

第 1 题 (课本习题 3.28) 得分: _____. 考虑两个自旋 $\frac{1}{2}$ 的粒子组成的一个系统. 观察者 A 专门测量其中一个离子的自旋分量 (s_{1z} , s_{1x} , 等等), 同时观察者 B 测量另一个粒子的自旋分量. 假定已知系统处在自旋单态, 即 $S_{\text{总}} = 0$.

- (a) 当观察者 B 不做任何测量时, 观察者 A 得到 $s_{1z} = \hbar/2$ 的概率是什么? 对于 $s_{1x} = \hbar/2$ 求解同样问题.
- (b) 观察者 B 肯定地确认粒子 2 的自旋处于 $s_{2z} = \hbar/2$ 态. 如果观察者 A (i) 测量 s_{1z} ; (ii) 测 s_{1x} , 则对观察者 A 的测量结果能给出的结论是什么? 解释你的答案.

解: (a) 这两个粒子组成的系统处于自旋单态, 可表为

$$|\psi\rangle = \frac{1}{\sqrt{2}}(\hat{z}+; \hat{z}-) + |\hat{z}-; \hat{z}+\rangle = \frac{1}{\sqrt{2}}(\hat{x}-; \hat{x}+) + |\hat{x}+; \hat{x}-\rangle. \quad (1)$$

当观察者 B 不做任何测量时, 观察者 A 得到 $s_{1z} = \hbar/2$ 的概率为 $\frac{1}{2}$;

当观察者 B 不做任何测量时, 观察者 A 得到 $s_{1x} = \hbar/2$ 的概率为 $\frac{1}{2}$.

- (b) 当观察者 B 肯定地确认粒子 2 的自旋处于 $s_{2z} = \hbar/2$ 后, 系统的状态塌缩至

$$|\psi'\rangle = |\hat{z}-; \hat{z}+\rangle = |\hat{z}-\rangle \otimes \frac{1}{\sqrt{2}}(|\hat{x}+\rangle - |\hat{x}-\rangle). \quad (2)$$

(i) 如果观察者测量 s_{1z} , 则有 100% 的概率得到 $s_{1z} = -\hbar/2$.

(ii) 如果观察者测量 s_{1x} , 则有 $\frac{1}{2}$ 的概率得到 $s_{1x} = \hbar/2$, 有 $\frac{1}{2}$ 的概率得到 $s_{1x} = -\hbar/2$.

□

第 2 题 (课本习题 3.30) 得分: _____. (a) 用两个不同的矢量 $\mathbf{U} = (U_x, U_y, U_z)$ 和 $\mathbf{V} = (V_x, V_y, V_z)$ 构造一个秩为 1 的球张量. 明确地用 $U_{x,y,z}$ 和 $V_{x,y,z}$ 写出 $T_{\pm 1,0}^{(1)}$.

- (b) 用两个不同的矢量 \mathbf{U} 和 \mathbf{V} 构造一个秩为 2 的球张量. 明确地用 $U_{x,y,z}$ 和 $V_{x,y,z}$ 写出 $T_{\pm 2, \pm 1, 0}^{(2)}$.

解: (a) 由

$$Y_{l=1}^{m=0}(\mathbf{U}) = \sqrt{\frac{3}{4\pi}} \cos \theta = \sqrt{\frac{3}{4\pi}} \frac{U_z}{|\mathbf{U}|}, \quad (3)$$

$$Y_{l=1}^{m=\pm 1}(\mathbf{U}) = \mp \sqrt{\frac{3}{8\pi}} \sin \theta e^{\pm i\theta} = \mp \sqrt{\frac{3}{4\pi}} \frac{U_x \pm iU_y}{|\mathbf{U}|}, \quad (4)$$

