Introduction to Quantum Computing

Lecture slides for the Isogeny-based Cryptography School 2021

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30th August 2021

Background

Quantum computers build on the principles of quantum mechanics. It can solve some problems much faster than the traditional computers.

A famous example is the Shor's integer factorization algorithm.





The widely used cryptosystem, RSA, relies on factoring being impossible for large integers. But Shor's algorithm shows that this problem is easy for a quantum computer.

Backgrounds

To study quantum computing, don't worry if you don't know too much about quantum mechanics. What you need to know is linear algebra in this lecture.

In this lecture, I will introduce some fundamental concepts and results. Hope to help you better understand other lectures this week.

I will not introduce the definitions in a very formal way because you can find it in many textbooks. I prefer to use examples to explain the concepts.

Preliminaries

The Deutsch-Jozsa algorithm

Simon's algorithm

Quantum Fourier transform

Grover's algorithm

Further readings

Qubits

Qubit (Quantum bit): $\alpha|0\rangle + \beta|1\rangle$, where $|\alpha|^2 + |\beta|^2 = 1$ and

$$|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad |1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}.$$

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2 qubit state:

$$\alpha_{00}|0\rangle\otimes|0\rangle + \alpha_{01}|0\rangle\otimes|1\rangle + \alpha_{10}|1\rangle\otimes|0\rangle + \alpha_{11}|1\rangle\otimes|1\rangle$$

where $|\alpha_{00}|^2 + |\alpha_{01}|^2 + |\alpha_{10}|^2 + |\alpha_{11}|^2 = 1$.

We will simply write $|i\rangle \otimes |j\rangle$ as $|i\rangle |j\rangle$, $|i,j\rangle$ or $|ij\rangle$.

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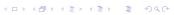
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n qubit state:

$$\sum_{x \in \{0,1\}^n} \alpha_x |x\rangle$$

where $\sum_{x \in \{0,1\}^n} |\alpha_x|^2 = 1$.



Dirac notation or bra-ket notation

For a unit column vector $\mathbf{v} = (v_0, \dots, v_{n-1})^T$, in quantum computing, we denote it as ("ket notation")

$$|\mathbf{v}\rangle = \sum_{j=0}^{n-1} v_j |j\rangle \leftrightarrow \begin{pmatrix} v_0 \\ \vdots \\ v_{n-1} \end{pmatrix}$$

where $\{|0\rangle, \dots, |n-1\rangle\}$ corresponds to the standard basis of \mathbb{C}^n .

Its conjugate transpose is denoted as ("bra notation")

$$\langle \mathbf{v} | = \sum_{j=0}^{n-1} \bar{v}_j \langle j | \leftrightarrow (\bar{v}_0 \quad \cdots \quad \bar{v}_{n-1})$$

It is a row vector.

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Examples:

Hadamard gate:
$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$$
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$$\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \sigma_y = \begin{pmatrix} 0 & i \\ -i & 0 \end{pmatrix}, \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}.$$

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Control gate: $|0\rangle\langle 0|\otimes U_0+|1\rangle\langle 1|\otimes U_1$, where U_0,U_1 are unitary operators. It means if the first qubit is $|i\rangle$, then we apply U_i to the second state.

$$\alpha_0|0\rangle|\phi_0\rangle + \alpha_1|1\rangle|\phi_1\rangle \mapsto \alpha_0|0\rangle U_0|\phi_0\rangle + \alpha_1|1\rangle U_1|\phi_1\rangle.$$

The matrix form
$$\begin{pmatrix} U_0 & \\ & U_1 \end{pmatrix}$$
 .



Measurements

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$$|\phi\rangle = \frac{1}{2}|00\rangle + \frac{1}{2}|01\rangle - \frac{1}{\sqrt{2}}|11\rangle.$$

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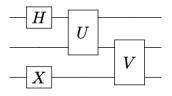
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We can do partial measurement. For $|\phi\rangle$, if we only measure the first qubit, then with probability 1/2, we obtain $|0\rangle$. The state associated to $|0\rangle$ is $(|0\rangle + |1\rangle)/\sqrt{2}$.

