS (Tohoku Univ.)	0.8	www.etl.go.jp/etl/linac/e/
	TERAS (ElectroTech. Lab.)	0.6	www.etl.go.jp/etl/linac/e/
	NIJI II (ElectroTech. Lab.)	0.5	www.etl.go.jp/etl/linac/e/
	NIJI IV (ElectroTech. Lab.)	2.5–3	pfwww.kek.jp/
	Photon Factory (KEK)	6	pfwww.kek.jp/
	Accumulator Ring (KEK)		

Korea

Pohang	Pohang Light Source	2	pal.postech.ac.kr/
Seoul	CESS (Seoul Nat. Univ.)	0.1	—

Middle East

SESAME	1	www.weizmann.ac.il/home/sesame/
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Russia

Moscow	Siberia I (Kurchatov Inst.)	0.45	—
	Siberia II (Kurchatov Inst.)	2.5	—
Dubna	DELSY (JINR)	0.6–1.2	—
Novosibirsk	VEPP-2M (BINP)	0.7	ssrc.inp.nsk.su/
	VEPP-3 (BINP)	2.2	ssrc.inp.nsk.su/
	VEPP-4 (BINP)	5–7	ssrc.inp.nsk.su/english/load.pl?right=vepp.html
Zelenograd	Siberia-SM (BINP)	0.8	ssrc.inp.nsk.su/
	TNK (F.V. Lukin Inst.)	1.2–1.6	—

Singapore

Helios2 (Univ. Singapore)	0.7	www.nus.edu.sg/NUSinfo/SSLS/
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Location	Ring (Institution)	Energy (GeV)	Internet address
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<i>Spain</i>				
Barcelona	Catalonia SR Lab	2.5–3		www.lls.ific.es/report/report.html
<i>Sweden</i>				
Lund	MAX I (Univ. Lund)	0.55		www.maxlab.lu.se/welcome.html
	MAX II (Univ. Lund)	1.5		www.maxlab.lu.se/welcome.html
	New Ring (Univ. Lund)	0.7		—
<i>Switzerland</i>				
Villigen	SLS (Paul Scherrer Inst.)	2.4		sls.web.psi.ch/
<i>Taiwan (ROC)</i>				
Hsinchu	SRRC (Synch. Rad. Res. Ctr.)	1.3–1.5		www.srrc.gov.tw/en/main2000.htm
<i>Thailand</i>				
Nakhon Ratchasima	SIAM (Suranaree Univ. Tech.)	1.0		—
<i>Ukraine</i>				
Kharkov	Pulse Stretcher/Synch. Rad.	0.75–2		—
Kiev	ISI-800 (UNSC)	0.7–1.0		—
<i>USA</i>				
Argonne, IL	APS (Argonne Nat. Lab.)	7		epics.aps.anl.gov/welcome.html
Baton Rouge, LA	CAMD (Louisiana State Univ.)	1.4		www.camd.lsu.edu/
Berkeley, CA	ALS (Lawrence Berkeley Nat. Lab.)	1.5–1.9		www-als.lbl.gov/als/
Durham, NC	ALS (Lawrence Berkeley Nat. Lab.)	1–1.3		monk.fel.duke.edu/
Gaithersburg, MD	FELL (Duke Univ.)	0.4		physics.nist.gov/MajResFac/SURF/SURF.html
	SURF III (NIST)	5.5		www.chess.cornell.edu/
Ithaca, NY	CESR (CHESS/Cornell Univ.)	2.5		—
Raleigh, NC	CESR (CHESS/Cornell Univ.)	3		www-ssrl.slac.stanford.edu/welcome.html
Stanford, CA	NC STAR (N. Carolina State Univ.)	3		www-ssrl.slac.stanford.edu/welcome.html
	NC STAR (N. Carolina State Univ.)	0.8–1		www.src.wisc.edu/
		0.80		www.nsls.bnl.gov/
Stoughton, WI	SPEAR2 (SSRL/SLAC)	2.5–2.8		www.nsls.bnl.gov/
Upton, NY	SPEAR3 (SSRL/SLAC)			
	Aladdin (Synch. Rad. Ctr.)			
	NSLS I (Brookhaven Nat. Lab.)			
	NSLS II (Brookhaven Nat. Lab.)			

X-Ray Data Booklet Table 2-1. Storage ring synchrotron radiation sources both planned and operating (October 1999).

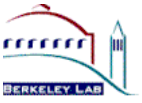
Location	Ring (Institution)	Energy (GeV)	Internet address
<i>Australia</i>	Boomerang	3	—
<i>Brazil</i>			
Campinas	LNLS-1	1.35	—
	LNLS-2	2	—
<i>Canada</i>			
Saskatoon	CLS (Canadian Light Source)	2.5–2.9	cls.usask.ca/
<i>China (PRC)</i>			
Beijing	BEPC (Inst. High Energy Phys.)	1.5–2.8	www.friends-partners.org/~china/ins/IHEP/bsrf/bsrf.html
	BLS (Inst. High Energy Phys.)	2.2–2.5	—
Hefei	NSRL (Univ. Sci. Tech. China)	0.8	—
Shanghai	SSRF (Inst. Nucl. Res.)	3.5	—
<i>Denmark</i>			
Aarhus	ASTRID (ISA)	0.6	www.isa.au.dk/
	ASTRID II (ISA)	1.4	www.isa.au.dk/
<i>England</i>			
Daresbury	SRS (Daresbury)	2	srs.dl.ac.uk/index.htm
	DIAMOND (Daresbury/Appleton)	3.0	srs.dl.ac.uk/top/diamond.html
	SINBAD (Daresbury)	0.6	—
<i>France</i>			
Grenoble	ESRF	6	www.esrf.fr/
Orsay	DCI (LURE)	1.8	www.lure.u-psud.fr/
	SuperACO (LURE)	0.8	—
	SOLEIL	2.5–2.75	www.sol.cnrs-gif.fr/
<i>Germany</i>			
Berlin	BESSY I	0.8	www.bessy.de/
	BESSY II	1.7–1.9	www.bessy.de/BII/
Bonn	ELSA (Bonn Univ.)	1.5–3.5	www-elsa.physik.uni-bonn.de/elsahome.html
Dortmund	DELTA (Dortmund Univ.)	1.5	www.delta.uni-dortmund.de/home_e.html
Hamburg	DORIS III (HASYLAB/DESY)	4.5–5.3	www-hasylab.desy.de/
	PETRA II (HASYLAB/DESY)	7–14	www-hasylab.desy.de/
Karlsruhe	ANKA (FZK)	2.5	www.fzk.de/anka/english/welcome.html

Table 2-1. Storage ring synchrotron radiation sources (continued).

Location	Ring (Institution)	Energy (GeV)	Internet address
<i>India</i>			
Indore	INDUS-I (Ctr. Adv. Tech.)	0.45	www.ee.ualberta.ca/~naik/cataccel.html
	INDUS-II (Ctr. Adv. Tech.)	2.5	www.ee.ualberta.ca/~naik/cataccel.html
<i>Italy</i>			
Frascati	DAFNE	0.51	www.lnf.infn.it/acceleratori/
Trieste	ELETTRA (Synch. Trieste)	1.5–2	waxa.elettra.trieste.it/ELETTRA.html
<i>Japan</i>			
Hiroshima	HISOR (Hiroshima Univ.)	0.7	www.hiroshima-u.ac.jp/Organization/src.html
Ichihara	Nano-hana (Japan SOR Inc.)	1.5–2	www.ijnet.or.jp/NANO-HANA/en.html
Kashiwa	VSX (Univ. of Tokyo–ISSP)	2–2.5	www.issp.u-tokyo.ac.jp/labs/sor/vsxs/index.html
Kusatsu	AURORA (Ritsumei Univ.)	0.6	www.ritsumei.ac.jp/se/d11/index-e.html
Kyoto	KSR (Kyoto University)	0.3	www.kuicr.kyoto-u.ac.jp/www/Buttons-e.imagemap?112,64
Nishi Harima	SPRING-8 (JASRI)	8	www.spring8.or.jp/
	Subaru (Himeji Inst. Tech.)	1–1.5	www.lasti.himeji-tech.ac.jp/NS/Index.html
Okasaki	UVSOR (Inst. Mol. Science)	0.75	www.ims.ac.jp/about_ims/facilities.html
	UVSOR-II (Inst. Mol. Science)	1.0	—
Sendai	TLS (Tohoku Univ.)	1.5	—
Tsukuba	TERAS (ElectroTech. Lab.)	0.8	www.etl.go.jp/etl/linac/e/
	NIJI II (ElectroTech. Lab.)	0.6	www.etl.go.jp/etl/linac/e/
	NIJI IV (ElectroTech. Lab.)	0.5	www.etl.go.jp/etl/linac/e/
	Photon Factory (KEK)	2.5–3	pfwww.kek.jp/
	Accumulator Ring (KEK)	6	pfwww.kek.jp/
<i>Korea</i>			
Pohang	Pohang Light Source	2	pal.postech.ac.kr/
Seoul	CESS (Seoul Nat. Univ.)	0.1	—
<i>Middle East</i>	SESAME	1	www.weizmann.ac.il/home/sesame/
<i>Russia</i>			
Moscow	Siberia I (Kurchatov Inst.)	0.45	—
	Siberia II (Kurchatov Inst.)	2.5	—
Dubna	DELSY (JINR)	0.6–1.2	—
Novosibirsk	VEPP-2M (BINP)	0.7	ssrc.inp.nsk.su/
	VEPP-3 (BINP)	2.2	ssrc.inp.nsk.su/
	VEPP-4 (BINP)	5–7	ssrc.inp.nsk.su/english/load.pl?right=vepp.html
	Siberia-SM (BINP)	0.8	ssrc.inp.nsk.su/
Zelenograd	TNK (F.V. Lukin Inst.)	1.2–1.6	—
<i>Singapore</i>	Helios2 (Univ. Singapore)	0.7	www.nus.edu.sg/NUSinfo/SSLS/

Table 2-1. Storage ring synchrotron radiation sources (continued).

Location	Ring (Institution)	Energy (GeV)	Internet address
<i>Spain</i>			
Barcelona	Catalonia SR Lab	2.5–3	www.lls.ifae.es/report/report.html
<i>Sweden</i>			
Lund	MAX I (Univ. Lund)	0.55	www.maxlab.lu.se/welcome.html
	MAX II (Univ. Lund)	1.5	www.maxlab.lu.se/welcome.html
	New Ring (Univ. Lund)	0.7	—
<i>Switzerland</i>			
Villigen	SLS (Paul Scherrer Inst.)	2.4	www1.psi.ch/www_sls_hn/
<i>Taiwan (ROC)</i>			
Hsinchu	SRRC (Synch. Rad. Res. Ctr.)	1.3–1.5	www.srrc.gov.tw/en/main2000.htm
<i>Thailand</i>			
Nakhon Ratchasima	SIAM (Suranaree Univ. Tech.)	1.0	—
<i>Ukraine</i>			
Kharkov	Pulse Stretcher/Synch. Rad.	0.75–2	—
Kiev	ISI-800 (UNSC)	0.7–1.0	—
<i>USA</i>			
Argonne, IL	APS (Argonne Nat. Lab.)	7	epics.aps.anl.gov/welcome.html
Baton Rouge, LA	CAMD (Louisiana State Univ.)	1.4	www.camd.lsu.edu/
Berkeley, CA	ALS (Lawrence Berkeley Nat. Lab.)	1.5–1.9	www-als.lbl.gov/als/
Durham, NC	FELL (Duke Univ.)	1–1.3	monk.fel.duke.edu/
Gaithersburg, MD	SURF III (NIST)	0.4	physics.nist.gov/MajResFac/SURF/SURF.html
Ithaca, NY	CESR (CHESS/Cornell Univ.)	5.5	www.chess.cornell.edu/
Raleigh, NC	NC STAR (N. Carolina State Univ.)	2.5	—
Stanford, CA	SPEAR2 (SSRL/SLAC)	3	www-ssrl.slac.stanford.edu/welcome.html
	SPEAR3 (SSRL/SLAC)	3	www-ssrl.slac.stanford.edu/welcome.html
Stoughton, WI	Aladdin (Synch. Rad. Ctr.)	0.8–1	www.src.wisc.edu/
Upton, NY	NSLS I (Brookhaven Nat. Lab.)	0.80	www.nsls.bnl.gov/
	NSLS II (Brookhaven Nat. Lab.)	2.5–2.8	www.nsls.bnl.gov/



3. Scattering Processes

Contents

- [Scattering of X-rays from Electrons and Atoms](#) - *Janos Kirz*
- [Low-Energy Electron Ranges in Matter](#) - *Piero Pianetta*

[CXRO Home](#) | [ALS Home](#) | [LBL Home](#)

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X-Ray Data Booklet

Section 3.1 SCATTERING of X-RAYS from ELECTRONS and ATOMS

Janos Kirz

A. COHERENT, RAYLEIGH, OR ELASTIC SCATTERING

Scattering from single electrons (Thomson scattering) has a total cross section

$$\sigma_T = 8\pi r_e^2 / 3 = 6.652 \times 10^{-29} \text{ m}^2, \quad (1)$$

where r_e is the classical radius of the electron, $e^2 / mc^2 = 2.818 \times 10^{-15}$ meter. The angular distribution for unpolarized incident radiation is proportional to $(1 + \cos^2 \theta)$, where θ is the scattering angle. For polarized incident radiation, the cross section vanishes at 90° in the plane of polarization.

Scattering from atoms involves the cooperative effect of all the electrons, and the cross section becomes

$$\sigma_R = \pi r_e^2 \int_{-1}^1 |f(\theta)|^2 (1 + \cos^2 \theta) d(\cos \theta), \quad (2)$$

where $f(\theta)$ is the (complex) atomic scattering factor, tabulated in Section 2.7 of this booklet. Up to about 2 keV, the scattering factor is approximately independent of scattering angle, with a real part that represents the effective number of electrons that participate in the scattering. At higher energies, the scattering factor falls off rapidly with scattering angle. For details see Ref. 1.

B. COMPTON SCATTERING

In relativistic quantum mechanics, the scattering of x-rays by a free electron is given by the Klein-Nishina formula. If we assume unpolarized x-rays and unaligned electrons, this formula can be approximated as follows for x-ray energies below 100 keV:

$$d\sigma_{\text{KN}} / d\Omega \cong \frac{r_e^2 (1 + \cos^2 \theta)}{2[1 + k(1 - \cos \theta)]^2} , \quad (3)$$

where $k = E / mc^2$, the photon energy measured in units of the electron rest energy. The total cross section is approximately

$$\sigma_{\text{KN}} \cong 8\pi r_e^2 \frac{(1 + 2k + 1.2k^2)}{3(1 + 2k)^2} . \quad (4)$$

Note that for very low energies ($k \rightarrow 0$), we recover the Thomson cross section. The real difference comes when we deal with atoms. In that case, if the scattering leaves the atom in the ground state, we deal with coherent scattering (see above), whereas if the electron is ejected from the atoms, the scattering is (incoherent) Compton scattering. At high energies, the total Compton cross section approaches $Z\sigma_{\text{KN}}$. At low energies and small scattering angles, however, binding effects are very important, the Compton cross section is significantly reduced, and coherent scattering dominates (see Figs. 3-1 and 3-2). For details see Refs. 1 and 2.

The scattered x-ray suffers an energy loss, which (ignoring binding effects) is given by

$$E' / E = 1 / [1 + k(1 - \cos \theta)] \quad (5)$$

or, in terms of the wavelength shift,

$$\lambda' - \lambda = \lambda_c (1 - \cos \theta) , \quad (6)$$

where $\lambda_c = h / mc = 2.426 \times 10^{-12}$ meter . The kinetic energy of the recoil electron is just the energy lost by the photon in this approximation:

$$E_e = E \frac{k(1 - \cos \theta)}{1 + k(1 - \cos \theta)} \quad (7)$$

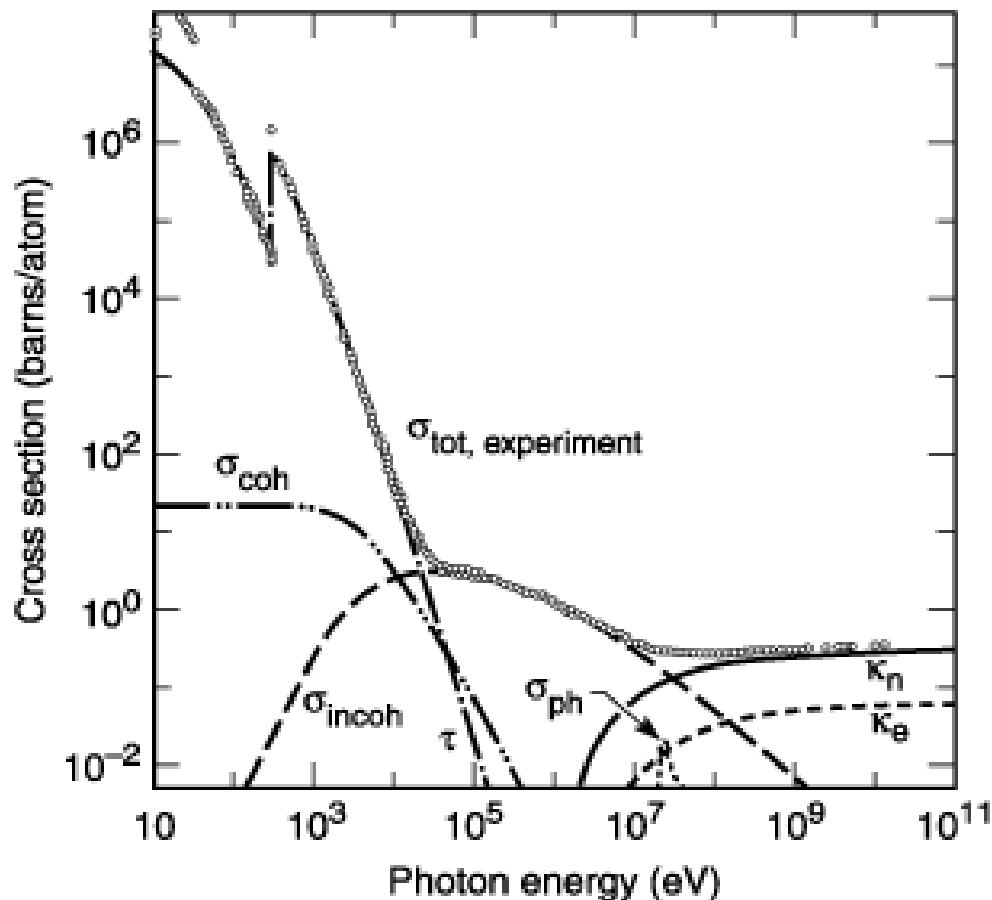


Fig. 3-1. Total photon cross section σ_{tot} in carbon, as a function of energy, showing the contributions of different processes: τ , atomic photo-effect (electron ejection, photon absorption); σ_{coh} , coherent scattering (Rayleigh scattering—atom neither ionized nor excited); σ_{incoh} , incoherent scattering (Compton scattering off an electron); κ_n , pair production, nuclear field; κ_e , pair production, electron field; σ_{ph} , photonuclear absorption (nuclear

absorption, usually followed by emission of a neutron or other particle). (From Ref. 3; figure courtesy of J. H. Hubbell.)

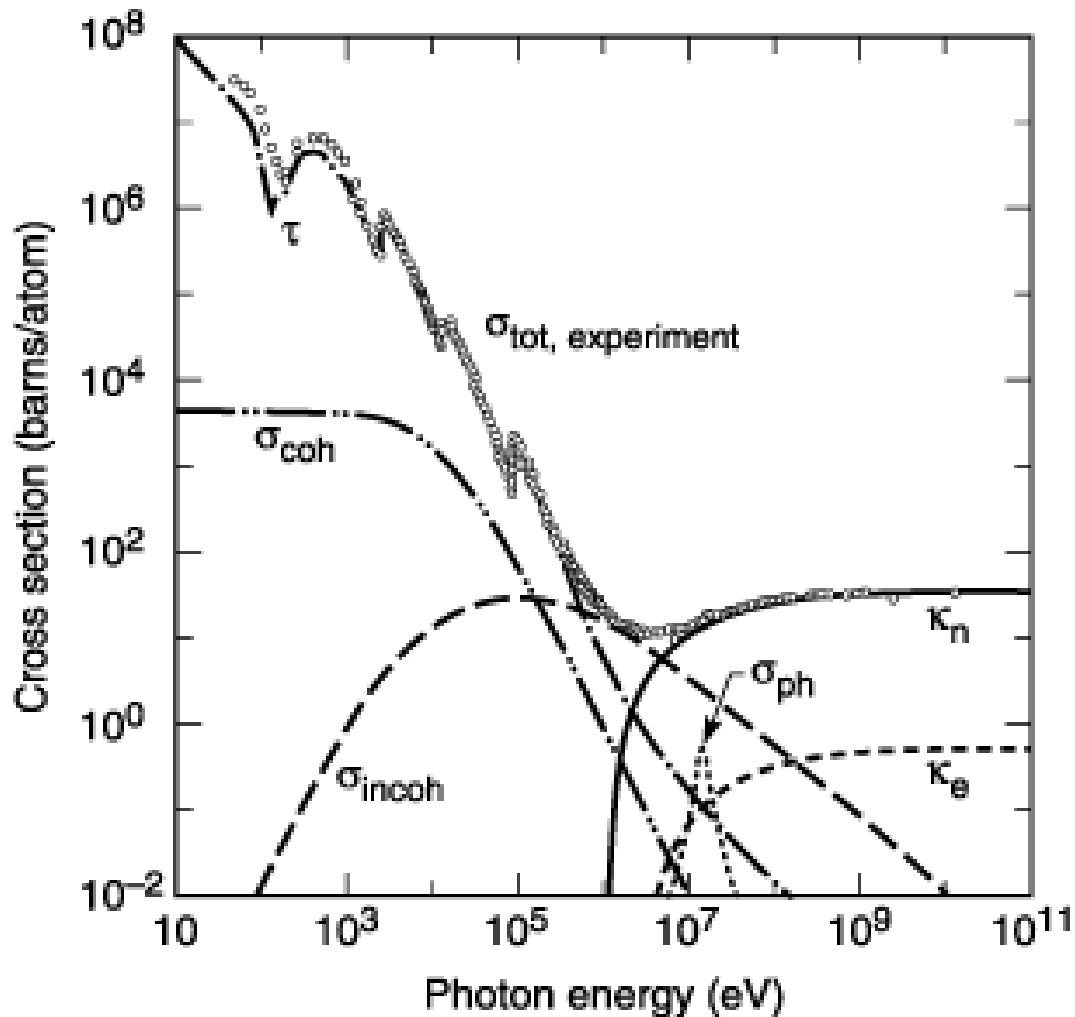


Fig. 3-2. Total photon cross section σ_{tot} in lead, as a function of energy. See Fig. 3-1. (From Ref. 3; figure courtesy of J. H. Hubbell.)

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X-Ray Data Booklet

Section 3.2 LOW-ENERGY ELECTRON RANGES IN MATTER

Piero Pianetta

The electron range is a measure of the straight-line penetration distance of electrons in a solid [1]. Electrons with energies in the kilo-electron volt range, traveling in a solid, are scattered inelastically in collisions with the electrons in the material. For low-*Z* materials, such as organic insulators, scattering from the valence electrons is the major loss mechanism for incident electron energies from 10 eV to 10 keV. The core levels contribute less than 10% to the electron's energy dissipation for energies between 1 keV and 10 keV [2]. A [PDF version of this section](#) is also available.

A. CSDA RANGES

For electron energies below 5 keV, the usual Bethe-Bloch formalism is inadequate for calculating the electron energy loss in a solid, and an approach using the dielectric response of the material is used [3]. The complex dielectric function $\epsilon(k, \omega)$ describes the response of a medium to a given energy transfer $\hbar\omega$ and momentum transfer $\hbar k$. The dielectric function contains contributions from both valence and core electrons. References 4 and 5 describe the steps for calculating $\epsilon(k, \omega)$ for insulators and metals, respectively. For an electron of energy E , the probability of an energy loss ω per unit distance is given by [2]

$$\tau(E, \hbar\omega) = \frac{1}{\pi a_0 E} \int_{k_-}^{k_+} \frac{dk}{k} \operatorname{Im} \left[\frac{-1}{\epsilon(k, \omega)} \right] , \quad (1)$$

where $\hbar k_{\pm} = \sqrt{2m}(\sqrt{E} \pm \sqrt{E - \hbar\omega})$ and $a_0 = \hbar^2 / me^2$. The quantity $\tau(E, \hbar\omega)$ is also known as the differential inverse mean free path, because by integrating it over all allowed energy transfers, the inelastic mean free path (IMFP) is obtained.

Furthermore, an integration of $\hbar\omega\tau(E, \hbar\omega)$ over all allowed energy transfers gives the energy loss per unit path length, or stopping power $S(E)$. The stopping power can then be used to calculate the distance it takes to slow an electron down to a given energy. This distance is called the continuous slowing down approximation range, or CSDA range, because the calculation assumes that the electron slows down continuously from the initial energy E to the final energy, which is usually taken to be 10 eV [2]. The CSDA range $R_0(E)$ is given by

$$R_0(E) = \int_{10\text{eV}}^E \frac{dE'}{S(E')} \quad (2)$$

The calculations for IMFP and stopping power have been carried out down to 10 eV for a number of materials, including SiO_2 [3]; polystyrene [2]; polyethylene [6]; collodion [7]; and silicon, aluminum, nickel, copper, and gold [5]. The CSDA ranges from 15 eV to 6 keV were then calculated for polystyrene, silicon, and gold by integrating Eq. (2) and are shown in Fig. 3-3. These curves can be used with confidence down to 100 eV. However, comparisons of different available calculations with the meager experimental data below 100 eV indicate that errors as large as 100% may occur at 10 eV. An example of this is shown in the figure, where experimental range data for collodion are given. It is clear that the agreement between the collodion and polystyrene data starts to become reasonable above 100 eV. The differences below 100 eV could equally well be due to problems with the theory or to the increased difficulty of the measurement. Stopping-power calculations for polymethyl methacrylate (PMMA) have been carried out only from 100 eV, so that the CSDA range as defined above could not be calculated [4]. However, data on effective electron ranges of photoelectrons in PMMA at several energies can be found in Ref. 8.

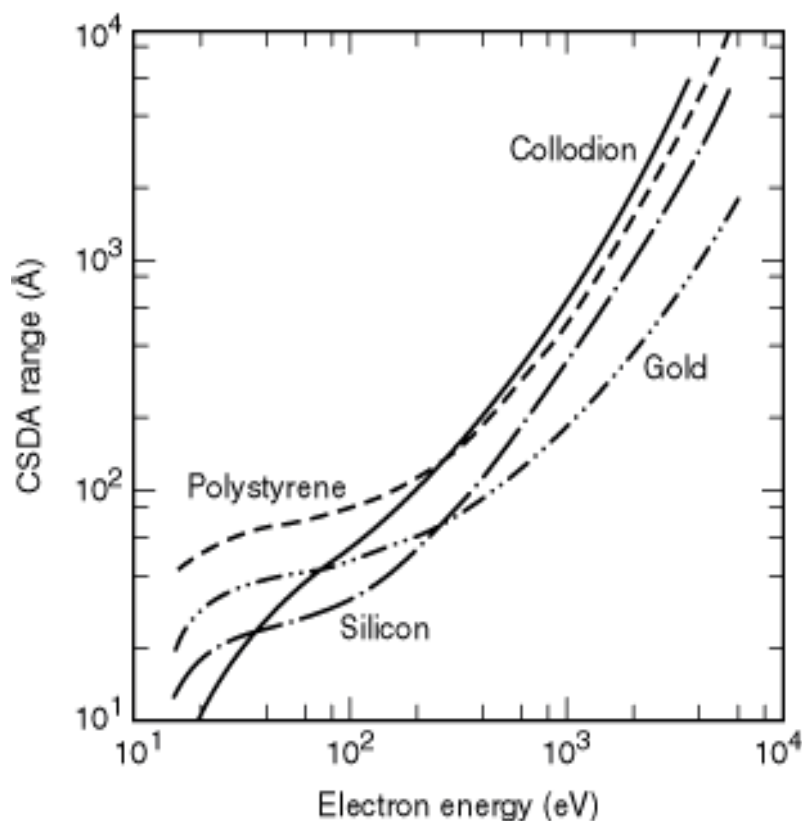


Fig. 3-3. Plot of the CSDA range, as a function of energy, for gold and silicon [5] and for polystyrene, $(C_8H_8)_n$, with a density of 1.05 g/cm^3 [2]. The measured electron range in collodion with a density of 1 g/cm^3 is also plotted [7].

B. ELECTRON INELASTIC MEAN FREE PATHS

A very important aspect of photoelectron spectroscopy, especially with synchrotron radiation, is the ability to effectively tune the surface sensitivity from a few angstroms or a few tens of angstroms in core-level photoemission measurements to a few hundred angstroms in total-electron-yield surface EXAFS experiments. This variation arises from the fact that the IMFP of the photoemitted electrons is a strong function of the electron kinetic energy, which can be tuned by the appropriate choice of photon energy. The definition of the IMFP [9] is the average distance traveled by an electron between inelastic collisions. Although the exact relationship between the IMFP and kinetic energy depends on the detailed electronic structure of the element or compound of interest, the general features are similar for all elements, starting at large values for kinetic energies

below 10–15 eV, dropping to a minimum value of 5–10 Å at kinetic energies between 30 and 100 eV, and then rising monotonically above 100 eV.

Since the surface sensitivity is determined by the depth *perpendicular* to the surface from which electrons can escape, it is best defined using the mean escape depth (MED), which is related to the IMFP by

$$\Delta = \lambda_i \cos \alpha \quad , \quad (3)$$

where Δ is the MED, λ_i is the IMFP and α is the emission angle of the electrons relative to the surface normal. However, it should be noted that *elastic* scattering effects within the solid could increase the MED as much as a factor of two at electron emission angles greater than 60°, depending on the angle of incidence of the incoming x-rays and the particular core level being studied [9,10]. Therefore, the standard technique of increasing the surface sensitivity by working at glancing emission angles using Eq. (3) must be qualified to take these effects into account. In addition, both angle-dependent cross sections and photoelectron diffraction effects can result in anisotropic emission from the solid that can also cause errors in the interpretation of the MEDs in solids. Because of these complications, graphs of the IMFPs, rather than the MEDs, versus electron kinetic energy will be presented here to give a measure of the surface sensitivity. The reader is referred to Ref. 9 when more complicated experimental conditions need to be considered.

Using the formalism developed by Penn that uses optical data to determine the IMFP of a material [11], Tanuma et al. have calculated the IMFPs for a large number of elements and compounds for kinetic energies up to 2000 eV [12–14]. Figure 3-4 shows IMFP curves for Ag, Al, Na, PMMA, Si, and SiO₂.

