

1. a) b)

$$(2) ds^2 = dr^2 + r^2 d\theta^2 \Rightarrow g_{\mu\nu} = \begin{pmatrix} 1 & 0 \\ 0 & r^2 \end{pmatrix} \text{ and } g^{\mu\nu} = (g_{\mu\nu})^{-1} = \begin{pmatrix} 1 & 0 \\ 0 & \frac{1}{r^2} \end{pmatrix} \quad \mu, \nu = 1, 2 (r, \theta)$$

$$\text{then } g_{rr} = 1, g_{r\theta} = 0, g_{\theta\theta} = r^2, g^{rr} = 1, g^{r\theta} = 0, g^{\theta\theta} = \frac{1}{r^2}$$

(3) since the space is locally flat, so we have  $g_{\mu\nu;\lambda} = 0$ ,  $g_{\mu\lambda;\mu} = 0$  and  $g_{\lambda\mu;\nu} = 0$

$$\Rightarrow \begin{cases} g_{\mu\nu,\lambda} - \Gamma_{\mu\lambda}^\alpha g_{\alpha\nu} - \Gamma_{\nu\lambda}^\alpha g_{\mu\alpha} = 0 & (1) \\ g_{\mu\lambda,\mu} - \Gamma_{\mu\mu}^\alpha g_{\alpha\lambda} - \Gamma_{\lambda\mu}^\alpha g_{\mu\alpha} = 0 & (2) \\ g_{\lambda\mu,\nu} - \Gamma_{\lambda\nu}^\alpha g_{\alpha\mu} - \Gamma_{\mu\nu}^\alpha g_{\lambda\alpha} = 0 & (3) \end{cases}$$

$$\frac{1}{2}((2)+(3)-(1)) \Rightarrow \Gamma_{\lambda\mu\nu} = g_{\lambda\alpha} \Gamma_{\mu\nu}^\alpha = \frac{1}{2}(g_{\mu\lambda,\nu} + g_{\nu\lambda,\mu} - g_{\mu\nu,\lambda})$$

$$\Rightarrow \Gamma_{\mu\nu}^\alpha = \frac{1}{2} g^{\alpha\lambda} (g_{\mu\lambda,\nu} + g_{\nu\lambda,\mu} - g_{\mu\nu,\lambda})$$

$$\text{so } \Gamma_{\theta\theta}^r = \frac{1}{2} g^{rr} (g_{r\theta,\theta} + g_{\theta r,\theta} - g_{\theta\theta,r}) = \frac{1}{2} \cdot 1 \cdot (-\frac{\partial r^2}{\partial r}) = -r$$

$$\Gamma_{r\theta}^\theta = \frac{1}{2} g^{\theta\theta} (g_{r\theta,\theta} + g_{\theta\theta,r} - g_{\theta r,\theta}) = \frac{1}{2} \frac{1}{r^2} \cdot 2r = \frac{1}{r}$$

$$\Gamma_{\theta r}^\theta = \frac{1}{2} g^{\theta\theta} (g_{\theta\theta,r} + g_{r\theta,\theta} - g_{\theta r,\theta}) = \frac{1}{2} \frac{1}{r^2} \cdot 2r = \frac{1}{r}$$

2. (1) based on the conditions given,  $\Gamma_{\alpha\beta}^\mu = \frac{1}{2} \eta^{\mu\rho} (h_{\rho\alpha,\beta} + h_{\rho\beta,\alpha} - h_{\rho\alpha,\beta})$  (1)

then rewrite the geodesic equation into 2 equations:

$$\begin{cases} \frac{d^2 x^0}{d\tau^2} = 0 & (2) \\ \frac{d^2 x^i}{d\tau^2} + \Gamma_{00}^i \left(\frac{dx^0}{d\tau}\right)^2 = 0 & (3) \end{cases}$$

$$(2) \Rightarrow x^0 = \alpha\tau + \text{const.}, \text{ where } \alpha \text{ is a constant} \quad (4)$$

$$(3) \Rightarrow \frac{d^2 x^i}{d\tau^2} + \Gamma_{00}^i = 0 \quad (5)$$

$$\text{also because of (1), } \Gamma_{00}^i = \frac{1}{2} \eta^{ii} (h_{i0,0} + h_{i0,0} - h_{00,i}) = \frac{1}{2} \cdot 1 \cdot (-h_{00,i}) = -\frac{1}{2} h_{00,i}$$

