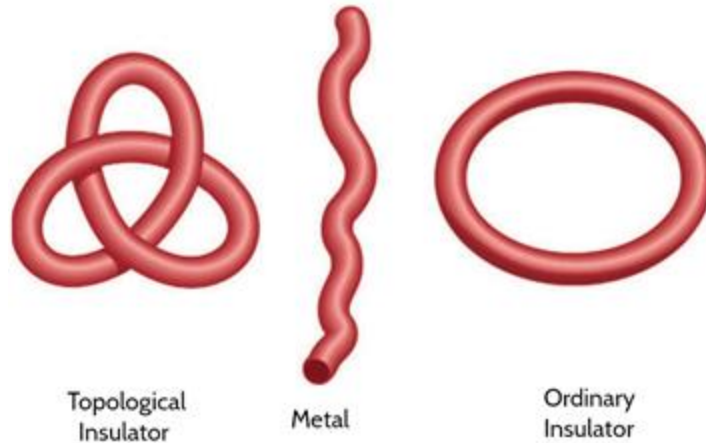


Topological Insulators and Topological Phase Transitions

Chaitali Shah

What are topological insulators?

- An insulator that always has a metallic boundary when placed next to an ordinary insulator
- Origin of metallic boundaries - topological invariants that cannot change as long as material is insulating



The Role of Symmetry

- Symmetry in **momentum** - $H(k)$
- Chiral Symmetry (S)

$$S H(k) S^{-1} = -H(k)$$

- Particle-Hole Symmetry (Charge Conjugation) Symmetry (C)

$$C H(k) C^{-1} = -H(k)$$

- Inversion (Time-Reversal) Symmetry (T)

$$T H(k) T^{-1} = H(k)$$

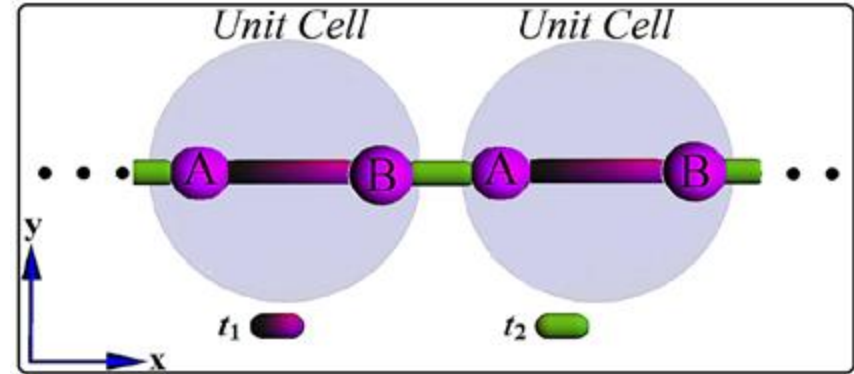
- 10 discrete symmetry classes

Symmetry Class	Time reversal symmetry	Particle hole symmetry	Chiral symmetry
A	No	No	No
AIII	No	No	Yes
AI	Yes, $T^2 = 1$	No	No
BDI	Yes, $T^2 = 1$	Yes, $C^2 = 1$	Yes
D	No	Yes, $C^2 = 1$	No
DIII	Yes, $T^2 = -1$	Yes, $C^2 = 1$	Yes
AII	Yes, $T^2 = -1$	No	No
CII	Yes, $T^2 = -1$	Yes, $C^2 = -1$	Yes
C	No	Yes, $C^2 = -1$	No
CI	Yes, $T^2 = 1$	Yes, $C^2 = -1$	Yes

Atomic Picture

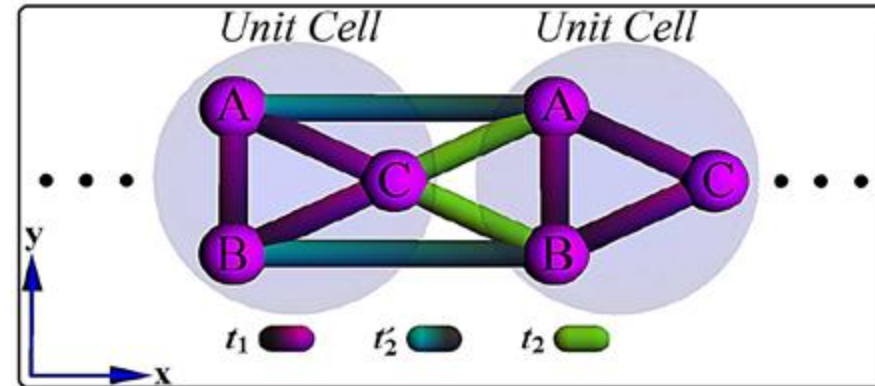
SSH Model (1D)

$$H = \sum_{n=1}^N t_1 (A_n^\dagger B_n) + \sum_{n=1}^N t_2 (B_n^\dagger A_{n+1})$$



