

第 1 题 得分: \_\_\_\_\_. 计算 Maxwell 分布下的最可几速度, 平均速度和速度分布的宽度, 即

$$\Delta v \equiv \sqrt{\langle (v - \langle v \rangle)^2 \rangle} \quad (1)$$

将结果用绝对温度和分子质量表达, 并算出氢气和氧气以上各量在室温下的数值.

解: 麦克斯韦速度分布律为

$$f(v) = \left(\frac{m}{2\pi kT}\right)^{3/2} 4\pi v^2 e^{-\frac{mv^2}{2kT}}. \quad (2)$$

令上式的导数等于零,

$$f'(v_0) = \left(\frac{m}{2\pi kT}\right)^{3/2} 4\pi v_0 e^{-\frac{mv_0^2}{2kT}} \left(2 - \frac{mv_0^2}{kT}\right) = 0, \quad (3)$$

得最可几速度为

$$v_0 = \sqrt{\frac{2kT}{m}}. \quad (4)$$

平均速度为

$$\langle v \rangle = \int_0^\infty dv v f(v) = \left(\frac{m}{2\pi kT}\right)^{3/2} 4\pi \int_0^\infty dv v^3 e^{-\frac{mv^2}{2kT}} = \left(\frac{m}{2\pi kT}\right)^{3/2} 4\pi \times \frac{2k^2 T^2}{m^2} = \sqrt{\frac{8kT}{\pi m}}.$$

速度分布的宽度可表为

$$\Delta v = \sqrt{\langle v^2 \rangle - \langle v \rangle^2}, \quad (5)$$

其中速率的平方平均为

$$\langle v^2 \rangle = \int_0^\infty dv v^2 f(v) = \left(\frac{m}{2\pi kT}\right)^{3/2} 4\pi \int_0^\infty dv v^4 e^{-\frac{mv^2}{2kT}} = \left(\frac{m}{2\pi kT}\right)^{3/2} 4\pi \times \frac{3\sqrt{\pi}}{8 \left(\frac{m}{2kT}\right)^{5/2}} = \frac{3kT}{m}, \quad (6)$$

故速度分布的宽度为

$$\Delta v = \sqrt{\left(3 - \frac{8}{\pi}\right) \frac{kT}{m}}. \quad (7)$$

氢气分子的质量为  $\frac{2 \times 10}{6.02 \times 10^{23}}$  kg, 在室温 300 K 下最可几速度为

$$v_{0, \text{H}_2} = 1.58 \times 10^3 \text{ m} \cdot \text{s}^{-1} = 1.58 \text{ km} \cdot \text{s}^{-1}, \quad (8)$$

平均速度为

$$\langle v_{\text{H}_2} \rangle = 1.78 \times 10^3 \text{ m} \cdot \text{s}^{-1} = 1.78 \text{ km} \cdot \text{s}^{-1}, \quad (9)$$

速度分布的宽度为

$$\Delta v_{\text{H}_2} = 752 \text{ m} \cdot \text{s}^{-1}. \quad (10)$$

氧气分子的质量为  $\frac{32 \times 10^{-3}}{6.02 \times 10^{23}}$  kg, 在室温下最可几速度为

$$v_{0, \text{O}_2} = 395 \text{ m} \cdot \text{s}^{-1}, \quad (11)$$

平均速度为

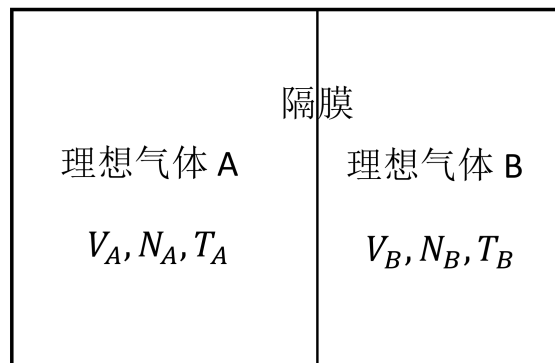
$$\langle v_{\text{O}_2} \rangle = 445 \text{ m} \cdot \text{s}^{-1}, \quad (12)$$