Quantum circuit

A quantum circuit can be drawn as a diagram by associating each qubit with a horizontal "wire", and drawing each gate as a box across the wires corresponding to the qubits on which it acts.

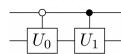


The above circuit corresponds to the unitary operator

$$(I_2 \otimes V)(U \otimes I_2)(H \otimes I_2 \otimes X)$$

on 3 qubits.

For control gate $|0\rangle\langle 0|\otimes U_0+|1\rangle\langle 1|\otimes U_1$, the quantum circuit is



Let $f:\{0,1\}^m \to \{0,1\}^n$ be a function, it becomes a unitary operator by the following trick

$$f': \{0,1\}^m \times \{0,1\}^n \to \{0,1\}^m \times \{0,1\}^n$$

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 $(x,y) \to (x,y \oplus f(x)).$

In a quantum computer, we denote it as

$$O_f: \{0,1\}^m \times \{0,1\}^n \rightarrow \{0,1\}^m \times \{0,1\}^n$$
$$|x\rangle|y\rangle \rightarrow |x\rangle|y \oplus f(x)\rangle.$$

It is called an oracle to query functions.

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We can implement U_f from O_f .

More precisely, denote $|-\rangle = \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$, then

$$|x\rangle|-\rangle \xrightarrow{O_f} \frac{1}{\sqrt{2}}|x\rangle(|f(x)\rangle-|1\oplus f(x)\rangle).$$

If f(x)=0, the result is $|x\rangle|-\rangle$; If f(x)=1, the result is $-|x\rangle|-\rangle$. In summary, the result is

$$(-1)^{f(x)}|x\rangle|-\rangle.$$

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A set of quantum gates is called universal if any unitary operator can be approximately represented as a circuit the gates in the set.

For example, the set $\{H, T, C\}$ with

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}, \quad T = \begin{pmatrix} 1 & 0 \\ 0 & e^{-\pi i/4} \end{pmatrix}, \quad C = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}.$$

Complexity

Gate complexity: The number of elementary gates used in the universal set.

Up to some poly-log terms, the gate complexity does not change if universal set varies.

Query complexity: The number of evaluations to the given function, i.e., the number of O_f (or U_f) used in the quantum circuit.

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Deutsch-Jozsa problem

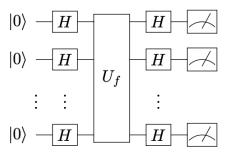
The Deutsch-Jozsa algorithm was the first to show a separation between the quantum and classical difficulty of a problem.

Definition 1 (Deutsch-Jozsa problem)

Let $f:\{0,1\}^n \to \{0,1\}$. It is promised to be constant or balanced (i.e., $|f^{-1}(0)| = |f^{-1}(1)| = 2^{n-1}$). The goal is to decide which is the case by making as few function evaluations as possible.

Classically, it requires $2^{n-1}+1$ function evaluations. However, the Deutsch-Jozsa algorithm only uses 1 function evaluation.

The circuit of Deutsch-Jozsa algorithm is very simple:



The last step means measurement.

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- 2. In the first step, we apply $H^{\otimes n}$, then we obtain

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3. In the second step, we apply U_f which gives

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4. Finally, we apply $H^{\otimes n}$ again

$$\frac{1}{2^n} \sum_{z \in \{0,1\}^n} \left(\sum_{y \in \{0,1\}^n} (-1)^{f(y) + y \cdot z} \right) |z\rangle.$$

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So final state is $|0\rangle^{\otimes n}$. If we perform measurement, we always obtain $|0\rangle^{\otimes n}$.

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So final state is $|0\rangle^{\otimes n}$. If we perform measurement, we always obtain $|0\rangle^{\otimes n}$.

▶ If f is balanced, then the coefficient of $|0\rangle^{\otimes n}$

$$\frac{1}{2^n} \sum_{y \in \{0,1\}^n} (-1)^{f(y)} = 0.$$

We therefore never obtain $|0\rangle^{\otimes n}$ by measuring the final state.

