

Comprehensive optical and data management infrastructure for high-throughput light-sheet microscopy of whole mouse brains

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Abstract. The comprehensive mapping and quantification of neuronal projections in the central nervous system requires high-throughput imaging of large volumes with microscopic resolution. To this end, we have developed a confocal light-sheet microscope that has been optimised for the 3D imaging of structurally intact clarified whole-mount mouse brains. Here we describe the optical and electro-mechanical arrangement of the microscope and give details on the organisation of the microscope management software. The software orchestrates all components of the microscope, coordinates critical timing and synchronisation and has been written in a versatile and modular structure using the LabVIEW language. It can easily be adapted and integrated to other microscope systems and has been made freely available to the light-sheet community. The tremendous amount of data routinely generated by light-sheet microscopy further requires novel strategies for data handling and storage. To complete the full imaging pipeline of our high-throughput microscope we further elaborate on big data management from streaming of raw images up to stitching of 3D data sets. The meso-scale neuroanatomy imaged at micron-scale resolution in those datasets allows characterization and quantification of neuronal projections in unsectioned mouse brains.

Keywords: Light-sheet microscopy, whole brain imaging, data management, software management, imaging pipeline, high-throughput imaging.

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1 Introduction

The highly ambitious project of mapping and understanding each and every neuronal projection in the whole brain has been moved to feasible reality by the recent advent of light-sheet fluorescent microscopy (LSFM).¹⁻³ With this technique 3D data sets can be acquired with a resolution that is high enough to identify neurons and their dendritic and axonal features in time scales which are no longer the bottle neck of high-throughput acquisition. In LSFM, the sample is illuminated with a thin sheet of light confined into the focal plane of the detection objective, which collects the fluorescence emission along an axis perpendicular to the illumination plane.⁴ This technique drastically reduces the imaging acquisition time by recording millions of pixels in parallel due to its wide-field detection and affords intrinsic optical sectioning due to the light-sheet illumination. Consequently fluorophores outside the light-sheet are neither bleached nor contribute blurring out-of-focus noise. Together these attributes cause a high spatio-temporal resolution and an excellent signal to noise ratio in LSFM.

Several challenges remain to be overcome, however, to allow fast and, most of all, systematic production of reliable datasets and their meaningful interpretation to further our understanding of neuronal networks. Those challenges include fast, cheap and reproducible sample preparation and mounting, automated image acquisition and, most crucially, storage, interpretation and analysis of

the vast data sets light-sheet microscopy routinely produces.^{5,6} The mapping and understanding of this “big data” is an immense task that requires the expertise of computer scientists to employ fully automated, scalable post-processing, for example, to do blood vessel segmentation or cell counting.⁷ On the other hand, the imaging of large, intrinsically opaque samples in LSFM necessitates clearing protocols based on refractive index matching which render the tissue transparent.^{8,9} Therefore technological advances in LSFM need to be matched by novel development in the area of information- and biotechnology.

Here we will present a state-of-the-art light-sheet microscope, as it is implemented in our lab, that is especially well suited to acquire 3D data sets of clarified and structurally intact mouse brains. The LSFM features double-sided illumination with a digitally scanned light-sheet and a sample chamber which has been specifically designed for the the imaging of large ($> 1\text{cm}^3$) immersed and freely movable samples. After briefly introducing the optical setup we explain how to prepare and mount the samples for stable 3D imaging for several days. In addition, a full description of the control hardware and software is presented. The microscope software coordinates every microscope component and ensures critical timing and synchronisation between those components. We further give a systematic and robust approach to handle and store light-sheet microscopy data of several terabytes. To conclude the comprehensive and high-throughput imaging pipeline we furthermore elaborate on stitching of large 3D data sets. The imaging capabilities of our light-sheet microscope are exemplified with tomographies of whole transgenic mouse brains.

2 Optical setup

2.1 Laser unit

The custom-made confocal light-sheet microscope is equipped with five linearly polarised, continuous-wave lasers for fluorescence excitation (Fig. 1,A). The wavelengths were chosen to excite the most common fluorophores (see Table 2 for the manufacturer and specifications of microscope components). The beam from each laser is first collimated and expanded with a telescope made of singlet lenses (f140, f200) and then combined into a common path with a mirror and a long-pass filter (Semrock, LaserMUXTM series). From here on all lenses are achromatic doublets coated for the visible spectrum. An acousto-optical tunable filter (AOTF) acts as fast (μs) electronically tunable filter and intensity modulator. Due its nonlinear response we measured the AOTF light transmittance for each wavelength as a function of radio frequency amplitude and determined a look-up table to linearise the output. The zero order light is blocked by an iris. A 1:1 telescope (f150,f150) is used to re-collimate the beam after the AOTF. An electro-optical laser modulator (EOM) acts like a wave plate with electronically controlled retardation and rapidly rotates (few hundreds ns) the input polarisation of the excitation light by 90° . A digital line is used to switch between the voltage levels, allowing control of the frequency with which the polarisation is changed (see also section 3.2 on timing).

2.2 Illumination unit

The light impinges onto a polarising beam splitter cube which splits the excitation light depending on its polarisation into one of the two identical excitation arms. The beam diameter is reduced with a telescope (f350, f200) whose telecentric plane coincides with the mirrored surface of a galvanometric scanner (galvo). Both galvos are mounted on a custom made optical breadboard

