



INTRODUCTION TO UNIVERSALITY

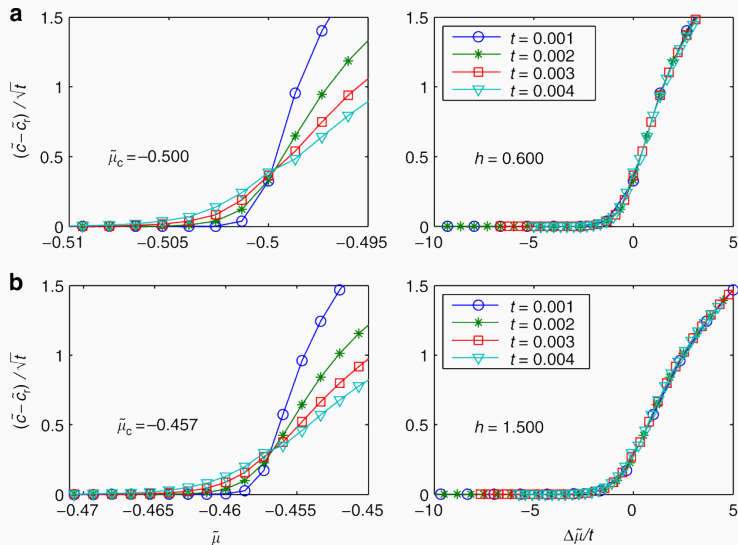
Phase Transitions in Random Graphs

Alberto Pezzotta

July 3, 2015

UNIVERSALITY IN PHASE TRANSITIONS

AN EXAMPLE FROM COLD ATOMS PHYSICS



(Chen et al., “Critical behaviours of contact near phase transitions”, 2014)

RANDOM GRAPHS

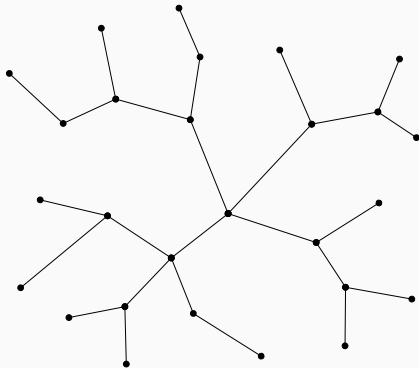
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- **Model real networks:**
 - possibility of tailoring the random graph;
 - sacrifice some features (e.g., clustering coefficient);
 - interpolating between random graphs and networks (e.g., WS and NW small world)

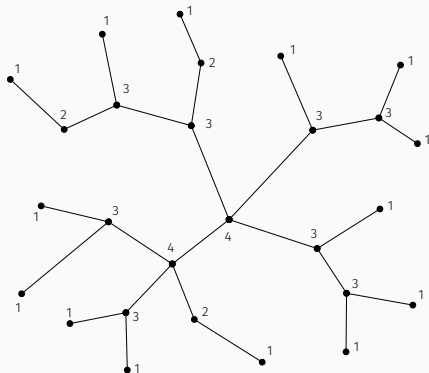
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- **Model real networks:**
 - possibility of tailoring the random graph;
 - sacrifice some features (e.g., clustering coefficient);
 - interpolating between random graphs and networks (e.g., WS and NW small world)
- Make connection with **statistical mechanics**:
 - universal properties?
 - phase transitions \rightsquigarrow giant component



Newman et al., 2001

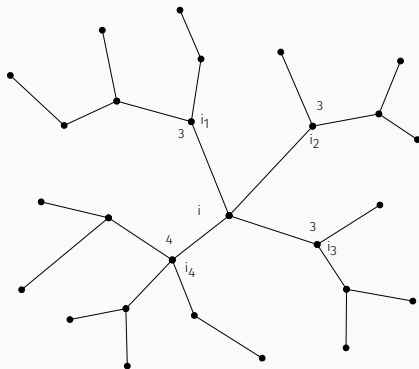
DEGREE DISTRIBUTION AND ...



