

**Q5.12**

Since  $\epsilon(t') = \epsilon_0 \sin \Omega t$ , place it into  $\sigma(t)$ :

$$\sigma(t) = \int_{-\infty}^t dt' G(t - t') \frac{d\epsilon_0 \sin \Omega t'}{dt'} = \int_{-\infty}^t dt' G(t - t') \epsilon_0 \Omega \cos \Omega t$$

Let  $\tau = t - t'$ , so

$$\begin{aligned} \sigma(t) &= \epsilon_0 \Omega \int_{-\infty}^0 d\tau (-G(\tau)) \cos(\Omega(t - \tau)) \\ &= \epsilon_0 \Omega \int_0^{\infty} d\tau G(\tau) \cos(\Omega(t - \tau)) \\ &= \epsilon_0 \Omega \int_0^{\infty} d\tau G(\tau) (\cos \Omega t \cos \Omega \tau + \sin \Omega t \sin \Omega \tau) \\ &= \epsilon_0 \left[ \left( \Omega \int_0^{\infty} d\tau G(\tau) \sin \Omega \tau \right) \sin \Omega t + \left( \Omega \int_0^{\infty} d\tau G(\tau) \cos \Omega \tau \right) \cos \Omega t \right] \\ &= \epsilon_0 (G'(\Omega) \sin \Omega t + G''(\Omega) \cos \Omega t) \end{aligned}$$

**Q6.1**

$\mathbf{A}, \mathbf{B}$  is unitary matrix, so  $\mathbf{A}^\dagger \mathbf{A} = I = \mathbf{A}^{-1} \mathbf{A}$  and  $\mathbf{B}^\dagger \mathbf{B} = I = \mathbf{B}^{-1} \mathbf{B}$

(a).  $\mathbf{A}^\dagger \mathbf{A} = I$

(b).  $\mathbf{A}^\dagger \mathbf{A}^{-1} = \mathbf{A}^\dagger \mathbf{A}^\dagger = \mathbf{A}^{\dagger^2}$

(c).  $(c\mathbf{A}\mathbf{A}^\dagger)^\dagger = (c\mathbf{I})^\dagger = \bar{c}\mathbf{I}$

(d).  $\begin{pmatrix} \mathbf{A} & \mathbf{0} \\ \mathbf{0} & \mathbf{B} \end{pmatrix}^{-1} = \begin{pmatrix} \mathbf{A}^{-1} & \mathbf{0} \\ \mathbf{0} & \mathbf{B}^{-1} \end{pmatrix} = \begin{pmatrix} \mathbf{A}^\dagger & \mathbf{0} \\ \mathbf{0} & \mathbf{B}^\dagger \end{pmatrix}$

### Q6.3

$\langle \mathbf{x}, \mathbf{y} \rangle \mathbf{z} = \mathbf{x}^\dagger \mathbf{y} \mathbf{z}$ , where  $\mathbf{x}^\dagger \mathbf{y}$  is multiplication of  $1 \times N$  and  $N \times 1$  matrix and the result is scalar  $m$ . Therefore,

$$\langle \mathbf{x}, \mathbf{y} \rangle \mathbf{z} = m \mathbf{z} = \mathbf{z} m = \mathbf{z} \mathbf{x}^\dagger \mathbf{y}$$

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Need to expand into elements details

$$\begin{aligned} \langle \mathbf{x}, \mathbf{y} \rangle \mathbf{z} &= \langle \mathbf{x}, \mathbf{y} \rangle \mathbf{z}[n] \\ &= \sum_{m=0}^{M-1} x^*[m] y[m] z[n] \\ &= \sum_{m=0}^{M-1} x_m^* y_m z_n \\ &= \sum_{m=0}^{M-1} z_n x_m^* y_m \\ &= \sum_{m=0}^{M-1} [z x^*]_{nm} y_m \\ &= [\mathbf{z} \mathbf{x}^\dagger \mathbf{y}]_n \end{aligned}$$