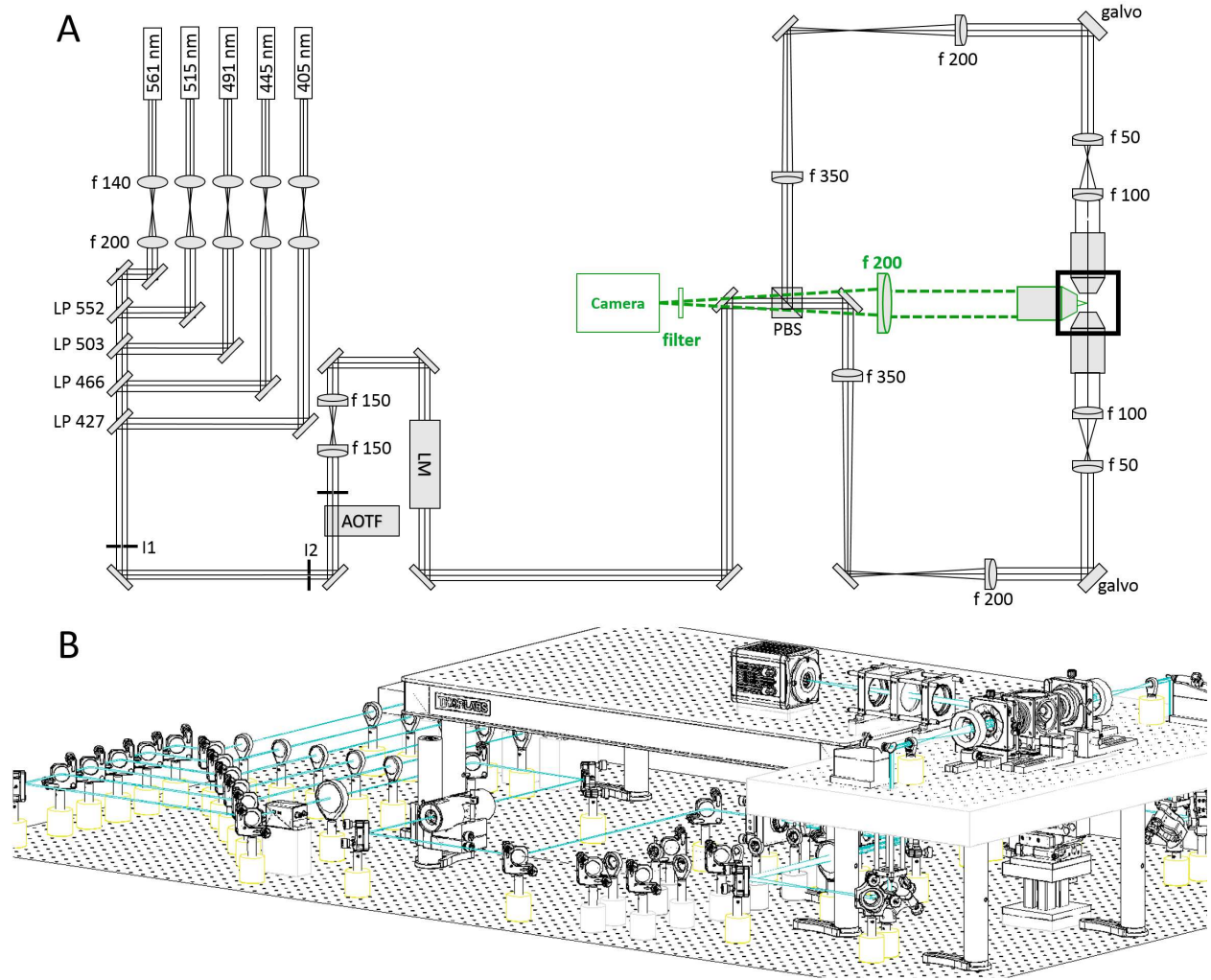


Fig 1: (A) Schematic of the light path. The galvo scanners are mounted above periscopes. LP: long-pass filter, I: iris, AOTF: acousto-optical tunable filter, LM: laser modulator, PBS: polarisation beam splitter. Green dotted: detection path. (B) Oblique view of the microscope. A custom-made breadboard serves to mount the sample chamber and objectives at an elevated height and features two circular holes at the edges for the periscopes and a large central cut-out for the translation stages. A second breadboard is used for the camera.

which features two circular holes to pass the periscopic beams and a large central cut-out for the sample chamber and motor stages (Fig. 1,B).

Each scan mirror surface is re-imaged with a telescope (f50, f100) onto the back aperture of a long working distance low magnification objective (Nikon, 10x 0.3NA WD 17.5mm). The focal lengths of the doublets making up the telescope have been chosen for space considerations and to achieve a $1/e^2$ beam radius of 0.73mm at the back focal plane resulting in an effective NA of 0.025 (see also section 2.5). The galvos are in planes conjugate with the back focal plane of their respective excitation objective to ensure that a tilt angle of the galvo converts into a lateral translation of the Gaussian beam across the specimen. The two excitation objectives are designed for air immersion but are immersed into the clearing solution ($n=1.45$) within the sample chamber. A coverslip glued to the front housing edge serves the dual purpose of maintaining the first diffractive surface between the front optical element and the air and to protect the front lens elements from the clearing agent. A guide to align objectives to a common focus can be found in the supplementary information. The light-sheet is generated digitally^{10,11} by scanning the excitation beam across the focal plane of the detection objective. This generates incoherent illumination resulting in fewer artefacts. Additionally each line in the specimen is illuminated with the same intensity creating a homogeneous light-sheet which is particularly advantageous for the quantitative and qualitative investigation of large samples in their entirety.

2.3 Detection unit

Fluorescence is collected with an objective that is specifically designed for immersion in clearing solutions. The objective is equipped with a correction collar allowing for immersion in media with refractive indices ranging from 1.41 to 1.52 (Olympus, XLSLPLN25XGMP, 25x, 1.0NA, effective focal length 7.2 mm, working distance 8mm). A tube lens of 200mm creates an image on a sCMOS camera (Orca Flash4.0, Hamamatsu) with a chip of over 4 megapixels. The cell size of the OrcaFlash is $p^2 = (6.5\mu m)^2$ over an active area of $r^2 = (13.3mm)^2$. With a 200mm tube lens the field of view (FOV) in the sample is $(480\mu m)^2$ (see also Table 3 in the supplementary information for relevant formulas).

In the camera's rolling shutter acquisition mode only a subset of adjacent horizontal pixel lines is simultaneously exposed and this active detection region is moved across the image sensor.¹³ With the delay between the exposure of two adjacent lines set to the minimum time required to read out a single line ($t_{\text{shutter}} = 9.7\mu s$) the entire frame of 2048 horizontal lines is activated within 19.86ms. Setting the line exposure time to $t_{\text{exp}} = 0.3ms$, for instance, results in $l = t_{\text{exp}}/t_{\text{shutter}} \approx 31$ lines being simultaneously exposed at any time. This detection area is swept from the top to the bottom of the chip and acts like a moving, virtual confocal slit corresponding to a width of $s_{\text{object}} = l p / M_{\text{eff}} = 7.2\mu m$ in the sample space, where M_{eff} is the effective magnification of the detection objective. With the given values of line exposure and line read out time, the camera acquires images at a full-frame rate of $R_{\text{frame}} = 1/(\#_{\text{horiz. lines}} * t_{\text{shutter}} + t_{\text{exp}}) \approx 50Hz$.

The camera was run as master in the internal trigger mode and a synchronisation pulse generated with the exposure of each new line was used as an output trigger for a DAQ board (NI PCIe-6353, National Instruments) which was operated as slave. The saw-tooth driving signal for the galvo mirrors was generated by the DAQ and synchronised with the pixel line reset pulses from the camera to achieve confocal line detection (see also section 3.2 on synchronisation).

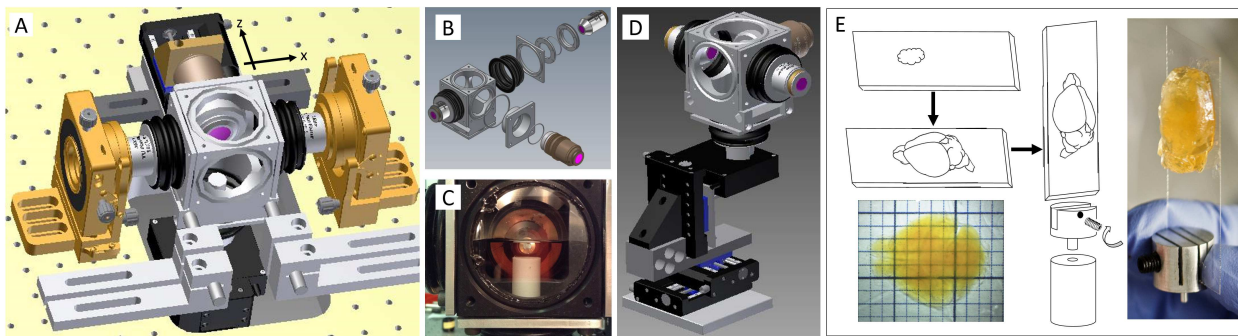


Fig 2: Schematic of sample mounting and motion. (A) Fluorescence excitation (along x axis) and detection (along z axis) are operated on independent perpendicular light paths where the excitation light-sheet and the detection focal plane overlap. (B) The custom-made sample chamber is assembled with silicone bellows and seal rings. The clarified fluorescently-labelled brain is mounted on a Teflon cylinder in the centre of the watertight chamber filled with clearing solution (C) and can be translated and rotated freely with piezo motors (D). (E) Mounting of a clarified fluorescently labelled brain. The brain is glued onto a coverslip and inserted into an adapter that slides into the Teflon cylinder.