速度分布的宽度为

$$\Delta v_{\text{O}_2} = 187 \text{ m} \cdot \text{s}^{-1}. \quad (13)$$

□

**第 2 题 得分:** \_\_\_\_\_. 由同一种分子组成的理想气体被隔膜分成两部分. 每部分的体积, 粒子数和温度如右图所示. 假设隔膜左右温度相同, 但密度不同, 且系统与外界热绝缘. 证明隔膜撤掉后气体的熵增加.



**证:** 隔膜撤掉后气体的熵的改变量为

$$\begin{aligned}\Delta S &= S_2 - S_1 = S_2 - S_{1A} - S_{1B} = (N_A + N_B)k \ln \frac{V_A + V_B}{N_A + N_B} - N_A k \ln \frac{V_A}{N_A} - N_B k \ln \frac{V_B}{N_B} \\ &= (N_A + N_B)k \left[ \ln \frac{V_A + V_B}{N_A + N_B} - \frac{N_A}{N_A + N_B} \ln \frac{V_A}{N_A} - \frac{N_B}{N_A + N_B} \ln \frac{V_B}{N_B} \right].\end{aligned}\quad (14)$$

函数

$$f(x) = \ln x \quad (15)$$

的二阶导

$$f''(x) = -\frac{1}{x^2} < 0, \quad (16)$$

故  $f(x)$  为上凸函数. 式 (14) 中括号内的部分可写作

$$\begin{aligned}& \ln \frac{V_A + V_B}{N_A + N_B} - \frac{N_A}{N_A + N_B} \ln \frac{V_A}{N_A} - \frac{N_B}{N_A + N_B} \ln \frac{V_B}{N_B} \\ &= f\left(\frac{N_A}{N_A + N_B} \frac{V_A}{N_A} + \frac{N_B}{N_A + N_B} \frac{V_B}{N_B}\right) - \frac{N_A}{N_A + N_B} f\left(\frac{V_A}{N_A}\right) - \frac{N_B}{N_A + N_B} f\left(\frac{V_B}{N_B}\right).\end{aligned}\quad (17)$$

其中  $\frac{N_A}{N_A + N_B} + \frac{N_B}{N_A + N_B} = 1$ . 根据上凸函数的性质,

$$f\left(\frac{N_A}{N_A + N_B} \frac{V_A}{N_A} + \frac{N_B}{N_A + N_B} \frac{V_B}{N_B}\right) > \frac{N_A}{N_A + N_B} f\left(\frac{V_A}{N_A}\right) + \frac{N_B}{N_A + N_B} f\left(\frac{V_B}{N_B}\right), \quad (18)$$

$$\Rightarrow \ln \frac{V_A + V_B}{N_A + N_B} - \frac{N_A}{N_A + N_B} \ln \frac{V_A}{N_A} - \frac{N_B}{N_A + N_B} \ln \frac{V_B}{N_B} \geq 0. \quad (19)$$

从而

$$\Delta S > 0, \quad (20)$$

即隔膜撤掉后气体的熵增加. □

**第 3 题 得分:** \_\_\_\_\_. 在高温条件下, 即

$$\varepsilon \equiv \frac{\hbar^2}{2IkT} \ll 1 \quad (21)$$

证明双原子气体的转动配分函数

$$q_r = \omega \left[ \frac{1}{\varepsilon} + \frac{1}{3} + \frac{1}{15}\varepsilon + O(\varepsilon^2) \right] \quad (22)$$

和每个分子平均转动动能

$$u_r = kT \left[ 1 - \frac{1}{3}\varepsilon - \frac{1}{45}\varepsilon^2 + O(\varepsilon^3) \right] \quad (23)$$

其中  $\omega$  为核自旋的简并度.

$$\omega = \begin{cases} (2s_A + 1)(2s_B + 1), & AB\text{型} \\ \frac{1}{2}(2s_A + 1)^2, & AA\text{型} \end{cases} \quad (24)$$

提示: 可用 Euler-Maclaurin 公式计算修正项.

证: 对于  $AB$  型双原子分子, 转动配分函数为

$$q_r = (2s_A + 1)(2s_B + 1) \sum_{j=0}^{\infty} (2j + 1) e^{-\varepsilon j(j+1)}. \quad (25)$$