得单由 \mathbf{U} 构造的秩为 1 的球张量

$$U_0^{(1)} = U_z, \quad (5)$$

$$U_{\pm 1}^{(1)} = \mp \frac{U_x \pm iU_y}{\sqrt{2}}. \quad (6)$$

对 \mathbf{V} 同理有

$$V_0^{(1)} = V_z, \quad (7)$$

$$V_{\pm 1}^{(1)} = \mp \frac{V_x \pm iV_y}{\sqrt{2}}. \quad (8)$$

利用定理

$$T_q^{(k)} = \sum_{q_1} \sum_{q_2} \langle k_1 k_2; q_1 q_2 | k_1 k_2; k q \rangle X_{q_1}^{(k_1)} Z_{q_2}^{(k_2)}, \quad (9)$$

有

$$T_q^{(1)} = \sum_{q_1} \sum_{q_2} \langle 11; q_1 q_2 | 11; 1q \rangle U_{q_1}^{(1)} V_{q_2}^{(2)}, \quad (10)$$

从而

$$\begin{aligned} T_{+1}^{(1)} &= \langle 11; 0, +1 | 11; 1, +1 \rangle U_0^{(1)} V_{+1}^{(1)} + \langle 11; +1, 0 | 11; 1, +1 \rangle U_{+1}^{(1)} V_0^{(1)} \\ &= -\frac{1}{\sqrt{2}} U_z \frac{-V_x - iV_y}{\sqrt{2}} + \frac{1}{\sqrt{2}} \frac{-U_x - iU_y}{\sqrt{2}} V_z \\ &= \frac{1}{2} (U_z V_x + iU_z V_y - U_x V_z - iU_y V_z), \end{aligned} \quad (11)$$

$$\begin{aligned} T_0^{(1)} &= \langle 11; -1, +1 | 11; 10 \rangle U_{-1}^{(1)} V_{+1}^{(1)} + \langle 11; 00 | 11; 10 \rangle U_0^{(1)} V_0^{(1)} + \langle 11; +1, -1 | 11; 10 \rangle U_{+1}^{(1)} V_{-1}^{(1)} \\ &= -\frac{1}{\sqrt{2}} \frac{U_x - iU_y}{\sqrt{2}} \frac{-V_x - iV_y}{\sqrt{2}} + \frac{1}{\sqrt{2}} \frac{-U_x - iU_y}{\sqrt{2}} \frac{V_x - iV_y}{\sqrt{2}} \\ &= \frac{i}{\sqrt{2}} (U_x V_y - U_y V_x), \end{aligned} \quad (12)$$

$$\begin{aligned} T_{-1}^{(1)} &= \langle 11; -1, 0 | 11; 1, -1 \rangle U_{-1}^{(1)} V_0^{(1)} + \langle 11; 0, -1 | 11; 1, -1 \rangle U_0^{(1)} V_{-1}^{(1)} \\ &= -\frac{1}{\sqrt{2}} \frac{U_x - iU_y}{\sqrt{2}} V_z + \frac{1}{\sqrt{2}} U_z \frac{V_x - iV_y}{\sqrt{2}} \\ &= \frac{1}{2} (-U_x V_z + iU_y V_z + U_z V_x - iU_z V_y). \end{aligned} \quad (13)$$

(b) 利用定理

$$T_q^{(k)} = \sum_{q_1} \sum_{q_2} \langle k_1 k_2; q_1 q_2 | k_1 k_2; kq \rangle X_{q_1}^{(k_1)} Z_{q_2}^{(k_2)}, \quad (14)$$

有

$$T_q^{(2)} = \sum_{q_1} \sum_{q_2} \langle 11; q_1 q_2 | 11; 2q \rangle U_{q_1}^{(1)} V_{q_2}^{(1)}, \quad (15)$$

