

21st January 2025 (Tuesday) →

① Laplacian in spherical polar coordinate →

$$\left\{ r, \theta, \phi \right\}$$

$$r \in [0, \infty), \quad \theta \in [0, \pi], \quad \phi \in [0, 2\pi]$$

The Laplacian looks like →

$$\frac{1}{r} \frac{\partial^2}{\partial r^2} (r \psi) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 \psi}{\partial \phi^2} = 0$$

We may separate,

$$r \sin \theta \frac{\partial^2}{\partial r^2} (r \psi) + \sin \theta \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \psi}{\partial \theta} \right) + \underbrace{\frac{\partial^2 \psi}{\partial \phi^2}}_{-m^2} = 0$$

We replace,

$$r \sin \theta \frac{\partial^2}{\partial r^2} (r \psi) + \sin \theta \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \psi}{\partial \theta} \right) - m^2 = 0$$

$$\Rightarrow r \frac{\partial^2}{\partial r^2} (r \psi) + \underbrace{\frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \psi}{\partial \theta} \right)}_{f(\theta)} - \frac{m^2}{\sin \theta} = 0$$

$$\underbrace{f(r)}$$

$$\downarrow \alpha^2 = L(L+1)$$

$$\underbrace{-\alpha^2}_{f(\theta)}$$

Looking at the radial part with $\psi = R(r) P(\theta) \dots$

$$\frac{r}{R} \frac{\partial^2}{\partial r^2} (r R) = \alpha^2$$

$$\Rightarrow \frac{\partial^2}{\partial r^2} (r R) = \frac{\alpha^2}{r} R$$

Now, this is an eigenvalue problem as usual.

$$\therefore R = r^L$$

$$\therefore r \frac{\partial^2}{\partial r^2} (r^L) = (L+1)L r^L$$

$\therefore (L+1)L$ is the eigenvalue.

After quite a bit of fiddling, we find that the solutions are,

$$r^L \rightarrow L(L+1)$$

$$r^{-L-1} \rightarrow L(L+1)$$

$$\therefore \alpha^2 = L(L+1)$$

Now, we have,

$$\frac{1}{\sin \theta} \cdot \frac{1}{P} \frac{d}{d\theta} \left(\sin \theta \frac{dP}{d\theta} \right) + \left[L(L+1) - \frac{m^2}{\sin^2 \theta} \right] P = 0$$

We do a coordinate transform to make it recognizable.

$$x = \cos \theta \Rightarrow dx = -\sin \theta d\theta \Rightarrow \sin^2 \theta = 1 - x^2$$

$$\therefore \frac{d}{dx} \left[(1-x^2) \frac{dP}{dx} \right] + \left[L(L+1) - \frac{m^2}{(1-x^2)} \right] P = 0 \rightarrow \text{Associated Legendre diff'l eqn.}$$

Standard form of this diff'l eqn,

$$\frac{d^2y}{dx^2} + \frac{Q(x)}{P(x)} \frac{dy}{dx} + \frac{R(x)}{P(x)} = 0 \quad \sum_{n=0}^{\infty} a_n x^n$$

We may use power series to solve this.

→ $P(x)$, $Q(x)$ is a regular point

Singular point \rightarrow regular singular
 \rightarrow irregular singular

→ $P(x)$ has pole of order -1 (Simple pole)

→ $Q(x)$ has pole of order -2

Now, we have,

$$y = \sum_{n=0}^{\infty} b_n x^{n+s} \rightarrow \text{Indicial eqn}$$

□ Review power series soln. of Associated Legendre diff'l eqn.

Note that in our case, $x \in \{-1, 1\}$ are singularities.

Moreover, they are regular singular pt's.

② Something about plates? (Charged ones, of course)

If ϕ is symmetric, we have azimuthal/axial symmetry.

Setting $m=0$, we have,

$$\frac{d}{dx} \left[(1-x^2) \frac{dP_L}{dx} \right] + L(L+1) P_L(x) = 0$$

□ Just look up the solution and do!, the class is not really going on very coherently.

There is a one-one correspondence b/w Legendre polynomials and powers of x .

$$\begin{aligned} 1 &\rightarrow P_0 \\ x &\rightarrow P_1 \\ x^2 &\rightarrow P_2 \dots \end{aligned}$$

So,

Any L^2 function in the interval $[-1, 1]$ can be expressed as,

$$f(x) = \sum_n a_n P_n(x)$$

□ Look up orthogonality of $P_L(x)$.