### Q6.4

$$(a). \mathbf{y} = \mathbf{A}\mathbf{x} = \sum_{n=1}^N A[m, n]x[n]$$

$$(b). \mathbf{x} = [\mathbf{A}^\dagger \mathbf{y}] = \sum_{m=1}^M [A^\dagger]_{nm} y_m = \sum_{m=1}^M A^*[n, m] y[m]$$

$$(c). \mathbf{y} = \mathcal{A}(x(t')) = \int_{-\infty}^{\infty} dt' a_m(t') x(t')$$

$$(d). x(t') = \mathbf{A}^\dagger \mathbf{y}(t') = \sum_{m=1}^M a_m^*(t') y_m$$

$$(e). y(t) = \mathcal{A}\{x(t')\} = \int_{-\infty}^{\infty} dt' a(t, t') x(t')$$

$$(f). x(t') = \mathcal{A}^\dagger\{y(t)\} = \int_{-\infty}^{\infty} dt a^*(t, t') y(t)$$

Assume  $\mathbf{y}$  is  $m \times 1$  matrix, and  $\mathbf{x}$  is  $n \times 1$  matrix,

(a). It is a discrete-to-discrete transformation in transformation matrix form.  $\mathbf{A}$  is  $m \times n$  matrix.

$$\mathbf{y} = \mathbf{A}\mathbf{x} = \left( \sum_{i=1}^n \mathbf{A}_{ji} \mathbf{x}_i \right), j = (1, \dots, m)$$

If  $\mathbf{A}$  is fourier basis vectors, according to §6.8 and forward DFT,

$$\mathbf{y} = \frac{1}{\sqrt{N}} \sum_{n=0}^{N-1} \mathbf{x}[n] \exp(-i2\pi nk/N) = \mathbf{Y}[k]$$

(b). Similarly,  $\mathbf{A}^\dagger$  is  $n \times m$  transformation matrix.

$$\mathbf{x} = [\mathbf{A}^\dagger \mathbf{y}] = \left( \sum_{i=1}^m \mathbf{A}_{ji}^\dagger \mathbf{y}_i \right), j = (1, \dots, n)$$

Use §5.8 inverse DFT, since the unitary transformation constant is  $1/\sqrt{N}$ , then

$$\mathbf{x} = [\mathbf{A}^\dagger \mathbf{y}] = \frac{1}{\sqrt{N}} \sum_{k=0}^{N-1} \mathbf{y}[k] \exp(i2\pi kn/N) = \mathbf{X}[n]$$

(c). It's a continuous-to-discrete transformation.  $\mathcal{A}$  is a transformation

operator. According to §5.4 forward FS, so

$$\mathbf{y} = \mathcal{A}\{x(t')\} = \int_{-\infty}^{\infty} dt' \mathcal{A}(x(t')) = \frac{1}{T_0} \int_{-T_0/2}^{T_0/2} dt' x(t') \exp(-i2\pi kt'/T_0) = \mathbf{Y}[k]$$

(d). It's a discrete-to-continuous transformation. According to §5.4 inverse FS, so

$$x(t') = [\mathbf{A}^\dagger \mathbf{y}](t') = \sum_{k=-\infty}^{\infty} \mathbf{A}^\dagger[k] \mathbf{y}[k] = \sum_{k=-\infty}^{\infty} \mathbf{y}[k] \exp(i2\pi kt'/T_0)$$

(f). It's a forward CT-FT as §5.5. So

$$y(t) = \mathcal{A}\{x(t')\} = \int_{-\infty}^{\infty} dt' \mathcal{A}(x(t')) = \int_{-\infty}^{\infty} dt' x(t') \exp(-i2\pi tt')$$

(g). It's a inverse CT-FT as §5.5. So

$$x(t') = \mathcal{A}^\dagger\{y(t)\} = \int_{-\infty}^{\infty} dt \mathcal{A}^\dagger(y(t)) = \int_{-\infty}^{\infty} dt y(t) \exp(i2\pi t't)$$

### Q6.5

Way 1: (Thanks Joseph Tibbs)

Move  $e$  to the left of equation so that  $\mathbf{g} - \mathbf{e} = \mathbf{H}\mathbf{f}$ . Then apply  $\mathbf{Q}^\dagger$  to both side,

$$\mathbf{Q}^\dagger(\mathbf{g} - \mathbf{e}) = \mathbf{Q}^\dagger\mathbf{H}\mathbf{f}$$

$$\frac{1}{\sqrt{N}} \sum_{n=0}^{N-1} (g[n] - e[n]) \exp(-i2\pi nk/N) = \mathbf{Q}^\dagger\mathbf{H}\mathbf{f} = \mathbf{\Lambda}\mathbf{Q}^\dagger\mathbf{f}$$

$$\frac{1}{\sqrt{N}} \sum_{n=0}^{N-1} g'[n] \exp(-i2\pi nk/N) = \mathbf{\Lambda}\mathbf{Q}^\dagger\mathbf{f}, \text{ Let } g' = g - e$$

$$G'(k) = H(k)F(k), \text{ use equation 6.8}$$

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Way 2:

Apply forward fourier operater to  $\mathbf{g} = \mathbf{H}\mathbf{f} + \mathbf{e}$ , then

$$\mathbf{Q}^\dagger\mathbf{g} = \mathbf{Q}^\dagger\mathbf{H}\mathbf{f} + \mathbf{Q}^\dagger\mathbf{e}.$$

Apply equation 6.7 and 6.8,

$$\mathbf{Q}^\dagger\mathbf{g} = \mathbf{\Lambda}\mathbf{Q}^\dagger\mathbf{f} + \mathbf{Q}^\dagger\mathbf{e}$$

$$\implies G(k) = H(k)F[k] + \frac{1}{\sqrt{N}} \sum_{k=0}^{N-1} e[n] \exp(-i2\pi nk/N)$$

$$= H(k)F(k) + E(k)$$

Because  $\mathbf{e}$  exists constantly in the time domain, then as example 5.4.1 shows, the narrow functions in one domain imply broad functions in the other and vice versa. In this case, the  $E(k)$  would have non-zero values in a very narrow width in frequency domain while the left values are all zeros. Therefore,  $G(k) \approx H(k)F(k)$  which is the Fourier convolution theorem.

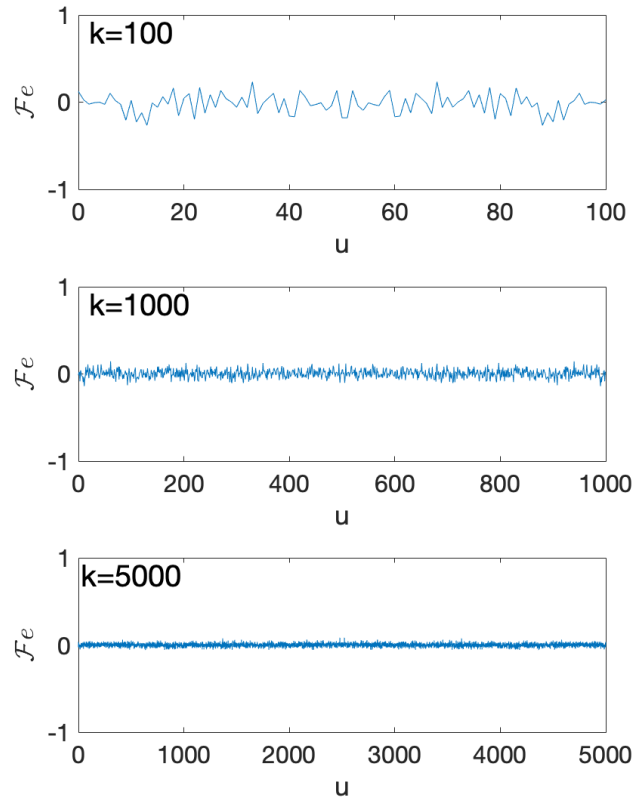


Figure 1:  $\mathcal{F}e$  in frequency domain with different  $k$

Code:

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1 hz=100;t=0:1/hz:1;uu=0:hz;a=2;
2 noi=randn(length(t),1)*a;
3 gu = fft(noi)/hz;

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Frequency  $k$  has impact on the transformed value. In theory,  $dt \rightarrow 0$ , which means  $k = 1/dt \rightarrow \infty$ . Then  $E(k) = \mathcal{F}e \rightarrow 0$