Newman et al., 2001

p_k = degree dist.

$g_0(z)$ = generating function of $p_k = \sum_k p_k z^k$



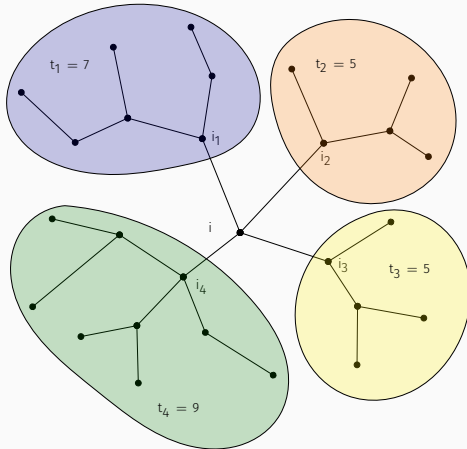
Newman et al., 2001

$$p_k \rightsquigarrow g_0(z)$$

$$q_k = \text{excess degree dist} = (k+1)p_{k+1}/\langle k \rangle_p$$

$$g_1(z) = \text{generating function of } q_k = \sum_k q_k z^k$$

DEGREE DISTRIBUTION AND ...



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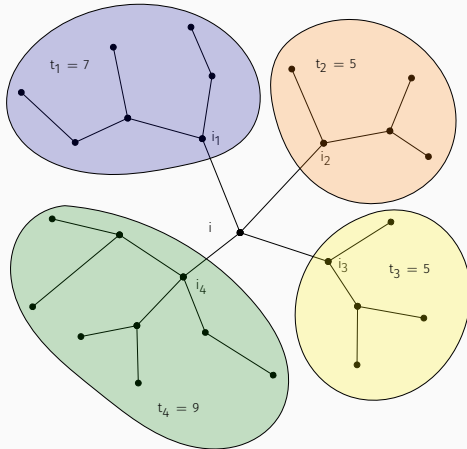
$$p_k \rightsquigarrow g_0(z)$$

$$q_k \rightsquigarrow g_1(z) = \frac{g'_0(z)}{g'_0(1)}$$

ρ_t = branch size dist

$h_1(z)$ = generating function of ρ_t

DEGREE DISTRIBUTION AND ...



Newman et al., 2001

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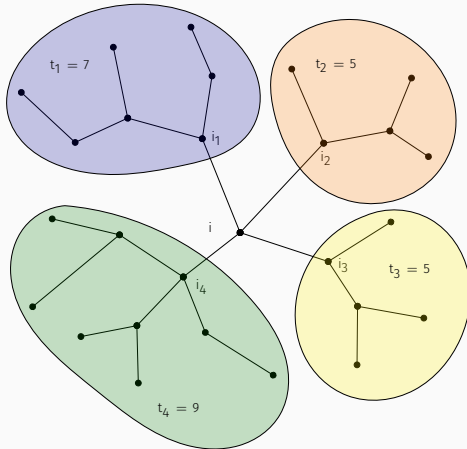
$$q_k \rightsquigarrow g_1(z) = \frac{g'_0(z)}{g'_0(1)}$$

$$\rho_t \rightsquigarrow h_1(z)$$

$$\pi_s = \text{component size dist} = \sum_k P(s|k) p_k$$

$$P(s|k) = \sum_{\{t_k\}} \delta(s-1, \sum_{m=1}^k t_m) \prod_{m=1}^k \rho_{t_m}$$