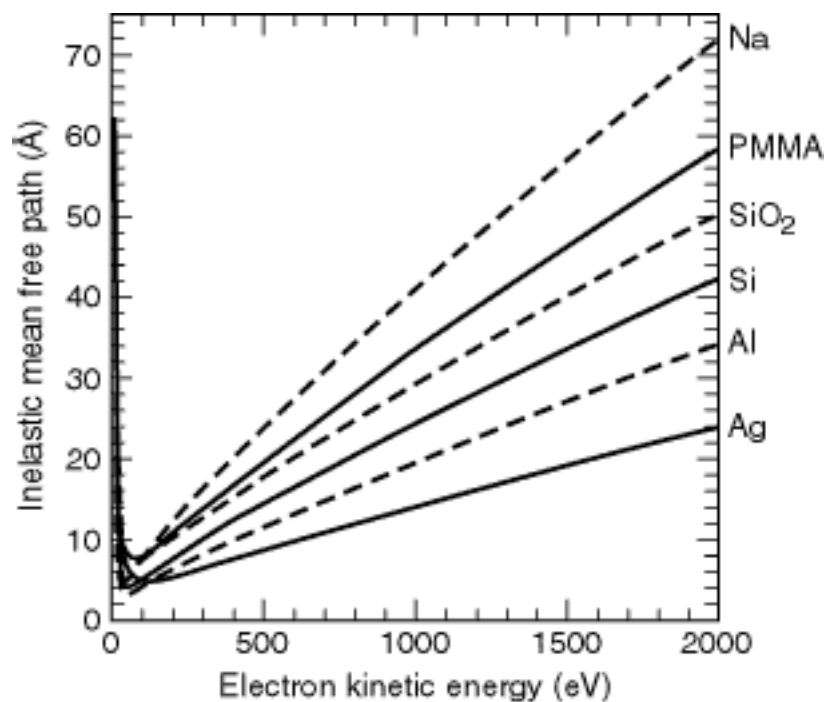


Fig. 3-4. *Inelastic mean free paths for electron kinetic energies up to 2000 eV, for Ag, Al, NA, PMMA, Si, and SiO₂.*

These materials are representative of a fairly wide variety of materials for kinetic energies between 200 and 2000 eV. For example, the IMFPs for Ni, Ta, W, Pt, and Au all hover around the values given here for Ag; Cr, Fe, and Cu fall between Al and Ag. Likewise, C falls between Si and SiO₂, whereas GaAs overlies the PMMA curve for much of this energy range. The behavior below 200 eV is more complex, because the IMFPs are strongly dependent on the details of the electronic structure. Figure 3-5 shows the region below 250 eV for Al, Ag, GaAs, NA, PMMA, and Si. Silicon dioxide is not shown here because it overlaps the PMMA curve in this range, whereas GaAs does not. Although the calculations below 50 eV may not be reliable, owing to limitations in the theory, the values are plotted at these low energies to show the general behavior of the IMFPs in this region, as well as the location of the minima for the different materials.

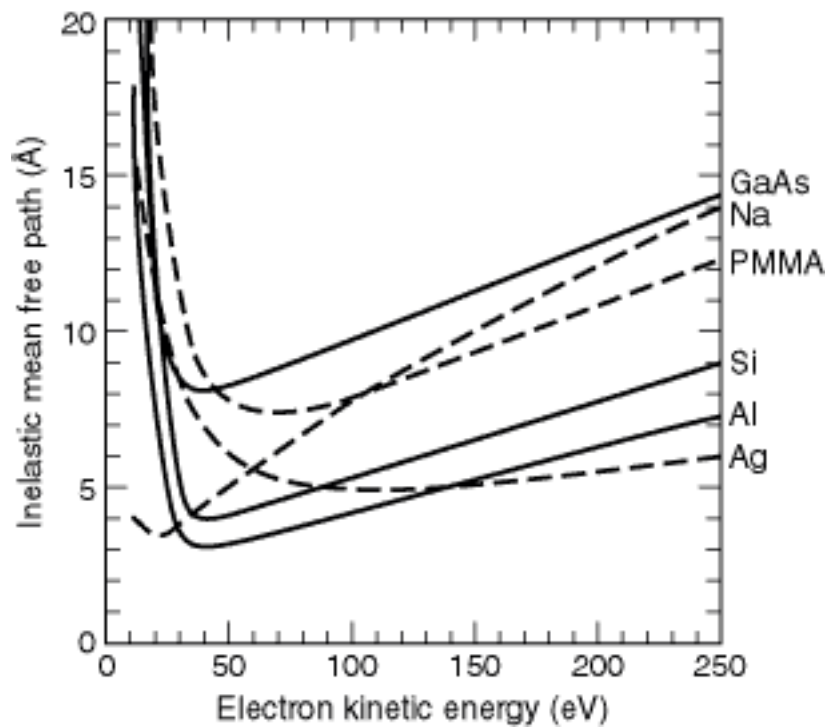


Fig. 3-5. Detail of the inelastic mean free paths in the kinetic energy range below 250 eV, for Ag, Al, GaAs, NA, PMMA, and Si.

Calculations for additional materials can be found in the literature as follows: (i) elements from C to Bi [12]; (ii) III-V and II-VI compound semiconductors, alkali halides, Si_3N_4 , and several oxides [13]; and (iii) organic compounds [14]. Calculations are being presented here because they provide the most complete and consistent set of values for the IMFPs. References 9 and 10 give the historical background for both the theory and the experimental work in this field and show that it is difficult to generalize much of the experimental data in the literature, owing to the experiment-specific effects described above, as well as uncertainties in sample preparation. Seah and Dench [15] were the first to classify the material dependence of the IMFPs and presented data for kinetic energies up to 10 keV. A good example of the care that is needed in determining IMFPs is given in Ref. 8, which is a study of the Si/SiO₂ system. Finally, it should be mentioned that spin-dependent effects on the IMFP have also been observed in ferromagnetic materials [17].

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X-Ray Data Booklet

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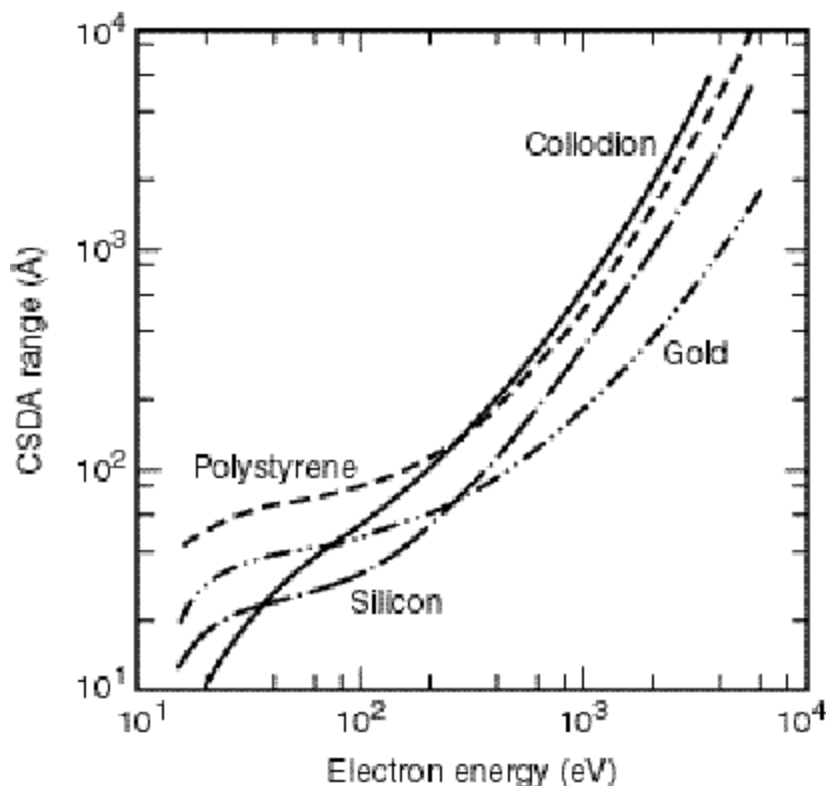


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$$\Delta = \lambda_i \cos \alpha, \quad (3)$$

where Δ is the MED, λ_i is the IMFP and α is the emission angle of the electrons relative to the surface normal. However, it should be noted that *elastic* scattering effects within the solid could increase the MED as much as a factor of two at electron emission angles greater than 60°, depending on the angle of incidence of the incoming x-rays and the particular core level being studied [9,10]. Therefore, the standard technique of increasing the surface sensitivity by working at glancing emission angles using Eq. (3) must be qualified to take these effects into account. In addition, both angle-dependent cross sections and photoelectron diffraction effects can result in anisotropic emission from the solid that can also cause errors in the interpretation of the MEDs in solids. Because of these complications, graphs of the IMFPs, rather than the MEDs, versus electron kinetic energy will be presented here to give a measure of the surface sensitivity. The reader is referred to Ref. 9 when more complicated experimental conditions need to be considered.

Using the formalism developed by Penn that uses optical data to determine the IMFP of a material [11], Tanuma et al. have calculated the IMFPs for a large number of elements and compounds for kinetic energies up to 2000 eV [12–14]. Figure 3-4 shows IMFP curves for Ag, Al, Na, PMMA, Si, and SiO₂. These materials are representative of a fairly wide variety of materials for kinetic energies between 200 and 2000 eV. For example, the IMFPs for Ni, Ta, W, Pt, and Au all hover around the values given here for Ag; Cr, Fe, and Cu fall between Al and Ag. Likewise, C falls between Si and SiO₂, whereas GaAs overlies the PMMA curve for much of this energy range. The behavior below 200 eV is more complex, because the IMFPs are strongly dependent on the details of the electronic structure. Figure 3-5 shows the region below 250 eV for Al, Ag, GaAs, Na, PMMA, and Si. Silicon dioxide is not shown here because it overlaps the PMMA curve in this range, whereas GaAs does not. Although the calculations below 50 eV may not be reliable, owing to limitations in the theory, the values are plotted at these low energies to show the general behavior of the IMFPs in this region, as well as the location of the minima for the different materials. Calculations for additional materials can be found in the literature as follows: (i) elements from C to Bi [12]; (ii) III-V and II-VI compound semiconductors, alkali halides, Si₃N₄, and several oxides [13]; and (iii) organic compounds [14]. Calculations are being presented here because they provide the most complete and consistent set of values for the IMFPs. References 9 and 10 give the historical background for both the theory and the experimental work in this field and show that it is difficult to generalize much of the experimental data in the literature, owing to the experiment-specific effects described above, as well as uncertainties in sample preparation. Seah and Dench [15] were the first to classify the material dependence of the IMFPs and presented data for kinetic energies up to 10 keV. A good example of the care that is needed in determining IMFPs is given in Ref. 8, which is a study of the Si/SiO₂ system. Finally, it should be mentioned that spin-dependent effects on the IMFP have also been observed in ferromagnetic materials [17].

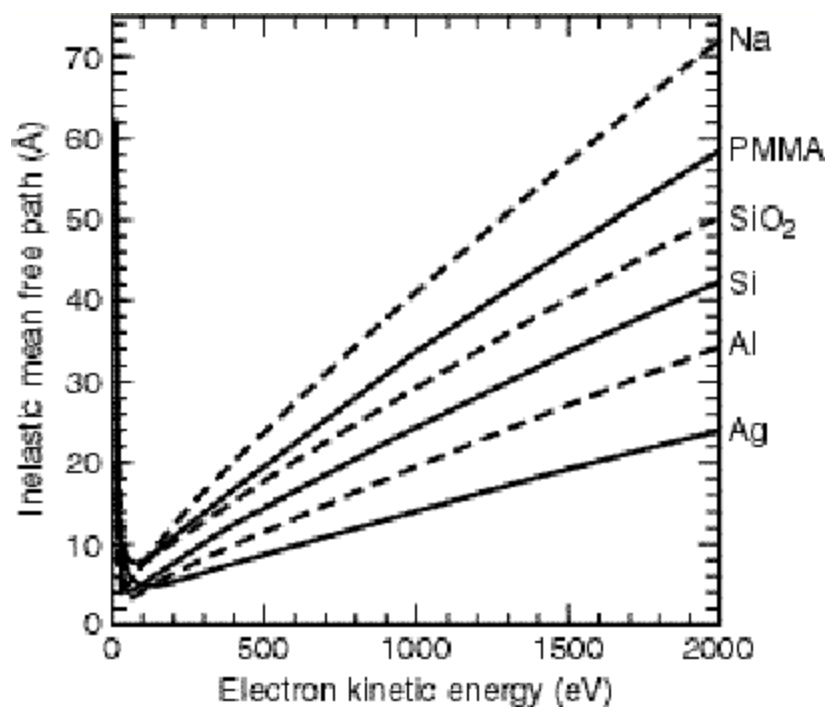


Fig. 3-4. Inelastic mean free paths for electron kinetic energies up to 2000 eV, for Ag, Al, Na, PMMA, Si, and SiO₂.

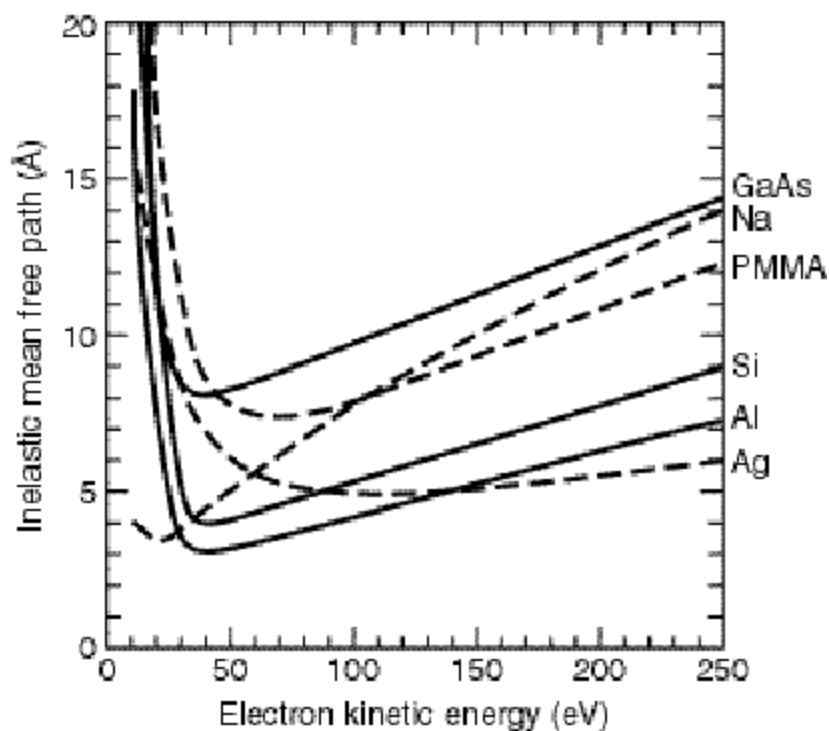
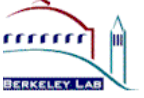


Fig. 3-5. Detail of the inelastic mean free paths in the kinetic energy range below 250 eV, for Ag, Al, GaAs, Na, PMMA, and Si.

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4. Optics

Contents

- [Crystal and Multilayer Elements](#) - *James H. Underwood*
- [Specular Reflectivities for Grazing-Incidence Mirrors](#) - *Eric M. Gullikson*
- [Gratings and Monochromators](#) - *Malcolm R. Howells*
- [Zone Plates](#) - *Janos Kirz and David T. Attwood*
- [X-Ray Detectors](#) - *Albert C. Thompson*

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X-Ray Data Booklet

Section 4.1 MULTILAYERS AND CRYSTALS

James H. Underwood

(This section is also available as a [PDF file](#).)

A. MULTILAYERS

By means of modern vacuum deposition technology, structures consisting of alternating layers of high- and low- Z materials, with individual layers having thicknesses of the order of nanometers, can be fabricated on suitable substrates. These structures act as multilayer interference reflectors for x-rays, soft x-rays, and extreme ultraviolet (EUV) light. Their high reflectivity and moderate energy bandwidth ($10 < E/\Delta E < 100$) make them a valuable addition to the range of optical components useful in instrumentation for EUV radiation and x-rays with photon energies from a few hundred eV to tens of keV. These multilayers are particularly useful as mirrors and dispersive elements on synchrotron radiation beamlines. They may be used to produce focal spots with micrometer-scale sizes and for applications such as fluorescent microprobing, microdiffraction, and microcrystallography. Multilayer reflectors also have a wide range of applications in the EUV region, where normal-incidence multilayer reflectors allow the construction of space telescopes and of optics for EUV lithography. Ordinary mirrors, operating at normal or near-normal angles of incidence, do not work throughout the EUV and x-ray regions. The reason is the value of the complex refractive index, $n = 1 - \delta - i\beta$, which can be close to unity for all materials in this region; δ may vary between 10^{-3} in the EUV region to 10^{-6} in the x-ray region. The Fresnel equations at normal incidence show that the reflectivity $R^2 = |(n-1)/(n+1)|^2$ is very small. However, the more general Fresnel equations show that x-rays and EUV radiation can be reflected by mirrors at large angles of incidence. Glancing (or grazing) incidence is the term used when the

rays make a small angle (a few degrees or less) with the mirror surface; the complement ($90^\circ - i$) of the optical angle of incidence i is called the glancing (or grazing) angle of incidence. Glancing-incidence mirrors are discussed in Section 4.2.

Although the normal-incidence intensity reflectivity R^2 of a surface might be 10^{-3} or 10^{-4} , the corresponding *amplitude* reflectivity $R = (n - 1)/(n + 1)$, the square root of the intensity reflectivity, will be 1/30 to 1/100. This implies that, if the re-reflections from 30–100 surfaces could be made to add in phase, a total reflectivity approaching unity could be obtained. This is the multilayer principle. As shown in Fig. 4-1, a multilayer reflector comprises a stack of materials having alternately high and low refractive indices. The thicknesses are adjusted so that the path length difference between reflections from successive layer pairs is equal to one wavelength. Hence, x-ray/EUV multilayers are approximately equivalent to the familiar “quarter-wave stacks” of visible-light coating technology. The equivalence is not exact, however, because of the absorption term β , which is usually negligible for visible-light multilayers. This absorption reduces the multilayer reflectivity below unity and requires the design to be optimized for highest reflectivity. Particularly critical is the value of Γ , which is the ratio of the thickness of the high- Z (high electron density, high $|n|$) layer to the total thickness of each layer pair. This optimization normally requires modeling or simulation of the multi-layer. Most thin-film calculation programs, even if designed for the visible region, will perform these calculations if given the right complex values of the refractive index. Alternatively, there are a number of on-line calculation programs available; links can be found at <http://www-cxro.lbl.gov/>.

Either elemental materials or compounds can be used to fabricate multilayer reflectors. The performance obtained from a multilayer depends largely on whether there exists a fortuitous combination of materials having the right refractive indices (good contrast, low absorption), whether these materials can be deposited in smooth thin layers, and whether they remain stable (low reactivity, low diffusion) in the deposited state. The roughness of the underlying substrate is also of prime importance; an rms roughness of the same order of magnitude as the layer

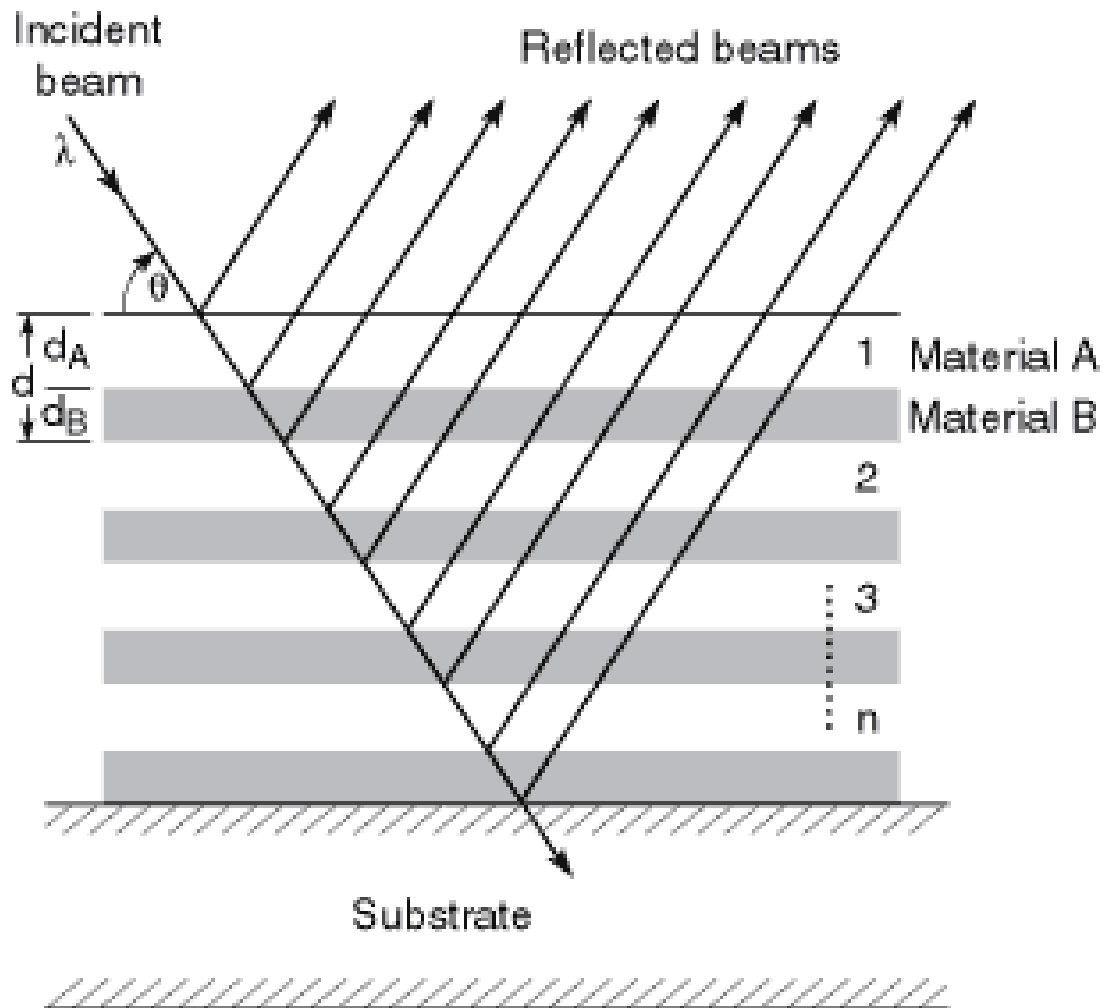


Fig. 4-1. Schematic of a multilayer reflector of n bilayer pairs. The parameters λ , θ , and d are chosen to satisfy the familiar Bragg equation, but the relative thicknesses of the high- and low- Z materials are also critical in optimizing reflectivity. The total reflectivity is the vector sum of the complex reflection coefficients at each interface, with the different path lengths taken into account.

thicknesses will spoil the performance of most coatings. “Superpolished” substrates, with roughness $\sigma \approx 0.1$ nm are preferred. On such substrates, the peak reflectivity of the coatings can approach 80% to 90% of theoretical predictions. A large variety of multilayers have been made and their reflectivities measured over the years. A database of reported reflectivities has been assembled from surveys taken at the biennial Physics of X-ray Multilayer Structures conferences

and can be found at the website listed above. For the EUV region around 100 eV, two remarkably successful combinations are molybdenum-silicon and molybdenum-beryllium. With Mo-Si, a normal-incidence reflectivity of 68% has been achieved at a wavelength of 13.4 nm; Mo-Be multilayers have achieved a reflectivity close to 70% at 11.4 nm (see Fig. 4-2). These relatively high reflectivities are the basis for current efforts in the field of EUV lithography.

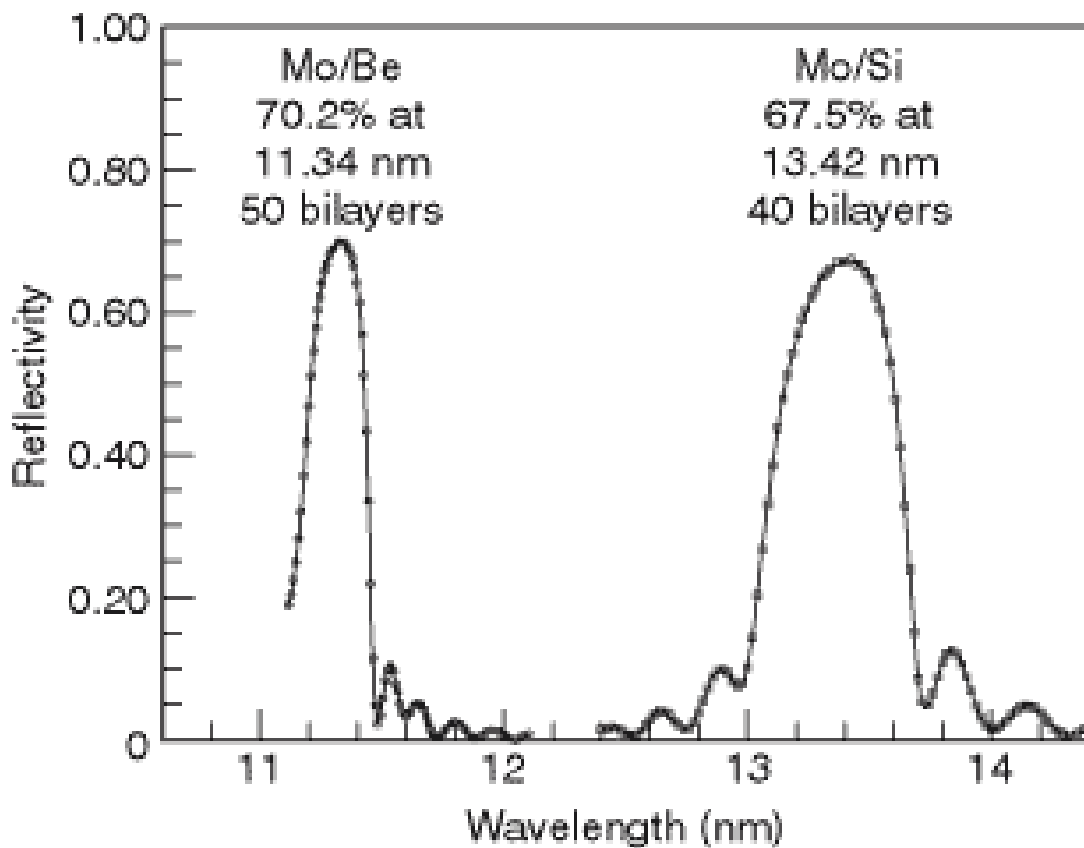


Fig. 4-2. The reflectivity of two multilayer reflectors at extreme ultraviolet wavelengths.

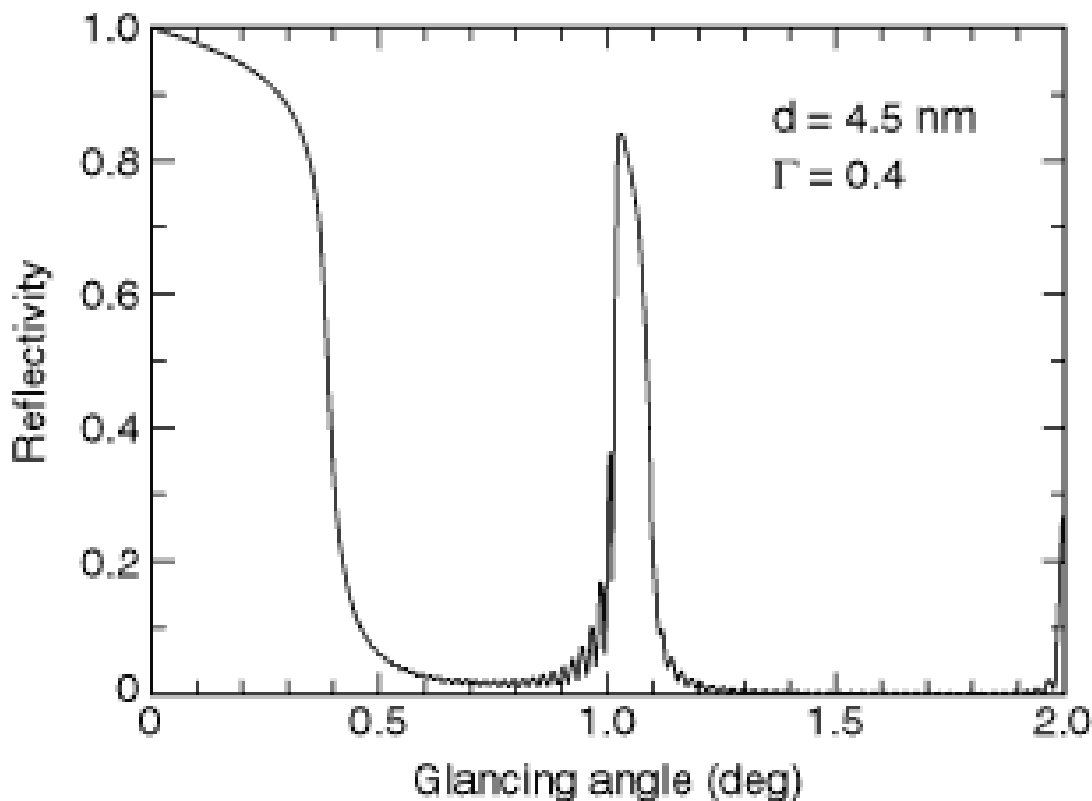


Fig. 4-3. *The reflectivity of a tungsten–boron carbide multilayer at 8048 eV. The parameters d and Γ are discussed in the text.*

At hard x-ray wavelengths near 10 keV, a commonly used multilayer is made with tungsten as the high- Z material and boron carbide (B_4C) as the low- Z material. A reflectivity of 84% has been achieved with this combination at 8048 eV, the energy of the Cu $K\alpha$ line (see Fig. 4-3).

B. CRYSTALS

Multilayers are examples of a periodic structure that can be used to analyze short-wavelength electromagnetic radiation. Such structures split the incident beam into a large number N of separate beams; between any beam i and the beam $i + 1$, the optical path difference is constant. After leaving the periodic structure, the beams are recombined and caused to interfere, whereupon the spectrum of the incident radiation is produced.

Dispersion of radiation by a periodic structure is thus formally equivalent to multiple-beam interferometry. Structures that are periodic across their surface and that produce the N interfering beams by division of the incident wave front

are called gratings and are treated in Section 4.3. Crystals and multilayer structures produce the N interfering beams by division of the incident *amplitude*. The spectrum of the incident radiation is dispersed in angle according to the Bragg equation $n\lambda = 2d \sin \theta$, where n is an integer representing the order of the reflection, λ is the wavelength of the incident radiation, d is the period of the multilayer or crystal structure, and θ is the angle of glancing incidence. For a crystal, d is the lattice spacing, the perpendicular distance between the successive planes of atoms contributing to the reflection. These planes are designated by their Miller indices (hkl) or, in the case of crystals belonging to the hexagonal group, $(hkil)$. (The value of $2d$ also represents the longest wavelength that the structure can diffract.)

For $2d$ values greater than about 25 Å, the choice of natural crystals is very limited, and those available (such as prochlorite) are likely to be small and of poor quality. Sputtered or evaporated multilayers can be used as dispersing elements at longer wavelengths. (Langmuir-Blodgett films have fallen into disfavor since the development of vacuum-deposited multilayers.)

Table 4-1 is a revision of one compiled by E. P. Bertin [1]. The crystals are arranged in order of increasing $2d$ spacing.