JH Model (quasi 1D)

$$H = \sum_{n=1}^N t_1 (A_n^\dagger B_n + B_n^\dagger C_n + A_n^\dagger C_n) + \sum_{n=1}^N t_2' (A_n^\dagger A_{n+1} + B_n^\dagger B_{n+1}) + \sum_{n=1}^N t_2 (C_n^\dagger A_{n+1} + C_n^\dagger B_{n+1})$$



Hamiltonian $H(k)$

SSH Model (1D)

$$\mathcal{H}(k) = \begin{pmatrix} 0 & t_1 + t_2 e^{ik} \\ t_1 + t_2 e^{-ik} & 0 \end{pmatrix}$$

$$E^\pm = \pm \sqrt{t_1^2 + t_2^2 + 2t_1 t_2 \cos(k)}$$

Also, note that

$$\mathcal{H}(k) = (t_1 + t_2 \cos(k))\sigma_x + (t_2 \sin(k))\sigma_y$$

where σ_x and σ_y are Pauli matrices

$$\mathcal{H}(k) = d_x \sigma_x + d_y \sigma_y$$

JH Model (quasi 1D)

$$\mathcal{H}(k) = \begin{pmatrix} 2t'_2 \cos(k) & t_1 & t_1 + t_2 e^{ik} \\ t_1 & 2t'_2 \cos(k) & t_1 + t_2 e^{ik} \\ t_1 + t_2 e^{-ik} & t_1 + t_2 e^{-ik} & 0 \end{pmatrix}$$

$$E^0 = -2t_1 + \eta$$

$$E^\pm = \frac{1}{2}(\eta \pm \sqrt{\eta^2 + 2(t_1^2 + 4t_2'^2 + 8t_1 t_2 \cos(k))})$$

$$\text{where } \eta = t_1 + 2t'_2 \cos(k)$$

Band closing and reopening at $k = (0, \pi)$

SSH Model

$$E^+ = E^- \text{ when}$$

$$t_1 = t_2$$

Topological Phase Transition is possible when coupling is same.

JH Model

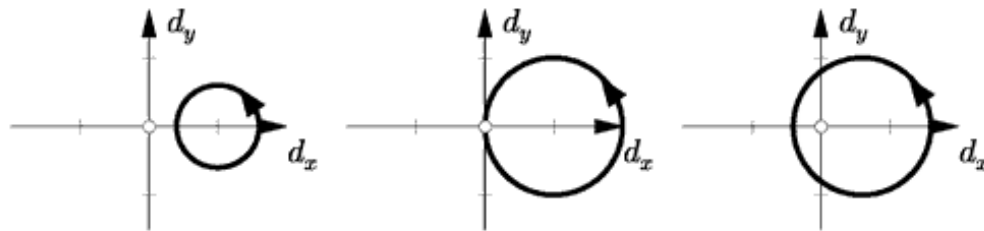
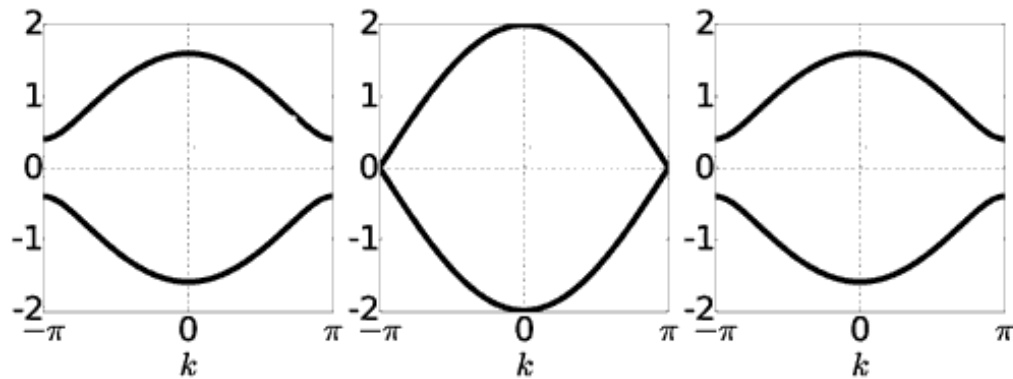
$$E^+ = E^- \text{ when}$$