2.4 Sample mounting and motion

We designed a cubic water-tight sample chamber (Fig. 2,A) that allows access from all six sides while maintaining the 3D integrity of large clarified and fluorescently labelled mouse brains. The sample chamber is tightly bolted to the optical breadboard while soft connections using silicone bellows allow for adjustable movements of the objectives and free 3D motion of the motor stages (Fig. 2,B). Connections are sealed with rubber rings and silicone caulk and additionally tightened with cable binders. In this way the objectives can be refocused and realigned without compromising the watertight seal of the chamber (Fig. 2,C). The clarified brains are imaged immersed in clearing solution composed of 63% 2,2'-thiodiethanol (TDE) in Phosphate buffered saline (PBS) and a refractive index of $n \approx 1.45$.¹² To fill the entire volume of the sample chamber approximately 150ml clearing solution are needed. A motorized x-, y-, z-, θ -stage (see Table 2) allows free 3D motion and rotation of a Teflon cylinder which reaches into the centre of the chamber (Fig. 2,D). Illumination and detection axes are horizontal while sample rotation occurs around the vertical axis.

To elucidate neuronal projections in structurally intact tissue it is paramount to be able to image centimeter-sized clarified samples, such as mouse brains, with high resolution in whole mount preparation. The problem of sample mounting in a light-sheet microscope arises from the fact that optical access to the sample is required from 3-4 planar sides leaving only two opposing sides to insert, fix and move the sample. Stable mounting is hereby a key concern as a whole-brain tomography can require image acquisition in excess of 24 hours and the effects of gravity, tissue shrinking/expansion and evaporation of the clearing solution over such time spans might have to be considered.

For the acquisition of whole-brain data sets, we fix a clarified and fluorescently-labelled mouse brain with super glue to a coverslip (Fig. 2,E). The brain is oriented along its axial orientation with the olfactory bulbs at the top and the cerebellum at the bottom of the coverslip. The coverslip is

slid into a bottom adaptor and tightened with a plastic-capped grub screw. The bottom adaptor is inserted into a Teflon cylinder with the coverslip being positioned on the far side of the detection objective. Three different slits in the adapter correspond to varying distances to the detection objective and give variability in sample thickness for example to allow also for the mounting of rat brain hemispheres.

2.5 Optical characterisation

2.5.1 Theoretical considerations

The optical properties of a microscope can be quantified by the point-spread function (PSF), which is the image of a point source in object space. The radial (i.e. perpendicular to the detection axis) and axial (i.e. parallel to the detection axis) resolution of the microscope are most commonly described by the full width at half-maximum (FWHM) of the PSF itself along the chosen direction (Fig. 3,A). In a light-sheet microscope axial and radial resolution are determined by the numerical aperture (NA) of the illumination and detection optics respectively. The radial PSF is a standard Airy function whose FWHM is given by:

$$\text{FWHM}_{\text{radial}} \simeq \frac{0.51\lambda_0}{\text{NA}_d}, \quad (1)$$

λ_0 being the wavelength in vacuum and NA_d the NA of the detection objective. The axial resolution is dominated by the thickness of the light sheet which for high NA detection objectives is usually larger than the depth of field of the detection objective. If Gaussian beams are used for illumination, the thickness of the light sheet can be calculated by the well-known formulas of Gaussian optics.¹⁴ The radius of the beam waist is given by:

$$w_0 = \frac{\lambda f_i^{\text{eff}}}{\pi w}, \quad (2)$$

where w is the $1/e^2$ radius of the beam entering the illumination objective, f_i^{eff} is the effective focal length of the illumination optics and λ is the wavelength in the medium ($\lambda = \lambda_0/n$, n being the refractive index). Note that for air objectives (as the one we use for illumination) the effective focal length is extended by a factor n when used inside a medium with refractive index different from unity.¹⁵ The axial FWHM is then given by:

$$\text{FWHM}_{\text{axial}} \simeq 1.177w_0. \quad (3)$$

The thickness of the light sheet is not constant throughout the FOV due to the limited depth of focus of the Gaussian beam. The latter can be quantified by the confocal parameter b , which is the distance along which the beam $1/e^2$ radius remains relatively constant and smaller than $\sqrt{2}w_0$:

$$b = \frac{2\pi w_0^2}{\lambda} = \frac{2\lambda(f_i^{\text{eff}})^2}{\pi w^2}. \quad (4)$$

The need to have an almost uniform resolution in the camera FOV imposes a lower bound to b and thus (according to the previous equations) a limit to the axial resolution. In our setup we adjusted the telescopes in the excitation path in order to have $w \simeq 0.73\text{mm}$, resulting in $b \simeq 346\mu\text{m}$ (corresponding to 72% of the FOV) and $\text{FWHM}_{\text{axial}} \simeq 5.13\mu\text{m}$ (assuming $\lambda_0 = 0.5\mu\text{m}$). An

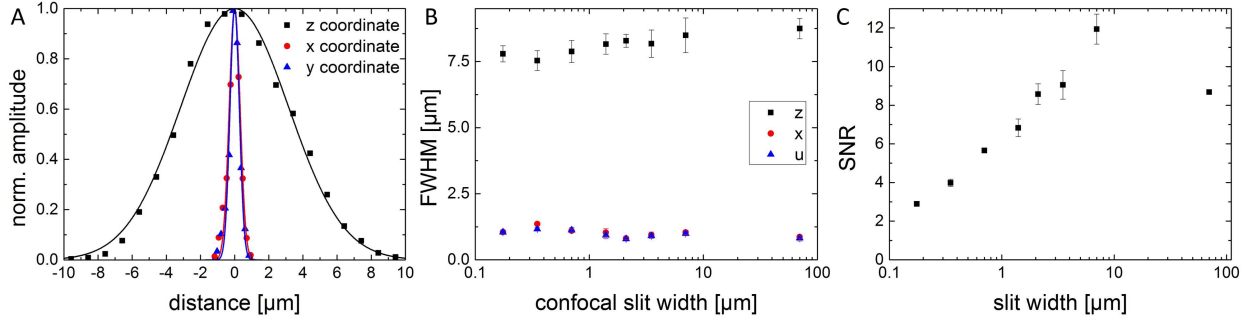


Fig 3: Optical characterisation. (A) Exemplary lateral and axial intensity profiles of nanometric beads and fits to Gaussian curves. (B) FWHM of the Gaussian fits with confocal slit width. Each data point is the average of 10 bead profiles. (C) Signal to noise ratio (SNR) as a function of slit width.

overview of the relevant formulas concerning the optical properties of a light-sheet microscope is given in Table 3. The reader is referred to this table as a quick guide to tailor their microscope to their specific needs.

2.5.2 Experimental characterisation

For optimal imaging it is crucial that the sample-emitted line of fluorescence is in the centre of the rolling shutter at the start of each image acquisition and furthermore travels with the same speed during stack acquisition. To this end the galvo scanners were calibrated using a fluorescein solution in the sample chamber and an image of the focal line was acquired for a set of mirror angles. By fitting each focal line profile to a Gaussian curve the position of the beam main axis was determined and plotted against the voltage applied to the galvo scanner. A linear fit to this data yielded the precise amplitude and offset for each galvo scanner. The finite width of the confocal slit was hereby taken into account such that maximum beam deflection corresponded to the central pixel line of the confocal slit coinciding with the first and last pixel row on the camera chip. From the same data the average FWHM of the Gaussian excitation line over the FOV of the camera was calculated as $(8.60 \pm 0.67) \mu m$ ($n = 14$). The disagreement with the theoretical value of $5.13 \mu m$ is likely attributed to the refractive index mismatch introduced in the excitation arm by the protective coverslip and by the optical path inside a medium with a refractive index different from air, both causing spherical aberrations.¹⁶