从而

$$T_{+2}^{(2)} = \langle 11; +1, +1 | 11; 2, +2 \rangle U_{+1}^{(1)} V_{+1}^{(1)} = \frac{-U_x - iU_y}{\sqrt{2}} \frac{-V_x - iV_y}{\sqrt{2}} = \frac{1}{2} (U_x V_x + iU_x V_y + iU_y V_x + U_y V_y), \quad (16)$$

$$\begin{aligned} T_{+1}^{(2)} &= \langle 11; 0, +1 | 11; 2, +1 \rangle U_0^{(1)} V_{+1}^{(1)} + \langle 11; +1, 0 | 11; 2, +1 \rangle U_{+1}^{(1)} V_0^{(1)} \\ &= \frac{1}{\sqrt{2}} U_z \frac{-V_x - iV_y}{\sqrt{2}} + \frac{1}{\sqrt{2}} \frac{-U_x - iU_y}{\sqrt{2}} V_z \\ &= -\frac{1}{2} [U_z (V_x + iV_y) + (U_x + iU_y) V_z], \end{aligned} \quad (17)$$

$$\begin{aligned} T_0^{(2)} &= \langle 11; -1, +1 | 11; 20 \rangle U_{-1}^{(1)} V_{+1}^{(1)} + \langle 11; 00 | 11; 20 \rangle U_0^{(1)} V_0^{(1)} + \langle 11; +1, -1 | 11; 20 \rangle U_{+1}^{(1)} V_{-1}^{(1)} \\ &= \frac{1}{\sqrt{6}} \frac{U_x - iU_y}{\sqrt{2}} \frac{-V_x - iV_y}{\sqrt{2}} + \frac{2}{\sqrt{6}} U_z V_z + \frac{1}{\sqrt{6}} \frac{-U_x - iU_y}{\sqrt{2}} \frac{V_x - iV_y}{\sqrt{2}} \\ &= \frac{1}{\sqrt{6}} (-U_x V_x - U_y V_y + 2U_z V_z), \end{aligned} \quad (18)$$

$$\begin{aligned} T_{-1}^{(2)} &= \langle 11; -1, 0 | 11; 2, -1 \rangle U_{-1}^{(1)} V_0^{(1)} + \langle 11; 0, -1 | 11; 2, -1 \rangle U_0^{(1)} V_{-1}^{(1)} \\ &= \frac{1}{\sqrt{2}} \frac{U_x - iU_y}{\sqrt{2}} V_z + \frac{1}{\sqrt{2}} U_z \frac{V_x - iV_y}{\sqrt{2}} \\ &= \frac{1}{2} (U_x V_z - iU_y V_z + U_z V_x - iU_z V_y), \end{aligned} \quad (19)$$

$$T_{-2}^{(2)} = \langle 11; -1, -1 | 11; 2, -2 \rangle U_{-1}^{(1)} V_{-1}^{(1)} = \frac{U_x - iU_y}{\sqrt{2}} \frac{V_x - iV_y}{\sqrt{2}} = \frac{1}{2} (U_x V_x - iU_x V_y - iU_y V_x - U_y V_y).$$

□

第 3 题 (课本习题 3.32) 得分: _____. (a) 把 xy , xz 和 $(x^2 - y^2)$ 写成一个秩为 2 的球 (不可约) 张量的分量.

(b) 期待值

$$Q = e \langle \alpha, j, m = j | (3z^2 - r^2) | \alpha, j, m = j \rangle$$

被称为 四极矩. 利用 Q 和适当的克莱布什-戈丹系数, 求

$$e \langle \alpha, j, m' | (x^2 - y^2) | \alpha, j, m = j \rangle,$$

其中 $m' = j, j-1, j-2, \dots$.

解: (a) 由

$$Y_2^0(\mathbf{r}) = \sqrt{\frac{5}{16\pi}} (3 \cos^2 \theta - 1) = \sqrt{\frac{5}{16\pi}} \left(3 \frac{z^2}{r^2} - 1 \right), \quad (20)$$

$$Y_2^{\pm 1}(\mathbf{r}) = \mp \sqrt{\frac{15}{8\pi}} \sin \theta \cos \theta e^{\pm i\phi} = \mp \sqrt{\frac{15}{8\pi}} \frac{(x \pm iy)z}{r^2}, \quad (21)$$

$$Y_2^{\pm 2}(\mathbf{r}) = \sqrt{\frac{15}{32\pi}} \sin^2 \theta e^{\pm 2i\phi} = \sqrt{\frac{15}{32\pi}} \frac{(x \pm iy)^2}{r^2}, \quad (22)$$

