

**第 1 题 得分:** \_\_\_\_\_. 试证明相对熵纠缠度量在纯态情况下和 Von Neumann 熵是等价的. (求任意给定纯态  $|\psi_{AB}\rangle$  和任意混合态  $\sum_i p_i \rho_i \otimes \sigma_i$  中的最小相对熵  $S(|\psi_{AB}\rangle\langle\psi_{AB}||\sum_i p_i \rho_i \otimes \sigma_i)$ ).

证: 记

$$\rho_{AB} = |\psi_{AB}\rangle\langle\psi_{AB}|. \quad (1)$$

纯态  $|\psi_{AB}\rangle$  和混合态  $\sum_i p_i \rho_i \otimes \sigma_i$  的相对熵为

$$\begin{aligned} S\left(\rho_{AB} \parallel \sum_i p_i \rho_i \otimes \sigma_i\right) &= \text{Tr} \left( \rho_{AB} \left( \log_2 \rho_{AB} - \log_2 \sum_i p_i \rho_i \otimes \sigma_i \right) \right) \\ &= \text{Tr}(\rho_{AB} \log_2 \rho_{AB}) - \text{Tr} \left( \rho_{AB} \log_2 \sum_i p_i \rho_i \otimes \sigma_i \right) \\ &= -S(\rho_{AB}) - \text{Tr} \left( \rho_{AB} \log_2 \sum_i p_i \rho_i \otimes \sigma_i \right). \end{aligned} \quad (2)$$

而  $|\psi_{AB}\rangle$  的 Von Neumann 熵为

$$S(\rho_{AB} \parallel \rho_A \otimes \rho_B) = S(\rho_A) + S(\rho_B) - S(\rho_{AB}). \quad (3)$$

要证  $\min_{\sum_i p_i \rho_i \otimes \sigma_i} S(|\psi_{AB}\rangle\langle\psi_{AB}||\sum_i p_i \rho_i \otimes \sigma_i) = S(|\psi_{AB}\rangle\langle\psi_{AB}|)$ , 即证  $\min \text{Tr}(\rho_{AB} \log_2 p_i \rho_i \otimes \sigma_i) = S(\rho_A) + S(\rho_B)$ .

对  $|\psi_{AB}\rangle$  进行 Schmidt 分解得

$$|\psi_{AB}\rangle = \sum_i a_i |\psi_i\rangle_A |\psi_i\rangle_B, \quad (4)$$

其中  $\sum_i |a_i|^2 = 1$ ,  $\{|\psi_i\rangle_A\}$  和  $\{|\psi_i\rangle_B\}$  分别是子系统 A 和 B 的正交归一基. □

**第 2 题 得分:** \_\_\_\_\_. 计算混合量子态  $\rho = p|\phi^+\rangle\langle\phi^+| + \frac{1-p}{4}I_{4 \times 4}$  的纠缠 concurrence, 其中  $0 \leq p \leq 1$ ,  $|\phi^+\rangle$  是 Bell 态.

解: 混合量子态为

$$\begin{aligned} \rho &= p|\phi^+\rangle\langle\phi^+| + \frac{1-p}{4}I_{4 \times 4} \\ &= p \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle) \frac{1}{\sqrt{2}}(\langle 00| + \langle 11|) + \frac{1-p}{4}I_{4 \times 4} \\ &= \frac{p}{2}(|00\rangle\langle 00| + |00\rangle\langle 11| + |11\rangle\langle 00| + |11\rangle\langle 11|) + \frac{1-p}{4}I_{4 \times 4} \\ &= \frac{p}{2} \begin{pmatrix} 1 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 1 \end{pmatrix} + \frac{1-p}{4} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} = \begin{pmatrix} \frac{1+p}{4} & 0 & 0 & \frac{p}{2} \\ 0 & \frac{1-p}{4} & 0 & 0 \\ 0 & 0 & \frac{1-p}{4} & 0 \\ \frac{p}{2} & 0 & 0 & \frac{1+p}{4} \end{pmatrix}. \end{aligned} \quad (5)$$

