

第 1 题 得分: _____. 计算下列系统内电子气体的 Fermi 能, Fermi 温度和 Fermi 速度.

- 1) 室温下的金属钠: 密度为 0.97 g/cm^3 , 每个原子贡献一个传导电子, 假设它们的能量和动量的关系为 $E = \frac{p^2}{2m}$, (即忽略晶格场对电子运动的影响).
- 2) 天狼星的伴星 (白矮星): 其质量约为太阳质量的 0.98 倍, 半径约为太阳半径的 0.0084 倍. 假设星体全部由氢组成.

解: 1) 电子属于 Fermion, 按照自由 Fermi 气体模型, 传导电子的数密度可用 Fermi 波矢表示,

$$\rho = \frac{\omega K_F^3}{6\pi^2}, \quad (1)$$

其中金属钠中的传导电子数密度 $\rho = \frac{N}{V} = \frac{N_A}{M_{\text{Na}}/\rho_m} = \frac{6.02 \times 10^{23} \text{ mol}^{-1}}{(22.99 \text{ g} \cdot \text{mol}^{-1})/(0.97 \times 10^6 \text{ g} \cdot \text{m}^{-3})} = 2.540 \times 10^{28} \text{ m}^{-3}$, 传导电子的各能级的简并度 $\omega = 2$, 故得 Fermi 波矢为

$$K_F = 9.09 \times 10^9 \text{ m}^{-1}. \quad (2)$$

Fermi 速度为

$$v_F = \frac{\hbar K_F}{m} = 6.62 \times 10^6 \text{ m} \cdot \text{s}^{-1}, \quad (3)$$

其中电子的质量 $m = 9.11 \times 10^{-31} \text{ kg}$. Fermi 能量为

$$\varepsilon_F = \frac{\hbar^2 K_F^2}{2m} = 5.05 \times 10^{-19} \text{ J} = 3.16 \text{ eV}. \quad (4)$$

Fermi 温度为

$$T_F = \frac{\varepsilon_F}{k} = 3.66 \times 10^4 \text{ K}. \quad (5)$$

- 2) 太阳的质量为 $m_\odot = 1.989 \times 10^{30} \text{ kg}$, 太阳的半径为 $r_\odot = 6.6934 \times 10^8 \text{ m}$, 该白矮星的质量密度为 $\rho_m = \frac{m}{V} = \frac{0.98 m_\odot}{\frac{4}{3}\pi(0.0084 r_\odot)^3} = 2.62 \times 10^9 \text{ kg} \cdot \text{m}^{-3}$. 假设白矮星中电子全部从原子内挤出, 每个氢原子挤出 $\chi = 2$ 个电子, 白矮星的电子数密度为 $\rho = \frac{\chi N_A}{M_{\text{He}}/\rho_m} = \frac{2 \times 6.02 \times 10^{23} \text{ mol}^{-1}}{(4.00 \times 10^{-3} \text{ kg} \cdot \text{mol}^{-1})/(2.62 \times 10^9 \text{ kg} \cdot \text{m}^{-3})} = 7.89 \times 10^{35} \text{ m}^{-3}$. 对于氢原子 Fermi 波矢为

$$K_F = \left(\frac{6\pi^2 \rho}{\omega} \right)^{1/3} = 2.86 \times 10^{12} \text{ m}^{-1}. \quad (6)$$

Fermi 能量为

$$\varepsilon_F = c\sqrt{p_F^2 + m^2 c^2} - mc^2 = c\sqrt{\hbar^2 k_F^2 + m^2 c^2} - mc^2 = 4.01 \times 10^{-14} \text{ J} = 2.51 \times 10^5 \text{ eV}. \quad (7)$$

Fermi 速度为

$$v_F = \frac{p_F c^2}{\varepsilon_F^2 + mc^2} = \frac{c^2 \hbar K_F}{\varepsilon_F + mc^2} = 2.22 \times 10^8 \text{ m} \cdot \text{s}^{-1}. \quad (8)$$