For quantitative analysis of spatial resolution and the signal to noise ratio (SNR) image stacks of nanometric fluorescent beads ($0.51 \mu m$ diameter, Bangs Laboratories) were acquired. The spatial resolution, in agreement with previous findings,^{17,18} was not affected by the confocal slit width on the camera (Fig. 3,B). The lateral and axial FWHMs were $(1.00 \pm 0.04) \mu m$ and $(8.13 \pm 0.14) \mu m$ respectively, averaged through stack depths and confocal slit widths. The experimentally determined lateral FWHM was larger than the theoretically calculated value of $0.26 \mu m$ by a factor of four. The discrepancy is likely caused by higher order aberrations introduced by the gel embedding the fluorescent beads exacerbated by the high NA of our detection objective.^{19,20}

Qualitatively, the SNR increased with increasing slit width up to a maximum around the $1/e^2$ illumination beam diameter of $8.6 \mu m$ and then decreased again with further increasing slit width

(Fig. 3,C) in agreement with previous observations.¹³

3 Microscope management software

3.1 General design

The many components of the microscope (camera, lasers, AOTF, stages, etc.) are orchestrated and operated through a custom-made software written in the LabVIEW language (National Instruments, Austin, US). Since the microscope should work mainly without human intervention for many hours (a whole-brain tomography can last a couple of days) without losing precise synchronization, stringent software design is a primary concern. The overall architecture of the software and its practical implementation were chosen to satisfy two main principles; firstly, operation in a distributed environment and secondly, robustness and ease of maintenance. This means, the apparatus should be able to deal with many independent hardware components and with several data streams (see below). Individual processing steps and/or data transfer can be computationally intensive, and may need to be distributed over several computers to perform properly at run-time. Any kind of hardware failure, from timing delays to complete system crash, must be automatically handled to avoid loss of data.

An object-oriented programming framework³⁰ has been used for software development, dividing the entire project into well-defined self-contained modules. This framework, besides helping defining the proper scope of each software component, guarantees robustness, ease of maintenance, scalability and reusability. Code was stored on GitHub (www.github.org) to help multiple users developing software simultaneously. The Murmex software development kit (see supplementary information 7.4) was used to operate the software efficiently and synchronously on different computers.

3.2 System timing

A schematic overview of the hardware architecture of our LSFM is shown in Fig. 4. One personal computer (Precision T5600, Dell, see also Table 2) controlled all the remote hardware components with the exception of the camera, which was controlled by a second, dedicated computer (Precision T7500, Dell) to handle and manage the enormous streams of data produced (see also section 4). Various instrumentations were addressed either through computer ports (as USB, RS232 and CameraLink) or via analog and/or digital signals generated by the DAQ board.

Synchronization of the camera rolling shutter to the scanning of the light sheet was achieved by letting the camera acquiring in free-run mode, and using a trigger output from the camera itself to time the generation of the waveforms sent to the galvo drivers. The HSYNC signal from the camera was used which provided true-type-logic (TTL) pulse every time the rolling shutter stepped from one line to the next one. This trigger signal was used as a clock by the DAQ board so that the line producing the light sheet moved step by step with the rolling shutter. The same HSYNC signal was also used to trigger the DAQ digital signal sent to the laser modulator. Since at each position of the scanning line the modulator should go through at least two states (vertical and horizontal polarization), HSYNC was used as a trigger while the much faster internal clock of the DAQ board served as clock.

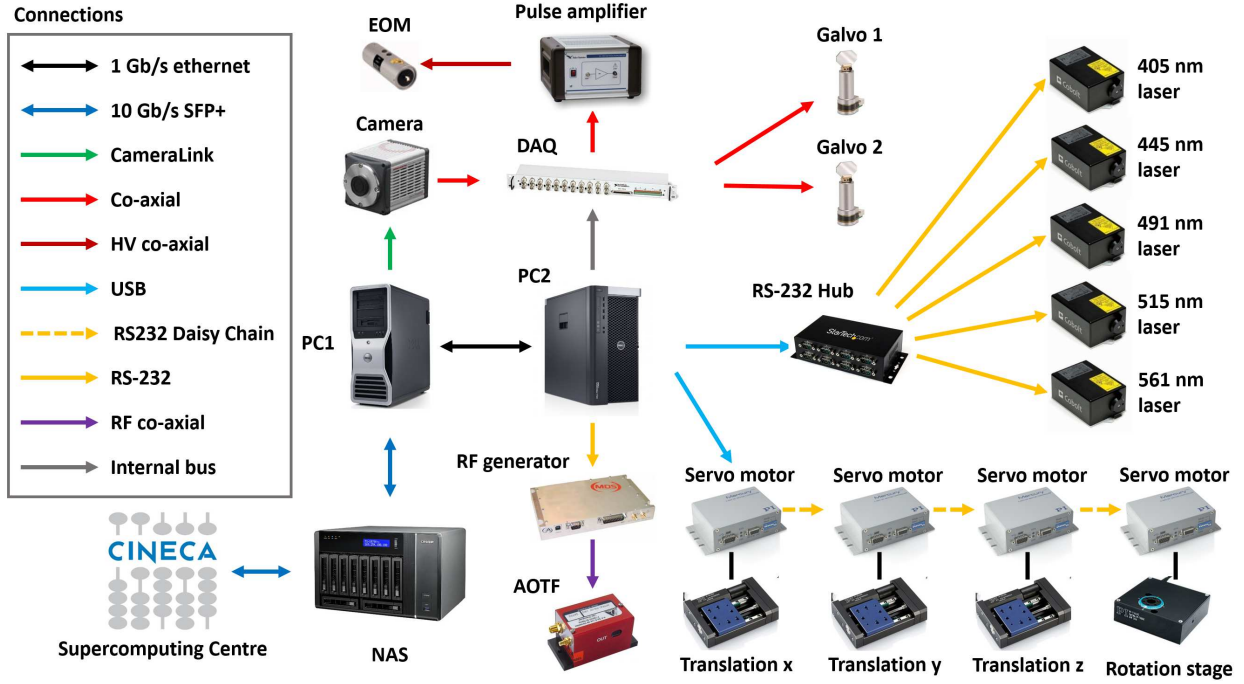


Fig 4: Schematic view of the hardware architecture of the microscope. Different kind of physical connections are depicted with different colours.

3.3 Software organisation

Each hardware component in the LSFM was managed through its own software component (see supplementary information 7.4). Since the galvo scanners and the electro-optical modulator were not addressed directly, but via analog and digital signals generated by the DAQ board, their software component just sent analog and digital waveforms to an additional intermediate component responsible for DAQ board management. Beyond hardware-related software, two additional software components, the Stack and Tomo components orchestrated the operation of the entire system.

The Stack component was used to acquire image stacks along the detection axis, here the z direction, and coordinates the Camera and the Stages components. Stacks are collected by continuously moving the sample along z while keeping the camera acquiring in free run. The translation speed was calculated such that in the time needed to collect a single frame the specimen was moved by a z step specified in the GUI. The Camera component used low-level libraries from Hamamatsu to save the data, producing one single *.cxd* file (the proprietary, non-compressed, tiff-like format from Hamamatsu) for each image stack. The Camera component must write data to disk with perfect synchronisation, otherwise the nominal z position of each frame is no more coincident with the real z position of the plane imaged. For this reason, the process executing the Camera code was set to highest priority. Additionally, the Stack component checked the timing of data production, in order to re-acquire the image stack if the camera had encountered a timing issue.

The Tomo component is a high-level component which coordinated the acquisition of adjacent, parallel, partially-overlapping stacks. The user can specify the volume to be imaged either by inserting the minimum and maximum x , y and z coordinates of a parallelepiped or by providing

by a dedicated 10 Gbit/s link provided by the Consortium GARR (the Public Institution managing research networks in Italy) to CINECA (the main Italian supercomputing centre). The data flow scheme is summarized in Fig. 5.