REFERENCE

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Table 4-1. Data for selected crystals used as dispersive elements in x-ray spectrometers and monochromators. The Miller indices $[(hkl)$, or $(hkil)$ for hexagonal crystals] are given for the diffracting planes parallel to the surface of the dispersive element. A question mark (?) indicates that the crystal is developmental and that the indices have not been ascertained. An asterisk following the indices indicates that literature references to this crystal without specification of (hkl) or $2d$ are likely to be references to this “cut.” The indicated useful wavelength region lies in the 2θ interval between 10° and 140° . The analyzer should be used outside this region in special cases only.

No.	Crystal	Miller indices	$2d$ (Å)	Chemical formula	Useful wavelength region (Å)	Applications, remarks
1	α -Quartz, silicon dioxide	(50 $\bar{5}$ 2)	1.624	SiO ₂	0.142–1.55	Shortest $2d$ of any practical crystal. Good for high-Z K -lines excited by 100-kV generators.
2	Lithium fluoride	(422)	1.652	LiF	0.144–1.58	Better than quartz (50 $\bar{5}$ 2) for the same applications.
3	Corundum, aluminum oxide	(146)	1.660	Al ₂ O ₃	0.145–1.58	Same applications as quartz (50 $\bar{5}$ 2)
4	Lithium fluoride	(420)	1.801	LiF	0.157–1.72	Similar to LiF (422).
5	Calcite, calcium carbonate	(633)	2.02	CaCO ₃	0.176–1.95	
6	α -Quartz, silicon dioxide	(22 $\bar{4}$ 3)	2.024	SiO ₂	0.177–1.96	
7	α -Quartz, silicon dioxide	(31 $\bar{4}$ 0)	2.3604	SiO ₂	0.205–2.25	Transmission-crystal optics.

8	α -Quartz, silicon dioxide	(22 $\bar{4}$ 0)	2.451	SiO ₂	0.213–2.37	
9	Topaz, hydrated aluminum fluorosilicate	(303)*	2.712	Al ₂ (F,OH) ₂ SiO ₄	0.236–2.59	Improves dispersion for V–Ni <i>K</i> -lines and rare earth <i>L</i> -lines.
10	Corundum, aluminum oxide, sapphire, alumina	(030)	2.748	Al ₂ O ₃	0.240–2.62	Diffacted intensity ~2–4X topaz (303) and quartz (203) with the same or better resolution.
11	α -Quartz, silicon dioxide	(20 $\bar{2}$ 3)	2.749	SiO ₂	0.240–2.62	Same applications as topaz (303) and LiF (220).

No.	Crystal	Miller indices	$2d$ (Å)	Chemical formula	Useful wavelength region (Å)	Applications, remarks
12	Topaz	(006)	2.795	Al ₂ (F,OH) ₂ SiO ₄	0.244–2.67	
13	Lithium fluoride	(220)	2.848	LiF	0.248–2.72	Same applications as topaz (303) and quartz (20 $\bar{2}$ 3), with 2–4X their diffracted intensity. Diffracted intensity ~0.4–0.8X LiF (200).
14	Mica, muscovite	(331)	3.00	K ₂ O·3Al ₂ O ₃ ·6SiO ₂ ·2H ₂ O	0.262–2.86	Transmission-crystal optics (Cauchois, DuMond types).
15	Calcite, calcium carbonate	(422)	3.034	CaCO ₃	0.264–2.93	
16	α -Quartz, silicon dioxide	(21 $\bar{3}$ 1)	3.082	SiO ₂	0.269–2.94	

17	α -Quartz, silicon dioxide	$(11\bar{2}2)$	3.636	SiO ₂	0.317–3.47	
18	Silicon	(220)	3.8403117	Si	0.335–3.66	Lattice period known to high accuracy.
19	Fluorite, calcium fluoride	(220)	3.862	CaF ₂	0.337–3.68	
20	Germanium	(220)	4.00	Ge	0.349–3.82	
21	Lithium fluoride	(200)*	4.027	LiF	0.351–384	Best general crystal for K <i>K</i> - to Lr <i>L</i> -lines. Highest intensity for largest number of elements of any crystal. Combines high intensity and high dispersion.
22	Aluminum	(200)	4.048	Al	0.353–3.86	Curved, especially doubly curved, optics.
23	α -Quartz, silicon dioxide	$(20\bar{2}0)$	4.246	SiO ₂	0.370–4.11	“Prism” cut.
24	α -Quartz, silicon dioxide	$(10\bar{1}2)$	4.564	SiO ₂	0.398–4.35	Used in prototype Laue multichannel spectrometer.
25	Topaz	(200)	4.638	Al ₂ (F,OH) ₂ SiO ₄	0.405–4.43	

No.	Crystal	Miller indices	$2d$ (Å)	Chemical formula	Useful wavelength region (Å)	Applications, remarks
26	Aluminum	(111)	4.676	Al	0.408–4.46	Curved, especially doubly curved, optics.
27	α -Quartz, silicon dioxide	$(11\bar{2}0)$	4.912	SiO ₂	0.428–4.75	

28	Gypsum, calcium sulfate dihydrate	(002)	4.990	$\text{CaSO}_4 \cdot 2\text{H}_2\text{O}$	0.435–4.76	Efflorescent: loses water in vacuum to become Plaster of Paris.
29	Rock salt, sodium chloride	(200)	5.641	NaCl	0.492–5.38	S $K\alpha$ and Cl $K\alpha$ in light matrixes. Like LiF (200), good general crystal for S K to Lr L .
30	Calcite, calcium carbonate	(200)	6.071	CaCO_3	0.529–5.79	Very precise wavelength measurements. Extremely high degree of crystal perfection with resultant sharp lines.
31	Ammonium dihydrogen phosphate (ADP)	(112)	6.14	$\text{NH}_4\text{H}_2\text{PO}_4$	0.535–5.86	
32	Silicon	(111)*	6.2712	Si	0.547–5.98	Very rugged and stable general-purpose crystal. High degree of perfection obtainable.
33	Sylvite, potassium chloride	(200)	6.292	KCl	0.549–6.00	
34	Fluorite, calcium fluoride	(111)	6.306	CaF_2	0.550–6.02	Very weak second order, strong third order.
35	Germanium	(111)*	6.532	Ge	0.570–6.23	Eliminates second order. Useful for intermediate- and low- Z elements where Ge $K\alpha$ emission is eliminated by pulse- height selection.
36	Potassium bromide	(200)	6.584	KBr	0.574–6.28	

37	α -Quartz, silicon dioxide	(10 $\bar{1}$ 0)	6.687	SiO ₂	0.583–6.38	P $K\alpha$ in low-Z matrixes, especially in calcium. Intensity for P–K K -lines greater than EDDT, but less than PET.
No.	Crystal	Miller indices	$2d$ (Å)	Chemical formula	Useful wavelength region (Å)	Applications, remarks
38	Graphite	(002)	6.708	C	0.585–6.40	P, S, Cl K -lines, P $K\alpha$ intensity > 5X EDDT. Relatively poor resolution but high integrated reflectivity.
39	Indium antimonide	(111)	7.4806	InSb	0.652–7.23	Important for K -edge of Si.
40	Ammonium dihydrogen phosphate (ADP)	(200)	7.5	NH ₄ H ₂ PO ₄	0.654–7.16	Higher intensity than EDDT.
41	Topaz	(002)	8.374	Al ₂ (F,OH) ₂ SiO ₄	0.730–7.99	
42	α -Quartz, silicon dioxide	(10 $\bar{1}$ 0)	*8.512	SiO ₂	0.742–8.12	Same applications as EDDT and PET; higher resolution, but lower intensity.
43	Pentaerythritol (PET)	(002)	8.742	C(CH ₂ OH) ₄	0.762–8.34	Al, Si, P, S, Cl $K\alpha$. Intensities ~1.5-2X EDDT, ~2.5X KHP. Good general crystal for Al–Sc $K\alpha$. Soft; deteriorates with age and exposure to x-rays.

44	Ammonium tartrate	(?)	8.80	$(\text{CHOH})_2(\text{COONH}_4)_2$	0.767–8.4	
45	Ethylenediamine- <i>d</i> -tartrate (EDDT, EDdT, EDT)	(020)	8.808		0.768–8.40	Same applications as PET, but lower intensity, substantially lower thermal expansion coefficient. Rugged and stable.
46	Ammonium dihydrogen phosphate (ADP)	(101)*	10.640	$\text{NH}_4\text{H}_2\text{PO}_4$	0.928–10.15	Mg $K\alpha$. Same applications as PET, EDDT, but lower intensity.
47	Na β -alumina	(0004)	11.24	$\begin{array}{c} \text{NaAl}_{11}\text{O}_{17} \\ \text{NH}_2 - \text{CH}_2 - \text{CH}_2 - \text{NH}_2 \\ \qquad \qquad \qquad \\ \text{COOH} - (\text{CHOH})_2 - \text{COOH} \end{array}$	0.980–10.87	
48	Oxalic acid dihydrate	(001)	11.92	$(\text{COOH})_2 \cdot 2\text{H}_2\text{O}$	1.04–11.37	
49	Sorbitol hexaacetate (SHA)	(110)	13.98	$\begin{array}{c} \text{CHOH} - \text{CO} - \text{CH}_3 \\ \\ (\text{COH} - \text{CO} - \text{CH}_3)_4 \\ \\ \text{CHOH} - \text{CO} - \text{CH}_3 \end{array}$	1.22–13.34	Applications similar to ADP (101) and gypsum (020). High resolution; stable in vacuum. Available in small pieces only.

No.	Crystal	Miller indices	$2d$ (Å)	Chemical formula	Useful wavelength region (Å)	Applications, remarks
50	Rock sugar, sucrose	(001)	15.12	$C_{12}H_{22}O_{11}$	1.32–14.42	
51	Gypsum, calcium sulfate dihydrate	(020)*	15.185	$CaSO_4 \cdot 2H_2O$	1.32–14.49	Na $K\alpha$. Inferior to KHP, RHP, and beryl. Poor in vacuum (efflorescent).
52	Beryl	(10 $\bar{1}$ 0)	15.954	$3BeO \cdot Al_2O_3 \cdot 6SiO_2$	1.39–15.22	Difficult to obtain. Good specimens have $\lambda/\delta\lambda \sim 2500$ –3000 at 12 Å. $2d$ may vary among specimens.
53	Bismuth titanate	(040)	16.40	$Bi_2(TiO_3)_3$	1.43–15.65	
54	Mica, muscovite	(002)*	19.84	$K_2O \cdot 3Al_2O_3 \cdot 6SiO_2 \cdot 2H_2O$	1.73–18.93	Easy to obtain. Easily bent: good for curved-crystal spectrometers, spectrographs.
55	Silver acetate	(001)	20.0	CH_3COOAg	1.74–19.08	
56	Rock sugar, sucrose	(100)	20.12	$C_{11}H_{22}O_{11}$	1.75–19.19	
57	Na β -alumina	(0002)	22.49	$NaAl_{11}O_{17}$	1.96–21.74	
58	Thallium hydrogen phthalate (THP, TIHP, TAP, TIAP)	(100)	25.9	$THC_8H_4O_4$	2.26–24.7	Same applications as KHP, RHP.

59	Rubidium hydrogen phthalate (RHP, RbHP, RAP, RbAP)	(100)	26.121	$\text{RbHC}_8\text{H}_4\text{O}_4$	2.28–24.92	Diffacted intensity ~3X KHP for Na, Mg, Al $K\alpha$ and Cu $L\alpha$; ~4X KHP for F $K\alpha$; ~8X KHP for O $K\alpha$
60	Potassium hydrogen phthalate (KHP, KAP)	(100)	26.632	$\text{KHC}_8\text{H}_4\text{O}_4$	2.32–25.41	Good general crystal for all low-Z elements down to O.
61	Octadecyl hydrogen maleate (OHM)	(?)	63.5	$\text{CH}_3(\text{CH}_2)_{17}\text{OOC}(\text{CH})_2\text{COOH}$	5.54–60.6	Ultralong-wavelength region down to C $K\alpha$

X-Ray Data Booklet

Section 4.1 MULTILAYERS AND CRYSTALS

James H. Underwood

A. MULTILAYERS

By means of modern vacuum deposition technology, structures consisting of alternating layers of high- and low- Z materials, with individual layers having thicknesses of the order of nanometers, can be fabricated on suitable substrates. These structures act as multilayer interference reflectors for x-rays, soft x-rays, and extreme ultraviolet (EUV) light. Their high reflectivity and moderate energy bandwidth ($10 < E/\Delta E < 100$) make them a valuable addition to the range of optical components useful in instrumentation for EUV radiation and x-rays with photon energies from a few hundred eV to tens of keV. These multilayers are particularly useful as mirrors and dispersive elements on synchrotron radiation beamlines. They may be used to produce focal spots with micrometer-scale sizes and for applications such as fluorescent microprobing, microdiffraction, and microcrystallography. Multilayer reflectors also have a wide range of applications in the EUV region, where normal-incidence multilayer reflectors allow the construction of space telescopes and of optics for EUV lithography.

Ordinary mirrors, operating at normal or near-normal angles of incidence, do not work throughout the EUV and x-ray regions. The reason is the value of the complex refractive index, $n = 1 - \delta - i\beta$, which can be close to unity for all materials in this region; δ may vary between 10^{-3} in the EUV region to 10^{-6} in the x-ray region. The Fresnel equations at normal incidence show that the reflectivity $R^2 = [(n - 1)/(n + 1)]^2$ is very small. However, the more general Fresnel equations show that x-rays and EUV radiation can be reflected by mirrors at large angles of incidence. Glancing (or grazing) incidence is the term used when the rays make a small angle (a few degrees or less) with the mirror surface; the complement ($90^\circ - i$) of the optical angle of incidence i is called the glancing (or grazing) angle of incidence. Glancing-incidence mirrors are discussed in Section 4.2.

Although the normal-incidence intensity reflectivity R^2 of a surface might be 10^{-3} or 10^{-4} , the corresponding *amplitude* reflectivity $R = (n - 1)/(n + 1)$, the square root of the intensity reflectivity, will be 1/30 to 1/100. This implies that, if the reflections from 30–100 surfaces could be made to add in phase, a total reflectivity approaching unity could be obtained. This is the multilayer principle. As shown in Fig. 4-1, a multilayer reflector comprises a stack of materials having alternately high and low refractive indices. The thicknesses are adjusted so that the path length difference between reflections from successive layer pairs is equal to one wavelength. Hence, x-ray/EUV multilayers are approximately equivalent to the familiar “quarter-wave stacks” of visible-light coating technology. The equivalence is not exact, however, because of the absorption term β , which is usually negligible for visible-light multilayers. This absorption reduces the multilayer reflectivity below unity and requires the design to be optimized for highest reflectivity. Particularly critical is the value of Γ , which is the ratio of the thickness of the high- Z (high electron density, high $|n|$) layer to the total thickness of each layer pair. This optimization

normally requires modeling or simulation of the multilayer. Most thin-film calculation programs, even if designed for the visible region, will perform these calculations if given the right complex values of the refractive index. Alternatively, there are a number of on-line calculation programs available; links can be found at <http://www-cxro.lbl.gov/>.

Either elemental materials or compounds can be used to fabricate multilayer reflectors. The performance obtained from a multilayer depends largely on whether there exists a fortuitous combination of materials having the right refractive indices (good contrast, low absorption), whether these materials can be deposited in smooth thin layers, and whether they remain stable (low reactivity, low diffusion) in the deposited state. The roughness of the underlying substrate is also of prime importance; an rms roughness of the same order of magnitude as the layer

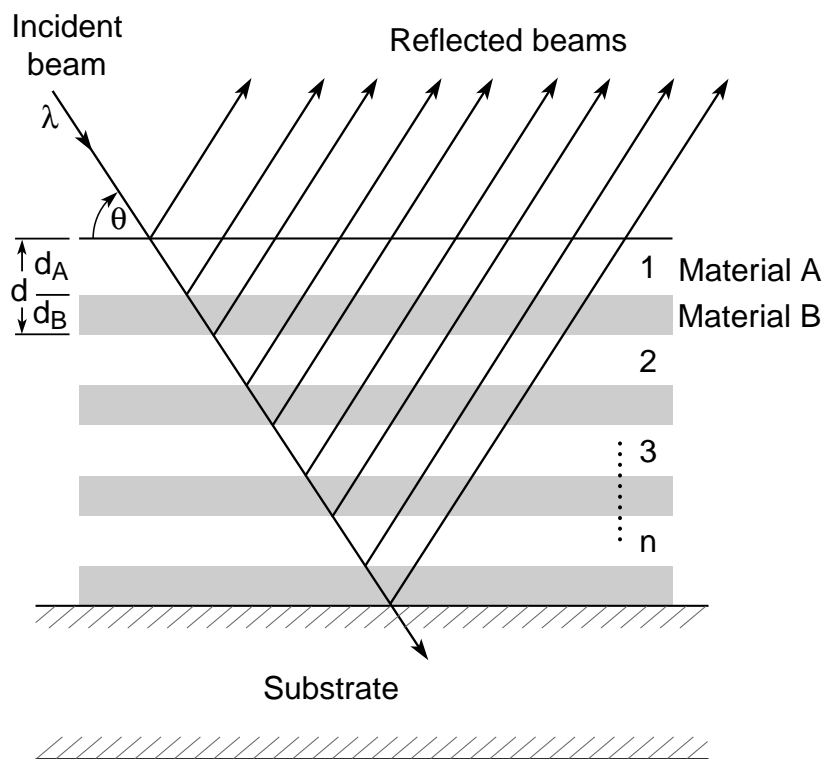


Fig. 4-1. Schematic of a multilayer reflector of n bilayer pairs. The parameters λ , θ , and d are chosen to satisfy the familiar Bragg equation, but the relative thicknesses of the high- and low- Z materials are also critical in optimizing reflectivity. The total reflectivity is the vector sum of the complex reflection coefficients at each interface, with the different path lengths taken into account.

thicknesses will spoil the performance of most coatings. “Superpolished” substrates, with roughness $\sigma \approx 0.1$ nm are preferred. On such substrates, the peak reflectivity of the coatings can approach 80% to 90% of theoretical predictions.

A large variety of multilayers have been made and their reflectivities measured over the years. A database of reported reflectivities has been assembled from surveys taken at the biennial Physics of X-ray Multilayer Structures conferences and can be found at the website listed above. For the EUV region around 100 eV, two remarkably successful combinations are molybdenum-silicon and molybdenum-beryllium. With Mo-Si, a normal-incidence reflectivity of 68% has been achieved at a wavelength of 13.4 nm; Mo-Be multilayers have achieved a reflectivity close

to 70% at 11.4 nm (see Fig. 4-2). These relatively high reflectivities are the basis for current efforts in the field of EUV lithography.

At hard x-ray wavelengths near 10 keV, a commonly used multilayer is made with tungsten as the high-Z material and boron carbide (B_4C) as the low-Z material. A reflectivity of 84% has been achieved with this combination at 8048 eV, the energy of the Cu $K\alpha$ line (see Fig. 4-3).

B. CRYSTALS

Multilayers are examples of a periodic structure that can be used to analyze short-wavelength electromagnetic radiation. Such structures split the incident beam into a large number N of separate beams; between any beam i and the beam $i + 1$, the optical path difference is constant. After leaving the periodic structure, the beams are recombined and caused to interfere, whereupon the spectrum of the incident radiation is produced.

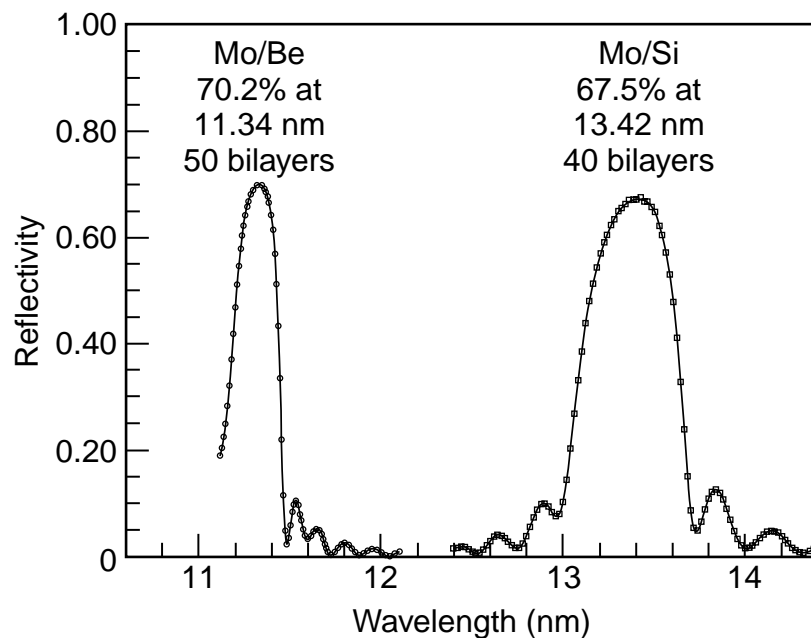


Fig. 4-2. The reflectivity of two multilayer reflectors at extreme ultraviolet wavelengths.

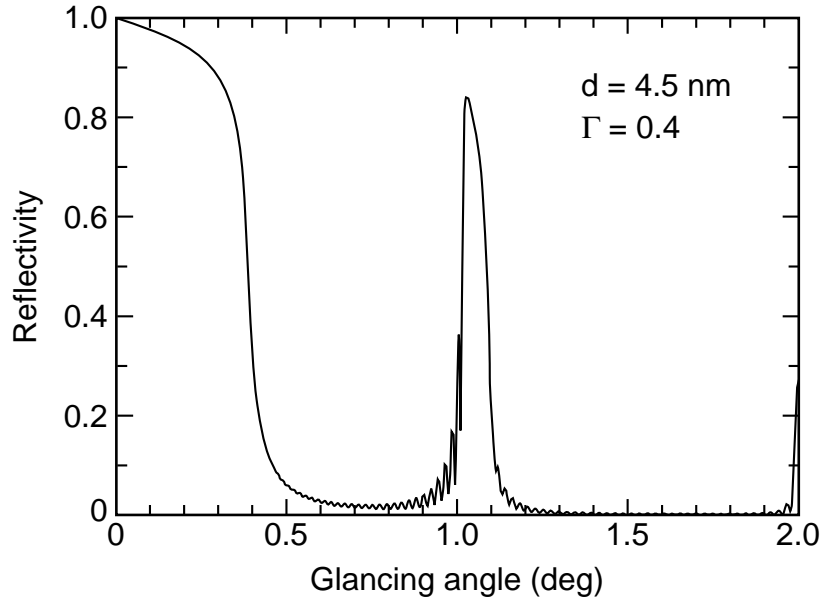


Fig. 4-3. The reflectivity of a tungsten–boron carbide multilayer at 8048 eV. The parameters d and Γ are discussed in the text.

Dispersion of radiation by a periodic structure is thus formally equivalent to multiple-beam interferometry. Structures that are periodic across their surface and that produce the N interfering beams by division of the incident wave front are called gratings and are treated in Section 4.3. Crystals and multilayer structures produce the N interfering beams by division of the incident *amplitude*. The spectrum of the incident radiation is dispersed in angle according to the Bragg equation $n\lambda = 2d \sin \theta$, where n is an integer representing the order of the reflection, λ is the wavelength of the incident radiation, d is the period of the multilayer or crystal structure, and θ is the angle of glancing incidence. For a crystal, d is the lattice spacing, the perpendicular distance between the successive planes of atoms contributing to the reflection. These planes are designated by their Miller indices $[hkl]$ or, in the case of crystals belonging to the hexagonal group, $[hkil]$. (The value of $2d$ also represents the longest wavelength that the structure can diffract.)

For $2d$ values greater than about 25 Å, the choice of natural crystals is very limited, and those available (such as prochlorite) are likely to be small and of poor quality. Sputtered or evaporated multilayers can be used as dispersing elements at longer wavelengths. (Langmuir-Blodgett films have fallen into disfavor since the development of vacuum-deposited multilayers.)

Table 4-1 is a revision of one compiled by E. P. Bertin [1]. The crystals are arranged in order of increasing $2d$ spacing.

REFERENCE

1. E. P. Bertin, "Crystals and Multilayer Langmuir-Blodgett Films Used as Analyzers in Wavelength-Dispersive X-Ray Spectrometers," in J. W. Robinson, Ed., *Handbook of Spectroscopy* (CRC Press, Cleveland, 1974), vol. 1, p. 238.

Table 4-1. Data for selected crystals used as dispersive elements in x-ray spectrometers and monochromators. The Miller indices $[hkl]$, or (hkl) for hexagonal crystals] are given for the diffracting planes parallel to the surface of the dispersive element. A question mark (?) indicates that the crystal is developmental and that the indices have not been ascertained. An asterisk following the indices indicates that literature references to this crystal without specification of (hkl) or $2d$ are likely to be references to this "cut." The indicated useful wavelength region lies in the 2θ interval between 10° and 140° . The analyzer should be used outside this region in special cases only.

No.	Crystal	Miller indices	$2d$ (Å)	Chemical formula	Useful wavelength region (Å)	Applications, remarks
1	α -Quartz, silicon dioxide	(50 $\bar{5}$ 2)	1.624	SiO ₂	0.142–1.55	Shortest $2d$ of any practical crystal. Good for high- Z K -lines excited by 100-KV generators.
2	Lithium fluoride	(422)	1.652	LiF	0.144–1.58	Better than quartz (50 $\bar{5}$ 2) for the same applications.
3	Corundum, aluminum oxide	(146)	1.660	Al ₂ O ₃	0.145–1.58	Same applications as quartz (50 $\bar{5}$ 2)
4	Lithium fluoride	(420)	1.801	LiF	0.157–1.72	Similar to LiF (422).
5	Calcite, calcium carbonate	(633)	2.02	CaCO ₃	0.176–1.95	
6	α -Quartz, silicon dioxide	(22 $\bar{4}$ 3)	2.024	SiO ₂	0.177–1.96	
7	α -Quartz, silicon dioxide	(31 $\bar{4}$ 0)	2.3604	SiO ₂	0.205–2.25	Transmission-crystal optics.
8	α -Quartz, silicon dioxide	(22 $\bar{4}$ 0)	2.451	SiO ₂	0.213–2.37	
9	Topaz, hydrated aluminum fluorosilicate	(303)*	2.712	Al ₂ (F,OH) ₂ SiO ₄	0.236–2.59	Improves dispersion for V–Ni K -lines and rare earth L -lines.
10	Corundum, aluminum oxide, sapphire, alumina	(030)	2.748	Al ₂ O ₃	0.240–2.62	Diffracted intensity ~2–4X topaz (303) and quartz (203) with the same or better resolution.
11	α -Quartz, silicon dioxide	(20 $\bar{2}$ 3)	2.749	SiO ₂	0.240–2.62	Same applications as topaz (303) and LiF (220).

Table 4-1. Selected data for crystals (continued).

No.	Crystal	Miller indices	2d (Å)	Chemical formula	Useful wavelength region (Å)	Applications, remarks
12	Topaz	(006)	2.795	Al ₂ (F,OH) ₂ SiO ₄	0.244–2.67	Same applications as topaz (303) and quartz (2023), with 2–4X their diffracted intensity. Diffracted intensity ~0.4–0.8X LiF (200).
13	Lithium fluoride	(220)	2.848	LiF	0.248–2.72	
14	Mica, muscovite	(331)	3.00	K ₂ O·3Al ₂ O ₃ ·6SiO ₂ ·2H ₂ O	0.262–2.86	
15	Calcite, calcium carbonate	(422)	3.034	CaCO ₃	0.264–2.93	Transmission-crystal optics (Cauchois, DuMond types).
16	α-Quartz, silicon dioxide	(21 $\bar{3}$ 1)	3.082	SiO ₂	0.269–2.94	Lattice period known to high accuracy.
17	α-Quartz, silicon dioxide	(11 $\bar{2}$ 2)	3.636	SiO ₂	0.317–3.47	
18	Silicon	(220)	3.8403117	Si	0.335–3.66	
19	Fluorite, calcium fluoride	(220)	3.862	CaF ₂	0.337–3.68	
20	Germanium	(220)	4.00	Ge	0.349–3.82	Best general crystal for K <i>K</i> - to Lr <i>L</i> -lines. Highest intensity for largest number of elements of any crystal. Combines high intensity and high dispersion.
21	Lithium fluoride	(200)*	4.027	LiF	0.351–3.84	
22	Aluminum	(200)	4.048	Al	0.353–3.86	
23	α-Quartz, silicon dioxide	(20 $\bar{2}$ 0)	4.246	SiO ₂	0.370–4.11	Curved, especially doubly curved, optics.
24	α-Quartz, silicon dioxide	(10 $\bar{1}$ 2)	4.564	SiO ₂	0.398–4.35	“Prism” cut.
25	Topaz	(200)	4.638	Al ₂ (F,OH) ₂ SiO ₄	0.405–4.43	Used in prototype Laue multichannel spectrometer.

Table 4-1. *Selected data for crystals (continued).*

No.	Crystal	Miller indices	2d (Å)	Chemical formula	Useful	Applications, remarks
					wavelength region (Å)	
26	Aluminum	(111)	4.676	Al	0.408–4.46	Curved, especially doubly curved, optics.
27	α -Quartz, silicon dioxide	(11 $\bar{2}$ 0)	4.912	SiO ₂	0.428–4.75	Efflorescent: loses water in vacuum to become Plaster of Paris.
28	Gypsum, calcium sulfate dihydrate	(002)	4.990	CaSO ₄ ·2H ₂ O	0.435–4.76	
29	Rock salt, sodium chloride	(200)	5.641	NaCl	0.492–5.38	S $K\alpha$ and Cl $K\alpha$ in light matrixes. Like LiF (200), good general crystal for S K to Lr L .
30	Calcite, calcium carbonate	(200)	6.071	CaCO ₃	0.529–5.79	Very precise wavelength measurements. Extremely high degree of crystal perfection with resultant sharp lines.
31	Ammonium dihydrogen phosphate (ADP)	(112)	6.14	NH ₄ H ₂ PO ₄	0.535–5.86	Very rugged and stable general-purpose crystal. High degree of perfection obtainable.
32	Silicon	(111)*	6.2712	Si	0.547–5.98	
33	Sylvite, potassium chloride	(200)	6.292	KCl	0.549–6.00	Very weak second order, strong third order.
34	Fluorite, calcium fluoride	(111)	6.306	CaF ₂	0.550–6.02	
35	Germanium	(111)*	6.532	Ge	0.570–6.23	Eliminates second order. Useful for intermediate- and low- Z elements where Ge $K\alpha$ emission is eliminated by pulse-height selection.
36	Potassium bromide	(200)	6.584	KBr	0.574–6.28	P $K\alpha$ in low- Z matrixes, especially in calcium. Intensity for P–K K -lines greater than EDDT, but less than PET.
37	α -Quartz, silicon dioxide	(10 $\bar{1}$ 0)	6.687	SiO ₂	0.583–6.38	

Table 4-1. Selected data for crystals (continued).