取该密度矩阵的复共轭, 并用泡利算符  $\sigma_y$  分别作用两个 qubit, 得

$$\begin{aligned} \tilde{\rho} &= (\sigma_y \otimes \sigma_y) \rho^* (\sigma_y \otimes \sigma_y) \\ &= \left( \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \otimes \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \right) \begin{pmatrix} \frac{1+p}{4} & 0 & 0 & \frac{p}{2} \\ 0 & \frac{1-p}{4} & 0 & 0 \\ 0 & 0 & \frac{1-p}{4} & 0 \\ \frac{p}{2} & 0 & 0 & \frac{1+p}{4} \end{pmatrix} \left( \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \otimes \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \right) \end{aligned}$$

$$\begin{aligned}
&= \begin{pmatrix} 0 & 0 & 0 & -1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ -1 & 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} \frac{1+p}{4} & 0 & 0 & \frac{p}{2} \\ 0 & \frac{1-p}{4} & 0 & 0 \\ 0 & 0 & \frac{1-p}{4} & 0 \\ \frac{p}{2} & 0 & 0 & \frac{1+p}{4} \end{pmatrix} \begin{pmatrix} 0 & 0 & 0 & -1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ -1 & 0 & 0 & 0 \end{pmatrix} \\
&= \begin{pmatrix} \frac{1+p}{4} & 0 & 0 & \frac{p}{2} \\ 0 & \frac{1-p}{4} & 0 & 0 \\ 0 & 0 & \frac{1-p}{4} & 0 \\ \frac{p}{2} & 0 & 0 & \frac{1+p}{4} \end{pmatrix}.
\end{aligned} \tag{6}$$

对原混合量子态的密度矩阵进行奇异值分解得

$$\rho = \begin{pmatrix} \frac{1+p}{4} & 0 & 0 & \frac{p}{2} \\ 0 & \frac{1-p}{4} & 0 & 0 \\ 0 & 0 & \frac{1-p}{4} & 0 \\ \frac{p}{2} & 0 & 0 & \frac{1+p}{4} \end{pmatrix} = \begin{pmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} & 0 & 0 \end{pmatrix} \begin{pmatrix} \frac{1+3p}{4} & 0 & 0 & 0 \\ 0 & \frac{1-p}{4} & 0 & 0 \\ 0 & 0 & \frac{1-p}{4} & 0 \\ 0 & 0 & 0 & \frac{1-p}{4} \end{pmatrix} \begin{pmatrix} \frac{1}{\sqrt{2}} & 0 & 0 & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & 0 & 0 & -\frac{1}{\sqrt{2}} \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \end{pmatrix} \tag{7}$$

故对其开方得

$$\sqrt{\rho} = \begin{pmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} & 0 & 0 \end{pmatrix} \begin{pmatrix} \frac{\sqrt{1+3p}}{2} & 0 & 0 & 0 \\ 0 & \frac{\sqrt{1-p}}{2} & 0 & 0 \\ 0 & 0 & \frac{\sqrt{1-p}}{2} & 0 \\ 0 & 0 & 0 & \frac{\sqrt{1-p}}{2} \end{pmatrix} \begin{pmatrix} \frac{1}{\sqrt{2}} & 0 & 0 & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & 0 & 0 & -\frac{1}{\sqrt{2}} \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \end{pmatrix} \tag{8}$$