Fermi 温度为

$$T_F = \frac{\varepsilon_F}{k} = 2.91 \times 10^9 \text{ K}. \quad (9)$$

□

第 2 题 得分: _____. 证明非相对论简并电子气体的热力学函数是

$$\begin{aligned} G &= N\mu = N\varepsilon_F \left[1 - \frac{1}{12}\pi^2 \left(\frac{kT}{\varepsilon_F} \right)^2 - \frac{1}{80}\pi^4 \left(\frac{kT}{\varepsilon_F} \right)^4 + \dots \right] \\ E &= \frac{3}{5}N\varepsilon_F \left[1 + \frac{5}{12}\pi^2 \left(\frac{kT}{\varepsilon_F} \right)^2 - \frac{1}{16}\pi^4 \left(\frac{kT}{\varepsilon_F} \right)^4 + \dots \right] \\ C_V &= \frac{1}{2}N\pi^2 \frac{k^2 T}{\varepsilon_F} \left[1 - \frac{3}{10}\pi^2 \left(\frac{kT}{\varepsilon_F} \right)^2 + \dots \right] \\ S &= \frac{1}{2}N\pi^2 \frac{k^2 T}{\varepsilon_F} \left[1 - \frac{1}{10} \left(\frac{kT}{\varepsilon_F} \right)^2 + \dots \right] \end{aligned}$$

其中 $N = \frac{V}{3\pi^2} \left(\frac{2m}{\hbar} \varepsilon_F \right)^{\frac{3}{2}}$ 为电子总数, \dots 代表 $\frac{kT}{\varepsilon_F}$ 的高阶项.

证:

Sommerfeld 方法:

$$\int_0^\infty d\varepsilon \frac{f(\varepsilon)}{e^{\beta(\varepsilon-\mu)} + 1} \approx 2 \sum_{n=0}^\infty F^{(2n)} c_n (kT)^{2n}, \quad (10)$$

其中 $F(\varepsilon)$ 为 $f(\varepsilon)$ 的原函数,

$$c_n = \frac{1}{(2n)!} \int_0^\infty \frac{\xi^{2n} e^\xi}{(e^\xi + 1)^2}. \quad (11)$$

$$c_0 = \frac{1}{2}, \quad c_1 = \frac{\pi^2}{12}, \quad c_2 = \frac{7\pi^2}{240}.$$

系统中的电子数为

$$N = V \int_0^\infty d\varepsilon \frac{D(\varepsilon)}{e^{\beta(\varepsilon-\mu)} + 1}, \quad (12)$$

其中对于非相对论自由电子, 态密度

$$D(\varepsilon) = \frac{\sqrt{2}m^{3/2}\varepsilon^{1/2}}{\pi^2\hbar^3}, \quad (13)$$

从而对于上面这一积分式,

$$f(\varepsilon) = D(\varepsilon), \quad (14)$$

$$\implies F(\varepsilon) = \frac{2\sqrt{2}m^{3/2}\varepsilon^{3/2}}{3\pi^2\hbar^3}, \quad (15)$$

$$F^{(2)}(\varepsilon) = D'(\varepsilon) = \frac{\sqrt{2}m^{3/2}\varepsilon^{-1/2}}{2\pi^2\hbar^3}, \quad (16)$$

$$F^{(4)}(\varepsilon) = \frac{3\sqrt{2}m^{3/2}\varepsilon^{-5/2}}{8\pi^2\hbar^3}. \quad (17)$$

从而利用 Sommerfeld 方法, 得

$$N = \frac{V}{3\pi^2} \left(\frac{2m\mu}{\hbar^2} \right)^{3/2} \left[1 + \frac{\pi^2}{8} \left(\frac{kT}{\mu} \right)^2 + \frac{7\pi^4}{640} \left(\frac{kT}{\mu} \right)^4 + \dots \right], \quad (18)$$

$$\implies \rho = \frac{N}{V} = \frac{1}{3\pi^2} \left(\frac{2m\mu}{\hbar^2} \right)^{3/2} \left[1 + \frac{\pi^2}{8} \left(\frac{kT}{\mu} \right)^2 + \frac{7\pi^4}{640} \left(\frac{kT}{\mu} \right)^4 + \dots \right]. \quad (19)$$

由于

$$\rho = \frac{\omega K_F^3}{6\pi^2} \stackrel{\omega=2}{=} \frac{K_F^3}{3\pi^2} = \frac{(2m\varepsilon_F/\hbar^2)^{3/2}}{3\pi^2}, \quad (20)$$