No.	Crystal	Miller indices	2d (Å)	Chemical formula	Useful wavelength region (Å)	Applications, remarks
38	Graphite	(002)	6.708	C	0.585–6.40	P, S, Cl <i>K</i> -lines, P <i>Kα</i> intensity > 5X EDDT. Relatively poor resolution but high integrated reflectivity.
39	Indium antimonide	(111)	7.4806	InSb	0.652–7.23	Important for <i>K</i> -edge of Si.
40	Ammonium dihydrogen phosphate (ADP)	(200)	7.5	NH ₄ H ₂ PO ₄	0.654–7.16	Higher intensity than EDDT.
41	Topaz	(002)	8.374	Al ₂ (F,OH) ₂ SiO ₄	0.730–7.99	Same applications as EDDT and PET; higher resolution, but lower intensity.
42	α-Quartz, silicon dioxide	(10 $\bar{1}$ 0)*	8.512	SiO ₂	0.742–8.12	
43	Pentaerythritol (PET)	(002)	8.742	C(CH ₂ OH) ₄	0.762–8.34	
44	Ammonium tartrate	(?)	8.80	(CHOH) ₂ (COONH ₄) ₂	0.767–8.4	Same applications as PET, but lower intensity, substantially lower thermal expansion coefficient. Rugged and stable.
45	Ethylenediamine- <i>d</i> -tartrate (EDDT, EDDT, EDT)	(020)	8.808	NH ₂ CH ₂ – CH ₂ – NH ₂ COOH – (CHOH) ₂ – COOH	0.768–8.40	
46	Ammonium dihydrogen phosphate (ADP)	(101)*	10.640	NH ₄ H ₂ PO ₄	0.928–10.15	
47	Na β-alumina	(0004)	11.24	NaAl ₁₁ O ₁₇	0.980–10.87	Applications similar to ADP (101) and gypsum (020). High resolution; stable in vacuum. Available in small pieces only.
48	Oxalic acid dihydrate	(001)	11.92	(COOH) ₂ ·2H ₂ O	1.04–11.37	
49	Sorbitol hexacetate (SHA)	(110)	13.98	CHOH – CO – CH ₃ (COH – CO – CH ₃) ₄ CHOH – CO – CH ₃	1.22–13.34	

Table 4-1. Selected data for crystals (continued).

No.	Crystal	Miller indices	$2d$ (Å)	Chemical formula	Useful wavelength region (Å)	Applications, remarks
50	Rock sugar, sucrose	(001)	15.12	$C_{12}H_{22}O_{11}$	1.32–14.42	
51	Gypsum, calcium sulfate dihydrate	(020)*	15.185	$CaSO_4 \cdot 2H_2O$	1.32–14.49	Na $K\alpha$. Inferior to KHP, RHP, and beryl. Poor in vacuum (efflorescent). Difficult to obtain. Good specimens have $\lambda/\delta\lambda \sim 2500$ –3000 at 12 Å. $2d$ may vary among specimens.
52	Beryl	(10 $\bar{1}$ 0)	15.954	$3BeO \cdot Al_2O_3 \cdot 6SiO_2$	1.39–15.22	
53	Bismuth titanate	(040)	16.40	$Bi_2(TiO_3)_3$	1.43–15.65	
54	Mica, muscovite	(002)*	19.84	$K_2O \cdot 3Al_2O_3 \cdot 6SiO_2 \cdot 2H_2O$	1.73–18.93	Easy to obtain. Easily bent: good for curved-crystal spectrometers, spectrographs.
55	Silver acetate	(001)	20.0	CH_3COOAg	1.74–19.08	
56	Rock sugar, sucrose	(100)	20.12	$C_{11}H_{22}O_{11}$	1.75–19.19	
57	Na β -alumina	(0002)	22.49	$NaAl_{11}O_{17}$	1.96–21.74	
58	Thallium hydrogen phthalate (THP, TIHP, TAP, TIAP)	(100)	25.9	$THC_8H_4O_4$	2.26–24.7	Same applications as KHP, RHP.
59	Rubidium hydrogen phthalate (RHP, RbHP, RAP, RbAP)	(100)	26.121	$RbHC_8H_4O_4$	2.28–24.92	Diffacted intensity $\sim 3X$ KHP for Na, Mg, Al $K\alpha$ and Cu $L\alpha$, $\sim 4X$ KHP for F $K\alpha$, $\sim 8X$ KHP for O $K\alpha$
60	Potassium hydrogen phthalate (KHP, KAP)	(100)	26.632	$KHC_8H_4O_4$	2.32–25.41	Good general crystal for all low- Z elements down to O.
61	Octadecyl hydrogen maleate (OHM)	(?)	63.5	$CH_3(CH_2)_{17}OOC(CH)_2COOH$	5.54–60.6	Ultralong-wavelength region down to C $K\alpha$

X-Ray Data Booklet

Section 4.2 Specular Reflectivities for Grazing-Incidence Mirrors

Eric M. Gullikson

The specular reflectivity of six common materials is given in Figs. 4-4 and 4-5 for photon energies between 30 eV and 30 keV. The reflectivity for a perfectly smooth surface and for s-polarization is

$$R = |r|^2 = \left| \frac{k_{iz} - k_{tz}}{k_{iz} + k_{tz}} \right|^2, \quad (1)$$

where

$$k_{iz} = \frac{2\pi}{\lambda} \sin \theta \quad \text{and} \quad k_{tz} = \frac{2\pi}{\lambda} \sqrt{n^2 - \cos^2 \theta}.$$

The grazing angle θ is measured from the plane of the mirror surface. The normal components of the incident and transmitted wave vectors are k_{iz} and k_{tz} , respectively. The complex index of refraction n is obtained from the average atomic scattering factor of the material, as described in Section 1.7 and in Ref. 1.

The effect of high-spatial-frequency roughness on the reflection coefficient of an interface can be approximated by the multiplicative factor

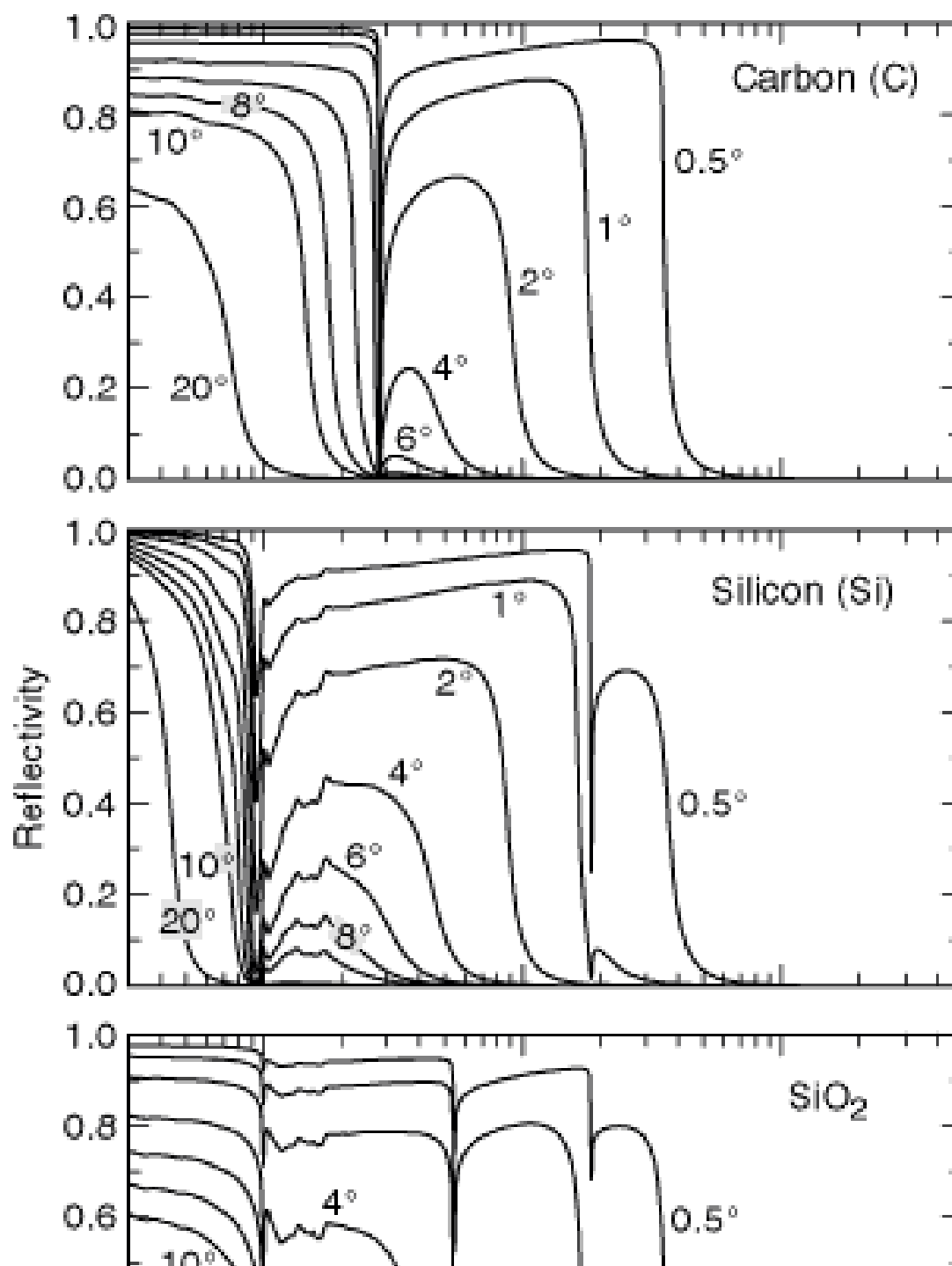
$$r = r_0 \exp(-2k_{iz}k_{tz}\sigma^2), \quad (2)$$

where r_0 is the complex reflection coefficient of a perfectly smooth interface and

σ is the rms roughness. For updated values of the atomic scattering factors and for on-line reflectivity calculations, see http://www-cxro.lbl.gov/optical_constants/.

REFERENCE

1. B. L. Henke, E. M. Gullikson, and J. C. Davis, "X-Ray Interactions: Photoabsorption, Scattering, Transmission, and Reflection at $E = 50\text{--}30,000$ eV, $Z = 1\text{--}92$," *At. Data Nucl. Data Tables* **54**, 181 (1993).



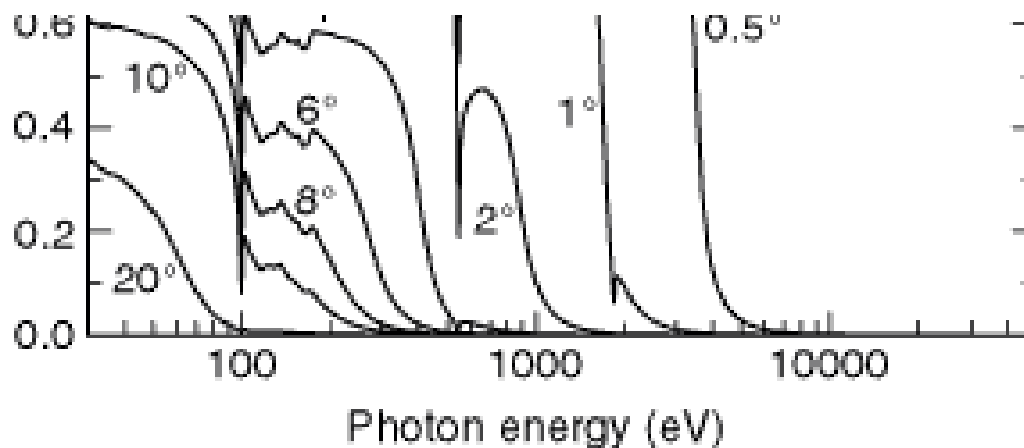


Fig. 4-4. Specular reflectivities of carbon ($\rho = 2.2 \text{ g/cm}^3$), silicon ($\rho = 2.33 \text{ g/cm}^3$), and silicon dioxide ($\rho = 2.2 \text{ g/cm}^3$). The reflectivity is calculated for *s*-polarization at grazing angles of 0.5, 1, 2, 4, 6, 8, 10, and 20 degrees.

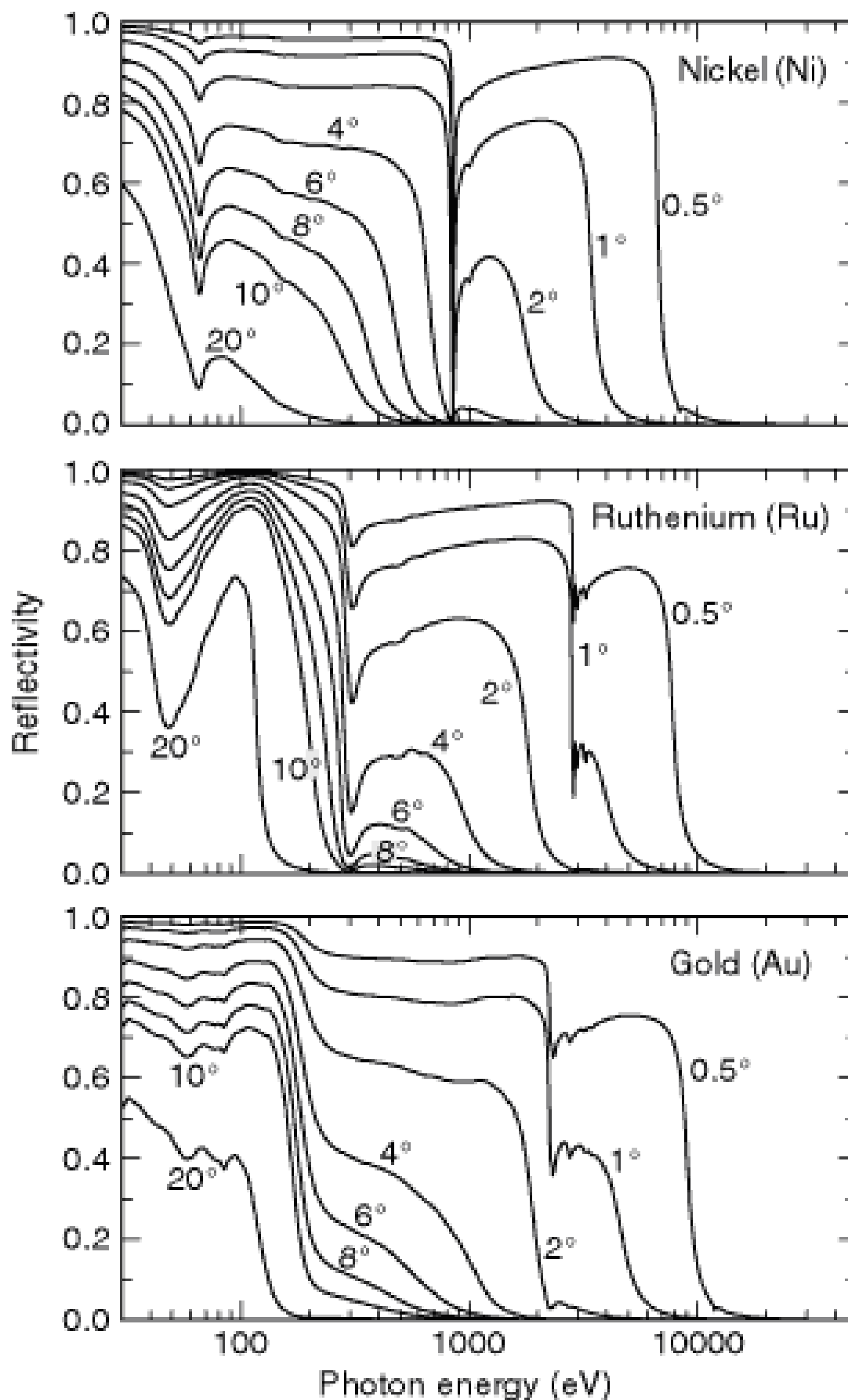


Fig. 4-5. Specular reflectivities of nickel ($\rho = 8.90 \text{ g/cm}^3$), ruthenium (ρ

= 12.41 g/cm³), and gold ($\rho = 19.3$ g/cm³). The reflectivity is calculated for s-polarization at grazing angles of 0.5, 1, 2, 4, 6, 8, 10, and 20 degrees.

X-Ray Data Booklet

Section 4.3 Gratings and Monochromators

Malcolm R. Howells

Both the [original version printed in the x-ray data booklet](#) and an [extended version with more details and references](#) of this section are available in PDF format.

X-Ray Data Booklet

Section 4.3 GRATINGS AND MONOCHROMATORS

Malcolm R. Howells

A. DIFFRACTION PROPERTIES

A.1 Notation and sign convention

We adopt the notation of Fig. 4-6, in which **a** and **b** have opposite signs if they are on opposite sides of the normal.

A.2 Grating equation

The grating equation may be written

$$m\lambda = d_0(\sin\alpha + \sin\beta) \quad . \quad (1)$$

The angles **a** and **b** are both arbitrary, so it is possible to impose various conditions relating them. If this is done, then for each **l**, there will be a unique **a** and **b**. The following conditions are used:

(i) *On-blaze condition:*

$$\alpha + \beta = 2\theta_B \quad , \quad (2)$$

where θ_B is the blaze angle (the angle of the sawtooth). The grating equation is then

$$m\lambda = 2d_0 \sin \theta_B \cos(\beta + \theta_B) \quad . \quad (3)$$

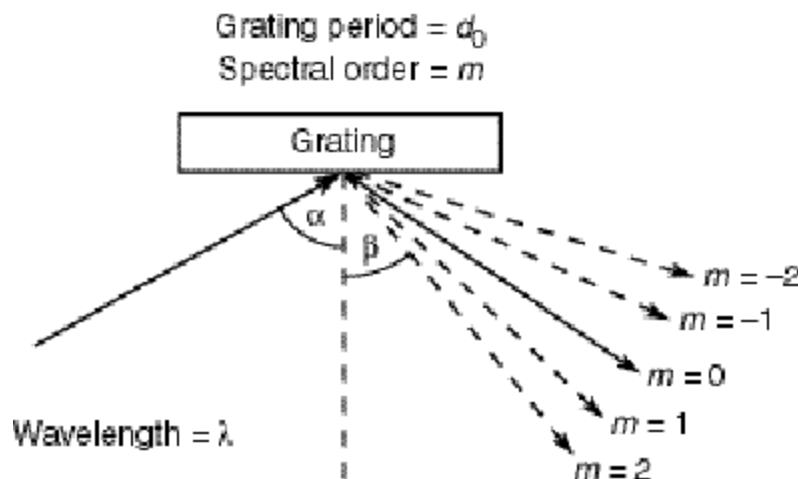


Fig. 4-6. Grating equation notation.

(ii) Fixed in and out directions:

$$\alpha - \beta = 2\theta \quad , \quad (4)$$

where 2θ is the (constant) included angle. The grating equation is then

$$m\lambda = 2d_0 \cos \theta \sin(\theta + \beta) \quad . \quad (5)$$

In this case, the wavelength scan ends when α or β reaches 90° , which occurs at the horizon wavelength $\lambda_H = 2d_0 \cos^2 \theta$.

(iii) *Constant incidence angle*: Equation (1) gives β directly.

(iv) *Constant focal distance (of a plane grating)*:

$$\frac{\cos \beta}{\cos \alpha} = \text{a constant } c_{ff} \quad , \quad (6)$$

leading to a grating equation

$$1 - \left(\frac{m\lambda}{d} - \sin \beta \right)^2 = \frac{\cos^2 \beta}{c_{ff}^2} \quad . \quad (7)$$

Equations (3), (5), and (7) give β (and thence α) for any λ . Examples of the above α - β relationships are (for references see <http://www-cxro.lbl.gov/>):

- (i) Kunz et al. plane-grating monochromator (PGM), Hunter et al. double PGM, collimated-light SX700 PGM
- (ii) Toroidal-grating monochromators (TGMs), spherical-grating monochromators (SGMs, “Dragon” system), Seya-Namioka, most aberration-reduced holographic SGMs, variable-angle SGM, PGMs
- (iii) Spectrographs, “Grasshopper” monochromator
- (iv) Standard SX700 PGM and most variants

B. FOCUSING PROPERTIES

The study of diffraction gratings (for references see <http://www-cxro.lbl.gov/>) goes back more than a century and has included plane, spherical [1], toroidal, and ellipsoidal surfaces and groove patterns made by classical (“Rowland”) ruling [2], holography [3,4], and variably spaced ruling [5,6]. In recent years the optical design possibilities of holographic groove patterns and variably spaced rulings have been extensively developed. Following normal practice, we provide an analysis of the imaging properties of gratings by means of the path function F [7]. For this purpose we use the notation of Fig. 4-7, in which the zeroth groove (of width d_0) passes through the grating pole O, while the n th groove passes through the variable point $P(\xi, w, l)$. The holographic groove pattern is taken to be made using two coherent point sources C and D with cylindrical polar coordinates (r_C, ϕ_C, z_C) , (r_D, ϕ_D, z_D) relative to O. The lower (upper) sign in Eq. (9) refers to C and D both real or both virtual (one real and one virtual), for which case the equiphase surfaces are confocal hyperboloids (ellipses) of revolution about

CD. Gratings with varied line spacing $d(w)$ are assumed to be ruled according to $d(w) = d_0(1 + v_1 w + v_2 w^2 + \dots)$.

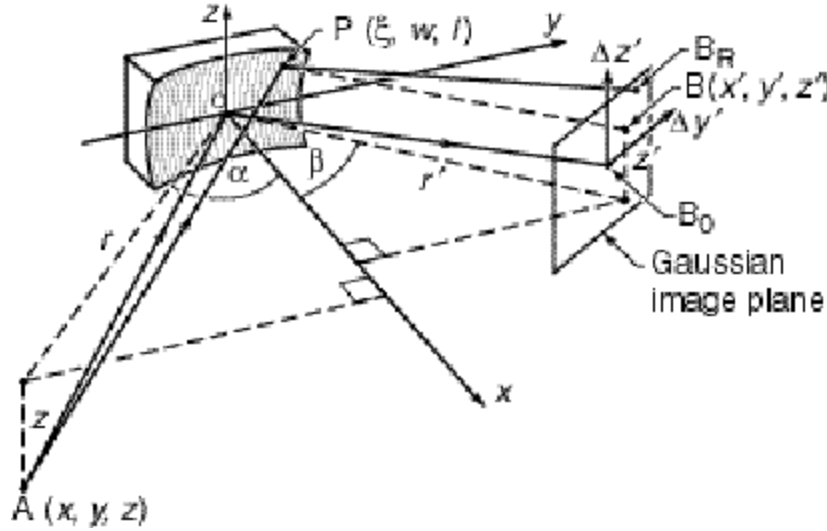


Fig. 4-7. *Focusing properties notation.*

We consider all the gratings to be ruled on the general surface

$$x = \sum_{ij} a_{ij} w^i l^j$$

and the a_{ij} coefficients are given below.

Ellipse coefficients a_{ij}

$$\begin{aligned} a_{20} &= \frac{\cos \theta}{4} \left(\frac{1}{r} + \frac{1}{r'} \right) & a_{12} &= \frac{a_{20} A}{\cos^2 \theta} \\ a_{30} &= a_{20} A & a_{22} &= \frac{a_{20} (2A^2 + C)}{2 \cos^2 \theta} \\ a_{40} &= \frac{a_{20} (4A^2 + C)}{4} & a_{04} &= \frac{a_{20} C}{8 \cos^2 \theta} \\ a_{02} &= \frac{a_{20}}{\cos^2 \theta} \end{aligned}$$

The other a_{ij} 's with $i + j \leq 4$ are zero. In the expressions above, r , r' , and θ are the object distance, image distance, and incidence angle to the normal, respectively, and

$$A = \frac{\sin \theta}{2} \left(\frac{1}{r} - \frac{1}{r'} \right), \quad C = A^2 + \frac{1}{rr'}, \quad .$$

Toroid coefficients a_{ij}

$$\begin{aligned} a_{20} &= \frac{1}{2R} & a_{22} &= \frac{1}{4\rho R^2} \\ a_{40} &= \frac{1}{8R^3} & a_{04} &= \frac{1}{8\rho^3} \\ a_{02} &= \frac{1}{2\rho} \end{aligned}$$

Other a_{ij} 's with $i + j \leq 4$ are zero. Here, R and ρ are the major and minor radii of the bicycle-tire toroid.

The a_{ij} 's for spheres; circular, parabolic, and hyperbolic cylinders; paraboloids; and hyperboloids can also be obtained from the values above by suitable choices of the input parameters r , r' , and q .

Values for the ellipse and toroid coefficients are given to sixth order at <http://www-cxro.lbl.gov/>.

B.1 Calculation of the path function F

F is expressed as

$$F = \sum_{ijk} F_{ijk} w^i l^j, \quad (8)$$

where

$$F_{ijk} = z^k C_{ijk}(\alpha, r) + z'^k C_{ijk}(\beta, r') + \frac{m\lambda}{d_0} f_{ijk}$$

and the f_{ijk} term, originating from the groove pattern, is given by one of the following expressions:

$$f_{ijk} = \begin{cases} 1 \text{ when } ijk = 100, 0 \text{ otherwise} & \text{Rowland} \\ \frac{d_0}{\lambda_0} \left[z_C^k C_{ijk}(\gamma, r_C) \pm z_D^k C_{ijk}(\delta, r_D) \right] & \text{holographic} \\ n_{ijk} & \text{varied line spacing} \end{cases} \quad (9)$$

The coefficient F_{ijk} is related to the strength of the i, j aberration of the wavefront diffracted by the grating. The coefficients C_{ijk} and n_{ijk} are given below, where the following notation is used:

$$T = T(r, \alpha) = \frac{\cos^2 \alpha}{r} - 2a_{20} \cos \alpha \quad (10a)$$

and

$$S = S(r, \alpha) = \frac{1}{r} - 2a_{02} \cos \alpha. \quad (10b)$$

Coefficients C_{ijk} of the expansion of F

$$\begin{aligned}
C_{011} &= -\frac{1}{r} & C_{020} &= \frac{S}{2} & C_{022} &= -\frac{S}{4r^2} - \frac{1}{2r^3} \\
C_{031} &= \frac{S}{2r^2} & C_{040} &= \frac{4a_{02}^2 - S^2}{8r} - a_{04} \cos \alpha \\
C_{100} &= -\sin \alpha & C_{102} &= \frac{\sin \alpha}{2r^2} \\
C_{111} &= -\frac{\sin \alpha}{r^2} & C_{120} &= \frac{S \sin \alpha}{2r} - a_{12} \cos \alpha \\
C_{200} &= \frac{T}{2} & C_{202} &= -\frac{T}{4r^2} + \frac{\sin^2 \alpha}{2r^3} \\
C_{211} &= \frac{T}{2r^2} - \frac{\sin^2 \alpha}{r^3} & C_{300} &= -a_{30} \cos \alpha + \frac{T \sin \alpha}{2r} \\
C_{220} &= -a_{22} \cos \alpha + \frac{1}{4r} (4a_{20}a_{02} - TS - 2a_{12} \sin 2\alpha) + \frac{S \sin^2 \alpha}{2r^2} \\
C_{400} &= -a_{40} \cos \alpha + \frac{1}{8r} (4a_{20}^2 - T^2 - 4a_{30} \sin 2\alpha) + \frac{T \sin^2 \alpha}{2r^2}
\end{aligned}$$

The coefficients for which $i \leq 4, j \leq 4, k \leq 2, i + j + k \leq 4, j + k = \text{even}$ are included here.

Coefficients n_{ijk} of the expansion of F

$$\begin{aligned}
n_{ijk} &= 0 \quad \text{for } j, k \neq 0 \\
n_{100} &= 1 & n_{300} &= \frac{v_1^2 - v_2}{3} \\
n_{200} &= \frac{-v_1}{2} & n_{400} &= \frac{-v_1^3 + 2v_1v_2 - v_3}{4}
\end{aligned}$$

Values for C_{ijk} and n_{ijk} are given to sixth order at <http://www-cxro.lbl.gov/>.

B.2 Determination of the Gaussian image point

By definition the principal ray AOB_0 arrives at the Gaussian image point $B_0(r'_0, \beta_0, z'_0)$ in Fig. 4-7. Its direction is given by Fermat's principal, which implies $[F/w]_{w=0, l=0} = 0$ and $[F/l]_{w=0, l=0} = 0$, from which

$$\frac{m\lambda}{d_0} = \sin \alpha + \sin \beta_0 \tag{11a}$$

and

$$\frac{z}{r} + \frac{z'_0}{r'_0} = 0 \quad , \quad (11b)$$

which are the grating equation and the law of magnification in the vertical direction. The tangential focal distance r'_0 is obtained by setting the focusing term F_{200} equal to zero and is given by

$$T(r, \alpha) + T(r'_0, \beta_0) = \begin{cases} 0 & \text{Rowland} \\ -\frac{m\lambda}{\lambda_0} [T(r_C, \gamma) \pm T(r_D, \delta)] & \text{holographic} \\ \frac{v_1 m \lambda}{d_0} & \text{varied line spacing} \end{cases} \quad (12)$$

Equations (11) and (12) determine the Gaussian image point B_0 and, in combination with the sagittal focusing condition ($F_{020} = 0$), describe the focusing properties of grating systems under the paraxial approximation. For a Rowland spherical grating the focusing condition, Eq. (12), is

$$\left(\frac{\cos^2 \alpha}{r} - \frac{\cos \alpha}{R} \right) + \left(\frac{\cos^2 \beta}{r'_0} - \frac{\cos \beta}{R} \right) = 0 \quad , \quad (13)$$

which has important special cases: (i) plane grating, $R = \infty$, implying

$$r'_0 = -r \cos^2 \mathbf{a} / \cos^2 \mathbf{b} = -r / c_{\text{ff}}^2$$

so that the focal distance and magnification are fixed if c_{ff} is held constant; (ii) object and image on the Rowland circle, i.e., $r = R \cos \alpha$, $r'_0 = R \cos \beta$, $M = 1$; and (iii) $\mathbf{b} = 90^\circ$ (Wadsworth condition). The focal distances of TGMs and SGMs, with or without moving slits, are also determined using Eq. (13).

B.3 Calculation of ray aberrations

In an aberrated system, the outgoing ray will arrive at the Gaussian image plane at a point B_R displaced from the Gaussian image point B_0 by the ray aberrations $\Delta y'$ and $\Delta z'$ (Fig. 4-7). The latter are given by

$$\Delta y' = \frac{r'_0}{\cos \beta_0} \frac{\partial F}{\partial w} \quad , \quad \Delta z' = r'_0 \frac{\partial F}{\partial l} \quad , \quad (14)$$

where F is to be evaluated for $A = (r, \mathbf{a}, z)$ and $B = (r'_0, \beta_0, z'_0)$. By means of the expansion of F , these equations allow the ray aberrations to be calculated separately for each aberration type:

$$\Delta y'_{ijk} = \frac{r'_0}{\cos \beta_0} F_{ijk} i w^{i-1} l^j \quad , \quad \Delta z'_{ijk} = r'_0 F_{ijk} w^i j l^{j-1} \quad . \quad (15)$$

Moreover, provided the aberrations are not too large, they are additive, so that they may either reinforce or cancel.

C. DISPERSION PROPERTIES

Dispersion properties can be summarized by the following relations.