$R$  矩阵为

$$\begin{aligned}
R &= \sqrt{\sqrt{\rho}\rho\sqrt{\rho}} \\
&= \sqrt{\begin{pmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} & 0 & 0 \end{pmatrix} \begin{pmatrix} \frac{\sqrt{1+3p}}{2} & 0 & 0 & 0 \\ 0 & \frac{\sqrt{1-p}}{2} & 0 & 0 \\ 0 & 0 & \frac{\sqrt{1-p}}{2} & 0 \\ 0 & 0 & 0 & \frac{\sqrt{1-p}}{2} \end{pmatrix} \begin{pmatrix} \frac{1}{\sqrt{2}} & 0 & 0 & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & 0 & 0 & -\frac{1}{\sqrt{2}} \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \end{pmatrix} \begin{pmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} & 0 & 0 \end{pmatrix} \begin{pmatrix} \frac{\sqrt{1+3p}}{2} & 0 & 0 & 0 \\ 0 & \frac{\sqrt{1-p}}{2} & 0 & 0 \\ 0 & 0 & \frac{\sqrt{1-p}}{2} & 0 \\ 0 & 0 & 0 & \frac{\sqrt{1-p}}{2} \end{pmatrix} \begin{pmatrix} \frac{1}{\sqrt{2}} & 0 & 0 & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & 0 & 0 & -\frac{1}{\sqrt{2}} \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \end{pmatrix}} \\
&= \sqrt{\begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} & 0 & 0 \end{pmatrix} \begin{pmatrix} (\frac{1+3p}{4})^2 & 0 & 0 & 0 \\ 0 & (\frac{1-p}{4})^2 & 0 & 0 \\ 0 & 0 & (\frac{1-p}{4})^2 & 0 \\ 0 & 0 & 0 & (\frac{1-p}{4})^2 \end{pmatrix} \begin{pmatrix} \frac{1}{\sqrt{2}} & 0 & 0 & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & 0 & 0 & -\frac{1}{\sqrt{2}} \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \end{pmatrix}} \\
&= \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} & 0 & 0 \end{pmatrix} \begin{pmatrix} \frac{1+3p}{4} & 0 & 0 & 0 \\ 0 & \frac{1-p}{4} & 0 & 0 \\ 0 & 0 & \frac{1-p}{4} & 0 \\ 0 & 0 & 0 & \frac{1-p}{4} \end{pmatrix} \begin{pmatrix} \frac{1}{\sqrt{2}} & 0 & 0 & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & 0 & 0 & -\frac{1}{\sqrt{2}} \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \end{pmatrix} \\
&= \begin{pmatrix} \frac{1+p}{4} & 0 & 0 & \frac{p}{2} \\ 0 & \frac{1-p}{4} & 0 & 0 \\ 0 & 0 & \frac{1-p}{4} & 0 \\ \frac{p}{2} & 0 & 0 & \frac{1+p}{4} \end{pmatrix}.
\end{aligned} \tag{9}$$

$R$  的本征值从大到小依次为  $\lambda_1 = \frac{1+3p}{4}$ ,  $\lambda_2 = \frac{1-p}{4}$ ,  $\lambda_3 = \frac{1-p}{4}$ ,  $\lambda_4 = \frac{1-p}{4}$ , 故 concurrence 为

$$C(\rho) = \max\{0, \lambda_1 - \lambda_2 - \lambda_3 - \lambda_4\} = \max\{0, \frac{3p-1}{2}\} = \begin{cases} 0, & 0 \leq p \leq \frac{1}{3}, \\ \frac{3p-1}{2}, & \frac{1}{3} < p \leq 1. \end{cases} \tag{10}$$

□

第 10 题 得分: \_\_\_\_\_. 证明  $S(\rho_A) + S(\rho_B) \leq S(\rho_{AC}) + S(\rho_{BC})$ .

证:<sup>1</sup>首先来证明一个引理:

<sup>1</sup>参考文献: Lieb, Elliott H., and Mary Beth Ruskai. "Proof of the strong subadditivity of quantum-mechanical entropy." *J. Math. Phys.* 19 (1973): 36-55.

映射  $f(\rho_{AC}) = S(\text{Tr}_C(\rho_{AC})) - S(\rho_{AC})$  在 A 与 C 的复合 Hilbert 空间  $H_A \otimes H_C$  上是凸的, 即

$$\begin{aligned} f(\rho_{AC}) &= S(\text{Tr}_C(\rho_{AC})) - S(\rho_{AC}) \\ &\leq \alpha f(\rho_{A'C}) + (1 - \alpha) f(\rho_{A''C}) = \alpha [S(\rho_{A'}) - S(\rho_{A'C})] + (1 - \alpha) [S(\rho_{A''}) - S(\rho_{A''C})], \end{aligned} \quad (11)$$

其中  $\rho_{AC} = \alpha \rho_{A'C} + (1 - \alpha) \rho_{A''C}$ .