$$\text{compare (5) to the Newton equation: } \frac{d^2 x^i}{dt^2} + \phi_{,i} = 0$$

$$\Rightarrow \Gamma_{00}^i = \phi_{,i}$$

$$\Rightarrow -\frac{1}{2} \frac{\partial h_{00}}{\partial x^i} = \frac{\partial \phi}{\partial x^i}$$

$$\Rightarrow h_{00} = -2\phi$$

$$(2) g_{00} = \eta_{00} + h_{00} = -(1 + 2\phi)$$

$$d\tau = \sqrt{-ds^2} = \sqrt{-g_{00} dx^0 dx^0} = \sqrt{-g_{00}} dt = \sqrt{1+2\phi} dt \Rightarrow dt = \frac{d\tau}{\sqrt{1+2\phi}}$$

$$\text{use the Taylor expansion: } (1+2\phi)^{-\frac{1}{2}} = 1 - \phi + \frac{3}{2}\phi^2 + \dots \approx 1 - \phi$$

$$\text{so } dt \approx (1 - \phi) d\tau$$

$$3. (1) ds^2 = -dt^2 + dr^2 + r^2 d\theta^2 \Rightarrow g_{\mu\nu} = \begin{pmatrix} -1 & & \\ & 1 & \\ & & r^2 \end{pmatrix}, g^{\mu\nu} = \begin{pmatrix} -1 & & \\ & 1 & \\ & & \frac{1}{r^2} \end{pmatrix} \mu, \nu = 0, 1, 2$$

rewrite the geodesic equation:

$$\begin{cases} \frac{d^2 x^0}{d\tau^2} = 0 \\ \frac{d^2 x^1}{d\tau^2} + \Gamma_{22}^1 \left( \frac{dx^2}{d\tau} \right)^2 = 0 \\ \frac{d^2 x^2}{d\tau^2} + \Gamma_{21}^2 \frac{dx^1}{d\tau} \frac{dx^2}{d\tau} + \Gamma_{12}^2 \frac{dx^1}{d\tau} \frac{dx^2}{d\tau} = 0 \end{cases} \Rightarrow \begin{cases} \frac{d^2 t}{d\tau^2} = 0 \quad (1) \\ \frac{d^2 r}{d\tau^2} - r \left( \frac{d\theta}{d\tau} \right)^2 = 0 \quad (2) \\ \frac{d^2 \theta}{d\tau^2} + \frac{2}{r} \frac{d\theta}{d\tau} \frac{dr}{d\tau} = 0 \quad (3) \end{cases}$$

$$(2) \quad (1) \Rightarrow t = \alpha\tau + \text{const.}, \text{ where } \alpha \text{ is a constant} \quad (4)$$

$$\text{use } (4), (2) (3) \Rightarrow \begin{cases} \frac{d^2 r}{dt^2} - r \left( \frac{d\theta}{dt} \right)^2 = 0 \quad (5) \\ \frac{d^2 \theta}{dt^2} + \frac{2}{r} \frac{d\theta}{dt} \frac{dr}{dt} = 0 \quad (6) \end{cases}$$

Lagrangian of a prime in polar coordinates is:  $L = \frac{m}{2}(\dot{r}^2 + r^2 \dot{\theta}^2)$

$$\text{Euler-Lagrangian equation: } \begin{cases} \frac{d}{dt} \frac{\partial L}{\partial \dot{r}} = \frac{\partial L}{\partial r} \\ \frac{d}{dt} \frac{\partial L}{\partial \dot{\theta}} = \frac{\partial L}{\partial \theta} \end{cases} \Rightarrow \begin{cases} \frac{d}{dt} \dot{r} = r \dot{\theta}^2 \\ \frac{d}{dt} (r^2 \dot{\theta}) = 0 \end{cases}$$

$$\Rightarrow \begin{cases} \frac{d^2 r}{dt^2} - r \left( \frac{d\theta}{dt} \right)^2 = 0 \\ r^2 \frac{d^2 \theta}{dt^2} + 2r \frac{d\theta}{dt} \frac{dr}{dt} = 0 \end{cases}$$

$$\Rightarrow \begin{cases} \frac{d^2 r}{dt^2} - r \left( \frac{d\theta}{dt} \right)^2 = 0, \text{ i.e. } (5) \\ \frac{d^2 \theta}{dt^2} + \frac{2}{r} \frac{d\theta}{dt} \frac{dr}{dt} = 0, \text{ i.e. } (6) \end{cases} \text{ which means they are equivalent.}$$