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Simon's problem

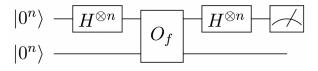
Simon's algorithm was the first quantum algorithm to show an exponential speed-up versus the best classical algorithm.

Definition 2 (Simon's problem)

Let $f:\{0,1\}^n \to \{0,1\}^n$. There is a unknown s such that f(x)=f(y) if and only if $y=x\oplus s$. The goal is to find s.

The classical algorithm needs at least $2^{n/2}$ queries to f. While Simon's algorithm only uses O(n) queries.

The circuit of Simon's algorithm is very similar to the circuit of Deutsch-Jozsa algorithm:



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3. In the second step, we apply O_f

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4. Finally we apply $H^{\otimes n} \otimes I$ again

$$\frac{1}{2^n} \sum_{z \in \{0,1\}^n} \sum_{y \in \{0,1\}^n} (-1)^{y \cdot z} |z\rangle |f(y)\rangle.$$

Recall that f(x)=f(y) iff $y=x\oplus s$, so we can split $\{0,1\}^n$ into $A\cup (A\oplus s)$. On $A,\,f$ is one-to-one.

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In the final state,

$$\sum_{y \in \{0,1\}^n} (-1)^{y \cdot z} |z\rangle = \sum_{y \in A} \left((-1)^{y \cdot z} + (-1)^{(y \oplus s) \cdot z} \right) |z\rangle
= \sum_{y \in A} (-1)^{y \cdot z} \left(1 + (-1)^{s \cdot z} \right) |z\rangle.$$

The coefficient is nonzero if $s \cdot z = 0$.

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If we run the above process n-1 times, we obtain z_1,\ldots,z_{n-1} such that $s\cdot z_i=0$ for all i. From this linear system, we can determine s.

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Quantum Fourier transform

Applications 1: quantum phase estimation

Applications 2: period finding

Grover's algorithm

Further readings

Quantum Fourier Transform (QFT)

Definition 3 (Quantum Fourier Transform (QFT))

Let N be a integer, $\omega=e^{2\pi i/N}$, the QFT is defined by

$$Q_N: \mathbb{Z}_N \to \mathbb{Z}_N |x\rangle \mapsto \frac{1}{\sqrt{N}} \sum_{y \in \mathbb{Z}_N} \omega^{xy} |y\rangle.$$

- ► A very important unitary operator in quantum information theory.
- ▶ It is the normalized discrete Fourier transform.

Quantum Fourier Transform (QFT)

In matrix form:

$$Q_N = \frac{1}{\sqrt{N}} \sum_{x,y \in \mathbb{Z}_N} \omega^{xy} |y\rangle \langle x|.$$

The inverse of QFT is

$$Q_N^{-1}: |x\rangle \mapsto \frac{1}{\sqrt{N}} \sum_{y \in \mathbb{Z}_N} \omega^{-xy} |y\rangle.$$

Check:

$$Q_N^{-1}Q_N|x\rangle = \frac{1}{N} \sum_{z \in \mathbb{Z}_N} \left(\sum_{y \in \mathbb{Z}_N} \omega^{y(x-z)} \right) |z\rangle = |x\rangle.$$

Quantum Fourier Transform (QFT)

Example 4

$$Q_2 = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}, \quad Q_3 = \frac{1}{\sqrt{3}} \begin{pmatrix} 1 & 1 & 1 \\ 1 & e^{2\pi i/3} & e^{-2\pi i/3} \\ 1 & e^{-2\pi i/3} & e^{2\pi i/3} \end{pmatrix},$$

$$Q_4 = \frac{1}{2} \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & i & -1 & -i \\ 1 & -1 & 1 & -1 \\ 1 & -i & -1 & i \end{pmatrix}.$$

It can be implemented using $O(\log^2 N)$ elementary quantum gates:

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}, \quad R_d = \begin{pmatrix} 1 & 0 \\ 0 & e^{\pi i/2^d} \end{pmatrix}.$$