这里我们定义

$$\Delta = \alpha \text{Tr}[\rho_{A'C}(-\log_2 \rho_{A'C} + \log_2 \rho_{A'} + \log_2 \rho_{AC} - \log_2 \rho_A)], \quad (12)$$

$$\Gamma = (1 - \alpha) \text{Tr}[\rho_{A''C}(-\log_2 \rho_{A''C} + \log_2 \rho_{A''} + \log_2 \rho_{AC} - \log_2 \rho_A)]. \quad (13)$$

要证上述引理, 即转化为证

$$\Delta + \Gamma \leq 0. \quad (14)$$

利用 Klein 定理, 得

$$\Delta + \Gamma \leq \alpha \text{Tr}[\exp(\log_2 \rho_{A'} + \log_2 \rho_{AC} - \log_2 \rho_A) - \rho_{A'C}] + (1 - \alpha) \text{Tr}[\exp(\log_2 \rho_{A''} + \log_2 \rho_{AC} - \log_2 \rho_A) - \rho_{A''C}], \quad (15)$$

再利用映射  $g(C) = \exp(K + \ln C)$  (其中  $K = \log_2 \rho_{AC} - \log_2 \rho_A$ ) 的上凸性 (即  $\alpha \text{Tr}_{AC}[\exp(\ln \rho_{A'} + \ln \rho_{AC} - \log_2 \rho_A)] + (1 - \alpha) \text{Tr}_{AC}[\exp(\log_2 \rho_{A''} + \log_2 \rho_{AC} - \log_2 \rho_A)] \leq \text{Tr}(\exp(\log_2 \rho_A + \log_2))$ ), 得

$$\Delta + \Gamma = \text{Tr}[\exp(\log_2 \rho_A + \log_2 \rho_{AB} - \log_2 \rho_A) - \rho_{AC}] = 0. \quad (16)$$

引理证毕.

我们定义

$$G(\rho_{ABC}) = S(\text{Tr}_{BC}(\rho_{ABC})) + S(\text{Tr}_{AC}(\rho_{ABC})) - S(\text{Tr}_B(\rho_{ABC})) - S(\text{Tr}_A(\rho_{ABC})), \quad (17)$$

利用上证得的引理, 这里  $S(\text{Tr}_{BC}(\rho_{ABC})) - S(\text{Tr}_B(\rho_{ABC}))$  和  $S(\text{Tr}_{AC}(\rho_{ABC})) - S(\text{Tr}_A(\rho_{ABC}))$  在  $H_A \otimes H_B \otimes H_C$  上也是凸的, 故  $G(\rho_{ABC})$  在  $H_A \otimes H_B \otimes H_C$  上是凸的. 利用 Araki, Huzihiro, and Elliott H. Lieb. "Entropy inequalities." *Inequalities*. Springer, Berlin, Heidelberg, 2002. 47-57. 中的引理 3 得

$$G(\rho_{ABC}) \leq 0, \quad (18)$$

故  $S(\rho_A) + S(\rho_B) \leq S(\rho_{AC}) + S(\rho_{BC})$ . □

**第 11 题 得分:** \_\_\_\_\_. 考虑 2-qubit 系统  $\rho_{AB} = \frac{1}{8}I \otimes I + \frac{1}{2}|\psi^-\rangle\langle\psi^-|$ , 分别沿  $\vec{n}, \vec{m}$  方向测 A, B 粒子的自旋. 其中  $\vec{m} \cdot \vec{n} = \cos \theta$ , 则测量结果均为向上的联合概率是多少? 由 Peres-Horodeski 判据, 确定  $\rho_{AB}$  是否为可分量子态.

**解:** 2-qubit 系统的密度矩阵为

$$\begin{aligned} \rho_{AB} &= \frac{1}{8}I \otimes I + \frac{1}{2}|\psi^-\rangle\langle\psi^-| \\ &= \frac{1}{8}I \otimes I + \frac{1}{2} \frac{1}{\sqrt{2}}(|01\rangle - |10\rangle) \frac{1}{\sqrt{2}}(\langle 01| - \langle 10|) \end{aligned}$$

$$\begin{aligned}
&= \frac{1}{8} I \times I + \frac{1}{4} (|01\rangle\langle 01| - |01\rangle\langle 10| - |10\rangle\langle 01| + |10\rangle\langle 10|) \\
&= \frac{1}{8} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} + \frac{1}{4} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 1 & -1 & 0 \\ 0 & -1 & 1 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} \\
&= \frac{1}{8} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 3 & -2 & 0 \\ 0 & -2 & 3 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}. \tag{19}
\end{aligned}$$