$$t_1 = e^{ik} \frac{2}{9} \left(t_2' + 4t_2 \pm \frac{\sqrt{-(-2t_2' + t_2)^2}}{2} \right)$$

$$t_2' = \frac{t_2}{2} \text{ and } t_1 = e^{ik} t_2$$



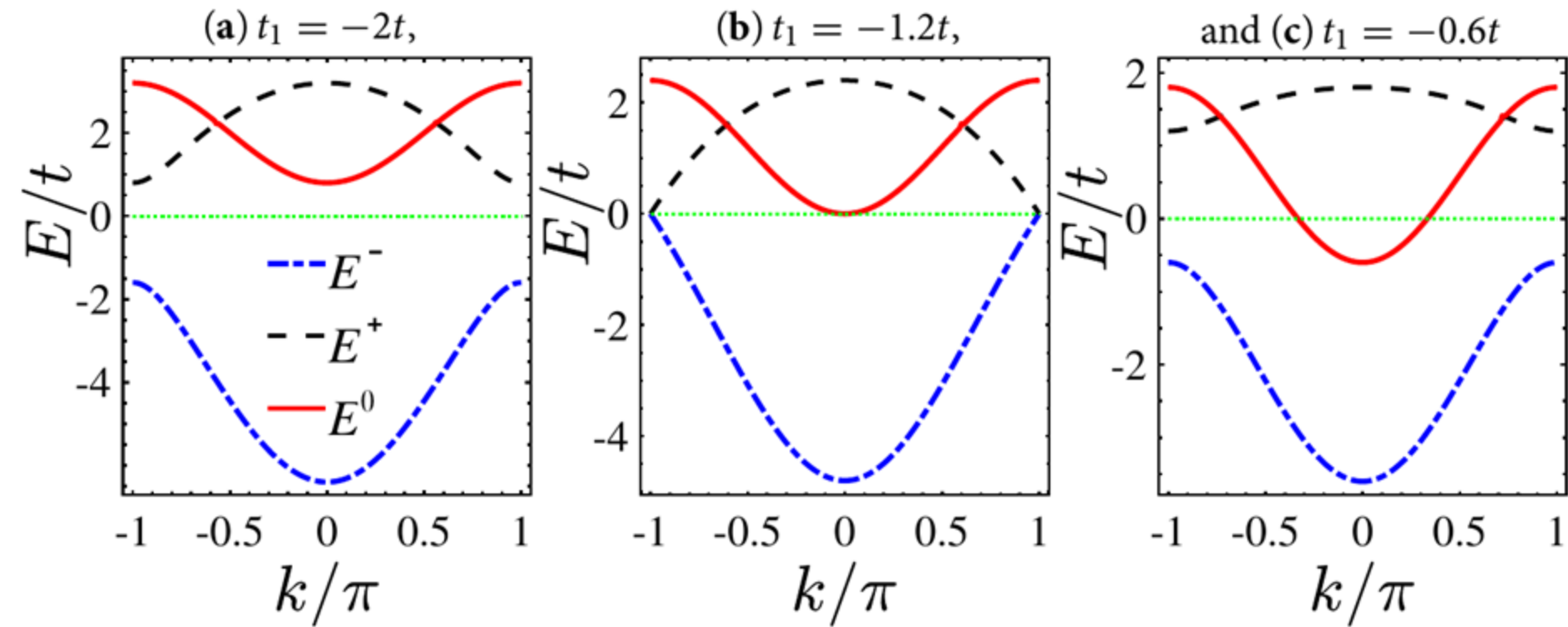
Transition in the SSH Model



Winding number

- Number of times the path encloses the origin
- Number of edge states
 - Zero for trivial states
 - Non-zero in topological states with edge modes

Phase Transition in the JH Model



Here, $t'_2 = t_2/2 = -0.6t$.