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Let's take a look at the case N=8:

$$\begin{aligned} &Q_8\big|x_0,x_1,x_2\big\rangle\\ &=& \frac{1}{\sqrt{8}}\sum_{y_0,y_1,y_2=0}^1 e^{2\pi i(x(y_0+2y_1+4y_2))/8}|y_0,y_1,y_2\rangle\\ &=& \frac{1}{\sqrt{8}}\left(\sum_{y_0=0}^1 e^{\pi ixy_0/4}|y_0\rangle\right)\left(\sum_{y_1=0}^1 e^{\pi ixy_1/2}|y_1\rangle\right)\left(\sum_{y_2=0}^1 e^{\pi ixy_2}|y_2\rangle\right) \end{aligned}$$

Note: $|x\rangle = |x_0, x_1, x_2\rangle$ and $x = x_0 + 2x_1 + 4x_2$ is the binary form.

$$\sum_{y_0=0}^1 e^{\pi i x y_0/4} |y_0\rangle \quad = \quad \sum_{y_0=0}^1 e^{\pi i x_0 y_0/4} e^{\pi i x_1 y_0/2} e^{\pi i x_2 y_0} |y_0\rangle$$

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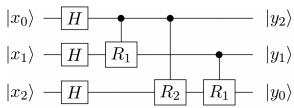
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$$\sum_{y_0=0}^{1} e^{\pi i x y_0/4} |y_0\rangle = \sum_{y_0=0}^{1} e^{\pi i x_0 y_0/4} e^{\pi i x_1 y_0/2} e^{\pi i x_2 y_0} |y_0\rangle$$

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$$\sum_{y_2=0}^{1} e^{\pi i x y_2} |y_2\rangle = \sum_{y_2=0}^{1} e^{\pi i x_0 y_2} |y_2\rangle.$$



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3. Apply control gate $\sum_{x} |x\rangle\langle x| \otimes U^x$:

$$\frac{1}{\sqrt{2^n}} \sum_{x=0}^{2^n-1} e^{2\pi i x \theta} |x\rangle |\psi\rangle.$$

An important subroutine of many quantum algorithms.

Input: a unitary operator U and an eigenvector $|\psi\rangle$.

Output: $\theta \in [0, 2\pi)$ such that $U|\psi\rangle = e^{2\pi i\theta}|\psi\rangle$.

- 1. Prepare the initial state $|0^n\rangle|\psi\rangle$.
- 2. Apply Hadamard gates $H^{\otimes n}$ to the first register:

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4. Apply QFT⁻¹ to $|x\rangle$:

$$\frac{1}{2^n} \sum_{y=0}^{2^n-1} \left(\sum_{x=0}^{2^n-1} e^{2\pi i x (\theta - y/2^n)} \right) |y\rangle |\psi\rangle.$$

Denote $\delta_y = \theta - y/2^n$. The coefficient of $|y\rangle|\psi\rangle$ is

$$\frac{1}{2^n} \left| \sum_{x=0}^{2^{n-1}} e^{2\pi i \delta_y x} \right| = \frac{1}{2^n} \left| \frac{e^{2\pi i \delta_y 2^n} - 1}{e^{2\pi i \delta_y} - 1} \right| = \frac{1}{2^n} \left| \frac{\sin(\pi \delta_y 2^n)}{\sin(\pi \delta_y)} \right|.$$

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If $|\delta_y|2^n \leq 1/2$, then the above quantity is lower bounded by

$$\geq \frac{1}{2^n} \left| \frac{2\delta_y 2^n}{\pi \delta_y} \right| = \frac{2}{\pi}$$

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We can modify the algorithm to ensure the success probability is at least $1-\epsilon$ for arbitrary small ϵ .

One of the most important applications of the QFT, the key step of Shor's algorithm.

One of the most important applications of the QFT, the key step of Shor's algorithm.