## Q6.6

As equation 6.3, fourier operator matrix

$$\mathbf{Q}_{4 \times 4} = \frac{1}{\sqrt{4}} \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & \exp(i2\pi/4) & \exp(i2\pi * 2/4) & \exp(i2\pi * 3/4) \\ 1 & \exp(i2\pi 2 * 1/4) & \exp(i2\pi 2 * 2/4) & \exp(i2\pi 2 * 3/4) \\ 1 & \exp(i2\pi 3 * 1/4) & \exp(i2\pi 3 * 2/4) & \exp(i2\pi 3 * 3/4) \end{pmatrix}$$

$$= \frac{1}{2} \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & \exp(i0.5\pi) & \exp(i\pi) & \exp(i1.5\pi) \\ 1 & \exp(i\pi) & \exp(i2\pi) & \exp(i3\pi) \\ 1 & \exp(i1.5\pi) & \exp(i3\pi) & \exp(i4.5\pi) \end{pmatrix}$$

$\mathbf{Q}_{4 \times 4}$  is symmetric. Use euler's equation, then

$$\mathbf{Q}_{4 \times 4} = \frac{1}{2} \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & \cos(0.5\pi) + i\sin(0.5\pi) & \cos(\pi) + i\sin(\pi) & \cos(1.5\pi) + i\sin(1.5\pi) \\ 1 & \cos(\pi) + i\sin(\pi) & \cos(2\pi) + i\sin(2\pi) & \cos(3\pi) + i\sin(3\pi) \\ 1 & \cos(1.5\pi) + i\sin(1.5\pi) & \cos(3\pi) + i\sin(3\pi) & \cos(4.5\pi) + i\sin(4.5\pi) \end{pmatrix}$$

$$= \frac{1}{2} \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & i & -1 & -i \\ 1 & -1 & 1 & -1 \\ 1 & -i & -1 & i \end{pmatrix}$$

$$\mathbf{Q}^\dagger \mathbf{Q} = \frac{1}{2} \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & -i & -1 & i \\ 1 & -1 & 1 & -1 \\ 1 & i & -1 & -i \end{pmatrix} * \frac{1}{2} \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & i & -1 & -i \\ 1 & -1 & 1 & -1 \\ 1 & -i & -1 & i \end{pmatrix}$$

$$= \frac{1}{4} \begin{pmatrix} 4 & 0 & 0 & 0 \\ 0 & 4 & 0 & 0 \\ 0 & 0 & 4 & 0 \\ 0 & 0 & 0 & 4 \end{pmatrix} = I_{4 \times 4} = \mathbf{Q} \mathbf{Q}^\dagger$$

So  $\mathbf{Q}_{4 \times 4}$  is unitary.



### Q6.7

According to equation 5.30 in §5.9:

$$\begin{aligned} F(u, v) &= \mathcal{F}f(x, y) = \int_{-\infty}^{\infty} dy \int_{-\infty}^{\infty} dx f(x, y) \exp(-i2\pi(ux + vy)) \\ &= \int_{-\infty}^{\infty} dx A \operatorname{rect} \frac{x - x_0}{X_0} \exp(-i2\pi ux) \int_{-\infty}^{\infty} dy B \operatorname{rect} \frac{y - y_0}{Y_0} \exp(-i2\pi vy) \end{aligned}$$

Let  $x' = x - x_0$  and use fourier shift theorem, then

$$\begin{aligned} &\int_{-\infty}^{\infty} dx A \operatorname{rect} \frac{x - x_0}{X_0} \exp(-i2\pi ux) \\ &= A \exp(-i2\pi ux_0) \int_{-\infty}^{\infty} dx' \operatorname{rect} \frac{x'}{X_0} \exp(-i2\pi ux') \\ &= A \exp(-i2\pi ux_0) \int_{-X_0/2}^{X_0/2} dx' \exp(-i2\pi ux') \\ &= A \exp(-i2\pi ux_0) \frac{(\exp(-i2\pi uX_0/2) - \exp(-i2\pi u(-X_0/2)))}{-i2\pi u} \\ &= AX_0 \exp(-i2\pi ux_0) \frac{\sin(\pi uX_0)}{\pi uX_0} \\ &= AX_0 \exp(-i2\pi ux_0) \operatorname{sinc}(uX_0) \quad \text{when } \operatorname{sinc}(t) = \sin(\pi t)/(\pi t) \end{aligned}$$

Do the same thing for the  $y$  part, then

$$F(u, v) = AX_0 \exp(-i2\pi ux_0) \operatorname{sinc}(uX_0) BY_0 \exp(-i2\pi vy_0) \operatorname{sinc}(vY_0)$$