(i) *Angular dispersion:*

$$\left(\frac{\partial \lambda}{\partial \beta} \right)_{\alpha} = \frac{d \cos \beta}{m} . \quad (16)$$

(ii) *Reciprocal linear dispersion:*

$$\left(\frac{\partial \lambda}{\partial (\Delta y')} \right)_{\alpha} = \frac{d \cos \beta}{mr'} \equiv \frac{10^{-3} d [\text{\AA}] \cos \beta}{mr' [\text{m}]} \text{\AA/mm} . \quad (17)$$

(iii) *Magnification:*

$$M(\lambda) = \frac{\cos \alpha}{\cos \beta} \frac{r'}{r} . \quad (18)$$

(iv) *Phase-space acceptance (ϵ):*

$$\epsilon = N \Delta \lambda_{S1} = N \Delta \lambda_{S2} \quad (\text{assuming } S_2 = MS_1) , \quad (19)$$

where N is the number of participating grooves.

D. RESOLUTION PROPERTIES

The following are the main contributions to the width of the instrumental line spread function. An estimate of the total width is the vector sum.

(i) *Entrance slit (width S_1):*

$$\Delta \lambda_{S1} = \frac{S_1 d \cos \alpha}{mr} . \quad (20)$$

(ii) *Exit slit (width S_2):*

$$\Delta \lambda_{S2} = \frac{S_2 d \cos \beta}{mr'} . \quad (21)$$

(iii) *Aberrations (of perfectly made grating):*

$$\Delta \lambda_A = \frac{\Delta y' d \cos \beta}{mr'} = \frac{d}{m} \left(\frac{\partial F}{\partial w} \right) . \quad (22)$$

(iv) *Slope error ΔF (of imperfectly made grating):*

$$\Delta\lambda_{\text{SE}} = \frac{d(\cos\alpha + \cos\beta)\Delta\phi}{m} . \quad (23)$$

Note that, provided the grating is large enough, diffraction at the entrance slit always guarantees a coherent illumination of enough grooves to achieve the slit-width-limited resolution. In such case a diffraction contribution to the width need not be added to the above.

E. EFFICIENCY

The most accurate way to calculate grating efficiencies is by the full electromagnetic theory [8]. However, approximate scalar-theory calculations are often useful and, in particular, provide a way to choose the groove depth h of a laminar grating. According to Bennett, the best value of the groove-width-to-period ratio r is the one for which the usefully illuminated area of the groove bottom is equal to that of the top. The scalar-theory efficiency of a laminar grating with $r = 0.5$ is given by Franks et al. as

$$E_0 = \frac{R}{4} \left[1 + 2(1 - P) \cos\left(\frac{4\pi h \cos\alpha}{\lambda}\right) + (1 - P)^2 \right]$$

$$E_m = \begin{cases} \left[\frac{R}{m^2\pi^2} [1 - 2\cos Q^+ \cos(Q^- + \delta)] \right. \\ \quad \left. + \cos^2 Q^+ \right] & m = \text{odd} \\ \left[\frac{R}{m^2\pi^2} \cos^2 Q^+ \right] & m = \text{even} \end{cases} \quad (24)$$

where

$$P = \frac{4h \tan\alpha}{d_0} ,$$

$$Q^\pm = \frac{m\pi h}{d_0} (\tan\alpha \pm \tan\beta) ,$$

$$\delta = \frac{2\pi h}{\lambda} (\cos\alpha + \cos\beta) ,$$

and R is the reflectance at grazing angle $\sqrt{\alpha_G \beta_G}$, where

$$\alpha_G = \frac{\pi}{2} - |\alpha| \text{ and } \beta_G = \frac{\pi}{2} - |\beta| .$$

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X-Ray Data Booklet

4.3 GRATINGS AND MONOCHROMATORS

Malcolm R. Howells

A. DIFFRACTION PROPERTIES

A.1 Notation and sign convention

We adopt the notation of Fig. 4.6 in which α and β have opposite signs if they are on opposite sides of the normal.

A.2 Grating equation

The grating equation may be written

$$m\lambda = d_0(\sin \alpha + \sin \beta). \quad (1)$$

The angles α and β are both arbitrary, so it is possible to impose various conditions relating them. If this is done, then for each λ , there will be a unique α and β . The following conditions are used:

(I) *ON-BLAZE CONDITION:*

$$\alpha + \beta = 2\theta_B, \quad (2)$$

where θ_B is the blaze angle (the angle of the sawtooth). The grating equation is then

$$m\lambda = 2d_0 \sin \theta_B \cos(\beta + \theta_B). \quad (3)$$

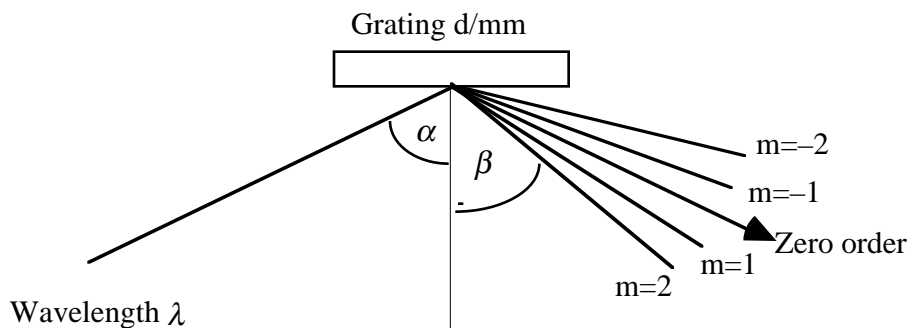


Fig. 4-6. Grating equation notation.

(II) *FIXED IN AND OUT DIRECTIONS:*

$$\alpha - \beta = 2\theta, \quad (4)$$

where 2θ is the (constant) included angle. The grating equation is then

$$m\lambda = 2d_0 \cos\theta \sin(\theta + \beta). \quad (5)$$

In this case, the wavelength scan ends when α or β reaches 90° , which occurs at the horizon wavelength $\lambda_H = 2d_0 \cos^2 \theta$

(III) *CONSTANT INCIDENCE ANGLE: EQUATION (1) GIVES β DIRECTLY.*(IV) *CONSTANT FOCAL DISTANCE (OF A PLANE GRATING):*

$$\frac{\cos\beta}{\cos\alpha} = \text{a constant } c_{ff}, \quad (6)$$

leading to a grating equation

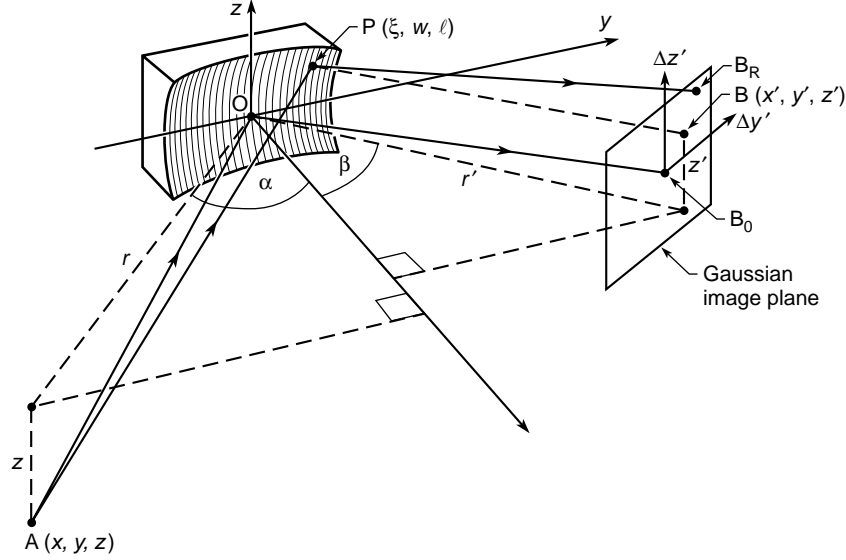
$$1 - \left(\frac{m\lambda}{d_0} - \sin\beta \right)^2 = \frac{\cos^2\beta}{c_{ff}^2} \quad (7)$$

Equations (3), (5), and (7) give β (and thence α) for any λ . Examples of the above α - β relationships are as follows:

- (i) Kunz et al. plane-grating monochromator (PGM) [1], Hunter et al. double PGM[2], collimated-light SX700 [3]
- (ii) Toroidal-grating monochromators (TGMs) [4, 5], spherical-grating monochromators (SGMs, “Dragon” system) [6], Seya-Namioka [7, 8] most aberration-reduced holographic SGMs [9], variable-angle SGM[10], PGMs [11, 12, 13]
- (iii) Spectrographs, “Grasshopper” monochromator [14]
- (iv) SX700 PGM [15] and variants [10, 16, 3]

B. FOCUSING PROPERTIES [17]

The study of diffraction gratings[18, 19] goes back more than a century and has included plane, spherical [20, 21, 22], toroidal [23] and ellipsoidal[24] surfaces and groove patterns made by classical (“Rowland”) ruling [25], holography [26, 27, 28] and variably-spaced ruling [29]. In recent years the optical design possibilities of holographic groove patterns [30, 31, 32] and variably-spaced rulings [13] have been extensively developed. Following normal practice, we provide an analysis of the imaging properties of gratings by means of the path function F [32]. For this purpose we use the notation of Fig. 4.7, in which the zeroth groove (of width d_0) passes through the grating pole O, while the n th groove passes through the variable point P(w, l). The holographic groove pattern is supposed to be made using two coherent point sources C and D with cylindrical polar coordinates (r_C, γ, z_C) , (r_D, δ, z_D) relative to O. The lower (upper) sign in eq. (9) refers



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Fig. 4-7*. Focusing properties notation.

to C and D both real or both virtual (one real and one virtual) for which case the equiphas surfaces are confocal hyperboloids (ellipses) of revolution about CD. The grating with varied line spacing $d(w)$ is assumed to be ruled according to $d(w) = d_0(1 + v_1 w + v_2 w^2 + \dots)$. We consider all the gratings to be ruled on the general surface $x = \sum_{ij} a_{ij} w^i l^j$ and the a_{ij} coefficients[33] are given for the important cases in Tables 1 and 2.

B.1 Calculation of the path function F

F is expressed as

$$F = \sum_{ijk} F_{ijk} w^i l^j \quad (8)$$

where $F_{ijk} = z^k C_{ijk}(\alpha, r) + z'^k C_{ijk}(\beta, r') + \frac{m\lambda}{d_0} f_{ijk}$.

and the f_{ijk} term, originating from the groove pattern, is given by one of the following expressions.

$$f_{ijk} = \begin{cases} \delta_{(i-1)jk} & \text{Rowland} \\ \frac{d_0}{\lambda_0} \{ z_C^k C_{ijk}(\gamma, r_C) \pm z_D^k C_{ijk}(\delta, r_D) \} & \text{holographic} \\ n_{ijk} & \text{varied line spacing} \end{cases} \quad (9)$$

The coefficient F_{ijk} is related to the strength of the i, j aberration of the wavefront diffracted by the grating. The coefficients C_{ijk} and n_{ijk} are given up to sixth order in Tables 3 and 4 in which the following notation is used:

$$T = T(r, \alpha) = \frac{\cos^2 \alpha}{r} - 2a_{20} \cos \alpha \quad S = S(r, \alpha) = \frac{1}{r} - 2a_{02} \cos \alpha \quad (10)$$

Table 1: Ellipse coefficients Q_{ij} from which the a_{ij} 's are obtained^a[33]

j	0	1	2	3	4	5	6
i							
0	0	0	1	0	$C/4$	0	$C^2/8$
1	0	0	A	0	$3AC/4$	0	*
2	1	0	$(2A^2 + C)/2$	0	$3C(4A^2 + C)/8$	0	*
3	A	0	$A(2A^2 + 3C)/2$	0	*	0	*
4	$(4A^2 + C)/4$	0	$(8A^4 + 24A^2C + 3C^2)/8$	0	*	0	*
5	$A(4A^2 + 3C)/4$	0	*	0	*	0	*
6	$(8A^4 + 12A^2C + C^2)/8$	0	*	0	*	0	*

^aIf r , r' and θ are the object distance, image distance, and incidence angle to the normal, respectively, then

$$a_{ij} = a_{20} \frac{Q_{ij}}{\cos^j \theta} \quad \text{where} \quad a_{20} = \frac{\cos \theta}{4} \left(\frac{1}{r} + \frac{1}{r'} \right), \quad A = \frac{\sin \theta}{2} \left(\frac{1}{r} - \frac{1}{r'} \right), \quad C = A^2 + \frac{1}{rr'}$$

The a_{ij} 's for spheres, circular cylinders, paraboloids and hyperboloids can also be obtained from Tables 1 and 2 by suitable choices of the input parameters r , r' and θ .

Table 2: Toroid^a a_{ij} 's [33]

j	0	1	2	3	4	5	6
i							
0	0	0	$1/(2\rho)$	0	$1/(8R^3)$	0	$1/(16\rho^5)$
1	0	0	0	0	0	0	*
2	$1/(2R)$	0	$1/(4\rho R^2)$	0	$(2\rho + R)/(16\rho^3 R^3)$	0	*
3	0	0	0	0	*	0	*
4	$1/(8R^3)$	0	$3/(16\rho R^4)$	0	*	0	*
5	0	0	*	0	*	0	*
6	$1/(16R^5)$	0	*	0	*	0	*

^a R and ρ are the major and minor radii of the bicycle-tire toroid we are considering.

Table 3: Coefficients C_{ijk} of the expansion of F^a

$C_{011} = -\frac{1}{r}$	$C_{020} = \frac{S}{2}$
$C_{022} = -\frac{S}{4r^2} - \frac{1}{2r^3}$	$C_{031} = \frac{S}{2r^2}$
$C_{040} = \frac{4a_{02}^2 - S^2}{8r} - a_{04} \cos \alpha$	$C_{042} = \frac{a_{04} \cos \alpha}{2r^2} + \frac{3S^2 - 4a_{02}^2}{16r^3} + \frac{3S}{4r^4}$
$C_{100} = -\sin \alpha$	$C_{102} = \frac{\sin \alpha}{2r^2}$
$C_{111} = -\frac{\sin \alpha}{r^2}$	$C_{120} = \frac{S \sin \alpha}{2r} - a_{12} \cos \alpha$
$C_{131} = -\frac{a_{12} \cos \alpha}{r^2} + \frac{3S \sin \alpha}{2r^3}$	$C_{122} = \frac{a_{12} \cos \alpha}{2r^2} - \frac{3S \sin \alpha}{4r^3} - \frac{3 \sin \alpha}{2r^4}$
$C_{200} = \frac{T}{2}$	$C_{140} = -a_{14} \cos \alpha + \frac{1}{2r} (2a_{02}a_{12} + a_{12}S \cos \alpha - a_{04} \sin 2\alpha) + \frac{\sin \alpha}{8r^2} (4a_{02}^2 - 3S^2)$
$C_{202} = -\frac{T}{4r^2} + \frac{\sin^2 \alpha}{2r^3}$	$C_{211} = \frac{T}{2r^2} - \frac{\sin^2 \alpha}{r^3}$
$C_{013} = \frac{1}{2r^3}$	$C_{300} = -a_{30} \cos \alpha + \frac{T \sin \alpha}{2r}$
$C_{220} = -a_{22} \cos \alpha + \frac{1}{4r} (4a_{20}a_{02} - TS - 2a_{12} \sin 2\alpha) + \frac{S \sin^2 \alpha}{2r^2}$	
$C_{222} = \frac{1}{2r^2} a_{22} \cos \alpha + \frac{1}{8r^3} (3ST - 4a_{02}a_{20} + 6a_{12} \sin 2\alpha) + \frac{3}{4r^4} (T - 2S \sin^2 \alpha) - \frac{3 \sin^2 \alpha}{r^5}$	
$C_{231} = -\frac{1}{r^2} a_{22} \cos \alpha + \frac{1}{4r^3} (-3ST + 4a_{02}a_{20} - 6a_{12} \sin 2\alpha) + \frac{3S \sin^2 \alpha}{r^4}$	
$C_{240} = -a_{24} \cos \alpha + \frac{1}{2r} (a_{12}^2 \sin^2 \alpha + 2a_{04}a_{20} + a_{22}S \cos \alpha + a_{04}T \cos \alpha - a_{14} \sin 2\alpha + 2a_{02}a_{22})$ $+ \frac{1}{16r^2} (-4a_{02}^2T - 8a_{02}a_{20}S + 12a_{12}S \sin 2\alpha + 3TS^2 + 16a_{02}a_{12} \sin \alpha - 8a_{04} \sin 2\alpha) + \frac{\sin^2 \alpha}{4r^3} (2a_{02}^2 - 3S^2)$	
$C_{302} = \frac{a_{30} \cos \alpha}{2r^2} - \frac{3T \sin \alpha}{4r^3} + \frac{\sin^3 \alpha}{2r^4}$	
$C_{311} = -\frac{a_{30} \cos \alpha}{r^2} + \frac{3T \sin \alpha}{2r^3} - \frac{\sin^3 \alpha}{r^4}$	

***Table 3:** Coefficients C_{ijk} of the expansion of F^a (continued)

$$\begin{aligned}
 C_{320} &= -a_{32} \cos \alpha + \frac{1}{2r} (2a_{20}a_{12} + 2a_{30}a_{02} + a_{30}S \cos \alpha + a_{12}T \cos \alpha - a_{22} \sin 2\alpha) \\
 &\quad + \frac{1}{4r^2} (4a_{20}a_{02} \sin \alpha - 3ST \sin \alpha - 4a_{12} \cos \alpha \sin^2 \alpha) + \frac{S \sin^3 \alpha}{2r^3} \\
 C_{400} &= -a_{40} \cos \alpha + \frac{1}{8r} (4a_{20}^2 - T^2 - 4a_{30} \sin 2\alpha) + \frac{T \sin^2 \alpha}{2r^2} \\
 C_{402} &= -\frac{1}{16r^3} (4a_{20}^2 + 3T^2 + 12a_{30} \sin 2\alpha) + \frac{a_{40} \cos \alpha}{2r^2} - \frac{3T \sin^2 \alpha}{2r^4} + \frac{\sin^4 \alpha}{2r^5} \\
 C_{411} &= -\frac{a_{40} \cos \alpha}{r^2} + \frac{1}{8r^3} (4a_{20}^2 - 3T^2 - 12a_{30} \sin 2\alpha) + \frac{3T \sin^2 \alpha}{r^4} - \frac{\sin^4 \alpha}{r^5} \\
 C_{420} &= -a_{42} \cos \alpha + \frac{1}{2r} (2a_{20}a_{22} + 2a_{12}a_{30} \sin^2 \alpha + 2a_{02}a_{40} - a_{32} \sin 2\alpha + a_{40}S \cos \alpha + a_{22}T \cos \alpha) \\
 &\quad + \frac{1}{16r^2} (-4a_{20}^2S - 8a_{02}a_{20}T + 3ST^2 + 12 \sin 2\alpha (a_{30}S + a_{12}T) + 8 \sin \alpha (2a_{02}a_{30} - 2a_{22} \sin 2\alpha + 2a_{12}a_{20})) \\
 &\quad + \frac{1}{2r^3} (2a_{02}a_{20} \sin^2 \alpha - 3ST \sin^2 \alpha - 2a_{12} \cos \alpha \sin^3 \alpha) + \frac{S \sin^4 \alpha}{2r^4} \\
 C_{500} &= -a_{50} \cos \alpha + \frac{1}{2r} (2a_{20}a_{30} + a_{30}T \cos \alpha - a_{40} \sin 2\alpha) + \frac{\sin \alpha}{2r^2} (a_{20}^2 - a_{30} \sin 2\alpha) - \frac{3T^2 \sin \alpha}{8r^2} + \frac{T \sin^3 \alpha}{2r^3} \\
 C_{600} &= -a_{60} \cos \alpha + \frac{1}{2r} (a_{30}^2 \sin^2 \alpha + 2a_{20}a_{40} + a_{40}T \cos \alpha - a_{50} \sin 2\alpha) \\
 &\quad + \frac{1}{16r^2} (-4a_{20}^2T + T^3 + 16a_{20}a_{30} \sin \alpha + 12a_{30}T \sin 2\alpha - 16a_{40} \cos \alpha \sin^2 \alpha) \\
 &\quad + \frac{1}{4r^3} (2a_{20}^2 \sin^2 \alpha - 3T^2 \sin^2 \alpha - 4a_{30} \cos \alpha \sin^3 \alpha) + \frac{T \sin^4 \alpha}{2r^4}
 \end{aligned}$$

^aThe coefficients for which $i \leq 6, j \leq 4, k \leq 2, i + j + k \leq 6, j + k = \text{even}$ are included in this table. The only addition to those is C_{013} , which has some interest, because, when the system is specialized to be symmetrical about the x axis, it represents a Seidel aberration, namely distortion.

Table 4: Coefficients n_{ijk} of the expansion of F for a grating with variable line spacing

$n_{ijk} = 0 \quad \text{for } j, k \neq 0$	
$n_{100} = 1$	$n_{400} = (-v_1^3 + 2v_1v_2 - v_3)/4$
$n_{200} = -v_1/2$	$n_{500} = (v_1^4 - 3v_1^2v_2 + v_2^2 + 2v_1v_3 - v_4)/5$
$n_{300} = (v_1^2 - v_2)/3$	$n_{600} = (-v_1^5 + 4v_1^3v_2 - 3v_1v_2^2 - 3v_1^2v_3 + 2v_2v_3 + 2v_1v_4 - v_5)/6$

B.2 Determination of the Gaussian image point

By definition the principal ray AOB_0 arrives at the Gaussian image point $B_0(r'_0, \beta_0, z'_0)$ (Fig. 4.7) and its direction is given by Fermat's principal which implies

$$[dF/dw]_{w=0, l=0} = 0, \quad [dF/dl]_{w=0, l=0} = 0 \quad \text{whence}$$

$$\frac{m\lambda}{d_0} = \sin \alpha + \sin \beta_0, \quad \frac{z}{r} + \frac{z'_0}{r'_0} = 0, \quad (11)$$

The tangential focal distance r'_0 is obtained by setting the focusing term F_{200} equal to zero and is given by

$$T(r, \alpha) + T(r'_0, \beta_0) = \begin{cases} 0 & \text{Rowland} \\ -\frac{m\lambda}{\lambda_0} \{T(r_C, \gamma) \pm T(r_D, \delta)\} & \text{holographic} \\ \frac{v_1 m \lambda}{d_0} & \text{varied line spacing} \end{cases} \quad (12)$$

Equations (11) and (12) determine the Gaussian image point B_0 , and in combination with the sagittal focusing condition ($F_{020}=0$), describe the focusing properties of grating systems under the paraxial approximation.

For a Rowland spherical grating the focusing condition (Eq. (12)) is

$$\left(\frac{\cos^2 \alpha}{r} - \frac{\cos \alpha}{R} \right) + \left(\frac{\cos^2 \beta}{r'_0} - \frac{\cos \beta}{R} \right) = 0 \quad (13)$$

which has important special cases. (i) plane grating, $R = \infty$ implying $r'_0 = -r \cos^2 \alpha / \cos^2 \beta$, (ii) object and image on the Rowland circle, or $r = R \cos \alpha$, $r'_0 = R \cos \beta$ and $M = 1$ and (iii) $\beta=90^\circ$ (Wadsworth condition). The focal distances of TGMs and SGMs are also determined by eq. (13).

B.3 Calculation of ray aberrations

In an aberrated system, the outgoing ray will arrive at the Gaussian image plane at a point B_R displaced from the Gaussian image point B_0 by the ray aberrations $\Delta y'$ and $\Delta z'$ (Fig. 4.7). The latter are given by [34, 35, 36]

$$\Delta y' = \frac{r'_0}{\cos \beta_0} \frac{\partial F}{\partial w}, \quad \Delta z' = r'_0 \frac{\partial F}{\partial l}, \quad (14)$$

where F is to be evaluated for $A = (r, \alpha, z)$, $B = (r'_0, \beta_0, z'_0)$. By means of the expansion of F , these equations allow the ray aberrations to be calculated separately for each aberration type.

$$\Delta y'_{ijk} = \frac{r'_0}{\cos \beta_0} F_{ijk} i w^{i-1} l^j, \quad \Delta z'_{ijk} = r'_0 F_{ijk} w^i j l^{j-1}. \quad (15)$$

Moreover, provided the aberrations are not too large, they are additive, so that they may either reinforce or cancel.

C. DISPERSION PROPERTIES

C.1 Angular dispersion

$$\left(\frac{\partial \lambda}{\partial \beta} \right)_\alpha = \frac{d \cos \beta}{m} \quad (16)$$

C.2 Reciprocal linear dispersion

$$\left(\frac{\partial \lambda}{\partial (\Delta y')} \right)_\alpha = \frac{d \cos \beta}{m r'} \equiv \frac{10^{-3} d [\text{\AA}] \cos \beta}{m r' [\text{m}]} \text{\AA}/\text{mm}, \quad (17)$$

C.3 Magnification (M)

$$M(\lambda) = \frac{\cos \alpha}{\cos \beta} \frac{r'}{r}. \quad (18)$$

C.4 Phase-space acceptance (ε)

$$\varepsilon = N \Delta \lambda_{S_1} = N \Delta \lambda_{S_2} \quad (\text{assuming } S_2 = M S_1) \quad (19)$$

where N is the number of participating grooves.

D. RESOLUTION PROPERTIES

The following are the main contributions to the width of the instrumental line spread function. The actual width is the vector sum.

(I) *ENTRANCE SLIT (WIDTH S_1):*

$$\Delta\lambda_{S_1} = \frac{S_1 d \cos \alpha}{mr}. \quad (20)$$

(II) *EXIT SLIT (WIDTH S_2):*

$$\Delta\lambda_{S_2} = \frac{S_2 d \cos \beta}{mr'}. \quad (21)$$

(III) *ABERRATIONS (OF PERFECTLY MADE GRATING):*

$$\Delta\lambda_A = \frac{\Delta y' d \cos \beta}{mr'} = \frac{d}{m} \left(\frac{\partial F}{\partial w} \right). \quad (22)$$

(IV) *SLOPE ERROR $\Delta\phi$ (OF IMPERFECTLY MADE GRATING):*

$$\Delta\lambda_{SE} = \frac{d(\cos \alpha + \cos \beta)\Delta\phi}{m}, \quad (23)$$

Note that, provided the grating is large enough, diffraction at the entrance slit always guarantees a coherent illumination of enough grooves to achieve the slit-width limited resolution and a diffraction contribution to the width need not be added to the above.

1. EFFICIENCY

The most accurate way to calculate grating efficiencies is by the full electromagnetic theory [37, 38] for which a code is available from Nevier. However, approximate scalar-theory calculations are often useful and, in particular, provide a way to choose the groove depth (h) of a laminar grating. According to Bennett [39], the best value of the groove-width-to-period ratio (r) is the one for which the usefully illuminated groove area is equal to the land area. The scalar theory efficiency of a laminar grating with $r=0.5$ is given by [40]

$$E_0 = \frac{R}{4} \left\{ 1 + 2(1-P) \cos\left(\frac{4\pi h \cos \alpha}{\lambda}\right) + (1-P)^2 \right\}$$

$$E_m = \begin{cases} \frac{R}{m^2 \pi^2} \left\{ 1 - 2 \cos Q^+ \cos(Q^- + \delta) + \cos^2 Q^+ \right\} & m = \text{odd} \\ \frac{R}{m^2 \pi^2} \cos^2 Q^+ & m = \text{even} \end{cases}$$

where

$$P = \frac{4h \tan \alpha}{d_0}, \quad Q^\pm = \frac{m\pi h}{d_0} (\tan \alpha \pm \tan \beta), \quad \delta = \frac{2\pi h}{\lambda} (\cos \alpha + \cos \beta) \quad (24)$$

and R is the reflectance at angle $\sqrt{\alpha\beta}$.

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X-Ray Data Booklet

Section 4.4 ZONE PLATES

Janos Kirz and David Attwood

A zone plate is a circular diffraction grating. In its simplest form, a transmission Fresnel zone plate lens consists of alternate transparent and opaque rings. The radii of the zone plate edges are given by

$$r_n^2 = n f \lambda + n^2 \lambda^2 / 4 \quad , \quad (1)$$

where n is the zone number (opaque and transparent zones count separately), λ is the wavelength, and f is the first-order focal length. The zone plate lens can be used to focus monochromatic, uniform plane wave (or spherical wave) radiation to a small spot, as illustrated in Fig. 4.8, or can be used for near-axis point-by-point construction of a full-field image, as illustrated in Fig. 4.9. When used in imaging applications, it obeys the thin-lens formula

$$\frac{1}{p} + \frac{1}{q} = \frac{1}{f} \quad , \quad (2)$$

where p and q are the object and image distances, respectively. Note that the second term on the right-hand side of Eq. (1) is correct only for $p \ll q$ or $p \gg q$, which are the common cases in zone plate applications. Descriptions of the diffractive properties of zone plate lenses, their use in x-ray microscopes, and extensive references to the literature are given elsewhere [1–3].

A zone plate lens is fully specified by three parameters. Most applications are dominated by the choice of photon energy $\hbar\omega$, and thus by λ . Resolution is set in large part by λ and the outer zone width, $\Delta r \equiv r_N - r_{N-1}$. To avoid chromatic blurring, the number of zones, N , must be less than the inverse relative spectral bandwidth, $\lambda/\Delta\lambda$. Thus, for many applications, λ , Δr , and N constitute a natural

set of zone plate–defining parameters. In terms of these three basic parameters, other zone plate parameters are given by [3]

$$D = 4N\Delta r \quad (3)$$

$$f = 4N(\Delta r)^2 / \lambda \quad (4)$$

$$F^\# = \Delta r / \lambda \quad (5)$$

$$NA = \lambda / 2\Delta r \quad (6)$$

To avoid chromatic blurring requires that

$$N < \lambda / \Delta \lambda \quad , \quad (7)$$

where $\Delta \lambda$ is the spectral width of the illuminating radiation. For uniform plane-wave illumination, as in Fig. 4-8, the Rayleigh criterion sets the diffraction-limited (λ, NA) spatial resolution of a perfect lens as

$$Res = \frac{0.610\lambda}{NA} = 1.22 \Delta r \quad , \quad (8)$$

where $NA = \sin\theta$ is the numerical aperture of the lens in vacuum and θ is the half-angle of the focused radiation. This resolution is obtained over an axial depth of focus given by

$$DOF = \pm \frac{\lambda}{2(NA)^2} \quad (9)$$

Spatial resolution in the full-field case can be improved somewhat from that given in Eq. (8), depending on both the object itself and the partial coherence of the illumination [4], as set by the parameter

$$\sigma = \frac{(NA)_c}{(NA)_o} \quad (10)$$

where $(NA)_c$ refers to the illumination numerical aperture of the condenser, and $(NA)_o$ is that of the zone plate objective lens, as given in Eq. (6).

The efficiency of the simple zone plate in the first order is ideally π^{-2} , or about 10%. The remainder of the radiation is absorbed (50%) or diffracted in other orders—zero order (25%), negative orders (12.5%), and higher positive orders (2.5%). If opaque zones are replaced by transparent but phase-shifting zones, efficiencies can be substantially improved [5]. To isolate the first order from unwanted orders, zone plate lenses are often made with a central stop, used in conjunction with a somewhat smaller collimating aperture near the focal region. This is particularly effective in scanning x-ray microscopes [6,7], where the zone plate is used to focus the radiation to a small spot, as in Fig. 4-8. For full-field x-ray microscopes, as suggested by Fig. 4-9, first-order imaging is assisted in a similar manner by stopping the central portion of the illuminating radiation [8].