Imagine we are given access to an oracle O_f function $f: \mathbb{Z}_N \to \mathbb{Z}_M$, for some integers N and M, such that:

- ▶ f is periodic: there exists r such that r divides N and f(x+r) = f(x) for all $x \in \mathbb{Z}_N$;
- ▶ f is one-to-one on each period: for all pairs (x,y) such that $|x-y| < r, f(x) \neq f(y)$.

Our task is to determine r.

Recall: $O_f: |x\rangle |y\rangle = |x\rangle |y \oplus f(x)\rangle$.

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$$\frac{1}{\sqrt{N}} \sum_{x=0}^{N-1} |x\rangle |f(x)\rangle = \frac{1}{\sqrt{N}} \sum_{y=0}^{r-1} \left(\sum_{j=0}^{N/r-1} |y+jr\rangle \right) |f(y)\rangle$$

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▶ Measure the second register: obtain a random *y*

$$\frac{\sqrt{r}}{\sqrt{N}} \sum_{j=0}^{N/r-1} |y+jr\rangle$$

• Apply Q_N to the first register: $\omega = e^{2\pi i/N}$

$$\frac{\sqrt{r}}{N} \sum_{z=0}^{N-1} \omega^{yz} \left(\sum_{j=0}^{N/r-1} \omega^{jrz} \right) |z\rangle.$$

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Note that if $\omega^{rz} \neq 1$, i.e., $rz \not\equiv 0 \mod N$, then

$$\sum_{j=0}^{N/r-1} \omega^{jrz} = \frac{\omega^{rzN} - 1}{\omega^{rz} - 1} = 0.$$

So the state is

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Measure it we obtain a random s (unknown) and z (known) such that z/N = s/r. If s is coprime to r, then we can determine r by simplify z/N. This happens with probability at least $1/\log\log r$.

Input: Integer N

Output: integers p,q such that ${\cal N}=pq$

¹gcd = greatest common divisor.

²order: the minimal r>0 s.t. $a^r\equiv 1\mod N$.

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Output: integers p, q such that N = pq

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Input: Integer N

- 1. Choose 1 < a < N uniformly at random.
- 2. Compute $b = \gcd(a, N)$. If b > 1 output b and stop. ¹

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Step 3 is technical, it relates to period finding. Consider

$$f(x) = a^x \mod N$$
.

We can check that f is periodic with period r and one-to-one on each period.

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Preliminaries

The Deutsch-Jozsa algorithm

Simon's algorithm

Quantum Fourier transform

Grover's algorithm

Further readings

A simple example of a problem that fits into the query complexity model is the unstructured search problem.

Definition 5 (Grover's search problem)

Given access to a function $f: \mathbb{Z}_N \to \{0,1\}$ with the promise that $f(x_0) = 1$ for a unique element x_0 . Our task is to determine x_0 .

Classical algorithm: N queries (i.e., N function evaluations to f). Quantum algorithm: $O(\sqrt{N})$ queries.

1. Prepare $|\phi\rangle = H^{\otimes n}|0^n\rangle$

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- 2. Repeat the following operations $O(\sqrt{N})$ times:
 - 2.1 Apply U_f
 - 2.2 Apply $D:=-H^{\otimes n}U_0H^{\otimes n}$, where U_0 maps $|0^n\rangle$ to $-|0^n\rangle$ and keeps all other basis states invariant.

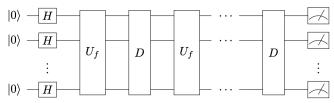
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- 3. Measure all the qubits and output the result.

Recall: $U_f|x\rangle=(-1)^{f(x)}|x\rangle$. This is a reflection.

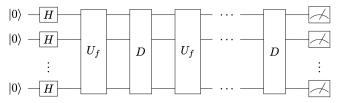
$$D=-H^{\otimes n}(I-2|0^n\rangle\langle 0^n|)H^{\otimes n}=2H^{\otimes n}|0^n\rangle\langle 0^n|H^{\otimes n}-I \text{ is another reflection}.$$

So DU_f is a rotation (need some analysis).