得

$$xy = \frac{1}{4i} \left[\sqrt{\frac{32\pi}{15}} r^2 Y_2^{+2}(\mathbf{r}) - \sqrt{\frac{32\pi}{15}} r^2 Y_2^{-2}(\mathbf{r}) \right] = i \sqrt{\frac{2\pi}{15}} r^2 [Y_2^{-2}(\mathbf{r}) - Y_2^{+2}(\mathbf{r})], \quad (23)$$

$$xz = \frac{1}{2} \left[\sqrt{\frac{8\pi}{15}} r^2 Y_2^{-1}(\mathbf{r}) - \sqrt{\frac{8\pi}{15}} r^2 Y_2^{+1}(\mathbf{r}) \right] = \sqrt{\frac{2\pi}{15}} r^2 [Y_2^{-1}(\mathbf{r}) - Y_2^{+1}(\mathbf{r})], \quad (24)$$

$$(x^2 - y^2) = \frac{1}{2} \left[\sqrt{\frac{32\pi}{15}} r^2 Y_2^{+2}(\mathbf{r}) + \sqrt{\frac{32\pi}{15}} r^2 Y_2^{-2}(\mathbf{r}) \right] = \sqrt{\frac{8\pi}{15}} r^2 [Y_2^{+2}(\mathbf{r}) + Y_2^{-2}(\mathbf{r})]. \quad (25)$$

(b) 利用前一小题的结论,

$$e \langle \alpha, j, m' | (x^2 - y^2) | \alpha, j, m = j \rangle = e \sqrt{\frac{8\pi}{15}} \langle \alpha, j, m' | r^2 [Y_2^{-2}(\mathbf{r}) + Y_2^{+2}(\mathbf{r})] | \alpha, j, m = j \rangle. \quad (26)$$

利用 Wigner-Eckart 定理,

$$\begin{aligned} e \langle \alpha, j, m' | (x^2 - y^2) | \alpha, j, m = j \rangle &= e \sqrt{\frac{8\pi}{15}} [\langle j2; j, -2 | j2; jm' \rangle + \langle j2; j, +2 | j2; jm' \rangle] \frac{\langle \alpha, j \| Y^{(2)} \| \alpha, j \rangle}{\sqrt{2j+1}} \\ &= e \sqrt{\frac{8\pi}{15}} \langle j2; j, -2 | j2; jm' \rangle \frac{\langle \alpha, j \| Y^{(2)} \| \alpha, j \rangle}{\sqrt{2j+1}} \end{aligned} \quad (27)$$

另一方面,

$$\begin{aligned} Q &= e \langle \alpha, j, m = j | (3z^2 - r^2) | \alpha, j, m = j \rangle = e \sqrt{\frac{16\pi}{5}} \langle \alpha, j, m' | Y_2^0(\mathbf{r}) | \alpha, j, m = j \rangle \\ &= e \sqrt{\frac{16\pi}{5}} \langle j2; j0 | j2; jj \rangle \frac{\langle \alpha, j \| Y^{(2)} \| \alpha, j \rangle}{\sqrt{2j+1}}. \end{aligned} \quad (28)$$

故

$$e\langle\alpha, j, m'|(x^2 - y^2)|\alpha, j, m = j\rangle = \frac{Q}{\sqrt{2}} \frac{\langle j^2; j, -2|j^2, jm'\rangle}{\langle j^2; j0|j^2, jj\rangle}. \quad (29)$$

□

第 4 题 (课本习题 3.10) 得分: _____. (a) 考虑全同制备的自旋 $\frac{1}{2}$ 系统的一个纯系综. 假定期待值 $\langle S_x \rangle$ 和 $\langle S_z \rangle$ 已知, 而 $\langle S_y \rangle$ 的符号也已知. 证明如何确定态矢量. 为什么不必知道 $\langle S_y \rangle$ 的大小?

(b) 考虑一个自旋 $\frac{1}{2}$ 系统的混合系综. 假定系综平均值 $[S_x], [S_y], [S_z]$ 都是已知的, 证明如何可以构造表征这个系综的 2×2 密度矩阵.