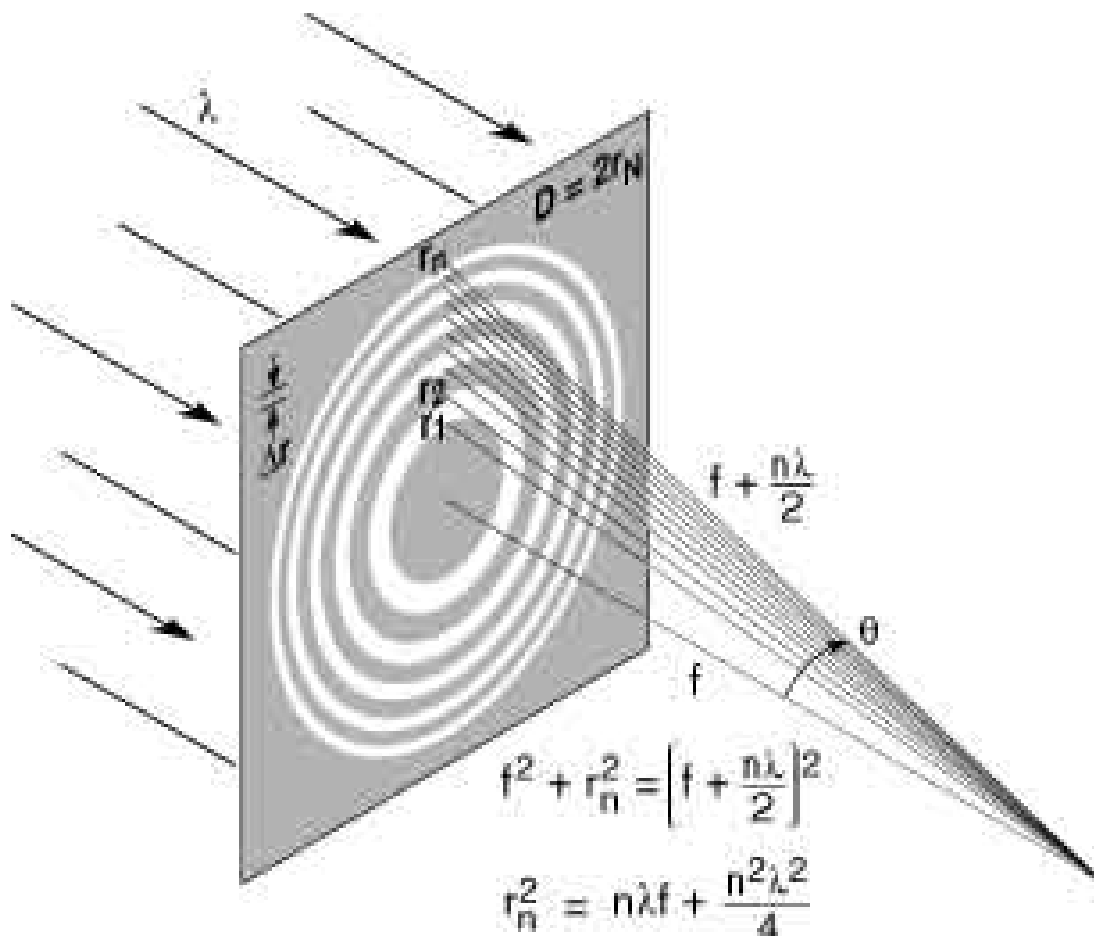


Fig. 4-8. A Fresnel zone plate lens with plane wave illumination, showing only the convergent (+1st) order of diffraction. Sequential zones of radius r_n are specified such that the incremental path length to the focal point is $n\lambda/2$. Alternate zones are opaque in the simple transmission zone plate. With a total number of zones, N , the zone plate lens is fully specified. Lens characteristics such as the focal length f , diameter D , and numerical aperture NA are described in terms of λ , N , and Δr , the outer zone width. [Courtesy of Cambridge University Press, Ref. 3.]

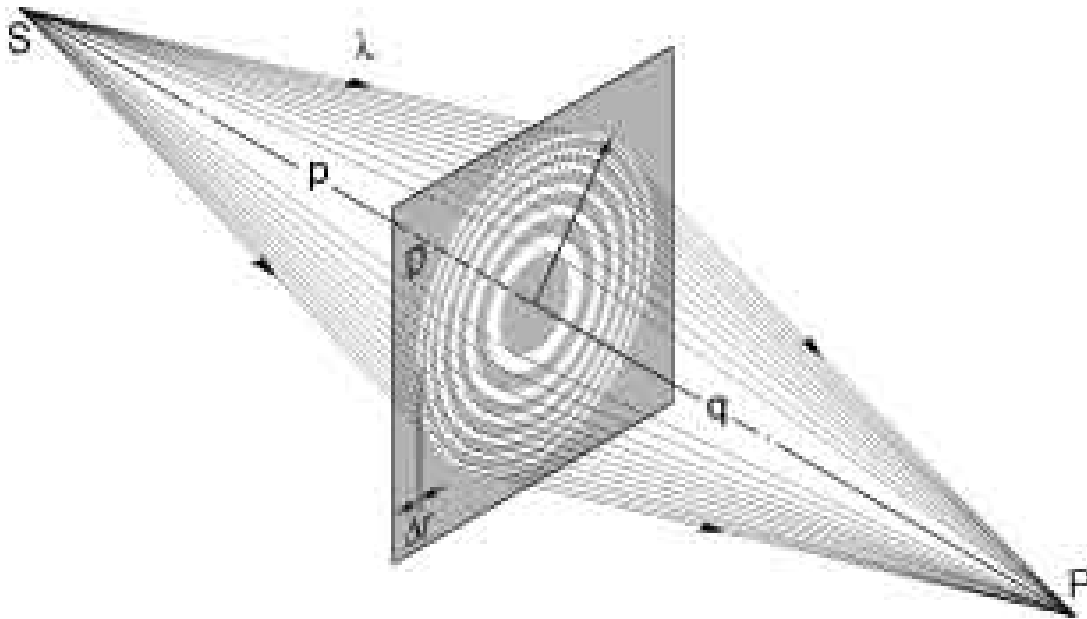


Fig. 4-9. A Fresnel zone plate used as a diffractive lens to form an x-ray image of a source point S in the image plane at P . The lens is shown as having a diameter D and outer zone width Δr . The object and image distances are p and q , respectively. A full-field image is formed concurrently in this manner. [Courtesy of Cambridge University Press, Ref. 3.]

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X-Ray Data Booklet

Section 4.5 X-RAY DETECTORS

Albert C. Thompson

A wide variety of x-ray detectors is available, some counting single photons, some providing only measurements of count rate or total flux, others measuring the energy, position, and/or incidence time of each x-ray. Table 4-2 provides typical values for useful energy range, energy resolution, dead time per event, and maximum count rate capability for common x-ray detectors. For special applications, these specifications can often be substantially improved.

Table 4-2. *Properties of common x-ray detectors;
 ΔE is measured as FWHM.*

Detector	Energy range (keV)	$\Delta E/E$ at 5.9 keV (%)	Dead time/event (μs)	Maximum count rate (s^{-1})
Gas ionization (current mode)	0.2–50	n/a	n/a	10^{11}a
Gas proportional	0.2–50	15	0.2	10^6
Multiwire and microstrip proportional	3–50	20	0.2	$10^6/\text{mm}^2$
Scintillation [NaI(Tl)]	3–10,000	40	0.25	2×10^6
Energy-resolving semiconductor	1–10,000	3	0.5–30	2×10^5
Surface-barrier (current mode)	0.1–20	n/a	n/a	10^8
Avalanche photodiode	0.1–50	20	0.001	10^8
CCD	0.1–70	n/a	n/a	n/a
Superconducting	0.1–4	< 0.5	100	5×10^3
Image plate	4–80	n/a	n/a	n/a

^a Maximum count rate is limited by space-charge effects to around 10^{11} photons/s per cm^2 .

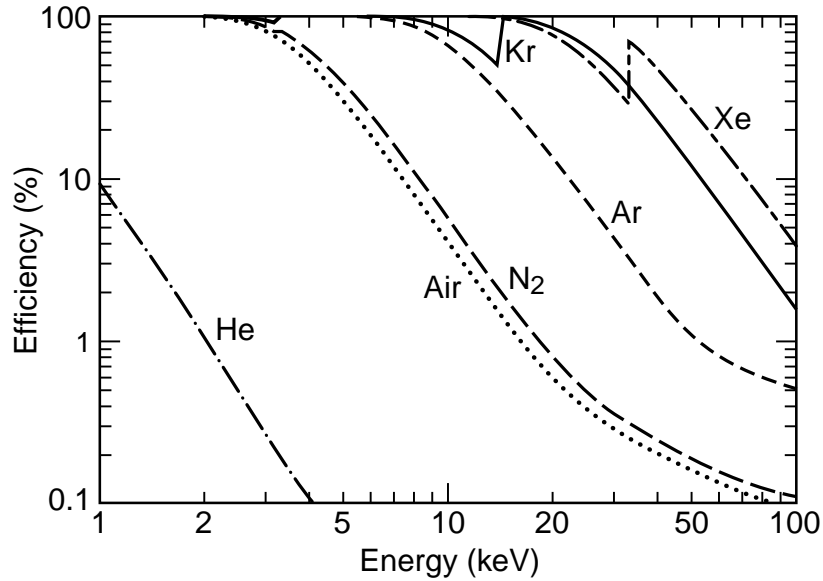


Fig. 4-10. Efficiency of a 10-cm-long gas ionization chamber as a function of energy, for different gases at normal pressure.

A. GAS IONIZATION DETECTORS

Gas ionization detectors are commonly used as integrating detectors to measure beam flux rather than individual photons. A typical detector consists of a rectangular gas cell with thin entrance and exit windows. Inside the detector, an electric field of about 100 V/cm is applied across two parallel plates. Some of the x-rays in the beam interact with the chamber gas to produce fast photoelectrons, Auger electrons, and/or fluorescence photons. The energetic electrons produce additional electron-ion pairs by inelastic collisions, and the photons either escape or are photoelectrically absorbed. The electrons and ions are collected at the plates, and the current is measured with a low-noise current amplifier. The efficiency of the detector can be calculated from the active length of the chamber, the properties of the chamber gas, and the x-ray absorption cross section at the appropriate photon energy. Figure 4-10 shows, for different gases at normal pressure, the efficiency of a 10-cm-long ion chamber as a function of energy. Once the efficiency is known, the photon flux can be estimated from chamber current and the average energy required to produce an electron-ion pair (Table 4-3).

Table 4-3. Average energy required to produce an electron-ion pair in several gases

Element	Energy (eV)
Helium	41
Nitrogen	36
Air	34.4
Neon	36.3
Argon	26
Krypton	24
Xenon	22

B. GAS PROPORTIONAL COUNTERS

Gas proportional detectors consist of a small-diameter anode wire in an enclosed gas volume. They are usually used to count single photon events. When a photon interacts in the gas, some gas atoms are ionized, and the electrons are attracted to the positive anode wire. Near the anode wire, the electrons are accelerated by the high electric field, producing a cascade of electrons that result in a large electrical pulse. The output is coupled to a low-noise preamplifier to give usable pulses. The pulse height resolution of the detector (about 20% at 6 keV) can be used for some energy discrimination, and the output counting rate can be as high as 10^6 counts per second.

C. MULTIWIRE AND MICROSTRIP PROPORTIONAL CHAMBERS

Multiwire and microstrip proportional chambers are widely used as position-sensitive detectors of both photons and charged particles. Multiwire chambers use a grid of fine wires spaced about 2 mm apart as the anode plane in a gas proportional chamber. Microstrip detectors use a patterned anode plane. The spatial resolution can be as good as $30\text{ }\mu\text{m}$. Recently gas electron multiplying (GEM) detectors have been developed that have improved spatial resolution and lower operating voltages [1].

D. SCINTILLATION DETECTORS

Scintillation detectors work by converting x-rays to optical photons in special materials and then detecting the light with a photomultiplier tube or a photodiode. The scintillator materials can be either organic scintillators, single crystals of thallium-activated sodium iodide [commonly referred to as NaI(Tl)], single crystals of sodium-activated cesium iodide [CsI(Na)], or single crystals of bismuth germanate (BGO). Since the light output is low (about 200–300 eV is required for each optical photon), the energy resolution is also low. Organic scintillators have very poor energy resolution, whereas the NaI(Tl), CsI(Na), and BGO crystals have energy resolutions of about 40% at 10 keV. These detectors can have a time resolution of better than 1 ns and a count rate capability up to 2×10^6 photons per second. For a scintillator thickness of more than 5 mm, for both NaI and CsI, the detection efficiency between 20 and 100 keV is essentially unity.

Gas scintillation detectors combine the operation of gas ionization chambers and photon detectors to give improved performance. Electrons generated from photon or charged-particle interactions in a gas (usually pure xenon or argon with 1% xenon) are accelerated in a high-field ($\sim 3\text{ kV/cm}$) region, where they produce UV scintillation light. This light is usually wave-shifted and then detected by a photomultiplier. These detectors have an energy resolution about two to three times better than conventional proportional chambers.

E. ENERGY-RESOLVING SEMICONDUCTOR DETECTORS

Silicon and germanium detectors can make excellent energy-resolving detectors of single photons (about 150 eV at 5.9 keV). They are basically large, reverse-biased $n^+ - i - p^+$ diodes. When a photon interacts in the intrinsic region, tracks of electron-hole pairs are produced (analogous to electron–positive ion pairs in a counting gas). In the presence of the electric field, these pairs separate and rapidly drift to the detector contacts. The average energy required to generate an electron-hole pair is 3.6 eV for silicon and 2.98 eV for germanium. To keep the leakage current low, the detector must have very few electrically active impurities. For example, germanium detectors are made from zone-refined crystals that have fewer than 10^{10} electrically

active impurities/cm³. They are usually cooled to reduce the thermal leakage current. The count rate capability with an energy resolution of <200 eV is limited to about 2×10^5 per second. To handle the high counting rates available at synchrotrons, multielement arrays of 4–30 elements have been developed for fluorescent EXAFS experiments.

F. CURRENT-MODE SEMICONDUCTOR DETECTORS

Semiconductor diodes are also used in current mode to measure x-ray flux. They have very linear responses and are available with thin entrance windows. Surface-barrier detectors are good beam monitors when used with low-noise current amplifiers. In addition, silicon avalanche detectors are now available in which the silicon is biased so that there is an internal avalanche of electron-hole pairs for each interacting photon. These devices can be used at lower beam intensities in a pulse-counting mode and in current mode at higher photon fluxes. They have excellent time resolution but limited energy resolution.

G. CCD DETECTORS

CCD detectors are now used in a variety of ways for x-ray imaging. They are available with up to 4096×4096 pixels, with pixel sizes of $12 \mu\text{m} \times 12 \mu\text{m}$ and readout times of less than 1 s. In most scientific applications, CCD detectors are cooled to below -30°C to reduce background noise. In most systems, a thin phosphor screen converts the incident x-rays into optical photons, which the CCD detects. A commonly used phosphor is $\text{Gd}_2\text{O}_2\text{S}(\text{Tb})$, which has a high efficiency and a light decay time of a few hundred microseconds. When used as a detector for macromolecular crystallography, a large phosphor screen (up to 300 mm^2) is usually coupled to the CCD with a tapered optical fiber [2]. On the other hand, for high-spatial-resolution x-ray imaging, a 5- to $20\text{-}\mu\text{m}$ -thick sapphire scintillation screen is optically coupled with a high-quality microscope lens to give a spatial resolution of around $1 \mu\text{m}$ [3]. For imaging with x-rays below 1 keV, direct exposure of back-thinned CCD detectors is used.

H. OTHER X-RAY DETECTORS

X-ray detectors operating at superconducting temperatures (0.1–4 K) have recently been developed; these devices achieve excellent energy resolution (12 eV at 700 eV). They are currently very small and very inefficient for x-rays above 1 keV, and they have maximum count rates of only about $5 \times 10^3 \text{ s}^{-1}$. With further development, however, they may make very useful x-ray spectrometers [4].

Microchannel plate detectors are compact, high-gain electron multipliers, which are often used as efficient electron or low-energy photon detectors. A typical MCP consists of about 10^7 closely packed lead-glass channels of equal diameter. Typically, the diameter of each channel, which acts as an independent, continuous dynode photomultiplier, is $\sim 10 \mu\text{m}$.

Image plate detectors are available that have many of the characteristics of film but with the advantage of excellent dynamic range, efficiency, and large area [5]. They are made with a plate containing a photosensitive material that on exposure to x-rays creates color centers that can be read out in a scanning mode with a laser as a digital image.

For high-speed x-ray imaging experiments, x-ray streak cameras have been developed that have time resolutions of around 350 fs.

Finally, photographic film is also available for quick x-ray–imaging experiments; however, because of the need for processing after exposure, it is no longer commonly used for scientific measurements. Special films are available that give improved efficiency, contrast, or resolution.

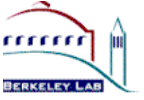
I. CALIBRATION OF X-RAY BEAM MONITORS

Measurement of the relative intensity of x-ray beams is usually done with a gas ionization chamber or a thin silicon diode in the beam path. Another technique is to place a thin foil (usually plastic) in the beam and to measure the scattered photons with a scintillation detector. The approximate efficiency of a gas ionization detector can be estimated from its active length and the properties of the chamber gas at the energy of the x-ray beam. However, calibration of these detectors to measure absolute x-ray intensity is more difficult. One calibration technique is to use a well-characterized single-photon detector as a standard and to establish the x-ray flux–to–detector current calibration of the beam monitor at a reduced beam flux where the single-photon counter response is linear.

At photon energies below 1000 eV, silicon photodiodes are available that can be used as absolute beam monitors [6]. At higher energies, avalanche photodiodes are now available with very wide dynamic ranges. Since they can be used in both single-photon and current-measuring modes, they can be easily calibrated once their efficiency as a function of energy is known (either from the device specifications or by direct measurement) [7].

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5. Miscellaneous

Contents

- [Physical constants](#)
- [Physical properties of the elements](#)
- [Electromagnetic Relations](#)
- [Radioactivity and Radiation Protection](#)
- [Useful formula](#)

[CXRO Home](#) | [ALS Home](#) | [LBL Home](#)

[Privacy and Security Notice](#)

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X-Ray Data Booklet

Section 5.1 PHYSICAL CONSTANTS

Table 5-1 was drawn from the recommendations of CODATA (the Committee on Data for Science and Technology). The full 1998 CODATA set of constants may be found at <http://physics.nist.gov/cuu/Constants/index.html>. A [PDF version of this table](#) is also available.

Table 5-1. Physical constants.

Quantity	Symbol, equation	Value	Uncert. (ppb)
speed of light	c (see note *)	$2.997\,924\,58 \times 10^8 \text{ m s}^{-1}$ ($10^{10} \text{ cm s}^{-1}$)	exact
Planck constant	h	$6.626\,068\,76(52) \times 10^{-34} \text{ J s}$ (10^{-27} erg s)	78
Planck constant, reduced	$\hbar = h/2\pi$	$1.054\,571\,596(82) \times 10^{-34} \text{ J s} = 6.582\,118\,89(26) \times 10^{-22} \text{ MeV s}$	78, 39
electron charge magnitude	e	$4.803\,204\,20(19) \times 10^{-10} \text{ esu} = 1.602\,176\,462(63) \times 10^{-19} \text{ C}$	39, 59
conversion constant	$\hbar c$	$197.326\,960\,1(78) \text{ MeV fm} (= \text{eV nm})$	39
electron mass	m_e	$0.510\,998\,902(21) \text{ MeV}/c^2 = 9.109\,381\,88(72) \times 10^{-31} \text{ kg}$	40, 79
proton mass	m_p	$938.271\,998(38) \text{ MeV}/c^2 = 1.672\,621\,58(13) \times 10^{-27} \text{ kg}$ $= 1.007\,276\,466\,88(13) \text{ u} = 1836.152\,667\,5(39) m_e$	40, 79 0.13, 2.1
deuteron mass	m_d	$1875.612\,762\,(75) \text{ MeV}/c^2$	40
unified atomic mass unit (u)	$(\text{mass } ^{12}\text{C atom})/12 = (1 \text{ g})/(N_A \text{ mol})$	$931.494\,013(37) \text{ MeV}/c^2 = 1.660\,538\,73(13) \times 10^{-27} \text{ kg}$	40, 79

Table 5-1

permittivity of free space	$\epsilon_0 = 1/(\mu_0 c^2)$	$8.854\,187\,817\dots\times 10^{-12}\text{ F m}^{-1}$	exact
permeability of free space	μ_0	$4\pi\times 10^{-7}\text{ N A}^{-2} = 12.566\,370\,614\dots\times 10^{-7}\text{ N A}^{-2}$	exact
fine-structure constant	$\alpha = e^2/4\pi\epsilon_0\hbar c$	$1/137.035\,999\,76(50)$	3.7
classical electron radius	$r_e = e^2/4\pi\epsilon_0 m_e c^2$	$2.817\,940\,285(31)\times 10^{-15}\text{ m}$	11
Bohr radius ($m_{\text{nucleus}} = \infty$)	$a_0 = 4\pi\epsilon_0\hbar^2/m_e e^2 = r_e\alpha^{-2}$	$0.529\,177\,208\,3(19)\times 10^{-10}\text{ m } (10^{-8}\text{ cm})$	3.7
Rydberg energy	$hcR_\infty = m_e e^4/2(4\pi\epsilon_0)^2\hbar^2 = m_e c^2\alpha^2/2$	$13.605\,691\,72(53)\text{ eV}$	39
Thomson cross section	$\sigma_T = 8\pi r_e^2/3$	$0.665\,245\,854(15)\text{ barn } (10^{-28}\text{ m}^2)$	22
Bohr magneton	$\mu_B = e\hbar/2m_e$	$5.788\,381\,749(43)\times 10^{-11}\text{ MeV T}^{-1}$	7.3
nuclear magneton	$\mu_N = e\hbar/2m_p$	$3.152\,451\,238(24)\times 10^{-14}\text{ MeV T}^{-1}$	7.6
electron cyclotron freq./field	$\omega_{\text{cycl}}^e/B = e/m_e$	$1.758\,820\,174(71)\times 10^{11}\text{ rad s}^{-1}\text{ T}^{-1}$	40
proton cyclotron freq./field	$\omega_{\text{cycl}}^p/B = e/m_p$	$9.578\,834\,08(38)\times 10^7\text{ rad s}^{-1}\text{ T}^{-1}$	40

Table 5-1. Physical constants(continued).

Quantity	Symbol, equation	Value	Uncert. (ppb)
Avogadro constant	N_A	$6.022\,141\,99(47)\times 10^{23}\text{ mol}^{-1}$	79
Boltzman constant	k	$1.380\,650\,3(24)\times 10^{-23}\text{ J K}^{-1} = 8.617\,342(15)\times 10^{-5}\text{ eV K}^{-1}$	1700

molar volume, ideal gas at STP	$N_A k (273.15 \text{ K}) / (101325 \text{ Pa})$	$22.413\,996(39) \times 10^{-3} \text{ m}^3 \text{ mol}^{-1}$	1700
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$$\pi = 3.141\,592\,653\,589\,793 \quad e = 2.718\,281\,828\,459\,045\,235 \quad \gamma = 0.577\,215\,664\,901\,532\,861\,238$$

The meter is the length of the path traveled by light in vacuum during a time interval of $1/299\,792\,458$ of a second.

1 in. 2.54 cm	=	1 newton	=	10^5 dyne	$1 \text{ eV}/c^2$ $662 \times 10^{-33} \text{ g}$	=	1.7821 coulomb $58 \times 10^9 \text{ esu}$	=	2.997 924	
1 Å 10^{-8} cm	=	1 joule	=	10^7 erg	$hc/(1 \text{ eV})$ $842 \mu\text{m}$	=	1.239	1 tesla	=	10^4 gauss
1 fm 10^{-13} cm	=	1 cal joule	=	4.184	$1 \text{ eV}/h$ $2.417\,989 \times 10^{14} \text{ Hz}$	=		1 atm $25 \times 10^6 \text{ dyne/cm}^2$	=	1.013
1 barn 10^{-24} cm^2	=	1 eV $176\,5 \times 10^{-12} \text{ erg}$	=	1.602	$1 \text{ eV}/k$ 604.5 K	=	11	0°C	=	273.15 K

X-Ray Data Booklet Table 5-1. Physical constants.

Quantity	Symbol, equation	Value	Uncert. (ppb)
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classical electron radius	$r_e = e^2/4\pi\epsilon_0 m_e c^2$	$2.817\,940\,285(31) \times 10^{-15} \text{ m}$	11
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Thomson cross section	$\sigma_T = 8\pi r_e^2/3$	$0.665\,245\,854(15) \text{ barn}$ (10^{-28} m^2)	22
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proton cyclotron freq./field	$\frac{e}{m_p} \text{ cycl}/B = e/m_p$	$9.578\,834\,08(38) \times 10^7 \text{ rad s}^{-1} \text{ T}^{-1}$	40

Table 5-1. Physical constants(continued).

Quantity	Symbol, equation	Value	Uncert. (ppb)
Avogadro constant	N_A	$6.022\,141\,99(47)\,10^{23}\,\text{mol}^{-1}$	79
Boltzman constant	k	$1.380\,650\,3(24)\,10^{-23}\,\text{J K}^{-1} = 8.617\,342(15)\,10^{-5}\,\text{eV K}^{-1}$	1700
molar volume, ideal gas at STP	$N_A k (273.15\,\text{K})/(101\,325\,\text{Pa})$	$22.413\,996(39)\,10^{-3}\,\text{m}^3\,\text{mol}^{-1}$	1700
$p = 3.141\,592\,653\,589\,793\,238$ $e = 2.718\,281\,828\,459\,045\,235$ $g = 0.577\,215\,664\,901\,532\,861$			
The meter is the length of the path traveled by light in vacuum during a time interval of 1/299 792 458 of a second.			
1 in. = 2.54 cm	1 newton = 10^5 dyne	$1\,\text{eV}/c^2 = 1.782\,662\,10^{-33}\,\text{g}$	1 coulomb = $2.997\,924\,58\,10^9\,\text{esu}$
$1\,\text{\AA} = 10^{-8}\,\text{cm}$	1 joule = 10^7 erg	$hc/(1\,\text{eV}) = 1.239\,842\,\text{m}$	1 tesla = 10^4 gauss
$1\,\text{fm} = 10^{-13}\,\text{cm}$	1 cal = 4.184 joule	$1\,\text{eV}/h = 2.417\,989\,10^{14}\,\text{Hz}$	1 atm = $1.013\,25\,10^6\,\text{dyne/cm}^2$
1 barn = $10^{-24}\,\text{cm}^2$	1 eV = $1.602\,176\,5\,10^{-12}\,\text{erg}$	$1\,\text{eV}/k = 11\,604.5\,\text{K}$	$0^\circ\text{C} = 273.15\,\text{K}$

X-Ray Data Booklet

Section 5.2 Physical Properties of the Elements

Table 5-2 lists several important properties of the elements. Data were taken mostly from D. R. Lide, Ed., *CRC Handbook of Chemistry and Physics*, 80th ed. (CRC Press, Boca Raton, Florida, 1999). Atomic weights apply to elements as they exist naturally on earth; values in parentheses are the mass numbers for the longest-lived isotopes. Some uncertainty exists in the last digit of each atomic weight. Specific heats are given for the elements at 25°C and a pressure of 100 kPa. Densities for solids and liquids are given as specific gravities at 20°C unless otherwise indicated by a superscript temperature (in °C); densities for the gaseous elements are given in g/cm³ for the liquids at their boiling points. The ionization energies were taken from <http://physics.nist.gov/PhysRefData/IonEnergy/ionEnergy.html>. A [PDF version of this table](#) is also available.

Table 5-2. Properties of the elements.