故有

$$\varepsilon_F^{3/2} = \mu^{3/2} \left[1 + \frac{\pi^2}{8} \left(\frac{kT}{\mu} \right)^2 + \frac{7\pi^4}{640} \left(\frac{kT}{\mu} \right)^4 + \cdots \right], \quad (21)$$

$$\begin{aligned} \Rightarrow \varepsilon_F &= \mu \left[1 + \frac{\pi^2}{8} \left(\frac{kT}{\mu} \right)^2 + \frac{7\pi^4}{640} \left(\frac{kT}{\mu} \right)^4 + \cdots \right]^{2/3} \\ &= \mu \left\{ 1 + \frac{2}{3} \left[\frac{\pi^2}{8} \left(\frac{kT}{\mu} \right)^2 + \frac{7\pi^4}{640} \left(\frac{kT}{\mu} \right)^4 + \cdots \right] + \frac{1}{2!} \frac{-1}{3} \frac{2}{3} \left[\frac{\pi^2}{8} \left(\frac{kT}{\mu} \right)^2 + \frac{7\pi^4}{640} \left(\frac{kT}{\mu} \right)^4 + \cdots \right]^2 + \cdots \right\} \\ &= \mu \left[1 + \frac{\pi^2}{12} \left(\frac{kT}{\mu} \right)^2 + \frac{\pi^4}{180} \left(\frac{kT}{\mu} \right)^4 + \cdots \right], \end{aligned} \quad (22)$$

$$\Rightarrow \mu = \varepsilon_F \left[1 + \frac{\pi^2}{12} \left(\frac{kT}{\mu} \right)^2 + \frac{\pi^4}{180} \left(\frac{kT}{\mu} \right)^4 + \cdots \right]^{-1}.$$

近似到 0 阶得

$$\mu = \varepsilon_F. \quad (23)$$

将上式回代，然后近似到 2 阶得

$$\mu = \varepsilon_F - \frac{\pi^2}{12\varepsilon_F} (kT)^2. \quad (24)$$

将上式回代，然后近似到 4 阶得

$$\mu = \varepsilon_F - \frac{\pi^2}{12\varepsilon_F \left[1 - \frac{\pi^2}{12} \left(\frac{kT}{\varepsilon_F} \right)^2 \right]} (kT)^2 - \frac{\pi^4}{180\varepsilon_F^4} (kT)^4 \approx \varepsilon_F - \frac{\pi^2}{12\varepsilon_F} (kT)^2 - \frac{1}{80\varepsilon_F^3} (kT)^4. \quad (25)$$

故 Gibbs 势为

$$G = N\mu = N\varepsilon \left[1 - \frac{1}{12} \pi^2 \left(\frac{kT}{\varepsilon_F} \right)^2 - \frac{1}{80} \pi^4 \left(\frac{kT}{\varepsilon_F} \right)^4 + \cdots \right]. \quad (26)$$

系统内能为

$$E = V \int_0^\infty d\varepsilon \frac{D(\varepsilon)}{e^{-\beta(\varepsilon-\mu)} + 1} \varepsilon. \quad (27)$$

从而对上面这一积分式，

$$f(\varepsilon) = \varepsilon D(\varepsilon) = \frac{\sqrt{2} m^{3/2} \varepsilon^{3/2}}{\pi^2 \hbar^3}, \quad (28)$$

$$\Rightarrow F(\varepsilon) = \frac{2\sqrt{2} m^{3/2} \varepsilon^{5/2}}{5\pi^2 \hbar^3}, \quad (29)$$

$$F^{(2)}(\varepsilon) = \frac{3\sqrt{2} m^{3/2} \varepsilon^{1/2}}{2\pi^2 \hbar^3}, \quad (30)$$

$$F^{(4)}(\varepsilon) = -\frac{3\sqrt{2} m^{3/2} \varepsilon^{-3/2}}{8\pi^2 \hbar^3}. \quad (31)$$

从而利用 Sommerfeld 方法，得

$$E = V \left(\frac{2m\mu}{\hbar^2} \right)^{3/2} \frac{1}{5\pi^2} \mu \left[1 + \frac{5\pi^2}{8} \left(\frac{kT}{\mu} \right)^2 - \frac{7\pi^4}{384} \left(\frac{kT}{\mu} \right)^4 + \cdots \right]. \quad (32)$$

