### **Differential flatness**

Consider a dynamical system with

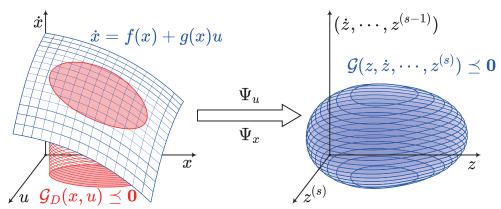
$$\dot{x}=f(x)+g(x)u, \ f \colon \mathbb{R}^n \mapsto \mathbb{R}^n, \ g \colon \mathbb{R}^n \mapsto \mathbb{R}^n, \ x \in \mathbb{R}^n, \ u \in \mathbb{R}^m, \ \mathrm{rank}(g)=m.$$

The system is differentially flat, if there exists  $z \in \mathbb{R}^m$  which can be determined by the state  $x \in \mathbb{R}^n$  and finite derivatives of  $u \in \mathbb{R}^m$ .

 $z \in \mathbb{R}^m$  is called the *flat output*, whose finite derivatives can uniquely determine all state and input:

$$egin{aligned} x &= \Psi_x(z,\dot{z},\cdots,z^{(s-1)}), \ u &= \Psi_u(z,\dot{z},\cdots,z^{(s)}). \end{aligned}$$

#### Differential flatness eliminates differential constraints.



#### **Multicopter State**

$$x=\{r,v,R,\omega\}\in\mathbb{R}^3{ imes}\mathrm{BO}(3) imes\mathbb{R}^3$$
 position, velocity, attitude, body rate

#### **Multicopter Control Input (After Input Mapping)**

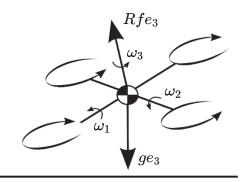
$$u=\{f, au\}\in\mathbb{R}_{\geq 0}{ imes}\mathbb{R}^3$$
 collective thrust, torque

#### **Multicopter Nonlinear Dynamics**

$$\left\{egin{aligned} \dot{r} &= v, \ m\dot{v} &= -mge_3 - RDR^{\mathrm{T}}\sigma(\|v\|)v + Rfe_3, \ \dot{R} &= R\hat{\omega}, \ M\dot{\omega} &= au - \omega imes M\omega - A(\omega) - B(R^{\mathrm{T}}v). \end{aligned}
ight.$$

#### **Multicopter Flat Output**

$$z=\{r,\psi\}\in\mathbb{R}^3 imes \mathrm{SO}(2)$$
 position, yaw heading



$$m\in \mathbb{R}_{\geq 0}$$

$$g\in \mathbb{R}_{\geq 0}$$

$$e_3 = (0,0,1)^{
m T}$$

$$D = \operatorname{Diag}(d_h, d_h, d_v)$$
 drag force coefficients

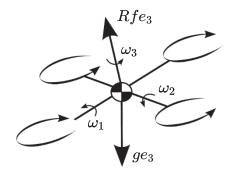
$$\sigma: \mathbb{R}_{\geq 0} \mapsto \mathbb{R}_{\geq 0}$$

$$M\in \mathbb{R}^{3 imes 3}_{>0}$$

$$A \colon \mathbb{R}^3 {\mapsto} \mathbb{R}^3$$

$$B: \mathbb{R}^3 \mapsto \mathbb{R}^3$$

#### **Flatness Transformation?**



Obviously, we have

$$r=r$$
  $v=\dot{r}$ 

Dot product the Newton equation by *X* and *Y* axes of the body frame:

$$egin{aligned} x_b = Re_1, \ y_b = Re_2 \ &(Re_i)^{ ext{T}} m \dot{v} = &(Re_i)^{ ext{T}} ig(-mge_3 - RDR^{ ext{T}} \sigma(\|v\|) v + Rfe_3ig), \ orall i \in \{1,2\} \end{aligned}$$

These yield

$$(Re_i)^{\mathrm{T}}(\dot{v}+rac{d_h}{m}\sigma(\|v\|)v+ge_3)=0,\ orall i\in\{1,2\}.$$

What does it mean?

$$(Re_i)^{\mathrm{T}}(\dot{v}+rac{d_h}{m}\sigma(\|v\|)v+ge_3)=0,\ orall i\in\{1,2\}.$$

Geometrically, we have

$$egin{aligned} x_b \perp (\dot{v} + rac{d_h}{m} \sigma(\|v\|) v + g e_3) \ y_b \perp (\dot{v} + rac{d_h}{m} \sigma(\|v\|) v + g e_3) \end{aligned}$$

Since the thrust force and body Z axis share the same direction,

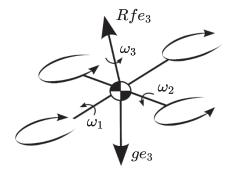
$$z_b = \mathcal{N}(\dot{v} + rac{d_h}{m}\sigma(\|v\|)v + ge_3), \ \mathcal{N}(x) := x/\|x\|_2.$$