22nd January 2025 (Wednesday) →

① Legendre's differential equation →

& shows up when dealing with Spherical Laplacian.

② Azimuthal symmetry: When potential does not depend on azimuth.

$$\frac{d}{dx} \left[(1-x^2) \frac{dP_L}{dx} \right] + L(L+1) P_L = 0$$

Whichever solution, $P_L(x)$ is a polynomial in x

Note: $x \in [-1, 1]$ here, as it stands for θ in spherical coord.

The differential equation looks like an eigenvalue eqn, thus suggesting that $P_L(x)$ are orthogonal eigenfunctions — allows us to represent any L^2 function on $[-1, 1]$.

$$\Rightarrow \int_{-1}^{+1} P_L(x) P_m(x) dx \propto \delta_{lm}$$

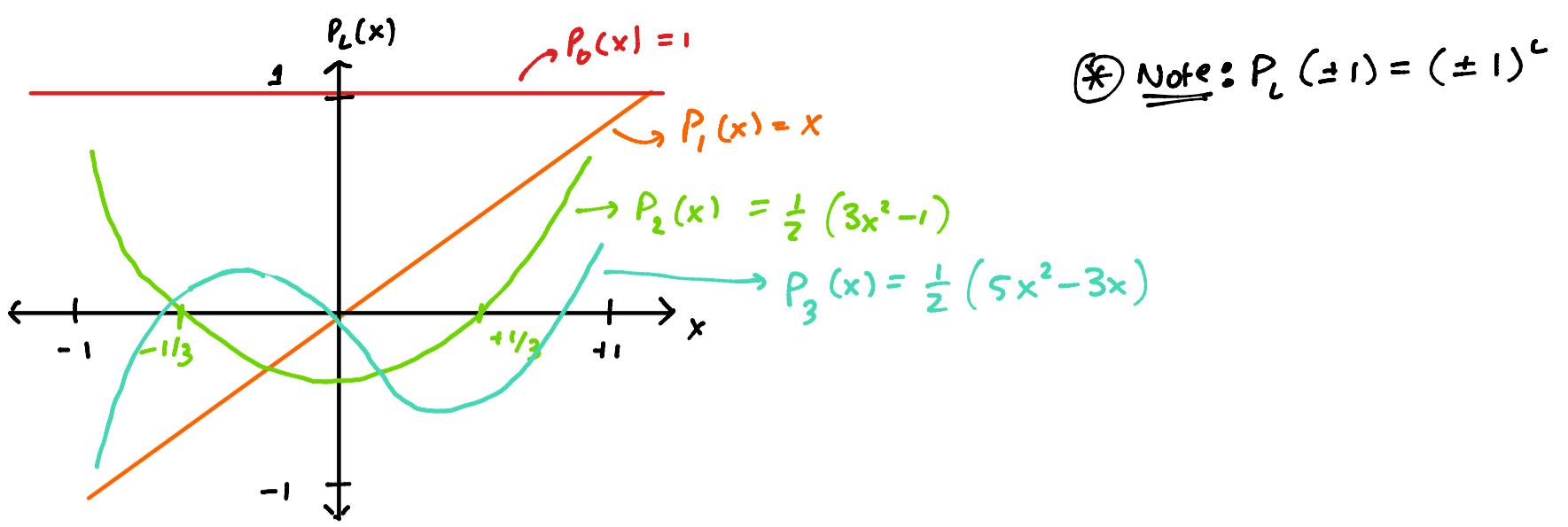
Also,

$$\int_{-1}^1 [P_L(x)]^2 dx = N_L = \frac{2L+1}{2}$$

③ Rodrigues' formulae →

$$P_L(x) = \frac{1}{2^L L!} \frac{d^L}{dx^L} [(x^2 - 1)^L]$$

We look at plots of $P_L(x)$



* Note: $P_L(\pm 1) = (\pm 1)^L$

Sort of like $\sin/\cos x$

If you switch off your rational brain, it is like
 \sin and \cos ?