利用第 9 题的结论,  $|\psi^-\rangle$  在  $U(\theta, \vec{n}) \otimes U(\theta, \vec{n})$  旋转变换下不变, 而  $\because (U(\theta, \vec{n}) \otimes U(\theta, \vec{n})) I \times I (U(\theta, \vec{n}) \otimes U(\theta, \vec{n}))^\dagger = (U(\theta, \vec{n}) \otimes U(\theta, \vec{n})) I \times I (U(\theta, \vec{n}) \otimes U(\theta, \vec{n}))^{-1} = I \otimes I$ ,  $\therefore I \times I$  也在  $U(\theta, \vec{n}) \otimes U(\theta, \vec{n})$  旋转变换下不变, 故我们不妨设  $\vec{n}$  沿  $|0\rangle$  方向,  $\vec{m}$  沿  $\cos\theta|0\rangle + e^{i\phi}\sin\theta|1\rangle$ . 分别沿  $\vec{n}$  和  $\vec{m}$  测量 A, B 粒子的自旋, 测量结果均向上对应的投影矩阵为

$$\begin{aligned}
P_{\uparrow\uparrow} &= |0\rangle\langle 0| \otimes \left( \cos\frac{\theta}{2}|0\rangle + e^{i\phi}\sin\frac{\theta}{2}|1\rangle \right) \left( \cos\frac{\theta}{2}\langle 0| + e^{-i\phi}\sin\frac{\theta}{2}\langle 1| \right) \\
&= \cos^2\frac{\theta}{2}|00\rangle\langle 00| + e^{-i\phi}\cos\frac{\theta}{2}\sin\frac{\theta}{2}|00\rangle\langle 01| + e^{i\phi}\cos\frac{\theta}{2}\sin\frac{\theta}{2}|01\rangle\langle 00| + \sin^2\frac{\theta}{2}|01\rangle\langle 01| \\
&= \begin{pmatrix} \cos^2\frac{\theta}{2} & e^{-i\phi}\cos\frac{\theta}{2}\sin\frac{\theta}{2} & 0 & 0 \\ e^{i\phi}\cos\frac{\theta}{2}\sin\frac{\theta}{2} & \sin^2\frac{\theta}{2} & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} \tag{20}
\end{aligned}$$

测量结果均为向上的联合概率为

$$\begin{aligned}
p_{\uparrow\uparrow} &= \text{Tr}(P_{\uparrow\uparrow}\rho_{AB}) = \text{Tr} \left( \begin{pmatrix} \cos^2\frac{\theta}{2} & e^{-i\phi}\cos\frac{\theta}{2}\sin\frac{\theta}{2} & 0 & 0 \\ e^{i\phi}\cos\frac{\theta}{2}\sin\frac{\theta}{2} & \sin^2\frac{\theta}{2} & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} \frac{1}{8} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 3 & -2 & 0 \\ 0 & -2 & 3 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \right) \\
&= \frac{1}{8} \text{Tr} \begin{pmatrix} \cos^2\frac{\theta}{2} & 3e^{-i\phi}\cos\frac{\theta}{2}\sin\frac{\theta}{2} & -2e^{-i\phi}\cos\frac{\theta}{2}\sin\frac{\theta}{2} & 0 \\ e^{i\phi}\cos\frac{\theta}{2}\sin\frac{\theta}{2} & 3\sin^2\frac{\theta}{2} & -2\sin^2\frac{\theta}{2} & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} = \frac{1}{8} \left( \cos^2\frac{\theta}{2} + 3\sin^2\frac{\theta}{2} \right) = \frac{1}{8} + \frac{1}{4}\sin^2\frac{\theta}{2}. \tag{21}
\end{aligned}$$

$\rho_{AB}$  的部分转置矩阵为

$$\sigma_{AB} = \rho_{AB}^{T_B} = \frac{1}{8} \begin{pmatrix} 1 & 0 & 0 & -2 \\ 0 & 3 & 0 & 0 \\ 0 & 0 & 3 & 0 \\ -2 & 0 & 0 & 1 \end{pmatrix}, \tag{22}$$

其特征值为  $\frac{3}{8}, \frac{3}{8}, \frac{3}{8}, -\frac{1}{8}$ , 其中有负特征值, 故  $\sigma_{AB}$  非半正定, 根据 Peres-Horodeski 判据,  $\rho_{AB}$  不是可分量子态.  $\square$