Dot product the Newton equation by Z axis of the body frame:

$$(Re_3)^{\mathrm{T}} m \dot{v} = \! (Re_3)^{\mathrm{T}} ig( -mge_3 - RDR^{\mathrm{T}} \sigma(\|v\|) v + Rfe_3 ig)$$

We obtain

$$f=z_b^{
m T}(m\dot{v}+d_v\sigma(\|v\|)v+mge_3)$$



Now that the yaw heading and body Z axis are both known, the yaw rotation quaternion is  $q_\psi = (\cos(\psi/2),0,0,\sin(\psi/2))^{\mathrm{T}}$ 

quaternion is 
$$q_{\psi}=(\cos(\psi/2),0,0,\sin(\psi/2))^{\mathrm{T}}$$
 The tilt rotation quaternion is  $q_z=rac{1}{\sqrt{2(1+z_b(3))}}(1+z_b(3),-z_b(2),z_b(1),0)^{\mathrm{T}}$ 

The tilt rotation has no component for inertial Z axis , since it is decomposed using Hopf fibration. The attitude quaternion of vehicle is thus

Expanding it yields 
$$q = \frac{1}{\sqrt{2(1+z_b(3))}} \begin{pmatrix} (1+z_b(3))\cos(\psi/2) \\ -z_b(2)\cos(\psi/2) + z_b(1)\sin(\psi/2) \\ z_b(1)\cos(\psi/2) + z_b(2)\sin(\psi/2) \\ (1+z_b(3))\sin(\psi/2), \end{pmatrix}$$

Thus the attitude rotation matrix is uniquely determined,  $R = \mathcal{R}_{quat}(q)$ 

Quaternion product, inverse, and rotation matrix formula are detailed by Vince.

Now the rotation is known.

$$\dot{R} = R\hat{\omega}$$

This implies

$$\omega = (R^{
m T}\dot{R})^ee$$

Equivalently,

$$\omega = 2(q_z \otimes q_\psi)^{-1} \otimes \left(\dot{q}_z {\otimes} q_\psi + q_z \otimes \dot{q}_\psi
ight)$$

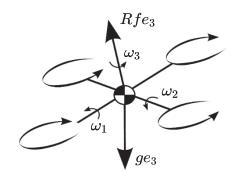
Substituting the tilt rotation quaternion and yaw rotation quaternion gives

$$\omega = egin{pmatrix} \dot{z}_b(1)\sin(\psi) - \dot{z}_b(2)\cos(\psi) - \dot{z}_b(3)(z_b(1)\sin(\psi) - z_b(2)\cos(\psi))/(1+z_b(3)) \ \dot{z}_b(1)\cos(\psi) + \dot{z}_b(2)\sin(\psi) - \dot{z}_b(3)(z_b(1)\cos(\psi) + z_b(2)\sin(\psi))/(1+z_b(3)) \ (z_b(2)\dot{z}_b(1) - z_b(1)\dot{z}_b(2))/(1+z_b(3)) + \dot{\psi}, \end{pmatrix}$$

Differentiate body Z axis gives

$$egin{aligned} \dot{z}_b &= rac{d_h}{m} \mathcal{DN}(\dot{v} + rac{d_h}{m} \sigma(\|v\|) v + g e_3)^{\mathrm{T}} (rac{m}{d_h} \ddot{v} + \sigma(\|v\|) \dot{v} + \dot{\sigma}(\|v\|) rac{v^{\mathrm{T}} \dot{v}}{\|v\|} v), \ \mathcal{DN}(x) &:= rac{1}{\|x\|} igg( \mathbf{I} - rac{x x^{\mathrm{T}}}{x^{\mathrm{T}} x} igg). \end{aligned}$$

Now we have expressed the body rate using finite derivatives of the flat output.

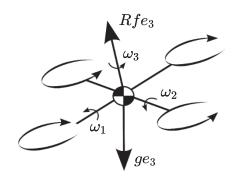


Now all states and inputs are known except the torque.

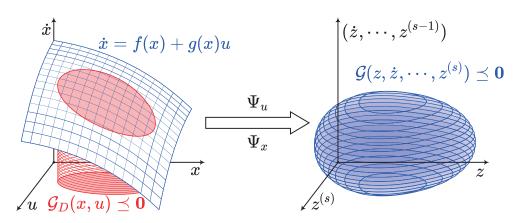
We trivially differentiate the body rate to get its expression using  $r^{(4)}, \ \psi^{(2)}$ 

Consequently,

$$au = M\dot{\omega} + \omega imes M\omega + A(\omega) + B(R^{
m T}v)$$



Finally, we have finished the derivation of **flatness transformation**.



Planning flat-output trajectories with high-order continuity suffices for the dynamics.