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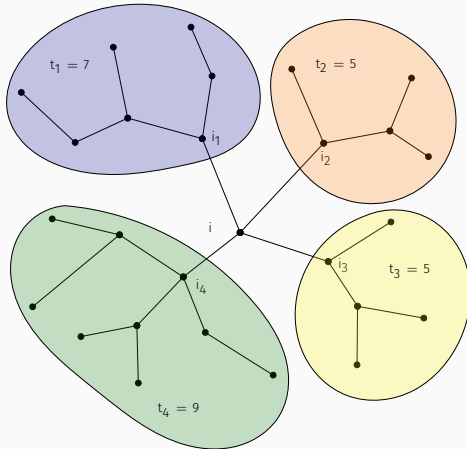
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$$\rho_t \rightsquigarrow h_1(z) = z g_1(h_1(z))$$

$$\pi_s \rightsquigarrow h_0(z) = z g_0(h_1(z))$$

TREE assumption

$$P(s|k) = \sum_{\{t_k\}} \delta(s-1, \sum_{m=1}^k t_m) \prod_{m=1}^k \rho_{t_m}$$

p_k normalized to 1: $g_0(1) = \sum_k p_k = 1$

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NOT the same for π_s , because

TREE \equiv SMALL components

$$h_0(1) = 1 - S$$

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$$\begin{cases} h_1(z) = z g_1(h_1(z)) \\ h_0(z) = z g_0(h_1(z)) \end{cases} \Rightarrow \begin{cases} h_1(1) = u = g_1(u) \\ h_0(1) = g_0(u) = 1 - S \end{cases}$$

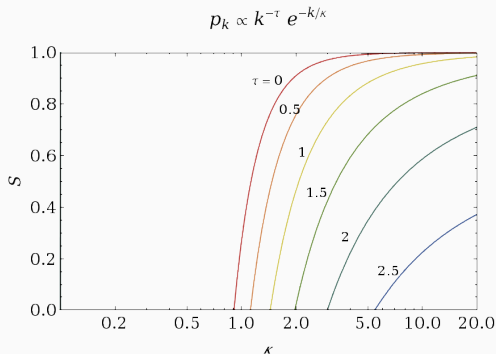
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Phase transition

Critical point at κ such that

$$g'_1(1) = 1 = \langle k \rangle_q$$



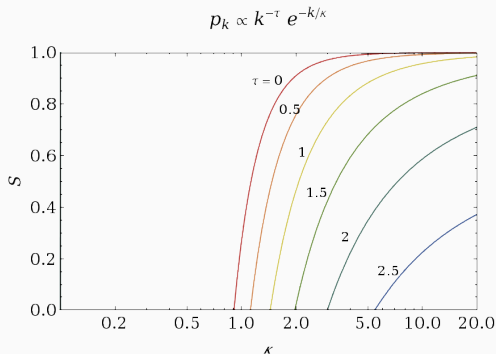
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Remarks:

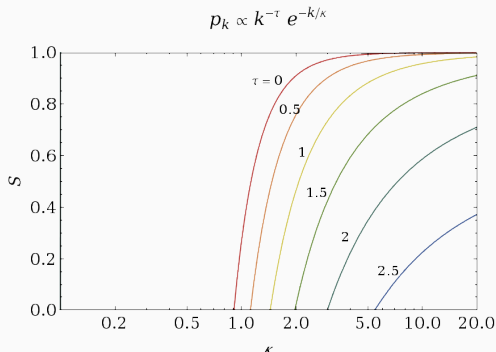


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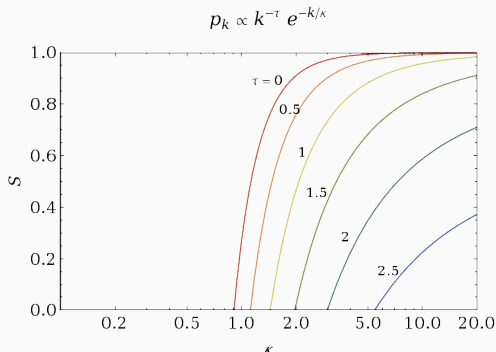
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Remarks:

- $\kappa \rightarrow \infty \equiv$ power law \rightsquigarrow always giant component
- linear increase at criticality \rightsquigarrow connection with percolation (??)