证: (a) 设该纯系综的状态为

$$|\alpha\rangle = \cos \frac{\beta}{2} |+\rangle + e^{i\alpha} \sin \frac{\beta}{2} |-\rangle. \quad (30)$$

期待值

$$\begin{aligned} \langle S_x \rangle &= \langle \alpha | S_x | \alpha \rangle = \left(\cos \frac{\beta}{2} \langle + | + e^{-i\alpha} \sin \frac{\beta}{2} \langle - | \right) \frac{\hbar}{2} (|+\rangle \langle -| + |-\rangle \langle +|) \left(\cos \frac{\beta}{2} |+\rangle + e^{i\alpha} \sin \frac{\beta}{2} |-\rangle \right) \\ &= \frac{\hbar}{2} \left(e^{i\alpha} \cos \frac{\beta}{2} \sin \frac{\beta}{2} + e^{-i\alpha} \sin \frac{\beta}{2} \cos \frac{\beta}{2} \right) = \frac{\hbar}{2} \cos \alpha \sin \beta, \end{aligned} \quad (31)$$

$$\begin{aligned} \langle S_z \rangle &= \langle \alpha | S_z | \alpha \rangle = \left(\cos \frac{\beta}{2} \langle + | + e^{-i\alpha} \sin \frac{\beta}{2} \langle - | \right) \frac{\hbar}{2} (|+\rangle \langle +| - |-\rangle \langle -|) \left(\cos \frac{\beta}{2} |+\rangle + e^{i\alpha} \sin \frac{\beta}{2} |-\rangle \right) \\ &= \frac{\hbar}{2} \left(\cos^2 \frac{\beta}{2} - \sin^2 \frac{\beta}{2} \right) = \frac{\hbar}{2} \cos \beta, \end{aligned} \quad (32)$$

$$\begin{aligned} \langle S_y \rangle &= \langle \alpha | S_y | \alpha \rangle = \langle \alpha | S_z | \alpha \rangle = \left(\cos \frac{\beta}{2} \langle + | + e^{-i\alpha} \sin \frac{\beta}{2} \langle - | \right) \frac{\hbar}{2} (-i|+\rangle \langle -| + i|-\rangle \langle +|) \left(\cos \frac{\beta}{2} |+\rangle + e^{i\alpha} \sin \frac{\beta}{2} |-\rangle \right) \\ &= \frac{\hbar}{2} \left(-ie^{i\alpha} \cos \frac{\beta}{2} \sin \frac{\beta}{2} + ie^{-i\alpha} \sin \frac{\beta}{2} \cos \frac{\beta}{2} \right) = \frac{\hbar}{2} \sin \alpha \sin \beta. \end{aligned} \quad (33)$$

因此由已知的 $\langle S_x \rangle$ 和 $\langle S_z \rangle$ 就可得

$$\langle S_y \rangle = \pm \sqrt{\left(\frac{\hbar}{2} \right)^2 - \langle S_x \rangle^2 - \langle S_z \rangle^2}. \quad (34)$$

当 $\langle S_y \rangle$ 符号已知, $\langle S_y \rangle$ 就可被完全确定下来. 此时, α 和 β 也被确定下来:

$$\alpha = \arctan \frac{\langle S_y \rangle}{\langle S_x \rangle}, \quad (35)$$

$$\beta = \arccos \frac{\langle S_z \rangle}{\hbar/2}, \quad (36)$$

即可确定态矢量 $|\alpha\rangle$.