While transferring data from the SSDs to the NAS, we perform on-the-fly conversion of files generated during sample acquisitions to a losslessly compressed multi-page TIFF format. Our conversion module leveraged the functionalities offered by the interface of the DCIMGAPI library, supplied by Hamamatsu, and the freely available LibTIFF library. The module read a single *.cxd* 3D image and wrote it in a variable number of multi-page TIFF files, depending on the user-defined limit of pages per file. Since the output of the Hamamatsu camera may exceed tens of GB, the image in proprietary format was read slice by slice to minimize memory requirements. The conversion module ran concurrently with the microscope controlling software and was synchronized to convert stack files as soon as they are completed. The module provided options to downsample the image or to save the input image at full or reduced resolution in order to adapt to storage requirements. Usually we downsampled the image in the XY plane by a factor 2 and left the Z sampling unaltered. This downsampling, in combination with the lossless compression offered by the TIFF format, reduced data size by about one order of magnitude. Thus the parallelepiped volume containing the whole mouse brain took up about 2-3 TB of disk space.

The process executing the file conversion was set to low priority to prevent interference with the Camera component, which needs to be perfectly timed with acquisition. Although file conversion drastically reduces data size, in the current implementation it was slower than data production. The camera flushed data to the SSDs with a rate R_{data} dependent on the frame rate R_{frame} and on the image size: $R_{\text{data}}[\text{Bytes}] = 2 \cdot R_{\text{frame}} \cdot \#_{\text{pixels}}$ where the factor of two takes into account the 16 bit grey depth of images. The volumetric imaging rate is given by $R_{\text{volume}} = 0.5 R_{\text{data}} \cdot \text{voxel size}$.

With a standard image size of 2048×2048 pixels, and $R_{\text{frame}} = 50$ Hz, data were produced at a rate $R_{\text{data}} = 400$ MB/s ≈ 1.37 TB/h, which resulted in a volumetric rate $R_{\text{volume}} \approx 2.3 \times 10^7 \mu\text{m}^3/\text{s} \approx 0.08 \text{ cm}^3/\text{h}$. The conversion module freed the SSDs with an average rate of ≈ 0.7 TB/h, meaning that the SSDs were actually filled up during imaging at a rate of ≈ 0.67 TB/h. Imaging sessions needed thus to be shorter than 16.3 hours which limited the maximum amount of raw data being acquired continuously to ≈ 22.3 TB. However, this was well above the disk space needed to image the parallelepiped volume encompassing a whole mouse brain. Larger specimens, as rat brain or portions of human brain, must be imaged in consecutive session until faster strategies for image conversion are devised.

After acquisition, the tiled raw images produced had to be stitched to a 3D data set. To perform this operation we used the TeraStitcher,²² a fully automated stitching tool designed to deal with tiled images of virtually unlimited size. TeraStitcher can run on machines with limited resources with a performance essentially determined only by the transfer rate towards secondary storage. TeraStitcher can be freely downloaded from <https://github.com/abria/TeraStitcher>.

5 Representative results

Some representative data produced by our microscope are reported in Fig. 6. With the apparatus described in this paper, whole mouse brains could be imaged in their entirety (Fig. 6,A), maintaining a high quality of the image also when going deep inside the specimen (Figs. 6,B). The resolution and contrast of the images allowed to distinguish axonal tracts (Fig. 6,C) in the brain

Table 1: Data production rates for typical camera operation.

Voxel size	$0.234 \times 0.234 \times 2$	μm^3
Frame rate	50	Hz
Frame size	2048×2048	pixels
Frame disk size	8	MB
Data production rate	1.37	TB/h
Data conversion rate	0.7	TB/h
Net SSD filling rate	0.67	TB/h
SSD RAID size	10.9	TB
Max acquisition time to fill up SSDs	16.3	h
Max raw data size that can be acquired continuously	22.3	TB
Max volume that can be imaged	1.04	cm^3
(considering 10% superposition between adjacent stacks)		

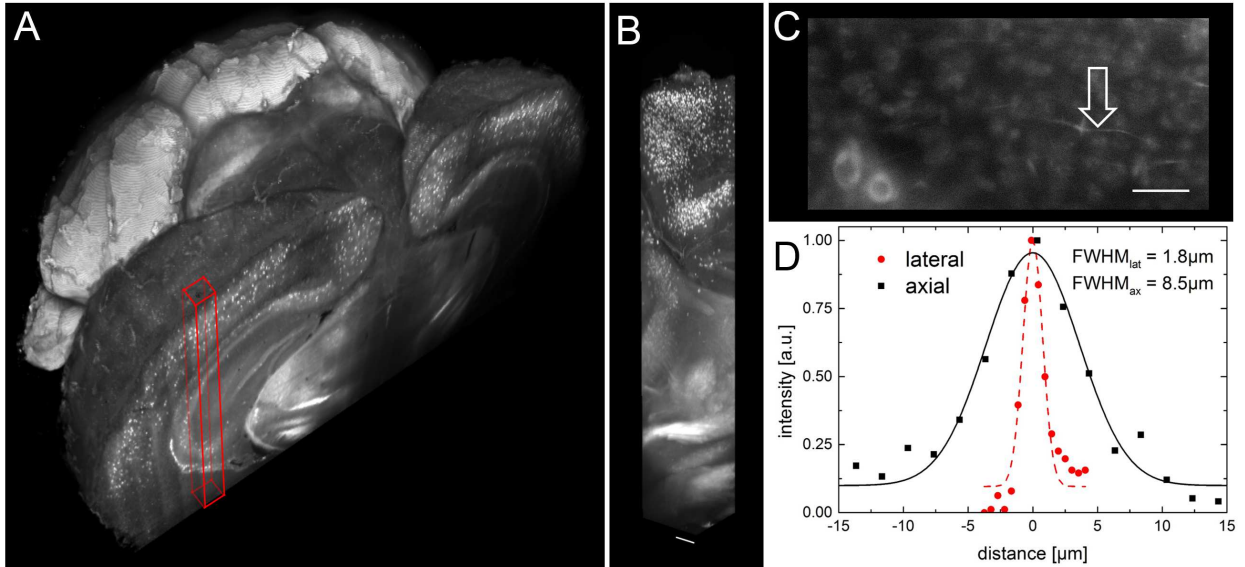


Fig 6: Whole mouse brain tomography. Imaging of whole transgenic mouse brains treated with CLARITY and cleared with TDE 63% imaged with Olympus, 25X objective. (A) 3D rendering of a parvalbumin-dTomato brain. (B) 3D rendering of a stack from PV-dTomato mouse brain, scale bar = $400\mu\text{m}$. (C) High resolution insert showing the axon (arrow) used for resolution estimation. Scale bar = $50\mu\text{m}$. (D) Lateral and axial FWHM of the intensity profiles through axon indicated in C. Modified with permission from.¹²

of a PV-cre-dTomato transgenic mouse, in which parvalbuminergic neurons are fluorescently labelled. GABAergic interneurons can be clearly identified using a GAD2-cre-dTomato mouse (Fig. 6D), while nuclear staining with propidium iodide (PI) permits to visualize all the cells (neuronal or not) in the brain (Fig. 6E). Such whole-brain datasets can be used to localize cell distributions in 3D using automatic cell-detection software.⁷

6 Discussion and conclusion

Light-sheet microscopy has already been a game changer for large-scale, high-throughput imaging by yielding data of high spatio-temporal scale. This measurement technique will continue to have a high impact in the field of whole-brain imaging due to its ability to record millions of images over the course of days or even weeks. Here we have presented a complete and detailed framework for whole mouse brain imaging. Starting with a full description of our double-sided light-sheet microscope and its constituent components we further gave details of the custom-written control software which coordinates the microscope's reliably synchronised operation. The data produced in our experiments easily amounts several TB per data set and needs to be compressed, stored, transferred, retrieved and processed necessitating the concurrent development of novel computational interface and analysis methods. Here we have presented a comprehensive, robust and fully automated pipeline of data management starting from the streaming of raw images up to the stitching of 3D data sets.