Hidden symmetry in the JH Model

$H(k)$ has exchange symmetry :

$$\Upsilon \psi_k = \psi'_k = \begin{pmatrix} B_k \\ A_k \\ C_k \end{pmatrix}$$

So it can be block-diagonalised in the basis of the operator Υ , under the transformation $\mathcal{U}^{-1} \mathcal{H}(k) \mathcal{U} = \mathcal{H}^{BD}(k)$

with

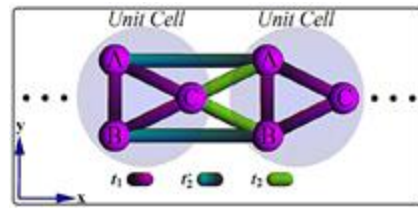
$$\mathcal{U} = \frac{1}{\sqrt{2}} \begin{pmatrix} -1 & 1 & 0 \\ 1 & 1 & 0 \\ 0 & 0 & \sqrt{2} \end{pmatrix}.$$

$$\mathcal{H}^{BD}(k) = \begin{pmatrix} h_1 & 0 \\ 0 & h_2 \end{pmatrix}$$

$$h_1 = -2t_1 + \eta,$$

$$h_2 = \begin{pmatrix} \eta & \sqrt{2}(t_1 + t_2 e^{ik}) \\ \sqrt{2}(t_1 + t_2 e^{-ik}) & 0 \end{pmatrix}$$

Generalised SSH Hamiltonian



Invariants for the JH Model

- Definition of an integer topological invariant : $\mathbb{Z} = |n_0 - n_\pi|$

where n_0 and n_π are the number of negative parities at $k_s = 0$ and $k_s = \pi$.

- Analytical expression of topological invariant \mathbb{Z} for the subsystem h_2 is

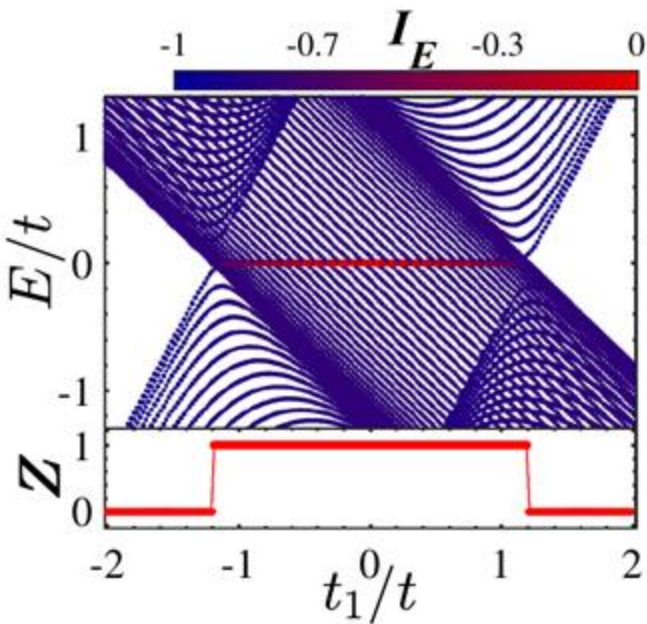
$$\mathbb{Z} = \begin{cases} 0, & \text{if } \text{sgn}(\eta(0)) = \text{sgn}(\eta(\pi)) \\ 1, & \text{if } \text{sgn}(\eta(0)) \neq \text{sgn}(\eta(\pi)) \end{cases},$$

- Definition of inverse participation ratio (IPR)

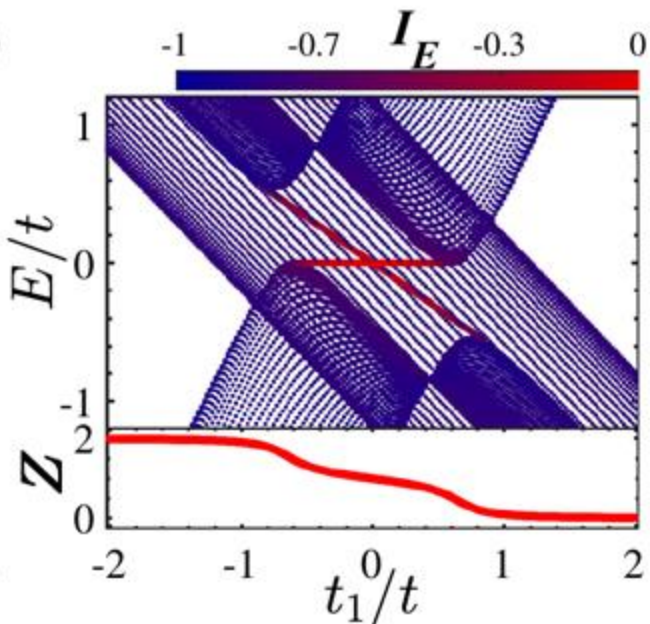
$$I_E = \frac{Ln \sum_j |\psi_E(j)|^4}{Ln 3N}.$$