In circuit diagram form, Grover's algorithm looks like this:



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Note that

$$|\phi\rangle = H^{\otimes n}|0^n\rangle = \frac{1}{\sqrt{N}} \sum_{x \in \mathbb{Z}_N} |x\rangle.$$

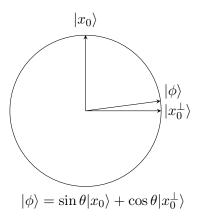
It formally equals

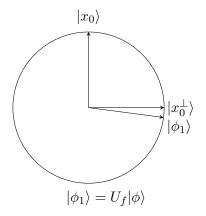
$$|\phi\rangle = \frac{1}{\sqrt{N}}|x_0\rangle + \frac{\sqrt{N-1}}{\sqrt{N}}|x_0^{\perp}\rangle,$$

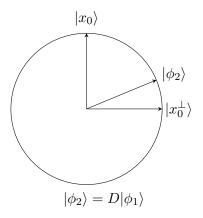
where

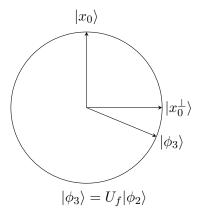
$$|x_0^{\perp}\rangle = \frac{1}{\sqrt{N-1}} \sum_{x \in \mathbb{Z}_N, x \neq x_0} |x\rangle.$$

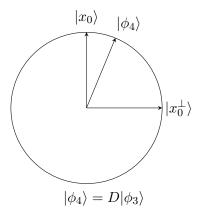












We denote $\sin\theta=1/\sqrt{N}$ and $\cos\theta=\sqrt{N-1}/\sqrt{N}$. In step 3, we can denote $U_0=I-2|0^n\rangle\langle 0^n|$ so that $D=-(I-2|\phi\rangle\langle\phi|)$.

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$$-\sin\theta|x_0\rangle + \cos\theta|x_0^{\perp}\rangle.$$

Then apply D to obtain

$$\sin \theta (I - 2|\phi\rangle\langle\phi|)|x_0\rangle - \cos \theta (I - 2|\phi\rangle\langle\phi|)|x_0^{\perp}\rangle$$

$$= \sin \theta (|x_0\rangle - 2\sin \theta (\sin \theta |x_0\rangle + \cos \theta |x_0^{\perp}\rangle))$$

$$- \cos \theta (|x_0^{\perp}\rangle - 2\cos \theta (\sin \theta |x_0\rangle + \cos \theta |x_0^{\perp}\rangle))$$

$$= \sin(3\theta)|x_0\rangle + \cos(3\theta)|x_0^{\perp}\rangle.$$

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As we have seen, DU_f is the product of two reflections in the plane spanned by $\{|x_0\rangle,|x_0^{\perp}\rangle\}$. So DU_f is a rotation of angle 2θ .

Hence, after T steps of iteration we obtain

$$\sin((2T+1)\theta)|x_0\rangle + \cos((2T+1)\theta)|x_0^{\perp}\rangle.$$

Hence, after T steps of iteration we obtain

$$\sin((2T+1)\theta)|x_0\rangle + \cos((2T+1)\theta)|x_0^{\perp}\rangle.$$

Since $\sin\theta=1/\sqrt{N}$, we have $\theta\approx 1/\sqrt{N}$. To make $\sin((2T+1)\theta)$ close to 1, we can choose T so that $(2T+1)\theta\approx\pi/2$. Namely, $T\approx\sqrt{N}\pi/4-1/2$.

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Further readings

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You may find the following lecture notes and books useful:

- ► Lecture Notes on Quantum Algorithms, Andrew Childs, University of Maryland http://www.cs.umd.edu/~amchilds/qa/ An excellent resource for more advanced topics on quantum algorithms.
- Quantum Computing: Lecture Notes, Ronald de Wolf, QuSoft, CWI and University of Amsterdam https://export.arxiv.org/abs/1907.09415 A comprehensive lecture note for more topics on quantum computing.
- Quantum Computation and Quantum Information, Nielsen and Chuang
 Cambridge University Press, 2001
 The Bible of quantum computing.