Z Element	Atomic weight	Density	Melting	Boiling	Ground-state configuration	Ground level	Ionization energy (eV)	Specific heat (J/g·K)
			point (°C)	point (°C)				
1 Hydrogen	1.00794	0.0708	-259.34	-252.87	1s	2S _{1/2}	13.598	14.304
2 Helium	4.002602	0.122	—	-268.93	1s ²	1S ₀	24.587	5.193
3 Lithium	6.941	0.534	180.50	1342	1s ² 2s	2S _{1/2}	5.392	3.582
4 Beryllium	9.012182	1.848	1287	2471	1s ² 2s ²	1S ₀	9.323	1.825
5 Boron	10.811	2.34	2075	4000	1s ² 2s ² 2p	2P° _{1/2}	8.298	1.026
6 Carbon	12.0107	1.9–2.3 (graphite)	4492 ^{10.3} MPa	3825 ^b	1s ² 2s ² 2p ²	3P ₀	11.260	0.709

Table 5-2

7 Nitrogen	14.00674	0.808	−210.00	−195.79	$1s^2 2s^2 2p^3$	$4S^{\circ}_{3/2}$	14.534	1.040
8 Oxygen	15.9994	1.14	−218.79	−182.95	$1s^2 2s^2 2p^4$	$3P_2$	13.618	0.918
9 Fluorine	18.9984032	1.50	−219.62	−188.12	$1s^2 2s^2 2p^5$	$2P^{\circ}_{3/2}$	17.423	0.824
10 Neon	20.1797	1.207	−248.59	−246.08	$1s^2 2s^2 2p^6$	$1S_0$	21.565	1.030
11 Sodium	22.989770	0.971	97.80	883	[Ne] $3s$	$2S_{1/2}$	5.139	1.228
12 Magnesium	24.3050	1.738	650	1090	[Ne] $3s^2$	$1S_0$	7.646	1.023
13 Aluminum	26.981538	2.6989	660.32	2519	[Ne] $3s^2 3p$	$2P^{\circ}_{1/2}$	5.986	0.897
14 Silicon	28.0855	2.33 ²⁵	1414	3265	[Ne] $3s^2 3p^2$	$3P_0$	8.152	0.705
15 Phosphorus	30.973761	1.82	44.15	280.5	[Ne] $3s^2 3p^3$	$4S^{\circ}_{3/2}$	10.487	0.769
16 Sulfur	32.066	2.07	119.6	444.60	[Ne] $3s^2 3p^4$	$3P_2$	10.360	0.710
17 Chlorine	35.4527	1.56 ^{−33.6}	−101.5	−34.04	[Ne] $3s^2 3p^5$	$2P^{\circ}_{3/2}$	12.968	0.479
18 Argon	39.948	1.40	−189.35	−185.85	[Ne] $3s^2 3p^6$	$1S_0$	15.760	0.520
19 Potassium	39.0983	0.862	63.5	759	[Ar] $4s$	$2S_{1/2}$	4.341	0.757
20 Calcium	40.078	1.55	842	1484	[Ar] $4s^2$	$1S_0$	6.113	0.647
21 Scandium	44.955910	2.989 ²⁵	1541	2836	[Ar] $3d 4s^2$	$2D_{3/2}$	6.562	0.568
22 Titanium	47.867	4.54	1668	3287	[Ar] $3d^2 4s^2$	$3F_2$	6.828	0.523
23 Vanadium	50.9415	6.11 ^{18.7}	1910	3407	[Ar] $3d^3 4s^2$	$4F_{3/2}$	6.746	0.489

Table 5-2

24 Chromium	51.9961	7.18–7.20	1907	2671	[Ar] $3d^5 4s$	7S_3	6.766	0.449
25 Manganese	54.938049	7.21–7.44	1246	2061	[Ar] $3d^5 4s^2$	$^6S_{5/2}$	7.434	0.479
26 Iron	55.845	7.874	1538	2861	[Ar] $3d^6 4s^2$	5D_4	7.902	0.449

		Melting		Boiling		Ionization			Specific
	Atomic		point	point	Ground-	Ground	energy	heat	
Z	Element	weight	Density	(°C)	(°C)	configuration	level	(eV)	(J/g·K)
27	Cobalt	58.933200	8.9	1495	2927	[Ar] $3d^7 4s^2$	$^4F_{9/2}$	7.881	0.421
28	Nickel	58.6934	8.902 ²⁵	1455	2913	[Ar] $3d^8 4s^2$	3F_4	7.640	0.444
29	Copper	63.546	8.96	1084.62	2562	[Ar] $3d^{10} 4s$	$^2S_{1/2}$	7.726	0.385
30	Zinc	65.39	7.133 ²⁵	419.53	907	[Ar] $3d^{10} 4s^2$	1S_0	9.394	0.388
31	Gallium	69.723	5.904 ^{29.6}	29.76	2204	[Ar] $3d^{10} 4s^2 4p$	$^2P^{\circ}_{1/2}$	5.999	0.371
32	Germanium	72.61	5.323 ²⁵	938.25	2833	[Ar] $3d^{10} 4s^2 4p^2$	3P_0	7.899	0.320
33	Arsenic	74.92160	5.73	817 ^{3.7 MPa}	603 ^b	[Ar] $3d^{10} 4s^2 4p^3$	$^4S^{\circ}_{3/2}$	9.789	0.329
34	Selenium	78.96	4.79	220.5	685	[Ar] $3d^{10} 4s^2 4p^4$	3P_2	9.752	0.321
35	Bromine	79.904	3.12	–7.2	58.8	[Ar] $3d^{10} 4s^2 4p^5$	$^2P^{\circ}_{3/2}$	11.814	0.226

Table 5-2

36 Krypton	83.80	2.16	157.38 ^{73.2} kPa	-153.22	[Ar] 3d ¹⁰ 4s ² 4p ⁶	1S ₀	14.000	0.248
37 Rubidium	85.4678	1.532	39.30	688	[Kr] 5s	2S _{1/2}	4.177	0.363
38 Strontium	87.62	2.54	777	1382	[Kr] 5s ²	1S ₀	5.695	0.301
39 Yttrium	88.90585	4.469 ²⁵	1522	3345	[Kr] 4d 5s ²	2D _{3/2}	6.217	0.298
40 Zirconium	91.224	6.506	1855	4409	[Kr] 4d ² 5s ²	3F ₂	6.634	0.278
41 Niobium	92.90638	8.57	2477	4744	[Kr] 4d ⁴ 5s	6D _{1/2}	6.759	0.265
42 Molybdenum	95.94	10.22	2623	4639	[Kr] 4d ⁵ 5s	7S ₃	7.092	0.251
43 Technetium	(98)	11.50 ^a	2157	4265	[Kr] 4d ⁵ 5s ²	6S _{5/2}	7.28	—
44 Ruthenium	101.07	12.41	2334	4150	[Kr] 4d ⁷ 5s	5F ₅	7.360	0.238
45 Rhodium	102.90550	12.41	1964	3695	[Kr] 4d ⁸ 5s	4F _{9/2}	7.459	0.243
46 Palladium	106.42	12.02	1554.9	2963	[Kr] 4d ¹⁰	1S ₀	8.337	0.246
47 Silver	107.8682	10.50	961.78	2162	[Kr] 4d ¹⁰ 5s	2S _{1/2}	7.576	0.235
48 Cadmium	112.411	8.65	321.07	767	[Kr] 4d ¹⁰ 5s ²	1S ₀	8.994	0.232
49 Indium	114.818	7.31	156.60	2072	[Kr] 4d ¹⁰ 5s ² 5p	2P ^o _{1/2}	5.786	0.233
50 Tin	118.710	7.31	231.93	2602	[Kr] 4d ¹⁰ 5s ² 5p ²	3P ₀	7.344	0.228

Table 5-2

51 Antimony	121.760	6.691	630.73	1587	[Kr] 4d ¹⁰ 5s ² 5p ³	4S° _{3/2}	8.608	0.207
52 Tellurium	127.60	6.24	449.51	988	[Kr] 4d ¹⁰ 5s ² 5p ⁴	3P ₂	9.010	0.202
53 Iodine	126.90447	4.93	113.7	184.4	[Kr] 4d ¹⁰ 5s ² 5p ⁵	2P° _{3/2}	10.451	0.145
54 Xenon	131.29	3.52	−111.79 ^{81.6} kPa	−108.12	[Kr] 4d ¹⁰ 5s ² 5p ⁶	1S ₀	12.130	0.158

Melting Boiling						IonizationSpecific			
Z	Element	Atomic weight	Density	point (°C)	point (°C)	Ground-state configuration	Ground level	energy (eV)	heat (J/g·K)
55	Cesium	132.90545	1.873	28.5	671	[Xe] 6s	2S _{1/2}	3.894	0.242
56	Barium	137.327	3.5	727	1897	[Xe] 6s ²	1S ₀	5.212	0.204
57	Lanthanum	138.9055	6.145 ²⁵	918	3464	[Xe] 5d 6s ²	2D _{3/2}	5.577	0.195
58	Cerium	140.116	6.770 ²⁵	798	3443	[Xe] 4f 5d 6s ²	1G ^o ₄	5.539	0.192
59	Praseodymium	140.90765	6.773	931	3520	[Xe] 4f ³ 6s ²	4I ^o _{9/2}	5.473	0.193
60	Neodymium	144.24	7.008 ²⁵	1021	3074	[Xe] 4f ⁴ 6s ²	5I ₄	5.525	0.190
61	Promethium	(145)	7.264 ²⁵	1042	3000	[Xe] 4f ⁵ 6s ²	6H ^o _{5/2}	5.582	—
62	Samarium	150.36	7.520 ²⁵	1074	1794	[Xe] 4f ⁶ 6s ²	7F ₀	5.644	0.197

Table 5-2

63 Europium	151.964	5.244 ²⁵	822	1529	[Xe] 4f ⁷ 6s ²	⁸ S° _{7/2}	5.670	0.182
64 Gadolinium	157.25	7.901 ²⁵	1313	3273	[Xe] 4f ⁷ 5d 6s ²	⁹ D° ₂	6.150	0.236
65 Terbium	158.92534	8.230	1356	3230	[Xe] 4f ⁹ 6s ²	⁶ H° _{15/2}	5.864	0.182
66 Dysprosium	162.50	8.551 ²⁵	1412	2567	[Xe] 4f ¹⁰ 6s ²	⁵ I ₈	5.939	0.170
67 Holmium	164.93032	8.795 ²⁵	1474	2700	[Xe] 4f ¹¹ 6s ²	⁴ I° _{15/2}	6.022	0.165
68 Erbium	167.26	9.066 ²⁵	1529	2868	[Xe] 4f ¹² 6s ²	³ H ₆	6.108	0.168
69 Thulium	168.93421	9.321 ²⁵	1545	1950	[Xe] 4f ¹³ 6s ²	² F° _{7/2}	6.184	0.160
70 Ytterbium	173.04	6.966	819	1196	[Xe] 4f ¹⁴ 6s ²	¹ S ₀	6.254	0.155
71 Lutetium	174.967	9.841 ²⁵	1663	3402	[Xe] 4f ¹⁴ 5d 6s ²	² D _{3/2}	5.426	0.154
72 Hafnium	178.49	13.31	2233	4603	[Xe] 4f ¹⁴ 5d ² 6s ²	³ F ₂	6.825	0.144
73 Tantalum	180.9479	16.654	3017	5458	[Xe] 4f ¹⁴ 5d ³ 6s ²	⁴ F _{3/2}	7.550	0.140
74 Tungsten	183.84	19.3	3422	5555	[Xe] 4f ¹⁴ 5d ⁴ 6s ²	⁵ D ₀	7.864	0.132
75 Rhenium	186.207	21.02	3186	5596	[Xe] 4f ¹⁴ 5d ⁵ 6s ²	⁶ S _{5/2}	7.834	0.137
76 Osmium	190.23	22.57	3033	5012	[Xe] 4f ¹⁴ 5d ⁶ 6s ²	⁵ D ₄	8.438	0.130

77 Iridium	192.217	22.42 ¹⁷	2446	4428	[Xe] 4f ¹⁴ 5d ⁷ 6s ²	4F _{9/2}	8.967	0.131
78 Platinum	195.078	21.45	1768.4	3825	[Xe] 4f ¹⁴ 5d ⁹ 6s	3D ₃	8.959	0.133
79 Gold	196.96655	~19.3	1064.18	2856	[Xe] 4f ¹⁴ 5d ¹⁰ 6s	2S _{1/2}	9.226	0.129
80 Mercury	200.59	13.546	−38.83	356.73	[Xe] 4f ¹⁴ 5d ¹⁰ 6s ²	1S ₀	10.438	0.140
81 Thallium	204.3833	11.85	304	1473	[Xe] 4f ¹⁴ 5d ¹⁰ 6s ² 6p	2P ^o _{1/2}	6.108	0.129
82 Lead	207.2	11.35	327.46	1749	[Xe] 4f ¹⁴ 5d ¹⁰ 6s ² 6p ²	3P ₀	7.417	0.129

Melting Boiling					IonizationSpecific				
Z	Element	Atomic weight	Density	point (°C)	point (°C)	Ground-state configuration	Ground level	energy (eV)	heat (J/g·K)
83	Bismuth	208.98038	9.747	271.40	1564	[Xe] 4f ¹⁴ 5d ¹⁰ 6s ² 6p ³	4S ^o _{3/2}	7.286	0.122
84	Polonium	(209)	9.32	254	962	[Xe] 4f ¹⁴ 5d ¹⁰ 6s ² 6p ⁴	3P ₂	8.417 ?	—
85	Astatine	(210)	—	302	—	[Xe] 4f ¹⁴ 5d ¹⁰ 6s ² 6p ⁵	2P ^o _{3/2}	—	—
86	Radon	(222)	—	−71	−61.7	[Xe] 4f ¹⁴ 5d ¹⁰ 6s ² 6p ⁶	1S ₀	10.748	0.094
87	Francium	(223)	—	27	—	[Rn] 7s	2S _{1/2}	4.073	—
88	Radium	(226)	—	700	—	[Rn] 7s ²	1S ₀	5.278	—

Table 5-2

89 Actinium	(227)	—	1051	3198	[Rn] 6d 7s ²	² D _{3/2}	5.17	0.120
90 Thorium	232.0381	11.72	1750	4788	[Rn] 6d ² 7s ²	³ F ₂	6.307	0.113
91 Protactinium	231.03588	15.37 ^a	1572	—	[Rn]5f ² (³ H ₄) 6d 7s ²	(4, ³ / ₂) _{11/2}	5.89	—
92 Uranium	238.0289	~18.95	1135	4131	[Rn]5f ³ (⁴ I° _{9/2}) 6d 7s ²	(⁹ / ₂ , ³ / ₂)° ₆	6.194	0.116
93 Neptunium	(237)	20.25	644	—	[Rn]5f ⁴ (⁵ I ₄) 6d 7s ²	(4, ³ / ₂) _{11/2}	6.266	—
94 Plutonium	(244)	19.84 ²⁵	640	3228	[Rn] 5f ⁶ 7s ²	⁷ F ₀	6.026	—
95 Americium	(243)	13.67	1176	2011	[Rn] 5f ⁷ 7s ²	⁸ S° _{7/2}	5.974	—
96 Curium	(247)	13.51 ^a	1345	3100	[Rn] 5f ⁷ 6d 7s ²	⁹ D° ₂	5.992	—
97 Berkelium	(247)	14 (est.)	1050	—	[Rn] 5f ⁹ 7s ²	⁶ H° _{15/2}	6.198	—
98 Californium	(251)	—	900	—	[Rn] 5f ¹⁰ 7s ²	⁵ I ₈	6.282	—
99 Einsteinium	(252)	—	860	—	[Rn] 5f ¹¹ 7s ²	⁴ I° _{15/2}	6.42	—
100 Fermium	(257)	—	1527	—	[Rn] 5f ¹² 7s ²	³ H ₆	6.50	—
101 Mendelevium	(258)	—	827	—	[Rn] 5f ¹³ 7s ²	² F° _{7/2}	6.58	—
102 Nobelium	(259)	—	827	—	[Rn] 5f ¹⁴ 7s ²	¹ S ₀	6.65	—

Table 5-2

103 Lawrencium	(262)	—	1627	—	[Rn] $5f^{14}$ $7s^2$ $7p$?	$2p_{1/2}^{\circ}$?	4.9 ?	—
104 Rutherfordium	(261)	—	—	—	[Rn] $5f^{14}$ $6d^2$ $7s^2$?	$3F_2$?	6.0 ?	—
105 Dubnium	(262)	—	—	—	—	—	—	—
106 Seaborgium	(266)	—	—	—	—	—	—	—
107 Bohrium	(264)	—	—	—	—	—	—	—
108 Hassium	(269)	—	—	—	—	—	—	—
109 Meitnerium	(268)	—	—	—	—	—	—	—

^aCalculated^bSublimes

X-Ray Data Booklet Table 5-2. Properties of the elements.

Z	Element	Atomic weight	Density	Melting point (°C)	Boiling point (°C)	Ground-state configuration	Ground level	Ionization energy (eV)	Specific heat (J/g·K)
1	Hydrogen	1.00794	0.0708	−259.34	−252.87	1s	² S _{1/2}	13.598	14.304
2	Helium	4.002602	0.122	—	−268.93	1s ²	¹ S ₀	24.587	5.193
3	Lithium	6.941	0.534	180.50	1342	1s ² 2s	² S _{1/2}	5.392	3.582
4	Beryllium	9.012182	1.848	1287	2471	1s ² 2s ²	¹ S ₀	9.323	1.825
5	Boron	10.811	2.34	2075	4000	1s ² 2s ² 2p	² P [°] _{1/2}	8.298	1.026
6	Carbon	12.0107	1.9–2.3 (graph)	4492 ^{10.3 MPa}	3825 ^b	1s ² 2s ² 2p ²	³ P ₀	11.260	0.709
7	Nitrogen	14.00674	0.808	−210.00	−195.79	1s ² 2s ² 2p ³	⁴ S [°] _{3/2}	14.534	1.040
8	Oxygen	15.9994	1.14	−218.79	−182.95	1s ² 2s ² 2p ⁴	³ P ₂	13.618	0.918
9	Fluorine	18.9984032	1.50	−219.62	−188.12	1s ² 2s ² 2p ⁵	² P [°] _{3/2}	17.423	0.824
10	Neon	20.1797	1.207	−248.59	−246.08	1s ² 2s ² 2p ⁶	¹ S ₀	21.565	1.030
11	Sodium	22.989770	0.971	97.80	883	[Ne] 3s	² S _{1/2}	5.139	1.228
12	Magnesium	24.3050	1.738	650	1090	[Ne] 3s ²	¹ S ₀	7.646	1.023
13	Aluminum	26.981538	2.6989	660.32	2519	[Ne] 3s ² 3p	² P [°] _{1/2}	5.986	0.897
14	Silicon	28.0855	2.33 ²⁵	1414	3265	[Ne] 3s ² 3p ²	³ P ₀	8.152	0.705
15	Phosphorus	30.973761	1.82	44.15	280.5	[Ne] 3s ² 3p ³	⁴ S [°] _{3/2}	10.487	0.769
16	Sulfur	32.066	2.07	119.6	444.60	[Ne] 3s ² 3p ⁴	³ P ₂	10.360	0.710
17	Chlorine	35.4527	1.56– ^{33.6}	−101.5	−34.04	[Ne] 3s ² 3p ⁵	² P [°] _{3/2}	12.968	0.479
18	Argon	39.948	1.40	−189.35	−185.85	[Ne] 3s ² 3p ⁶	¹ S ₀	15.760	0.520
19	Potassium	39.0983	0.862	63.5	759	[Ar] 4s	² S _{1/2}	4.341	0.757
20	Calcium	40.078	1.55	842	1484	[Ar] 4s ²	¹ S ₀	6.113	0.647
21	Scandium	44.955910	2.989 ²⁵	1541	2836	[Ar] 3d 4s ²	² D _{3/2}	6.562	0.568
22	Titanium	47.867	4.54	1668	3287	[Ar] 3d ² 4s ²	³ F ₂	6.828	0.523
23	Vanadium	50.9415	6.11 ^{18.7}	1910	3407	[Ar] 3d ³ 4s ²	⁴ F _{3/2}	6.746	0.489
24	Chromium	51.9961	7.18–7.20	1907	2671	[Ar] 3d ⁵ 4s	⁷ S ₃	6.766	0.449
25	Manganese	54.938049	7.21–7.44	1246	2061	[Ar] 3d ⁵ 4s ²	⁶ S _{5/2}	7.434	0.479
26	Iron	55.845	7.874	1538	2861	[Ar] 3d ⁶ 4s ²	⁵ D ₄	7.902	0.449

Table 5-2. Properties of the elements (continued).

Z	Element	Atomic weight	Density	Melting point (°C)	Boiling point (°C)	Ground-state configuration	Ground level	Ionization energy (eV)	Specific heat (J/g·K)
27	Cobalt	58.933200	8.9	1495	2927	[Ar] 3d ⁷ 4s ²	⁴ F _{9/2}	7.881	0.421
28	Nickel	58.6934	8.902 ²⁵	1455	2913	[Ar] 3d ⁸ 4s ²	³ F ₄	7.640	0.444
29	Copper	63.546	8.96	1084.62	2562	[Ar] 3d ¹⁰ 4s	² S _{1/2}	7.726	0.385
30	Zinc	65.39	7.133 ²⁵	419.53	907	[Ar] 3d ¹⁰ 4s ²	¹ S ₀	9.394	0.388
31	Gallium	69.723	5.904 ^{29.6}	29.76	2204	[Ar] 3d ¹⁰ 4s ² 4p	² P ^o _{1/2}	5.999	0.371
32	Germanium	72.61	5.323 ²⁵	938.25	2833	[Ar] 3d ¹⁰ 4s ² 4p ²	³ P ₀	7.899	0.320
33	Arsenic	74.92160	5.73	817 ^{3.7} MPa	603 ^b	[Ar] 3d ¹⁰ 4s ² 4p ³	⁴ S ^o _{3/2}	9.789	0.329
34	Selenium	78.96	4.79	220.5	685	[Ar] 3d ¹⁰ 4s ² 4p ⁴	³ P ₂	9.752	0.321
35	Bromine	79.904	3.12	−7.2	58.8	[Ar] 3d ¹⁰ 4s ² 4p ⁵	² P ^o _{3/2}	11.814	0.226
36	Krypton	83.80	2.16	157.38 ^{73.2} kPa	−153.22	[Ar] 3d ¹⁰ 4s ² 4p ⁶	¹ S ₀	14.000	0.248
37	Rubidium	85.4678	1.532	39.30	688	[Kr] 5s	² S _{1/2}	4.177	0.363
38	Strontium	87.62	2.54	777	1382	[Kr] 5s ²	¹ S ₀	5.695	0.301
39	Yttrium	88.90585	4.469 ²⁵	1522	3345	[Kr] 4d 5s ²	² D _{3/2}	6.217	0.298
40	Zirconium	91.224	6.506	1855	4409	[Kr] 4d ² 5s ²	³ F ₂	6.634	0.278
41	Niobium	92.90638	8.57	2477	4744	[Kr] 4d ⁴ 5s	⁶ D _{1/2}	6.759	0.265
42	Molybdenum	95.94	10.22	2623	4639	[Kr] 4d ⁵ 5s	⁷ S ₃	7.092	0.251
43	Technetium	(98)	11.50 ^a	2157	4265	[Kr] 4d ⁵ 5s ²	⁶ S _{5/2}	7.28	—
44	Ruthenium	101.07	12.41	2334	4150	[Kr] 4d ⁷ 5s	⁵ F ₅	7.360	0.238
45	Rhodium	102.90550	12.41	1964	3695	[Kr] 4d ⁸ 5s	⁴ F _{9/2}	7.459	0.243
46	Palladium	106.42	12.02	1554.9	2963	[Kr] 4d ¹⁰	¹ S ₀	8.337	0.246
47	Silver	107.8682	10.50	961.78	2162	[Kr] 4d ¹⁰ 5s	² S _{1/2}	7.576	0.235
48	Cadmium	112.411	8.65	321.07	767	[Kr] 4d ¹⁰ 5s ²	¹ S ₀	8.994	0.232
49	Indium	114.818	7.31	156.60	2072	[Kr] 4d ¹⁰ 5s ² 5p	² P ^o _{1/2}	5.786	0.233
50	Tin	118.710	7.31	231.93	2602	[Kr] 4d ¹⁰ 5s ² 5p ²	³ P ₀	7.344	0.228
51	Antimony	121.760	6.691	630.73	1587	[Kr] 4d ¹⁰ 5s ² 5p ³	⁴ S ^o _{3/2}	8.608	0.207
52	Tellurium	127.60	6.24	449.51	988	[Kr] 4d ¹⁰ 5s ² 5p ⁴	³ P ₂	9.010	0.202
53	Iodine	126.90447	4.93	113.7	184.4	[Kr] 4d ¹⁰ 5s ² 5p ⁵	² P ^o _{3/2}	10.451	0.145
54	Xenon	131.29	3.52	−111.79 ^{81.6} kPa	−108.12	[Kr] 4d ¹⁰ 5s ² 5p ⁶	¹ S ₀	12.130	0.158

Table 5-2. Properties of the elements (continued).

Z	Element	Atomic weight	Density	Melting point (°C)	Boiling point (°C)	Ground-state configuration		Ground level	Ionization energy (eV)	Specific heat (J/g·K)
55	Cesium	132.90545	1.873	28.5	671	[Xe]	6s	² S _{1/2}	3.894	0.242
56	Barium	137.327	3.5	727	1897	[Xe]	6s ²	¹ S ₀	5.212	0.204
57	Lanthanum	138.9055	6.145 ²⁵	918	3464	[Xe]	5d 6s ²	² D _{3/2}	5.577	0.195
58	Cerium	140.116	6.770 ²⁵	798	3443	[Xe]	4f 5d 6s ²	¹ G ^o ₄	5.539	0.192
59	Praseodymium	140.90765	6.773	931	3520	[Xe]	4f ³ 6s ²	⁴ I ^o _{9/2}	5.473	0.193
60	Neodymium	144.24	7.008 ²⁵	1021	3074	[Xe]	4f ⁴ 6s ²	⁵ I ₄	5.525	0.190
61	Promethium	(145)	7.264 ²⁵	1042	3000	[Xe]	4f ⁵ 6s ²	⁶ H ^o _{5/2}	5.582	—
62	Samarium	150.36	7.520 ²⁵	1074	1794	[Xe]	4f ⁶ 6s ²	⁷ F ₀	5.644	0.197
63	Europium	151.964	5.244 ²⁵	822	1529	[Xe]	4f ⁷ 6s ²	⁸ S ^o _{7/2}	5.670	0.182
64	Gadolinium	157.25	7.901 ²⁵	1313	3273	[Xe]	4f ⁷ 5d 6s ²	⁹ D ^o ₂	6.150	0.236
65	Terbium	158.92534	8.230	1356	3230	[Xe]	4f ⁹ 6s ²	⁶ H ^o _{15/2}	5.864	0.182
66	Dysprosium	162.50	8.551 ²⁵	1412	2567	[Xe]	4f ¹⁰ 6s ²	⁵ I ₈	5.939	0.170
67	Holmium	164.93032	8.795 ²⁵	1474	2700	[Xe]	4f ¹¹ 6s ²	⁴ I ^o _{15/2}	6.022	0.165
68	Erbium	167.26	9.066 ²⁵	1529	2868	[Xe]	4f ¹² 6s ²	³ H ₆	6.108	0.168
69	Thulium	168.93421	9.321 ²⁵	1545	1950	[Xe]	4f ¹³ 6s ²	² F ^o _{7/2}	6.184	0.160
70	Ytterbium	173.04	6.966	819	1196	[Xe]	4f ¹⁴ 6s ²	¹ S ₀	6.254	0.155
71	Lutetium	174.967	9.841 ²⁵	1663	3402	[Xe]	4f ¹⁴ 5d 6s ²	² D _{3/2}	5.426	0.154
72	Hafnium	178.49	13.31	2233	4603	[Xe]	4f ¹⁴ 5d ² 6s ²	³ F ₂	6.825	0.144
73	Tantalum	180.9479	16.654	3017	5458	[Xe]	4f ¹⁴ 5d ³ 6s ²	⁴ F _{3/2}	7.550	0.140
74	Tungsten	183.84	19.3	3422	5555	[Xe]	4f ¹⁴ 5d ⁴ 6s ²	⁵ D ₀	7.864	0.132
75	Rhenium	186.207	21.02	3186	5596	[Xe]	4f ¹⁴ 5d ⁵ 6s ²	⁶ S _{5/2}	7.834	0.137
76	Osmium	190.23	22.57	3033	5012	[Xe]	4f ¹⁴ 5d ⁶ 6s ²	⁵ D ₄	8.438	0.130
77	Iridium	192.217	22.42 ¹⁷	2446	4428	[Xe]	4f ¹⁴ 5d ⁷ 6s ²	⁴ F _{9/2}	8.967	0.131
78	Platinum	195.078	21.45	1768.4	3825	[Xe]	4f ¹⁴ 5d ⁹ 6s	³ D ₃	8.959	0.133
79	Gold	196.96655	~19.3	1064.18	2856	[Xe]	4f ¹⁴ 5d ¹⁰ 6s	² S _{1/2}	9.226	0.129
80	Mercury	200.59	13.546	–38.83	356.73	[Xe]	4f ¹⁴ 5d ¹⁰ 6s ²	¹ S ₀	10.438	0.140
81	Thallium	204.3833	11.85	304	1473	[Xe]	4f ¹⁴ 5d ¹⁰ 6s ² 6p	² P ^o _{1/2}	6.108	0.129
82	Lead	207.2	11.35	327.46	1749	[Xe]	4f ¹⁴ 5d ¹⁰ 6s ² 6p ²	³ P ₀	7.417	0.129

Table 5-2. Properties of the elements (continued).

Z	Element	Atomic weight	Density	Melting point (°C)	Boiling point (°C)	Ground-state configuration	Ground level	Ionization energy (eV)	Specific heat (J/g·K)
83	Bismuth	208.98038	9.747	271.40	1564	[Xe] 4f ¹⁴ 5d ¹⁰ 6s ² 6p ³	4S° _{3/2}	7.286	0.122
84	Polonium	(209)	9.32	254	962	[Xe] 4f ¹⁴ 5d ¹⁰ 6s ² 6p ⁴	3P ₂	8.417 ?	—
85	Astatine	(210)	—	302	—	[Xe] 4f ¹⁴ 5d ¹⁰ 6s ² 6p ⁵	2P° _{3/2}	—	—
86	Radon	(222)	—	−71	−61.7	[Xe] 4f ¹⁴ 5d ¹⁰ 6s ² 6p ⁶	1S ₀	10.748	0.094
87	Francium	(223)	—	27	—	[Rn] 7s	2S _{1/2}	4.073	—
88	Radium	(226)	—	700	—	[Rn] 7s ²	1S ₀	5.278	—
89	Actinium	(227)	—	1051	3198	[Rn] 6d 7s ²	2D _{3/2}	5.17	0.120
90	Thorium	232.0381	11.72	1750	4788	[Rn] 6d ² 7s ²	3F ₂	6.307	0.113
91	Protactinium	231.03588	15.37 ^a	1572	—	[Rn] 5f ² (³ H ₄) 6d 7s ²	(4, ³ / ₂) _{11/2}	5.89	—
92	Uranium	238.0289	~18.95	1135	4131	[Rn] 5f ³ (⁴ I° _{9/2}) 6d 7s ²	(⁹ / ₂ , ³ / ₂)° ₆	6.194	0.116
93	Neptunium	(237)	20.25	644	—	[Rn] 5f ⁴ (⁵ I ₄) 6d 7s ²	(4, ³ / ₂) _{11/2}	6.266	—
94	Plutonium	(244)	19.84 ²⁵	640	3228	[Rn] 5f ⁶ 7s ²	7F ₀	6.026	—
95	Americium	(243)	13.67	1176	2011	[Rn] 5f ⁷ 7s ²	8S° _{7/2}	5.974	—
96	Curium	(247)	13.51 ^a	1345	3100	[Rn] 5f ⁷ 6d 7s ²	9D° ₂	5.992	—
97	Berkelium	(247)	14 (est.)	1050	—	[Rn] 5f ⁹ 7s ²	6H° _{15/2}	6.198	—
98	Californium	(251)	—	900	—	[Rn] 5f ¹⁰ 7s ²	5I ₈	6.282	—
99	Einsteinium	(252)	—	860	—	[Rn] 5f ¹¹ 7s ²	4I° _{15/2}	6.42	—
100	Fermium	(257)	—	1527	—	[Rn] 5f ¹² 7s ²	3H ₆	6.50	—
101	Mendelevium	(258)	—	827	—	[Rn] 5f ¹³ 7s ²	2F° _{7/2}	6.58	—
102	Nobelium	(259)	—	827	—	[Rn] 5f ¹⁴ 7s ²	1S ₀	6.65	—
103	Lawrencium	(262)	—	1627	—	[Rn] 5f ¹⁴ 7s ² 7p ?	2P° _{1/2} ?	4.9 ?	—
104	Rutherfordium	(261)	—	—	—	[Rn] 5f ¹⁴ 6d ² 7s ² ?	3F ₂ ?	6.0 ?	—
105	Dubnium	(262)	—	—	—	—	—	—	—
106	Seaborgium	(266)	—	—	—	—	—	—	—
107	Bohrium	(264)	—	—	—	—	—	—	—
108	Hassium	(269)	—	—	—	—	—	—	—
109	Meitnerium	(268)	—	—	—	—	—	—	—

^aCalculated

^bSublimes

X-Ray Data Booklet

Section 5.3 ELECTROMAGNETIC RELATIONS

	Gaussian CGS	SI
Units and conversions:		
Charge	$2.997\,92 \times 10^9 \text{ esu}$	$= 1 \text{ C} = 1 \text{ A s}$
Potential	$(1/299.792) \text{ statvolt}$ $= (1/299.792) \text{ erg/esu}$	$= 1 \text{ V} = 1 \text{ J C}^{-1}$
Magnetic field	$10^4 \text{ gauss} = 10^4 \text{ dyne/esu}$	$= 1 \text{ T} = 1 \text{ N A}^{-1} \text{ m}^{-1}$
Electron charge	$e = 4.803\,204 \times 10^{-10} \text{ esu}$	$= 1.602\,176 \times 10^{-19} \text{ C}$
Lorentz force	$\mathbf{F} = q \left(\mathbf{E} + \frac{\mathbf{v}}{c} \times \mathbf{B} \right)$	$\mathbf{F} = q (\mathbf{E} + \mathbf{v} \times \mathbf{B})$
Maxwell equations	$\nabla \cdot \mathbf{D} = 4\pi\rho$	$\nabla \cdot \mathbf{D} = \rho$
	$\nabla \times \mathbf{E} + \frac{1}{c} \frac{\partial \mathbf{B}}{\partial t} = 0$	$\nabla \times \mathbf{E} + \frac{\partial \mathbf{B}}{\partial t} = 0$
	$\nabla \cdot \mathbf{B} = 0$	$\nabla \cdot \mathbf{B} = 0$
	$\nabla \times \mathbf{H} - \frac{1}{c} \frac{\partial \mathbf{D}}{\partial t} = \frac{4\pi}{c} \mathbf{J}$	$\nabla \times \mathbf{H} - \frac{\partial \mathbf{D}}{\partial t} = \mathbf{J}$
Linear media	$\mathbf{D} = \epsilon \mathbf{E}, \mathbf{B} = \mu \mathbf{H}$	$\mathbf{D} = \epsilon \mathbf{E}, \mathbf{B} = \mu \mathbf{H}$
Permittivity of free space	$\epsilon_{\text{vac}} = 1$	$\epsilon_{\text{vac}} = \epsilon_0$

Permeability of free space	$\mu_{\text{vac}} = 1$	$\mu_{\text{vac}} = \mu_0$
Fields from potentials	$\mathbf{E} = -\nabla V - \frac{1}{c} \frac{\partial \mathbf{A}}{\partial t}$ $\mathbf{B} = \nabla \times \mathbf{A}$	$\mathbf{E} = -\nabla V - \frac{\partial \mathbf{A}}{\partial t}$ $\mathbf{B} = \nabla \times \mathbf{A}$
Static potentials (coulomb gauge)	$V = \sum_{\text{charges}} \frac{q_i}{r_i}$ $\mathbf{A} = \frac{1}{c} \oint \frac{I d\mathbf{r}}{ \mathbf{r} - \mathbf{r}' }$	$V = \frac{1}{4\pi\epsilon_0} \sum_{\text{charges}} \frac{q_i}{r_i}$ $\mathbf{A} = \frac{\mu_0}{4\pi} \oint \frac{I d\mathbf{r}}{ \mathbf{r} - \mathbf{r}' }$
Relativistic transformations (\mathbf{v} is the velocity of primed system as seen in unprimed system)	$\mathbf{E}'_{\parallel} = \mathbf{E}_{\parallel}$ $\mathbf{E}'_{\perp} = \gamma \left(\mathbf{E}_{\perp} + \frac{1}{c} \mathbf{v} \times \mathbf{B} \right)$ $\mathbf{B}'_{\parallel} = \mathbf{B}_{\parallel}$ $\mathbf{B}'_{\perp} = \gamma \left(\mathbf{B}_{\perp} - \frac{1}{c} \mathbf{v} \times \mathbf{E} \right)$	$\mathbf{E}'_{\parallel} = \mathbf{E}_{\parallel}$ $\mathbf{E}'_{\perp} = \gamma (\mathbf{E}_{\perp} + \mathbf{v} \times \mathbf{B})$ $\mathbf{B}'_{\parallel} = \mathbf{B}_{\parallel}$ $\mathbf{B}'_{\perp} = \gamma \left(\mathbf{B}_{\perp} - \frac{1}{c^2} \mathbf{v} \times \mathbf{E} \right)$
$4\pi\epsilon_0 = \frac{1}{c^2} 10^7 \text{ A}^2 \text{ N}^{-1} = \frac{1}{8.987 55 \text{ K}} \times 10^{-9} \text{ F m}^{-1}$ $\frac{\mu_0}{4\pi} = 10^{-7} \text{ N A}^{-1} ; c = 2.997 924 58 \times 10^8 \text{ m s}^{-1}$		

Impedances (SI units)

ρ = resistivity at room temperature in $10^{-8} \Omega \text{ m}$:

~ 1.7 for Cu

~ 5.5 for W

~ 2.4 for Au

~ 73 for SS 304

~ 2.8 for Al

~ 100 for Nichrome

(Al alloys may have

double this value.)