The mouse brains imaged in our microscope are intrinsically opaque and need to be clarified before imaging. A successful optical clearing method not only renders the tissue transparent but also causes neither quenching of protein fluorescence nor tissue distortion and is compatible with repetitive immunostaining. A very promising method called CLARITY^{8,9} transforms tissue into a nanoporous, hydrogel-hybridized, lipid-free form. Crucially for the high-throughput imaging of large samples that we aim for, the protocol, however, is also fast, easy and cheap. We therefore developed a versatile, simple, rapid and inexpensive clearing method described elsewhere¹² which in combination with CLARITY allows for the optical clearing of entire mouse brains or portions of human tissue.

Even in refractive index-matched samples we have noticed that significant aberrations degrade image quality, especially at depths towards the end of the stack. Recently objectives specifically conceived for LSFM have become commercially available²³ which have been designed for use with clearing solutions and typically also feature a correction collar for the compensation of a certain amount of spherical aberration. A complete correction of system-induced aberrations, however, requires the ability to also correct for residual spherical aberration and higher-order aberrations introduced by the sample mount, immersion medium and slight misalignments of the optical system. Of course, aberrations introduced by the sample itself also degrade image quality and it is those aberrations that are most difficult to correct at run-time due to their modal complexity, large amplitude and high spatial variability.²⁶

Additionally another problem inherent to light-sheet microscopy degrades image quality: due to the uncoupled nature of excitation and detection in LSFM, the aberrations encountered in each arm are independent too and this can lead to a walk-off between the light sheet and the detection FOV when large samples are imaged at depths. Adaptive optics²⁴ provides a means to compensate for these aberrations and has recently been applied to improve LSFM.²⁵ The prospect of high quality images is especially essential for high through-put imaging methods for whole-brain 3D

reconstructions because the size of data sets are so large they warrant fully automated post processing for example for automated cell counting and identification.⁷ The effect of optical aberrations can severely obscure features even above the resolution limit and therefore hinder automated segmentation of the data. We have mounted the camera on a separate breadboard which leaves us with a lot of free, infinity-corrected space in the detection path. The implementation of adaptive optics technology into the fluorescence emission path of our LSFM in order to demonstrate real-time aberration corrected images for whole-brain reconstructions will be considered in the future.

Further improvements in image quality may be achieved with beam shaping in the illumination arm, for example by implementing Bessel beam illumination.²⁷⁻²⁹ Using these self-reconstructing, non-diffractive beams the optical sectioning capability can be improved over an extended FOV.

The streamlining of high-throughput imaging of entire mouse brains, starting from the live animal up to quantifiable data sets, requires tight integration of all aspects of the work flow. Even if the brains we image are transparent, their large size means that also comparatively small heterogeneities in the refractive index matching sum up to large optical path differences and therefore strong aberrations at large imaging depth. The microscope used for whole-brain imaging needs to be developed and adapted bearing those imaging circumstances in mind and sample preparation methods, imaging technology and data management schemes have to work as a unit in order to deliver on the highly anticipated, quantifiable data sets of whole-brain neuronal projections.

Table 2: Overview of microscope components.

Component	Manufacturer	Part#	Specifications
Lasers	Cobolt AB, Sweden	MLD	405nm, 80mW, s-polarised
		MLD	445nm, 50mW, s-polarised
		Calypso	491nm, 50mW, s-polarised
		Fandango	515nm, 50mW, s-polarised
		Jive	561nm, 50mW, s-polarised
AOTF	AA Opto-Electronic, France	AOTFnC-400.650-TN	> 90% diffraction efficiency, 3nm resolution, low cross talk between laser lines, high separation angle
Laser modulator	Qioptiq GmbH, Germany	LM 0202 VIS ADP	400-650nm , $\lambda/2$ -voltage (633nm): 210V
Pulse amplifier	Falco Systems, The Netherlands	WMA-300	50x amplification up to ± 150 V, DC to 5MHz signal bandwidth
Galvo scanner	Cambridge Technology, USA	6220H	small angle step response $200\mu s$
Objectives	Nikon; Japan	Plan Fluor EPI	10x0.3NA, WD 17.5mm, EFL 20mm (excitation)
	Olympus, Japan	LMPLFLN20X	20x0.4NA, WD 12mm, EFL 9mm (detection)
Motor stages	Physik Instrumente, Germany	C-863.11	DC servo-motor controller
		M-122	Travel range 25mm, $0.1\mu m$ resolution, max. velocity 20mm/s
		M-116	Precision Rotation Stage, $2.5\mu rad$, max. velocity $20^\circ/s$
Camera	Hamamatsu, Japan	Orca Flash4.0 V2.0	sCMOS sensor, 2048(H) x 2048(V), cell dim.: $6.5\mu m$, active area: 13.3mm x 13.3mm, 16bit images
DAQ board	National Instruments, USA	NI PCIe-6353	AI: 1 MS/s multichannel; 16-bit resolution, 10 V; AO: 2.86 MS/s, 16-bit resolution, 10 V; digital I/O lines (hardware-timed up to 10 MHz), 100MHz max counter frequency
Workstation 1	Dell, USA	T7500	12GB RAM, Intel Xeon Processor X5647 @ 2.93 GHz, OS Windows 7 64 bit
Workstation 2	Dell, USA	T5600	8GB RAM, Intel Xeon Processor E5-2620 @ 2 GHz, OS Windows 7 64 bit

Acknowledgments

The authors are grateful to Riccardo Ballerini and Ahmed Hajeb from the mechanics workshop at LENS for fabrication of the sample chamber. We further thank Marco De Pas from the electronic workshop for his expertise in the fabrication of the custom-made amplification electronics. We thank GARR for the 10GB optical fibre connection to CINECA and CINECA for hosting our data. The research leading to these results has received funding from the European Union Seventh Framework Program (FP7/2007-2013) under grant agreements no. 604102 (Human Brain Project) and n° 284464 (LASERLAB-EUROPE). The research has also been supported by the Italian Ministry for Education, University and Research in the framework of the Flagship Project NanoMAX, by “Ente Cassa di Risparmio di Firenze” (private foundation). Research activities were also supported by Regione Toscana in the program POR-CreO 2007-2013 (Linea di interventi 1.5.a - 1.6 - Bando Unico R&S 2012) under grant agreement (CUP) n° 6408.30122011.026000201.

Conflict of interest

M.v.H has a financial interest in Murmex by Distrio, Amsterdam the Netherlands.