Erdős–Rényi random graph $G_{n,p}$ (Erdős and Rényi, 1960):

- n vertices $\rightsquigarrow n(n-1)/2$ possible edges (links);
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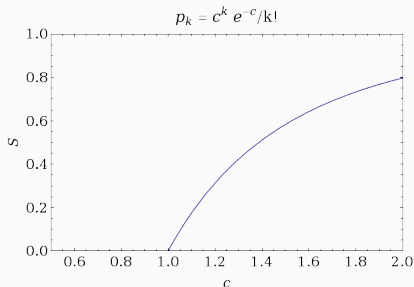
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A WELL UNDERSTOOD EXAMPLE

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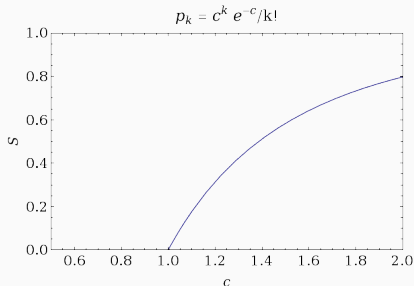


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Known properties:

- **equivalency** to Potts model (at the level of **Large Deviations**, Engel et al., 2004);
- ph. transition in the universality class of **percolation**, ($q \rightarrow 1$)-Potts

$\varphi(G)$ = # conn. comps in G/N ; c = average degree

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Large deviations: rate and cumulant generating functions

$$r(\varphi; c) = \sup_{q \geq 0} \{\varphi \ln q - y(q; c)\} \quad \leftrightarrow \quad y(q; c) = \sup_{\varphi \in [0,1]} \{\varphi \ln q - r(\varphi; c)\}$$

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Then **equivalence** is established:

$$Z_N(q; T = 1/c) \simeq e^{Nc/2} Y_N(q; c) \quad \text{or} \quad f(q; 1/c) \simeq -\frac{1}{2} - \frac{1}{c} y(q; c)$$

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Equilibrium solution encoding **symmetry breaking**:

$$x(0) = [1 + (q-1)s]/q; \quad x(\sigma > 0) = (1-s)/q$$

Solutions in the $q \rightarrow 1$ limit:

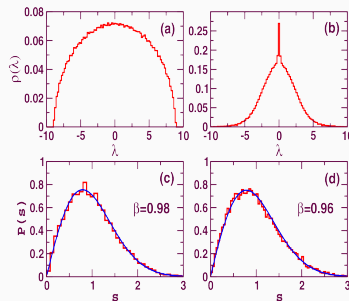
$$c \leq 1 : s^* = 0;$$

$$c > 1 : s^* = 1 - e^{-cs^*} \rightsquigarrow \text{Erdős-Renyi equation}$$

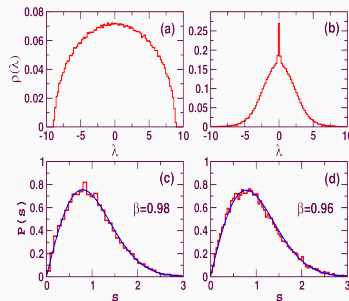
OPEN PROBLEMS

- Connection with **Random Matrix Theory**

- Connection with **Random Matrix Theory**, spectra of
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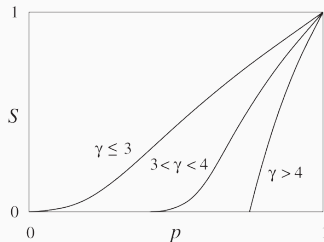
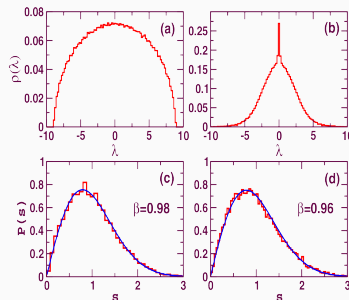


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- Universality of phase transitions in generic degree distribution random graphs (Dorogovtsev et al., 2008; Cohen et al., 2002).



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