(b) 设该混合系综的密度矩阵为

$$\rho = \begin{bmatrix} a & b \\ c & d \end{bmatrix}, \quad (37)$$

其满足

$$\rho = \rho^\dagger \implies b = c^*, \text{ and } a, d \in \mathbb{R}, \quad (38)$$

$$\text{Tr}(\rho) = a + d = 1. \quad (39)$$

系综平均值

$$[S_x] = \text{Tr}(\rho S_x) = \text{Tr} \left(\begin{bmatrix} a & b \\ c & d \end{bmatrix} \frac{\hbar}{2} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \right) = \frac{\hbar}{2}(b + c), \quad (40)$$

$$[S_y] = \text{Tr}(\rho S_y) = \text{Tr} \left(\begin{bmatrix} a & b \\ c & d \end{bmatrix} \frac{\hbar}{2} \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \right) = i\frac{\hbar}{2}(b - c), \quad (41)$$

$$[S_z] = \text{Tr}(\rho S_z) = \text{Tr} \left(\begin{bmatrix} a & b \\ c & d \end{bmatrix} \frac{\hbar}{2} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \right) = \frac{\hbar}{2}(a - d). \quad (42)$$

当 $[S_x], [S_y], [S_z]$ 已知, 则

$$a = \frac{1}{2} \left(1 + \frac{[S_z]}{\hbar/2} \right), \quad (43)$$

$$b = \frac{1}{\hbar} ([S_x] - i[S_y]), \quad (44)$$

$$c = \frac{1}{\hbar} ([S_x] + i[S_y]), \quad (45)$$

$$d = \frac{1}{2} \left(1 - \frac{[S_z]}{\hbar/2} \right), \quad (46)$$

由此可以构造出表征该混合系综的 2×2 密度矩阵.

□

第 5 题 (课本习题 3.11) 得分: _____. (a) 证明密度算符 (在薛定谔绘景中) 的时间演化由下式给定

$$\rho(t) = \mathcal{U}(t, t_0) \rho(t_0) \mathcal{U}^\dagger(t, t_0).$$

(b) 假定在 $t = 0$ 时有一个纯系统. 证明只要时间演化由薛定谔方程控制, 则它不可能演化成一个混合系统.

证: (a) 设 t_0 时刻密度算符为

$$\rho(t_0) = \sum_i w_i |\alpha^{(i)}, t_0\rangle \langle \alpha^{(i)}, t_0|. \quad (47)$$

在薛定谔绘景中, 任意纯态 $|\alpha^{(i)}\rangle$ 按照

$$|\alpha, t_0; t\rangle = \mathcal{U}(t, t_0) |\alpha^{(i)}, t_0\rangle \quad (48)$$

的形式演化, 其中 $\mathcal{U}(t, t_0)$ 为时刻 t_0 时刻至 t 时刻的演化算符. 故 t 时刻密度算符演化为

$$\begin{aligned} \rho(t) &= \sum_i w_i \mathcal{U}(t, t_0) |\alpha^{(i)}, t_0\rangle \langle \alpha^{(i)}, t_0| \mathcal{U}^\dagger(t, t_0) \\ &= \mathcal{U}(t, t_0) \left[\sum_i w_i |\alpha^{(i)}, t_0\rangle \langle \alpha^{(i)}, t_0| \right] \mathcal{U}^\dagger(t, t_0) \\ &= \mathcal{U}(t, t_0) \rho(t_0) \mathcal{U}^\dagger(t, t_0). \end{aligned} \quad (49)$$

(b) $t = 0$ 时刻该纯态系统的密度算符 $\rho(t = 0)$ 满足

$$\text{Tr}[\rho^2(t = 0)] = 1. \quad (50)$$

任意 t 时刻该系统的密度算符仍然满足

$$\begin{aligned}\mathrm{Tr}[\rho^2(t)] &= \mathrm{Tr}[U(t, 0)\rho(t=0)U^\dagger(t, 0)U(t, 0)\rho(t=0)U^\dagger(t, 0)] \\ &= \mathrm{Tr}[U(t, 0)\rho^2(t=0)U^\dagger(t, 0)] \\ &= \mathrm{Tr}[U^\dagger(t, 0)U(t, 0)\rho^2(t=0)] \\ &= \mathrm{Tr}[\rho^2(t=0)] \\ &= 1,\end{aligned}\tag{51}$$

即该系统仍然为一纯态系统. 因此只要时间演化由薛定谔方程控制, 则它不可能演化成一个混合系统.

□