- $I_E = 0$ for a localised eigenstate
- $I_E = -1$ for an extended eigenstate

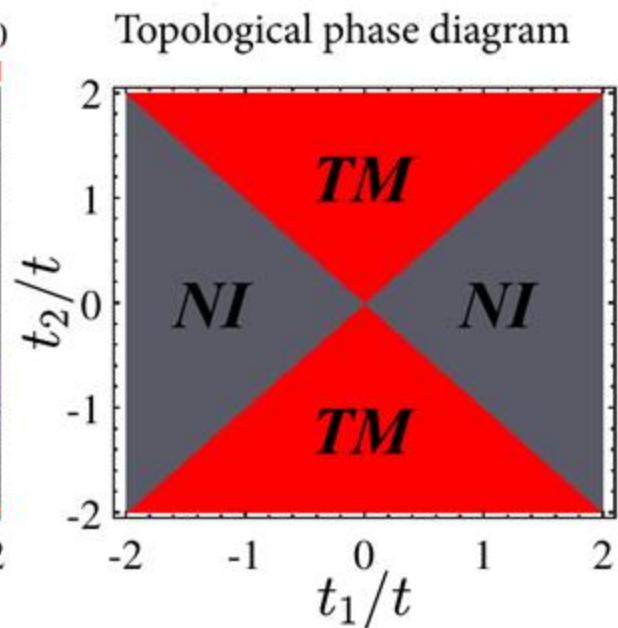
Numerical Results for JH Model



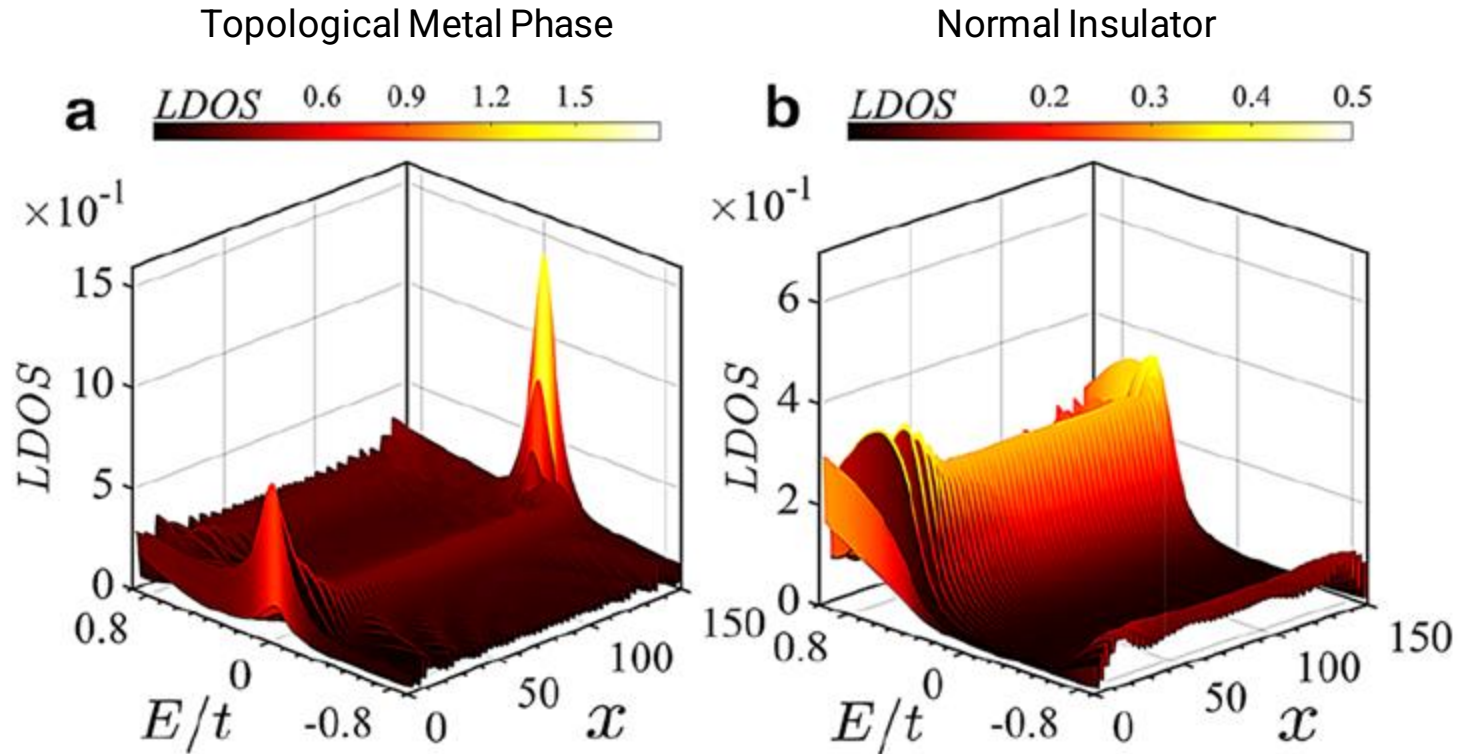
(a) $t'_2 = t_2/2$
 $t'_2 = -0.6t$



(b) $t'_2 \neq t_2/2 = -0.3$.
 $t'_2 = -0.6t$

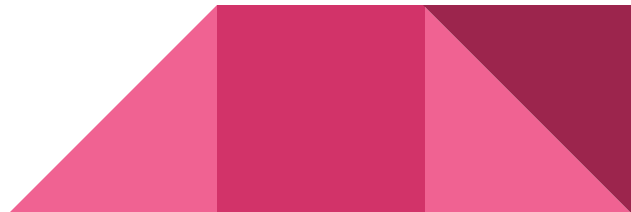


Numerical Results for JH Model



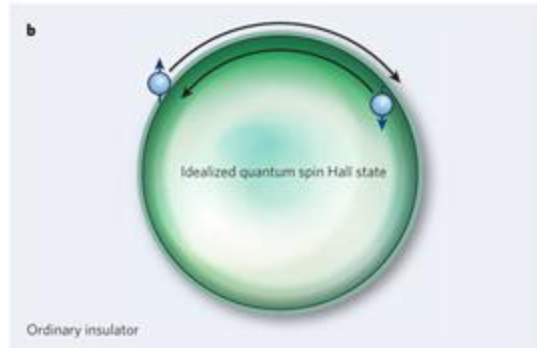
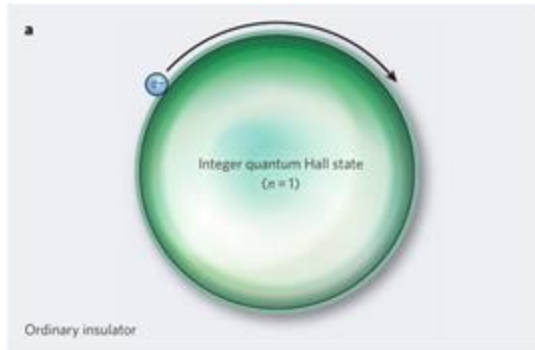
Possible experimental realisations of the JH Model

- Edge states created in the absence of chiral or particle-hole symmetry and *protected* by the hidden inversion (time-reversal) symmetry
- Experimentally, model can be realized by coupled acoustic resonators, topoelectrical circuits, optical lattices, photonic crystals and mechanical systems
- Can also simulate using cold atoms, to reveal the topological features employing density and momentum-distribution measurements
- LDOS can be probed using spatially resolved radio-frequency spectroscopy
- Can distinguish between topologically trivial and nontrivial edge states, using edge state transport in topological states of matter