— RKN

① Fourier Legendre Series →

If a function $f(x)$ is L^2 in interval $[-1, 1]$. Then it can be expressed

as,

$$f(x) = \sum_{l=0}^{\infty} A_l P_l(x) \quad \text{Probably in } \frac{2}{2l+1}$$

$$\text{Where, } A_l = \frac{2l+1}{2} \int_{-1}^1 f(x) P_l(x) dx$$

□ PL check all equations.

Now,

$$f(x) = \begin{cases} +1, & x > 0 \\ -1, & x < 0 \end{cases}$$

$$\begin{aligned} \frac{2}{2l+1} A_l &= \int_{-1}^{+1} f(x) P_l(x) dx \\ &= - \int_{-1}^0 P_l(x) dx + \int_0^1 P_l(x) dx \end{aligned}$$

□ Can be integrated using Rodrigues formula.

$$\rightarrow A_l = \left(\frac{-1}{2} \right)^{\frac{l-1}{2}} \frac{(2l+1)(l-2)!!}{2 \left(\frac{l+1}{2} \right)!}$$

$$A_1 = \frac{3}{2}, \quad A_2 = \frac{7}{8}, \quad A_3 = \frac{11}{16}, \dots$$

□ HW: Decompose step function P_c in interval $[-1, 1]$

- Establish relation between x^l and $P_l(x)$ for $l=1, 3, 5, \dots$
- Write step function as $\sum a_n x^n$ accurate upto 5th order.

Actual assignment - due next Monday - (10:00 A.M.)

- Also, plot with pen and paper.

If there is azimuthal symmetry, we just ignore z axis - we take projection on some x-y plane.

BVP with azimuthal symmetry →

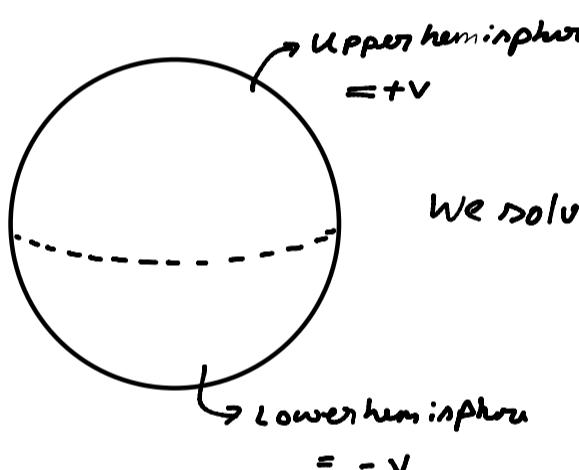
$$\nabla^2 \phi = 0$$

Something about "Object oriented"? Nevermind.

$$\phi(r, \theta) = \sum_{l=0}^{\infty} [A_l r^l + B_l r^{-(l+1)}] P_l(\cos\theta)$$

Again something about $r \rightarrow 0, r \rightarrow \infty$.

Problem →



We solve Laplace inside this sphere.

$$\nabla^2 \phi = 0$$

$$\Rightarrow \phi(r, \theta) = \sum_l A_l \left(\frac{r}{a}\right)^l P_l(\cos\theta)$$

We know, as $r \rightarrow a$,

$$\phi(r=a, \theta) = \begin{cases} +v, & 0 \leq \theta < \frac{\pi}{2} \\ -v, & \pi > \theta > \frac{\pi}{2} \end{cases}$$

$$\Rightarrow \phi(r=a, \theta) = \sum_l A_l P_l(\cos\theta)$$

$$A_1 = \frac{3}{2}, A_2 = \frac{7}{8}, A_3 = \frac{11}{16}, \dots$$

(from Legendre Series of Step function)

Associated Legendre →

$$\frac{d}{dx} \left[(1-x^2) \frac{dP_l^m}{dx} \right] + \left[(l+1)l - \frac{m^2}{(1-x^2)} \right] P_l^m(x) = 0$$

$$l = 0, 1, \dots \quad m = -l, \dots, +l$$

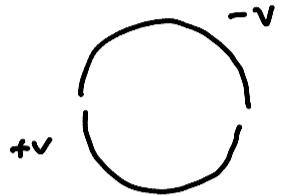
Very incoherent scribbling on polar plots of associated Legendre polynomials.

23rd January 2025 (Thursday) →

④ Slices: Representing functions on circles and spheres →

Plotting $\sin(x)$ vs. x is simple.

But in problems like →



→ Plotting potential on this boundary is problematic, moreover because it is on 2D polar.