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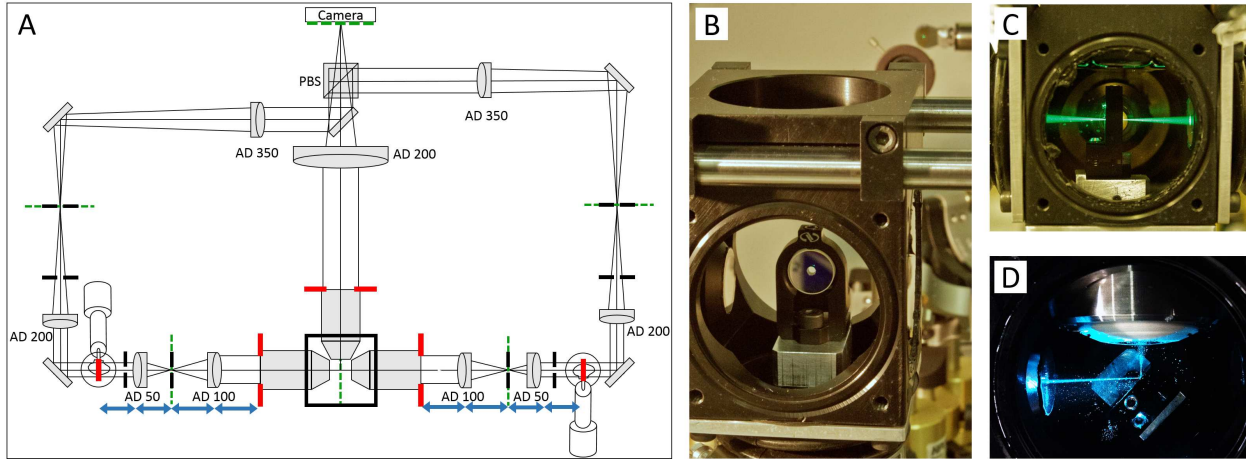


Fig 7: (A) Basic geometrical optics of a double-sided illumination light-sheet microscope. AD: achromatic doublet, PBS: polarising beam splitter, red: objective back focal planes and conjugated telecentric planes, green: image planes, blue: 4f telecentric lens system, black: alignment irises. (B) Alignment mirror with drilled hole for light transmission, mount for tip adjustment and adaptor for the Teflon tube. (C) With lateral movement the alignment mirror can be placed such that light is transmitted into the opposing excitation arm. (D) The alignment mirror can be rotated to precisely reflect by 90° .

7 Supplementary information

7.1 Alignment

For LSFM it is crucial that excitation and detection occur on perpendicular axes because any deviation from this geometry results in obscured images of reduced resolution and contrast. The objectives need to be perfectly confocal so that the fluorescence generated in the swept excitation beam also falls within the detection objective's focal plane. Telecentric imaging, that is the imaging with two lenses which are the sum of their focal lengths apart (Fig. 7,A), is used to re-image the galvo scanners onto the back apertures of the excitation objectives. In the detection path, the camera chip needs to be positioned in an image plane of the detection objective. Additionally, homogeneous illumination from both sides impose strict symmetry considerations on both illumination arms that have not only to be sufficiently aligned with respect to themselves but function as a pair with recursive dependence.

Our microscope was aligned with two tools, a shear plate to qualitatively assess collimation and a small mirror that can be mounted inside the sample chamber. The 0.5in mirror was first pierced with a drill using a ceramic drill bit to produce a hole roughly in its centre of the approximately the same diameter as the excitation beam inside the sample chamber (Fig. 7,B). Using a compact single axis adjustable mirror mount (V50-AX, Newport) attached to the Teflon cylinder allows to adjust the pitch of the reflection with the mount and the yaw with the rotation stage while at the same time providing a very space efficient mounting for the pierced 0.5in mirror. With this “sample mirror” three different position can be easily implemented, firstly, back reflection by hitting the reflective surface at 0° , secondly, transmission by laterally displacing the mirror until the light passes through the drilled hole (Fig. 7,C), and thirdly, reflection of the light by 90° by

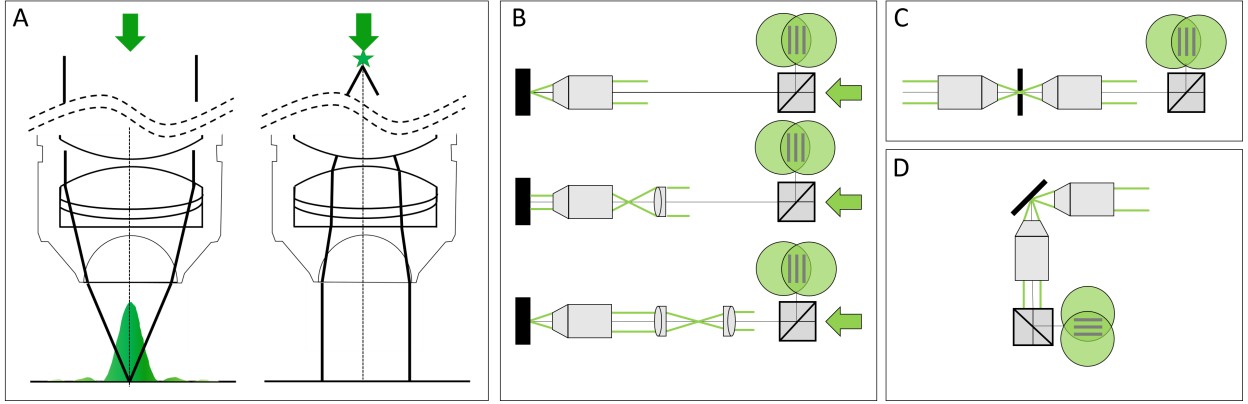


Fig 8: (A) Working principle of a microscope objective (left): a collimated input beam is converted into a spherical wavefront. Right: inverted light path, a focal spot in the back aperture creates a collimated output. (B) Recursive placement of lenses using the sample mirror in back-reflection mode, (C) in transmission mode and (D) in reflection mode.

precisely turning the mirror mount by 45° using the rotation stage (Fig. 7,D).

The excitation objectives are fixed on mounts which allow three dimensional translation plus pitch and yaw adjustment (LP-1A, Newport). In a first alignment step the beam paths of both excitation arms were brought to overlap through irises placed on the breadboard and the optical bench without any microscope objectives. We found it useful to use two mirrors in each periscope, one vertically mounted and one mounted at 45° for vertical deflection of the incoming beam. In this way full beam steering can be achieved to realign the periscope before hitting the galvo scanner. After this initial alignment the sample mirror was placed with its reflective surface in the centre of the central cut-out of the breadboard. The first excitation objective was placed into the beam path and brought to focus onto the sample mirror which was adjusted to reflect the light back into the same objective through the irises. Using a shear plate and a beam splitter cube the back reflected light was qualitatively adjusted for collimation using translation along the excitation axis of the objective mount (Fig. 8,B). The sample mirror was then moved laterally to allow the excitation beam to pass through the drilled hole and the second excitation objective was placed (Fig. 8,C). This time collimation was checked with the light going through both objectives and so their confocal placement was insured. Both excitation arms were then aligned recursively by starting at the putative confocal point between the objectives and placing successively lens after lens in the direction towards the light source, alternating between the two modalities illustrated in Fig. 8,A. Finally, the sample mirror was positioned in deflection mode and the detection objective was aligned (Fig. 8,D).

7.2 Sample preparation for optical characterisation

The fluorescent beads (CP01F, Bang Laboratories, $0.51\mu\text{m}$ diameter) used for quantitative analysis of spatial resolution and signal to noise ratio (SNR) were embedded in low- melting point agarose (4% agarose in 63%TDE in PBS). The agarose (A9793, Sigma-Aldrich) was dissolved then mixed with a diluted bead solution and left to polymerise over night on a coverslip. The coverslip was mounted inside the sample chamber as described in section 2.4. The embedded sub-diffraction

fluorescent beads were imaged with varying confocal slit widths ranging from $0.175\mu m - 70\mu m$ in the sample space. For each slit width one stack of $200\mu m$ depth ($1\mu m$ step size) was acquired and the one dimensional PSF in the lateral and axial directions were extracted manually with FIJI/ImageJ throughout the depth of the image stack. After fitting the lateral and axial profiles with Gaussian curves, the lateral and axial FWHM were averaged ($n=10$ data points) for different slit widths. The SNR was calculated according to $SNR = (y_0 + A)/y_0$ where y_0 and A are the offset and amplitude of the Gaussian fits.