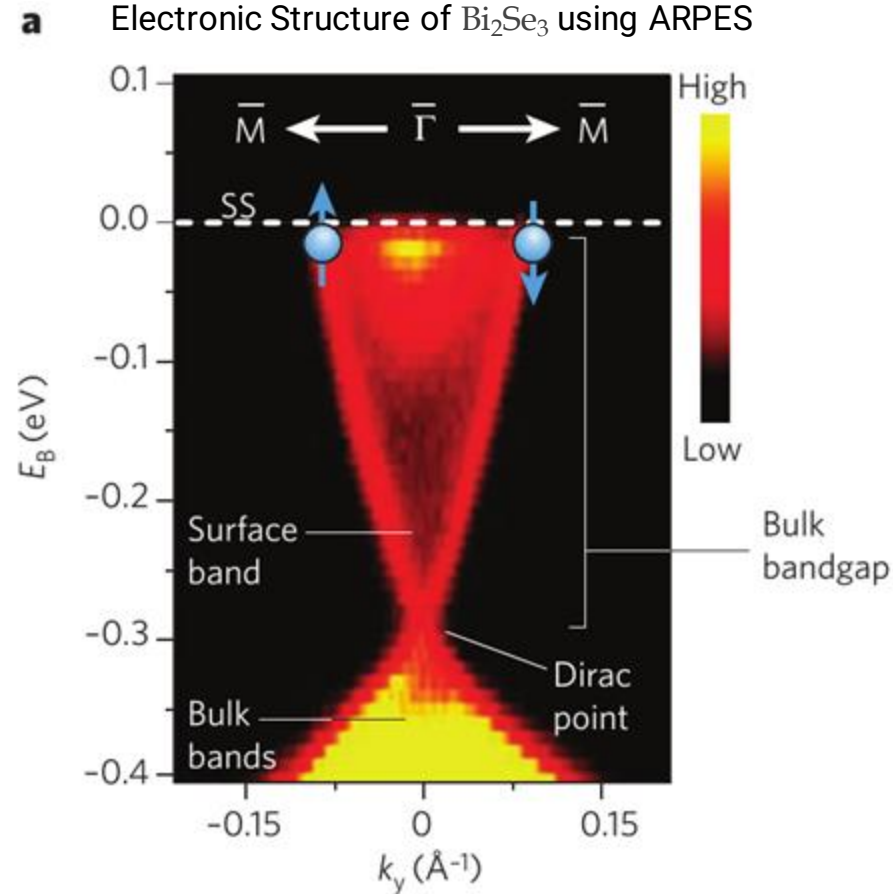
2D topological order

- Quantum Hall effect and Quantum spin Hall effect
- Theoretical Advances by Kane and Mele - topological invariant
- Prediction by Bernevig, Hughes and Zhang that 2D topological insulators with quantized charge conductance would be realised in (Hg,Cd)Te quantum wells



3D topological insulators

- 'Weak' and 'Strong' topological insulators
- The first topological insulator : alloy $\text{Bi}_x\text{Sb}_{1-x}$ using angle-resolved photoemission spectroscopy (ARPES) experiment
- 'Next-generation' topological insulators : Bi_2Se_3 and Bi_2Te_3



Applications and Future Scope

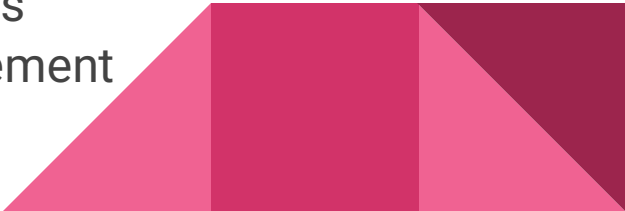
- Spin torque device for magnetic memory applications
 - Heterostructure - topological insulator + ferromagnet
 - Switching of ferromagnet by passing current in TI surface
- Magnetoelectric effect and Axion dynamics
 - Applied electric field generates a magnetic dipole and vice versa
 - Advantage of multiferroic materials - increase in speed and reproducibility without fatigue
- Emergent particles - Majorana fermions
 - Direct observation
 - Obey special kind of non-abelian quantum statistics
 - Topological quantum computers well protected from error



The Course - NE240

- Symmetry
 - Symmetry in Real space
 - Group Theory
 - 32 point and 230 space Groups
- Thermodynamics and Stat. Mech.
 - Atomic picture
 - Origin of various properties
- Phase transitions
 - Change in space group
 - Landau theory
 - Order Parameters
- Disorders
- Measurement

The Presentation

- Symmetry
 - Symmetry in k-space
 - Topology
 - 10 discrete symmetry classes
 - Band Theory
 - SSH lattice and quasi-1D (paper)
 - Spin-orbital coupling, Magnetic field
 - Phase transitions
 - Change in topological invariant
 - ~~Landau theory~~
 - Topological Invariant - Winding number
 - Disorders
 - Measurement
- 



Thank You