令

$$f(x) = (2x + 1) e^{-\varepsilon x(x+1)}, \quad (26)$$

则

$$\int_0^{\infty} f(x) dx = \int_0^{\infty} (2x + 1) e^{-\varepsilon x(x+1)} dx = -\frac{1}{\varepsilon} e^{-\varepsilon x(x+1)} \Big|_0^{\infty} = \frac{1}{\varepsilon}, \quad (27)$$

$$f(0) = 1, \quad (28)$$

$$f(\infty) = 0, \quad (29)$$

$$f^{(1)}(x) = [2 - \varepsilon(2x + 1)^2] e^{-\varepsilon x(x+1)}, \quad (30)$$

$$f^{(1)}(0) = 2 - \varepsilon, \quad (31)$$

$$f^{(1)}(\infty) = 0, \quad (32)$$

$$f^{(3)}(x) = [-\varepsilon^3(2x + 1)^4 + 12\varepsilon^2(2x + 1)^2 - 12\varepsilon] e^{-\varepsilon x(x+1)}, \quad (33)$$

$$f^{(3)}(0) = -\varepsilon^3 + 12\varepsilon^2 - 12\varepsilon, \quad (34)$$

$$f^{(3)}(\infty) = 0, \quad (35)$$

利用 Euler-Maclaurin 公式可得

$$\begin{aligned} \sum_{j=0}^{\infty} (2j + 1) e^{-\varepsilon j(j+1)} &= \int_0^{\infty} f(x) dx + \frac{1}{2}[f(0) - f(\infty)] + (-1)^1 \frac{1}{(2 \times 1)!} \frac{1}{6} [f^{(1)}(0) - f^{(1)}(\infty)] \\ &\quad + (-1)^2 \frac{1}{(2 \times 2)!} \frac{1}{30} [f^{(3)}(0) - f^{(3)}(\infty)] + O(\varepsilon^2) \\ &= \frac{1}{\varepsilon} + \frac{1}{3} + \frac{1}{15}\varepsilon. \end{aligned} \quad (36)$$

从而转动配分函数为

$$q_r = (2s_A + 1)(2s_B + 1) \left[ \frac{1}{\varepsilon} + \frac{1}{3} + \frac{1}{15}\varepsilon + O(\varepsilon^2) \right]. \quad (37)$$

对  $AA$  型双原子分子, 只需将其核自旋的简并度换为  $\frac{1}{2}(2s_A + 1)^2$ , 转动配分函数为

$$q_r = (2s_A + 1) \sum_{j=0}^{\infty} (2j + 1) e^{-\varepsilon j(j+1)} = \frac{1}{2} (2s_A + 1) \left[ \frac{1}{\varepsilon} + \frac{1}{3} + \frac{1}{15} \varepsilon + O(\varepsilon^2) \right]. \quad (38)$$

从而一般的转动配分函数可写为

$$q_r = \omega \left[ \frac{1}{\varepsilon} + \frac{1}{3} + \frac{1}{15} \varepsilon + O(\varepsilon^2) \right], \quad (39)$$

其中核自旋简并度

$$\omega = \begin{cases} (2s_A + 1)(2s_B + 1), & AB \text{型} \\ \frac{1}{2}(2s_A + 1)^2, & AA \text{型} \end{cases}. \quad (40)$$

每个分子的平均转动动能为

$$u_r = -\frac{\partial}{\partial \beta} \ln q_r, \quad (41)$$

其中

$$\ln q_r = \ln \omega \left[ \frac{1}{\varepsilon} + \frac{1}{3} + \frac{1}{15} \varepsilon + O(\varepsilon^2) \right] = \ln \omega - \ln \varepsilon + \ln \left( 1 + \frac{\varepsilon}{3} + \frac{1}{15} \varepsilon^2 + O(\varepsilon^3) \right) = \ln \omega - \ln \varepsilon + \frac{\varepsilon}{3} + \frac{\varepsilon^2}{90} + O(\varepsilon^3), \quad (42)$$

$$\frac{\partial}{\partial \beta} \ln q_r = \left[ -\frac{1}{\varepsilon} + \frac{1}{3} + \frac{\varepsilon}{45} + O(\varepsilon^2) \right] \frac{\partial \varepsilon}{\partial \beta}, \quad (43)$$

$$\frac{\partial \varepsilon}{\partial \beta} = \frac{\partial \left( \frac{\hbar^2 \beta}{2I} \right)}{\partial \beta} = \frac{\hbar^2}{2I}, \quad (44)$$

故每个分子的平均转动动能为

$$u_r = \frac{\hbar^2}{2I} \left[ \frac{1}{\varepsilon} - \frac{1}{3} - \frac{\varepsilon}{45} + O(\varepsilon^2) \right] = kT \left[ 1 - \frac{\varepsilon}{3} - \frac{\varepsilon^2}{45} + O(\varepsilon^3) \right]. \quad (45)$$

□