

Advanced Statistical Physics Homework

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I. Please calculate the average ensemble of a free particle placed in a cube with sides of length L .

Answer I. For this free particle, Hamiltonian is formed in $\hat{H} = -\frac{\hbar^2}{2m}\nabla^2$, and its eigenfunctions is defined by $\phi_{\mathbf{k}}(x, y, z) = \langle \mathbf{r} | \mathbf{k} \rangle = \frac{e^{i\mathbf{k}\cdot\mathbf{r}}}{L^{3/2}}$. Besides that $\phi_{\mathbf{k}}(x, y, z)$ shall satisfy periodic boundary condition, we can get eigenenergy by solving Schrödinger equation :

$$\hat{H}\phi_{\mathbf{k}} = E_{\mathbf{k}}\phi_{\mathbf{k}} \longrightarrow E_{\mathbf{k}} = \frac{\hbar^2 k^2}{2m}, \quad \mathbf{k} = (k_x + k_y + k_z) = \frac{2\pi}{L}(n_x + n_y + n_z) \quad (0.1)$$

which

$$\begin{aligned} E_{\mathbf{k}} &= \frac{\hbar^2 k^2}{2m} \\ \mathbf{k} &\equiv (k_x, k_y, k_z) = \frac{2\pi}{L}(n_x, n_y, n_z) \\ (n_x, n_y, n_z) &= 0, \pm 1, \pm 2, \dots \end{aligned}$$

Free particle partition function :

$$Z = \text{Tre}^{-\beta\hat{H}} = \sum_{\mathbf{k}} \int e^{-\beta E_{\mathbf{k}}} \phi_{\mathbf{k}}^*(\mathbf{r}) \phi_{\mathbf{k}}(\mathbf{r}) d\mathbf{r} = V \left(\frac{m}{2\pi\hbar^2\beta} \right)^{\frac{3}{2}} \quad (0.2)$$

Operational representation of canonical ensemble :

$$\hat{\rho} = \frac{1}{Z} e^{-\beta \hat{H}} = \sum_{\mathbf{k}} |\mathbf{k}\rangle \frac{1}{Z} e^{-\beta \hat{H}} \langle \mathbf{k}| = \sum_{\mathbf{k}} |\mathbf{k}\rangle \frac{1}{Z} e^{-\beta E_{\mathbf{k}}} \langle \mathbf{k}| \quad (0.3)$$

Under coordinate representation, we can describe matrix elements by :

$$\begin{aligned} \langle \mathbf{r} | \hat{\rho} | \mathbf{r}' \rangle &= \sum_{\mathbf{k}} \langle \mathbf{r} | \mathbf{k} \rangle \frac{1}{Z} e^{-\beta \hat{H}} \langle \mathbf{k} | \mathbf{r}' \rangle = \sum_{\mathbf{k}} \phi_{\mathbf{k}}^*(\mathbf{r}) \frac{1}{Z} e^{-\beta \hat{H}} \phi_{\mathbf{k}}(\mathbf{r}') \\ &= \frac{1}{V} \exp \left(-\frac{m(\mathbf{r} - \mathbf{r}')^2}{2\hbar^2 \beta} \right) \end{aligned} \quad (0.4)$$

The average ensemble of Hamiltonian is :

$$\begin{aligned} \langle \hat{H} \rangle &= Tr(\hat{\rho} \hat{H}) = -\frac{\partial}{\partial \beta} \ln [Tr(e^{-\beta \hat{H}})] = -\frac{\partial}{\partial \beta} \ln Z \\ &= \frac{3}{2} k_B T \end{aligned} \quad (0.5)$$

2. Please calculate the relationship between average population number c with temperature and specific heat capacity of Bose-Einstein condensation.

Answer 2. *The equation of state for ideal Boson gas :*

$$\begin{cases} \frac{P}{k_B T} = \frac{1}{\lambda^3} g_{5/2}(z) - \frac{1}{V} \ln(1 - z) \\ \frac{1}{v} = \frac{1}{\lambda^3} g_{3/2}(z) + \frac{1}{V} \frac{z}{1 - z} \end{cases} \quad (0.6)$$

For specific capacity is $v = \frac{V}{N}$, mean thermal wavelength is $\lambda = \sqrt{\frac{2\pi\hbar^2}{mk_B T}}$, and fugacity $z = e^{\beta\mu}$, which μ

is chemical potential. For Boson gas we have : $0 \leq z \leq 1$, it's obvious that $z \geq 0$, and we can confirm it by average population number $\langle n_0 \rangle = \frac{z}{1-z} \geq 0$. $g_n(z)$ is generated by :

$$g_n(z) = \sum_{l=1}^{\infty} \frac{z^l}{l^n} \quad (0.7)$$

when z valued from 0 to 1, $g_n(z)$ is positive monotone increasing and it's a limited function. If $n > 1$, $g_n(z)$ would turn into Riemann- ζ function :

$$g_n(1) = \sum_{l=1}^{\infty} \frac{1}{l^n} = \zeta(n) \quad (n > 1) \quad (0.8)$$

And $g_n(z)$ diverged when $n \leq 1$, so we have :

$$g_{3/2}(z) \leq g_{3/2}(1) = \zeta(3/2) = 2.612... \quad (0.9)$$

Now we'll continue to discuss conditions that produce condensation. First, we can rewrite the second formula of Equ.0.6 :

$$\lambda^3 \frac{\langle n_0 \rangle}{V} = \frac{\lambda^3}{v} - g_{3/2}(z) \quad (0.10)$$

When $\frac{\langle n_0 \rangle}{V} > 0$, the condensation achieved, and it must be true when $\frac{\lambda^3}{v} > g_{3/2}(1)$.

Discuss the critical condition :

$$\frac{\lambda_c^3}{v} = g_{3/2}(z) \quad (0.11)$$

then we will get critical temperature and critically specific heat capacity :

$$T_c = \frac{2\pi\hbar^2}{mk_B[v g_{3/2}(1)]^{2/3}} \quad v_c = \frac{\lambda^3}{g_{3/2}(1)} \quad (0.12)$$

When $T < T_c$ (v is constant) or $v < v_c$ (T is constant), Bose-Einstein condensation occurred. That is to say creation conditions of Bose-Einstein condensation are low temperature and high density.

Besides, when $V \rightarrow \infty$, fugacity z tends to

$$z = \begin{cases} 1 & \left(\frac{\lambda^3}{v} \geq g_{3/2}(1) \right) \\ \text{The roots of } g_{3/2}(z) = \frac{\lambda^3}{v} & \left(\frac{\lambda^3}{v} \leq g_{3/2}(1) \right) \end{cases} \quad (\text{O.I3})$$

And from Equ.0.10 , we have

$$\frac{\langle n_0 \rangle}{N} = \begin{cases} 1 - \left(\frac{T}{T_c} \right)^{\frac{3}{2}} = 1 - \frac{v}{v_c} & \left(\frac{\lambda^3}{v} \geq g_{3/2}(1) \right) \\ 0 & \left(\frac{\lambda^3}{v} \leq g_{3/2}(1) \right) \end{cases} \quad (\text{O.I4})$$