根据式 (18),

$$\begin{aligned}
 V \left(\frac{2m\mu}{\hbar^2} \right)^{3/2} &= 3\pi^2 N \left[1 + \frac{\pi^2}{8} \left(\frac{kT}{\mu} \right)^2 + \frac{7\pi^2}{640} \left(\frac{kT}{\mu} \right)^4 + \dots \right]^{-1} \\
 &= 3\pi^2 N \left\{ 1 - \left[\frac{\pi^2}{8} \left(\frac{kT}{\mu} \right)^2 + \frac{7\pi^2}{640} \left(\frac{kT}{\mu} \right)^4 \right] + \frac{1}{2!}(-1)(-2) \left[\frac{\pi^2}{8} \left(\frac{kT}{\mu} \right)^2 + \frac{7\pi^2}{640} \left(\frac{kT}{\mu} \right)^4 + \dots \right]^2 + \dots \right\} \\
 &= 3\pi^2 N \left[1 - \frac{\pi^2}{8} \left(\frac{kT}{\mu} \right)^2 + \frac{3\pi^2}{640} \left(\frac{kT}{\mu} \right)^4 + \dots \right]. \quad (33)
 \end{aligned}$$

将上式代入式 (32), 得

$$\begin{aligned}
 E &= \frac{3}{5} N \mu \left[1 + \frac{5\pi^2}{8} \left(\frac{kT}{\mu} \right)^2 - \frac{7\pi^4}{384} \left(\frac{kT}{\mu} \right)^4 + \dots \right] \left[1 - \frac{\pi^2}{8} \left(\frac{kT}{\mu} \right)^2 + \frac{3\pi^2}{640} \left(\frac{kT}{\mu} \right)^4 + \dots \right] \\
 &= \frac{3}{5} N \mu \left[1 + \frac{\pi^2}{2} \left(\frac{kT}{\mu} \right)^2 - \frac{19\pi^4}{384} \left(\frac{kT}{\mu} \right)^4 + \dots \right] \\
 &= \frac{3}{5} N \left[\mu + \frac{\pi^2}{2\mu} (kT)^2 - \frac{11\pi^4}{120\mu^3} (kT)^4 + \dots \right]. \quad (34)
 \end{aligned}$$

将式 (25) 代入上式, 并近似至 4 阶, 得

$$\begin{aligned}
 E &= \frac{3}{5} N \left\{ \varepsilon_F \left[1 - \frac{\pi^2}{12} \left(\frac{kT}{\varepsilon_F} \right)^2 - \frac{\pi^4}{80} \left(\frac{kT}{\varepsilon_F} \right)^4 \right] + \frac{\pi^2}{2\varepsilon_F \left[1 - \frac{\pi^2}{12} \left(\frac{kT}{\varepsilon_F} \right)^2 - \dots \right]} (kT)^2 - \frac{11\pi^4}{120\varepsilon_F^3 [1 - \dots]^3} (kT)^4 \right\} \\
 &= \frac{3}{5} N \left[\varepsilon_F + \frac{5\pi^2}{12\varepsilon_F} (kT)^2 - \frac{\pi^4}{16\varepsilon_F^3} (kT)^4 + \dots \right] \\
 &= \frac{3}{5} N \varepsilon_F \left[1 + \frac{5\pi^2}{12} \left(\frac{kT}{\varepsilon_F} \right)^2 - \frac{\pi^4}{16} \left(\frac{kT}{\varepsilon_F} \right)^4 + \dots \right]. \quad (35)
 \end{aligned}$$

系统的定容热容量为

$$C_V = \left(\frac{\partial E}{\partial T} \right)_V = \frac{1}{2} N \pi^2 \frac{k^2 T}{\varepsilon_F} \left[1 - \frac{3}{10} \left(\frac{kT}{\varepsilon_F} \right)^2 + \dots \right] \quad (36)$$

对于非相对论电子气体,

$$PV = \frac{2}{3} E. \quad (37)$$

系统的自由能为

$$F = G - PV = G - \frac{2}{3} E = N \varepsilon_F \left[\frac{3}{5} - \frac{\pi^2}{4} \left(\frac{kT}{\varepsilon_F} \right)^2 + \frac{\pi^4}{80} \left(\frac{kT}{\varepsilon_F} \right)^4 + \dots \right]. \quad (38)$$

系统的熵为

$$S = - \frac{\partial G}{\partial T} = \frac{\pi^2}{2} N \frac{k^2 T}{\varepsilon_F} \left[1 - \frac{1}{10} \left(\frac{kT}{\varepsilon_F} \right)^2 + \dots \right]. \quad (39)$$

□