For alternating currents, instantaneous current I , voltage V , angular frequency ω :

$$V = V_0 e^{j\omega t} = ZI .$$

Impedance of self-inductance L : $Z = j\omega L$.

Impedance of capacitance C : $Z = 1/j\omega C$.

Impedance of free space: $Z = \sqrt{\mu_0 / \epsilon_0} = 376.7 \Omega$.

High-frequency surface impedance of a good conductor:

$$Z = \frac{(1+j)\rho}{\delta} , \text{ where } \delta = \text{effective skin depth} ;$$

$$\delta = \sqrt{\frac{\rho}{\pi \nu \mu}} \cong \frac{6.6 \text{ cm}}{\sqrt{\nu [\text{Hz}]}} \text{ for Cu} .$$

Capacitance \hat{C} and inductance \hat{L} per unit length (SI units)

Flat rectangular plates of width w , separated by $d \ll w$ with linear medium (ϵ , μ) between:

$$\hat{C} = \epsilon \frac{w}{d} ; \quad \hat{L} = \mu \frac{d}{w} ;$$

$$\frac{\epsilon}{\epsilon_0} = 2 \text{ to } 6 \text{ for plastics; } 4 \text{ to } 8 \text{ for porcelain, glasses;}$$

$$\frac{\mu}{\mu_0} \cong 1 .$$

Coaxial cable of inner radius r_1 , outer radius r_2 :

$$\hat{C} = \frac{2\pi\epsilon}{\ln(r_2/r_1)} ; \quad \hat{L} = \frac{\mu}{2\pi} \ln(r_2/r_1) .$$

Transmission lines (no loss):

$$\text{Impedance: } Z = \sqrt{\hat{L}/\hat{C}} .$$

$$\text{Velocity: } v = 1/\sqrt{\hat{L}\hat{C}} = 1/\sqrt{\mu\epsilon} .$$

Motion of charged particles in a uniform, static magnetic field

The path of motion of a charged particle of momentum p is a helix of constant radius R and constant pitch angle λ , with the axis of the helix along **B**:

$$p[\text{GeV}/c] \cos \lambda = 0.29979 \, qB[\text{tesla}] \, R[\text{m}] ,$$

where the charge q is in units of the electronic charge. The angular velocity about the axis of the helix is

$$\omega[\text{rad s}^{-1}] = 8.98755 \times 10^7 \, qB[\text{tesla}] / E[\text{GeV}] ,$$

where E is the energy of the particle.

This section was adapted, with permission, from the 1999 web edition of the *Review of Particle Physics* (<http://pdg.lbl.gov>). See J. D. Jackson, *Classical Electrodynamics*, 2d ed. (John Wiley & Sons, New York, 1975) for more formulas and details. A [PDF version of this table](#) is also available.

X-Ray Data Booklet

Section 5.3 ELECTROMAGNETIC RELATIONS

	Gaussian CGS	SI
Units and conversions:		
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Lorentz force	$\mathbf{F} = q(\mathbf{E} + \frac{\mathbf{v}}{c} \times \mathbf{B})$	$\mathbf{F} = q(\mathbf{E} + \mathbf{v} \times \mathbf{B})$
Maxwell equations	$\nabla \cdot \mathbf{D} = 4\pi \rho$ $\nabla \times \mathbf{E} + \frac{1}{c} \frac{\partial \mathbf{B}}{\partial t} = 0$ $\nabla \cdot \mathbf{B} = 0$ $\nabla \times \mathbf{H} - \frac{1}{c} \frac{\partial \mathbf{D}}{\partial t} = \frac{4\pi}{c} \mathbf{J}$	$\nabla \cdot \mathbf{D} = \rho$ $\nabla \times \mathbf{E} + \frac{\partial \mathbf{B}}{\partial t} = 0$ $\nabla \cdot \mathbf{B} = 0$ $\nabla \times \mathbf{H} - \frac{\partial \mathbf{D}}{\partial t} = \mathbf{J}$
Linear media	$\mathbf{D} = \epsilon \mathbf{E}, \mathbf{B} = \mu \mathbf{H}$	$\mathbf{D} = \epsilon \mathbf{E}, \mathbf{B} = \mu \mathbf{H}$
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Permeability of free space	$\mu_{\text{vac}} = 1$	$\mu_{\text{vac}} = 4\pi \cdot 10^{-7}\text{ T A}^{-1}\text{ m}$
Fields from potentials	$\mathbf{E} = -\nabla V - \frac{1}{c} \frac{\partial \mathbf{A}}{\partial t}$ $\mathbf{B} = \nabla \times \mathbf{A}$	$\mathbf{E} = -\nabla V - \frac{\partial \mathbf{A}}{\partial t}$ $\mathbf{B} = \nabla \times \mathbf{A}$
Static potentials (coulomb gauge)	$V = \frac{1}{4\pi} \sum_{\text{charges}} \frac{q_i}{r_i}$ $\mathbf{A} = \frac{1}{c} \oint \frac{I d\mathbf{r}}{ \mathbf{r} - \mathbf{r}' }$	$V = \frac{1}{4\pi\epsilon_0} \sum_{\text{charges}} \frac{q_i}{r_i}$ $\mathbf{A} = \frac{\mu_0}{4} \oint \frac{I d\mathbf{r}}{ \mathbf{r} - \mathbf{r}' }$
Relativistic transformations (\mathbf{v} is the velocity of primed system as seen in unprimed system)	$E_{\parallel} = E_{\parallel}$ $E_{\perp} = \gamma(E_{\perp} + \frac{1}{c} \mathbf{v} \times \mathbf{B})$ $B_{\parallel} = B_{\parallel}$ $B_{\perp} = \gamma(B_{\perp} - \frac{1}{c} \mathbf{v} \times \mathbf{E})$	$E_{\parallel} = E_{\parallel}$ $E_{\perp} = \gamma(E_{\perp} + \mathbf{v} \times \mathbf{B})$ $B_{\parallel} = B_{\parallel}$ $B_{\perp} = \gamma(B_{\perp} - \frac{1}{c^2} \mathbf{v} \times \mathbf{E})$
	$\frac{1}{4\pi\epsilon_0} = \frac{1}{c^2} 10^7\text{ A}^2\text{ N}^{-1} = \frac{1}{8.987\,55 \cdot 10^{18}} 10^{-9}\text{ F m}^{-1}$ $\frac{\mu_0}{4\pi} = 10^{-7}\text{ N A}^{-1}; c = 2.997\,924\,58 \cdot 10^8\text{ m s}^{-1}$	

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$$= \sqrt{\frac{\omega \mu}{2\sigma}} = \frac{6.6 \text{ cm}}{\sqrt{[\text{Hz}]}} \text{ for Cu.}$$

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$$\frac{\epsilon}{\epsilon_0} = 2 \text{ to } 6 \text{ for plastics; } 4 \text{ to } 8 \text{ for porcelain, glasses;}$$

$$\frac{\mu}{\mu_0} = 1.$$

Coaxial cable of inner radius r_1 , outer radius r_2 :

$$\hat{C} = \frac{2\pi\epsilon}{\ln(r_2/r_1)}; \quad \hat{L} = \frac{\mu}{2\pi} \ln(r_2/r_1).$$

Transmission lines (no loss):

$$\text{Impedance: } Z = \sqrt{\hat{L}/\hat{C}}.$$

$$\text{Velocity: } v = 1/\sqrt{\hat{L}\hat{C}} = 1/\sqrt{\epsilon\mu}.$$

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The path of motion of a charged particle of momentum p is a helix of constant radius R and constant pitch angle θ , with the axis of the helix along \mathbf{B} :

$$p[\text{GeV} / c] \cos \theta = 0.29979 qB[\text{tesla}] R[\text{m}] ,$$

where the charge q is in units of the electronic charge. The angular velocity about the axis of the helix is

$$[\text{rad s}^{-1}] = 8.98755 \cdot 10^7 qB[\text{tesla}] / E[\text{GeV}] ,$$

where E is the energy of the particle.

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X-Ray Data Booklet

Section 5.4 RADIOACTIVITY AND RADIATION PROTECTION

The International Commission on Radiation Units and Measurements (ICRU) recommends the use of SI units. Therefore, we list SI units first, followed by cgs (or other common) units in parentheses, where they differ.

A. DEFINITIONS

Unit of activity = becquerel (curie):

$$1 \text{ Bq} = 1 \text{ disintegration s}^{-1} [= 1/(3.7 \times 10^{10}) \text{ Ci}]$$

Unit of absorbed dose = gray (rad):

$$\begin{aligned} 1 \text{ Gy} &= 1 \text{ J kg}^{-1} (= 10^4 \text{ erg g}^{-1} = 100 \text{ rad}) \\ &= 6.24 \times 10^{12} \text{ MeV kg}^{-1} \text{ deposited energy} \end{aligned}$$

Unit of exposure, the quantity of x- or γ -radiation at a point in space integrated over time, in terms of charge of either sign produced by showering electrons in a small volume of air about the point:

$$\begin{aligned} &= 1 \text{ C kg}^{-1} \text{ of air (roentgen; } 1 \text{ R} = 2.58 \times 10^{-4} \text{ C kg}^{-1}) \\ &= 1 \text{ esu cm}^{-3} (= 87.8 \text{ erg released energy per g of air}) \end{aligned}$$

Implicit in the definition is the assumption that the small test volume is embedded in a sufficiently large uniformly irradiated volume that the number of secondary electrons entering the volume equals the number leaving.

Unit of equivalent dose for biological damage = sievert. $1 \text{ Sv} = 100 \text{ rem}$ (roentgen equivalent for man). The equivalent dose in Sv = absorbed dose in grays $\times w_R$, where w_R is the radiation weighting factor (formerly the quality factor Q), which depends upon the type of radiation and other factors, as shown in Table 5-3. The equivalent dose expresses the long-term risk (primarily due to cancer and leukemia) from low-level chronic exposure.

B. RADIATION LEVELS

Natural annual background from all sources. In most of the world, the whole-body equivalent dose rate $\sim 0.4\text{--}4 \text{ mSv}$ (40–400 mrem). It can range up to 50 mSv (5 rem) in certain areas. The U.S. average $\sim 3.6 \text{ mSv}$, including about 2 mSv ($\sim 200 \text{ mrem}$) from inhaled natural radioactivity, mostly radon and radon daughters. This radon exposure value is for a typical house; radon exposure varies by more than an order of magnitude.

Cosmic ray background in counters (Earth's surface): $\sim 1 \text{ min}^{-1} \text{ cm}^{-2} \text{ sr}^{-1}$.

Man-made radiation dose: The greatest contribution to man-made radiation dose has been from irradiation from x-ray diagnostics in medicine, which accounts for about 20% of the average natural radiation dose.

Table 5-3. Radiation weighting factors.

Type of radiation	w_R
X- and γ -rays, all energies	1
Electrons and muons, all energies	1
Neutrons:	
< 10 keV	5
10–100 keV	10
0.1–2 MeV	20
2–20 MeV	10
> 20 MeV	5
Protons (other than recoils), > 2 MeV	5
Alphas, fission fragments, and heavy nuclei	20

Fluxes (per cm^2) to deposit one Gy, assuming uniform irradiation:

For photons:

$6.24 \times 10^9 I/Ef$, for photons of energy E [MeV], attenuation length I (g cm^{-2}), and fraction $f \ll 1$ expressing the fraction of the photon's energy deposited in a small volume of thickness $\ll I$ but large enough to contain the secondary electrons.
 2×10^{11} photons cm^{-2} for 1-MeV photons on carbon ($f \approx 0.5$).

For charged particles:

$6.24 \times 10^9/(dE/dx)$, where dE/dx [MeV $\text{g}^{-1} \text{cm}^2$], the energy loss per unit length, may be obtained from range-energy figures.

$3.5 \times 10^9 \text{ cm}^{-2}$ for minimum-ionizing singly-charged particles in carbon.

Quoted fluxes are good to about a factor of two for all materials.

Recommended exposure limits for radiation workers (whole-body dose):

ICRP: 20 mSv yr^{-1} averaged over 5 years, with the dose in any one year ≤ 50 mSv.

U.S.: 50 mSv yr^{-1} (5 rem yr^{-1}). Many laboratories in the U.S. and elsewhere set lower limits.

Lethal dose: Whole-body dose from penetrating ionizing radiation resulting in 50% mortality in 30 days (assuming no medical treatment), about 5 Gy (500 rads) as measured internally on body longitudinal center line. Surface dose varies owing to variable body attenuation and may be a strong function of energy.

This section was adapted, with permission, from the 1999 web edition of the *Review of Particle Physics* (<http://pdg.lbl.gov/>). For further information, see ICRP Publication 60, *1990 Recommendation of the International Commission on Radiological Protection* (Pergamon Press, New York, 1991) and E. Pochin, *Nuclear Radiation: Risks and Benefits* (Clarendon Press, Oxford, 1983).

X-Ray Data Booklet

Section 5.5 USEFUL EQUATIONS

The following pages include a number of equations useful to x-ray scientists, either expanding on subjects covered in this booklet or addressing topics not covered here. The equations have been drawn from D. T. Attwood, *Soft X-Rays and Extreme Ultraviolet Radiation: Principles and Applications* (Cambridge Univ. Press, Cambridge, 1999) [<http://www.coe.berkeley.edu/AST/sxreuv>], and the equation numbers refer to that volume, which should be consulted for further explanation and discussion. That reference also expands on the discussions in this booklet on zone plate optics, synchrotron radiation, and other topics.

General X-Ray Formulas

Wavelength and photon energy relationship:

$$\hbar\omega \cdot \lambda = hc = 1239.842 \text{ eV} \cdot \text{nm} \quad (1.1)$$

Number of photons required for 1 joule of energy:

$$1 \text{ joule} \Rightarrow 5.034 \times 10^{15} \lambda \text{ [nm] photons} \quad (1.2a)$$

X-Ray Scattering and Absorption

Thomson cross section for a free electron:

$$\sigma_e = \frac{8\pi}{3} r_e^2 \quad (2.45)$$

$$r_e = \frac{e^2}{4\pi\epsilon_0 mc^2} = 2.82 \times 10^{-13} \text{ cm} \quad (2.44)$$

and r_e is the classical electron radius.

Scattering cross section for a bound electron:

$$\sigma = \frac{8\pi}{3} r_e^2 \frac{\omega^4}{(\omega^2 - \omega_s^2)^2 + (\gamma\omega)^4} \quad (2.51)$$

Rayleigh cross section ($\omega^2 \ll \omega_s^2$):

$$\sigma_R = \frac{8\pi}{3} r_e^2 \left(\frac{\omega}{\omega_s} \right)^4 = \frac{8\pi}{3} r_e^2 \left(\frac{\lambda_s}{\lambda} \right)^4 \quad (2.52)$$

Scattering by a multi-electron atom:

$$\frac{d\sigma(\omega)}{d\Omega} = r_e^2 |f|^2 \sin^2 \Theta \quad (2.68)$$

$$\sigma(\omega) = \frac{8\pi}{3} |f|^2 r_e^2 \quad (2.69)$$

where the complex atomic scattering factor represents the electric field scattered by an atom, normalized to that of a single electron:

$$f(\Delta\mathbf{k}, \omega) = \sum_{s=1}^Z \frac{\omega^2 e^{-i\Delta\mathbf{k} \cdot \Delta\mathbf{r}_s}}{(\omega^2 - \omega_s^2 + i\gamma\omega)} \quad (2.66)$$

For forward scattering or long wavelength this reduces to

$$f^0(\omega) = \sum_{s=1}^Z \frac{\omega^2}{(\omega^2 - \omega_s^2 + i\gamma\omega)} = f_1^0 - if_2^0 \quad (2.72 \& 2.79)$$

Refractive index for x-ray radiation is commonly written * as

$$n(\omega) = 1 - \delta + i\beta = 1 - \frac{n_a r_e \lambda^2}{2\pi} (f_1^0 - if_2^0) \quad (3.9 \& 3.12)$$

where

$$\delta = \frac{n_a r_e \lambda^2}{2\pi} f_1^0(\omega) \quad (3.13a)$$

$$\beta = \frac{n_a r_e \lambda^2}{2\pi} f_2^0(\omega) \quad (3.13b)$$

Absorption length in a material:

$$\ell_{\text{abs}} = \frac{\lambda}{4\pi\beta} = \frac{1}{2n_a r_e \lambda f_2^0(\omega)} \quad (3.22 \& 3.23)$$

Mass-dependent absorption coefficient:

$$\mu = \frac{2r_e \lambda}{Am_u} f_2^0(\omega) \quad (3.26)$$

Atomic absorption cross section:

$$\sigma_{\text{abs}} = 2r_e \lambda f_2^0(\omega) = Am_u \mu(\omega) \quad (3.28a\&b)$$

Relative phase shift through a medium compared to a vacuum:

$$\Delta\phi = \left(\frac{2\pi\delta}{\lambda} \right) \Delta r \quad (3.29)$$

where Δr is the thickness or propagation distance.

* The choice of $+i\beta$ is consistent with a wave description $E = E_0 \exp[-i(\omega t - kr)]$. A choice of $-i\beta$ is consistent with $E = E_0 \exp[i(\omega t - kr)]$.

Snell's law:

$$\sin \phi' = \frac{\sin \phi}{n} \quad (3.38)$$

Critical angle for total external reflection of x-rays:

$$\theta_c = \sqrt{2\bar{\delta}} \quad (3.41)$$

$$\theta_c = \sqrt{2\bar{\delta}} = \sqrt{\frac{n_a r_e \lambda^2 f_1^0(\lambda)}{\pi}} \quad (3.42a)$$

Brewster's angle (or polarizing angle):

$$\phi_B \simeq \frac{\pi}{4} - \frac{\delta}{2} \quad (3.60)$$

Multilayer Mirrors

Bragg's law:

$$m\lambda = 2d \sin \theta \quad (4.6b)$$

Correction for refraction:

$$m\lambda = 2d \sin \theta \sqrt{1 - \frac{2\bar{\delta}}{\sin^2 \theta}} = 2d \sin \theta \left(1 - \frac{4\bar{\delta}d^2}{m^2\lambda^2} \right)$$

where $\bar{\delta}$ is the period-averaged real part of the refractive index.

$$\Gamma = \frac{\Delta t_H}{\Delta t_H + \Delta t_L} = \frac{\Delta t_H}{d} \quad (4.7)$$

Plasma Equations

Electron plasma frequency:

$$\omega_p^2 = \frac{e^2 n_e}{\epsilon_0 m} \quad (6.5)$$

Debye screening distance:

$$\lambda_D = \left(\frac{\epsilon_0 \kappa T_e}{e^2 n_e^2} \right)^{1/2} \quad (6.6)$$

No. of electrons in Debye sphere:

$$N_D = \frac{4\pi}{3} \lambda_D^3 n_e \quad (6.7)$$

Electron cyclotron frequency:

$$\omega_c = \frac{eB}{m} \quad (6.8)$$

Maxwellian velocity distribution for electrons characterized by a single-electron temperature κT_e :

$$f(v) = \frac{1}{(2\pi)^{3/2} v_e^3} e^{-v^2/2v_e^2} \quad (6.1)$$

where

$$v_e = \left(\frac{\kappa T_e}{m} \right)^{1/2} \quad (6.2)$$

Electron sound speed:

$$a_e = \left(\frac{\gamma \kappa T_e}{m} \right)^{1/2} \quad (6.79)$$

Critical electron density:

$$n_c \equiv \frac{\epsilon_0 m \omega^2}{e^2} = 1.11 \times 10^{21} \frac{\text{e/cm}^3}{\lambda^2(\mu\text{m})} \quad (6.112a \& b)$$

Refractive index of plasma is

$$n = \sqrt{1 - \frac{n_e}{n_c}} \quad (6.114b)$$

Ratio of electron energy in coherent oscillations to that in random motion:

$$\left| \frac{v_{os}}{v_e} \right|^2 = \frac{e^2 E^2}{m \omega^2 \kappa T_e} = \frac{I/c}{n_c \kappa T_e} \quad (6.131a)$$

$$\left| \frac{v_{os}}{v_e} \right|^2 = \frac{0.021 I (10^{14} \text{ W/cm}^2) \lambda^2 (\mu\text{m})}{\kappa T_e (\text{keV})} \quad (6.131b)$$

Spectral brightness of blackbody radiation within $\Delta\omega/\omega = 0.1\% \text{BW}$:

$$B_{\Delta\omega/\omega} = 3.146 \times 10^{11} \left(\frac{\kappa T}{eV} \right)^3 \frac{(\hbar\omega/\kappa T)^3}{(e^{\hbar\omega/\kappa T} - 1)} \frac{\text{photons/sec}}{(\text{mm})^2 (\text{mr})^2 (0.1\% \text{BW})} \quad (6.136b)$$

Photon energy at peak spectral brightness:

$$\hbar\omega|_{pk} = 2.822 \kappa T \quad (6.137)$$

where κ is the Boltzmann constant.

Stefan-Boltzmann radiation law (blackbody intensity at any interface):

$$I = \sigma T^4 \quad (6.141b)$$

where the Stefan-Boltzmann constant is

$$\sigma = \frac{\pi^2 \kappa^4}{60 c^2 \hbar^3} \quad (6.142)$$

With κT in eV:

$$I = \hat{\sigma} (\kappa T)^4 \quad (6.143a)$$

where $\hat{\sigma}$ is the modified Stefan-Boltzmann constant

$$\hat{\sigma} = \frac{\pi^2}{60 \hbar^3 c^2} = 1.027 \times 10^5 \frac{\text{watts}}{\text{cm}^2 (\text{eV})^4} \quad (6.143b)$$

Coherence

Longitudinal coherence length:*

$$\ell_{\text{coh}} = \lambda^2 / 2\Delta\lambda \quad (8.3)$$

Spatial or transverse coherence (rms quantities):

$$d \cdot \theta = \lambda / 2\pi \quad (8.5)$$

or in terms of FWHM values

$$d \cdot 2\theta|_{\text{FWHM}} = 0.44\lambda$$

Spatially coherent power within a relative spectral bandwidth $\lambda/\Delta\lambda = N$ for an undulator with N periods:

$$\bar{P}_{\text{coh},N} = \frac{e\lambda_u I}{8\pi\epsilon_0 d_x d_y \theta_x \theta_y \gamma^2} \cdot \left(\frac{\hbar\omega_0}{\hbar\omega} - 1 \right) f(\hbar\omega/\hbar\omega_0) \quad (8.7c)$$

where $\hbar\omega_0$ corresponds to $K = 0$, and where

$$f(\hbar\omega/\hbar\omega_0) = \frac{7}{16} + \frac{5}{8} \frac{\hbar\omega}{\hbar\omega_0} - \frac{1}{16} \left(\frac{\hbar\omega}{\hbar\omega_0} \right)^2 + \dots \quad (8.8)$$

When the undulator condition ($\sigma' \ll \theta_{\text{cen}}$) is satisfied, the coherent power within a relative spectral bandwidth $\Delta\lambda/\lambda < 1/N$, is

$$\bar{P}_{\text{coh},\lambda/\Delta\lambda} = \frac{e\lambda_u I \eta (\lambda/\Delta\lambda) N^2}{8\pi\epsilon_0 d_x d_y} \cdot \left(1 - \frac{\hbar\omega}{\hbar\omega_0} \right) f(\hbar\omega/\hbar\omega_0) \quad (\sigma'^2 \ll \theta_{\text{cen}}^2) \quad (8.10c)$$

where η is the combined beamline and monochrometer efficiency.

* The factor of two here is somewhat arbitrary and depends, in part, on the definition of $\Delta\lambda$. Equation (26) on page 2-14 omits this factor. See Attwood, op.cit., for further discussion.

Spatially coherent power available from a laser is

$$P_{\text{coh}} = \frac{(\lambda/2\pi)^2}{(d_x\theta_x)(d_y\theta_y)} P_{\text{laser}} \quad (8.11)$$

where P_{laser} is the total laser power.

Normalized degree of spatial coherence, or complex coherence factor:

$$\mu_{12} = \frac{\langle E_1(t) E_2^*(t) \rangle}{\sqrt{\langle |E_1|^2 \rangle} \sqrt{\langle |E_2|^2 \rangle}} \quad (8.12)$$

The van Cittert-Zernike theorem for the complex coherence factor is

$$\mu_{\text{OP}} = \frac{e^{-i\psi} \int \int I(\xi, \eta) e^{ik(\xi\theta_x + \eta\theta_y)} d\xi d\eta}{\int \int I(\xi, \eta) d\xi d\eta} \quad (8.19)$$

For a uniformly but incoherently illuminated pinhole

$$\mu_{\text{OP}}(\theta) = e^{-i\psi} \frac{2J_1(ka\theta)}{(ka\theta)} \quad (8.27)$$

which has its first null ($\mu_{\text{OP}} = 0$) at $ka\theta = 3.832$, which for $d = 2a$ corresponds to $d \cdot \theta = 1.22\lambda$.

EUV/Soft X-Ray Lasers

Growth of stimulated emission:

$$\frac{I}{I_0} = e^{GL} \quad (7.2)$$

where L is the laser length and G is the gain per unit length. For an upper-state ion density n_u and a density inversion factor $F(\leq 1)$

$$G = n_u \sigma_{\text{stim}} F \quad (7.4)$$

where the cross section for stimulated emission is

$$\sigma_{\text{stim}} = \frac{\lambda^3 A_{u\ell}}{8\pi c (\Delta\lambda/\lambda)} \quad (7.16)$$

$$\sigma_{\text{stim}} = \frac{\pi \lambda r_e}{(\Delta\lambda/\lambda)} \left(\frac{g_\ell}{g_u} \right) f_{\ell u} \quad (7.18)$$

where $A_{u\ell}$ is the spontaneous decay rate, $f_{\ell u}$ is the oscillator strength and g_ℓ/g_u is the ratio of degeneracy factors.

Laser wavelength scaling goes as $1/\lambda^4$:

$$\frac{P}{A} = \frac{16\pi^2 c^2 \hbar (\Delta\lambda/\lambda) GL}{\lambda^4} \quad (7.22)$$

Doppler-broadened linewidth:

$$\left. \frac{(\Delta\lambda)}{\lambda} \right|_{\text{FWHM}} = \frac{v_i}{c} = \frac{2\sqrt{2 \ln 2}}{c} \sqrt{\frac{\kappa T_i}{M}} \quad (7.19a)$$

where v_i is the ion thermal velocity, κT_i is the ion temperature, and M is the ion mass. With κT_i expressed in eV and an ion mass of $2m_p Z$

$$\left. \frac{(\Delta\lambda)}{\lambda} \right|_{\text{FWHM}} = 7.68 \times 10^{-5} \left(\frac{\kappa T_i}{2Z} \right)^{1/2} \quad (7.19b)$$

Lithography

Minimum printable line width:

$$L_w = k_1 \frac{\lambda}{NA} \tag{10.1}$$

where k_1 is a constant dominated by the optical system, but affected by pattern transfer processes.

Depth of focus:

$$DOF = \pm k_2 \frac{\lambda}{(NA)^2} \tag{10.2}$$

Degree of partial coherence:

$$\sigma = \frac{NA_{cond}}{NA_{obj}} \tag{10.3}$$

where the subscript cond refers to the condenser or illumination optics, and obj refers to the objective lens of the reduction optics.

International Technology Road Map for Semiconductors

	Years			
	2005	2008	2011	2014
1:1 lines (nm)	100	70	50	35
Isolated lines (nm)	65	45	30	20

<http://www.sematech.org>