7.3 Sample preparation for whole-brain imaging

Mouse brains were processed using a modified CLARITY protocol⁸ and cleared with a new refractive-index-matching medium.¹² Adult mice (post-natal day 56) were anaesthetized with isoflurane and transcardially perfused with 20 ml ice-cold phosphate-buffered saline (PBS) solution (pH 7.6) followed by 20 ml of a mixture of 4% (wt/vol) paraformaldehyde (PFA), 4% (wt/vol) acrylamide, 0.05% (wt/vol) bis-acrylamide, 0.25% (wt/vol) VA044 in PBS. We then extracted the brains from the skull and put them in the same PFA/acrylamide/bis-acrylamide/VA044/PBS solution at $4^\circ C$ for 3 days. Afterwards, we degassed the samples and increased their temperature to $37^\circ C$ to initiate polymerization. The embedded samples were then extracted from the gel and washed with clearing solution (sodium borate buffer (200 mM, pH 8.5) containing 4% (wt/vol) sodium dodecyl sulfate) at $37^\circ C$ through gentle shaking. To perform electrophoretic tissue clearing (ETC), hydrogel-embedded brains were placed in a custom-built organ-electrophoresis chamber. The same clearing solution was circulated through the chamber and a voltage of 20V was applied across the ETC chamber at $37^\circ C$ for several days. After clearing, brains were incubated in PBST0.1 (PBS and 0.1% Triton X-100, pH 7.6) at $37^\circ C$ for 2 days to remove the clearing solution. As a final step, murine brain were incubated in 50 ml of 30% and 63% (vol/vol) 2,2'-thiodiethanol in 0.01M PBS (TDE/PBS), each for 1 day respectively at $37^\circ C$ while gently shaking. The final 63% TDE/PBS solution has a refractive index $n \simeq 1.45$, allowing for refractive index matching and thus tissue transparency.

7.4 Murmex

To operate the software efficiently and synchronously on different machines, the [Murmex](#) software development kit was used. This kit integrates the object-oriented programming paradigm with a messaging system, and can be used to create communicating yet independent software modules within a standardized programming scheme.

Murmex creates distributed finite state machines, design patterns in which one component proceeds from a particular state to another based on the message it receives either from itself or from other components. For sending and receiving messages, Murmex uses the library [LabbitMQ](#), which is a wrapper of [RabbitMQ](#), a message-oriented middle-ware, however, these layers of abstraction are mostly hidden from the developers. Buffered, asynchronous and reliable communication between several software components is established by routing of messages by the broker to the specific software component pinpointed by its ID. Murmex has twelve pre-defined states (eg, Initialise, Configure, Start, Stop, Proceed, Update) and a generic state for implementing custom states. Usually, just a few of this states are actually used in standard application. An example of a standard implementation of the Murmex state machine for the camera is shown in Fig. 9. Each Murmex component streams information about its current status to its 'observers' which are

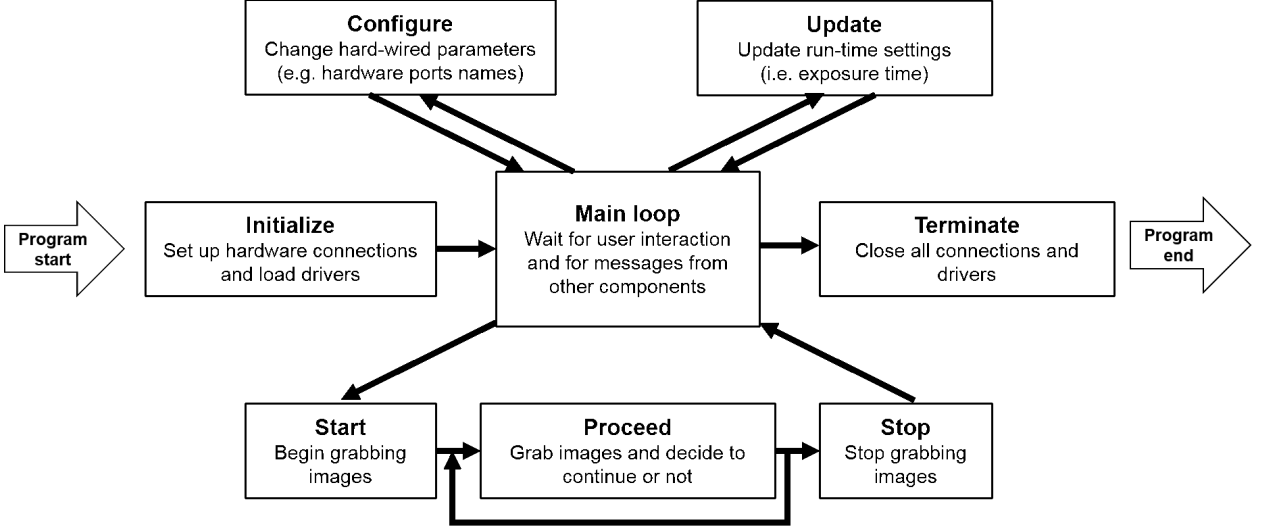


Fig 9: Schematic view of the work flow of a representative Murmex component operating a camera.

other components specified in its configuration file. Using this observer/observee hierarchy, it is possible to design a large and complex network of different components, running on various computers connected through a local connection or the internet while maintaining precise timing and synchronisation.

7.5 Formulas

In the table, f_{TL}^{design} is the focal length of the tube lens for which the objective has been designed (200 mm for Leica and Nikon, 180 mm for Olympus and 165 mm for Zeiss), f_{TL}^{real} the focal length of the tube lens actually employed, M the nominal magnification of the objective, n the refractive index of the immersion medium, λ_0 the wavelength in vacuum, λ the wavelength in the medium, w the $1/e^2$ radius of the Gaussian beam at the objective back focal plane, FN the field number of the objective, p the physical pixel size, r_{chip} the size of the camera chip, $t_{shutter}$ the pace of the rolling shutter (i.e. the time between exposure start in subsequent lines), t_{exp} the line exposure time. The subscripts i and d refers to illumination and detection, respectively.

Table 3: Overview of the main formulas related to light-sheet microscopy.

General formulas for microscope objectives			
Nominal focal length	f	$=$	$f_{\text{TL}}^{\text{design}}/M$
Effective focal length	f^{eff}	$=$	nf
Effective magnification	M^{eff}	$=$	$f_{\text{TL}}^{\text{real}}/f^{\text{eff}}$
Excitation (Gaussian)			
Beam waist	w_0	$=$	$\lambda f_i^{\text{eff}}/\pi w$
Axial resolution	$\text{FWHM}_{\text{axial}}$	\approx	$1.17w_0$
Confocal parameter	b	$=$	$2\pi\omega_0^2/\lambda$
		$=$	$2\lambda(f_i^{\text{eff}})^2/\pi w^2$
Detection			
Radial resolution	$\text{FWHM}_{\text{radial}}$	$=$	$0.5\lambda_0/\text{NA}_d$
Objective field size	S	$=$	FN/M
Camera			
Pixel size object space	p_{object}	$=$	p/M^{eff}
Camera FOV object space	$\text{FOV}_{\text{object}}$	$=$	$r_{\text{chip}}/M^{\text{eff}}$
		$=$	$p_{\text{object}} \cdot \#\text{pixels}$
Slit width on camera	s_{camera}	$=$	$p \cdot t_{\text{exp}}/t_{\text{shutter}}$
Slit width in object space	s_{object}	$=$	$s_{\text{camera}}/M^{\text{eff}}$
Frame rate	ν	$=$	$1/(\#\text{horiz. lines} \cdot t_{\text{shutter}} + t_{\text{exp}})$