第一次作业

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成绩:

第 1 题 得分: _____. 计算二元对称信道的信道容量.

 \mathbf{M} : 对二元对称信道, 设信源发送信号为 $X \in \{0,1\}$, 经信道传输后信宿接收信号为 $Y \in \{0,1\}$, 信道如图 1 所示:

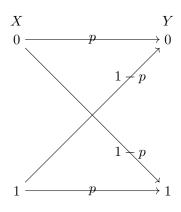


图 1: 二元对称信道

即

$$P(Y = 0 \mid X = 0) = P(Y = 1 \mid X = 1) = p,$$
(1)

$$P(Y = 1 \mid X = 0) = P(Y = 0 \mid X = 1) = 1 - p.$$
(2)

X 与 Y 的互信息量为

$$I(X;Y) = H(Y) - H(Y \mid X), \tag{3}$$

其中 Y 的信息熵为

$$\begin{split} H(Y) &= -P(Y=0)\log_2 P(Y=0) - P(Y=1)\log_2 P(Y=1) \\ &= -P(Y=0)\log_2 P(Y=0) - [1 - P(Y=0)]\log_2 [1 - P(Y=0)] \\ &\leq -\frac{1}{2}\log_2 \frac{1}{2} - \frac{1}{2}\log_2 \frac{1}{2} = 1, \end{split} \tag{4}$$

而 $Y \mid X$ 的条件熵

$$H(Y \mid X) = -P(X = 0)[P(Y = 0 \mid X = 0) \log_2 P(Y = 0 \mid X = 0) + P(Y = 1 \mid X = 0) \log_2 P(Y = 1 \mid X = 0)]$$

$$-P(X = 1)[P(Y = 0 \mid X = 1) \log_2 P(Y = 0 \mid X = 1) + P(Y = 1 \mid X = 1) \log_2 P(Y = 1 \mid X = 1)]$$

$$= -P(X = 0)[p \log_2 p + (1 - p) \log_2 (1 - p)] - P(X = 1)[(1 - p) \log_2 (1 - p) + p \log_2 p]$$

$$= -p \log_2 p - (1 - p) \log_2 (1 - p).$$
(5)

故二元对称信道的信道容量为

$$C = \max_{P(X=0)} I(X;Y) = 1 + p \log_2 p + (1-p) \log_2 (1-p).$$
(6)

第 2 题 得分: ______. 空间 H 中存在两组正交归一化态 $\{|\psi_i\rangle\}$ 、 $\{|\tilde{\psi}_i\rangle\}$,则存在幺正变换 U,使得 $U|\psi_i\rangle=|\tilde{\psi}_i\rangle$,试构造该 U 变换.

解: 构造

$$U = \sum_{i} |\tilde{\psi}_{i}\rangle\langle\psi_{i}|,\tag{7}$$

则 U 满足变换

$$U|\psi_i\rangle = \sum_{i} |\tilde{\psi}_j\rangle\langle\psi_j|\psi_i\rangle = \sum_{i} \delta_{ij}|\tilde{\psi}_j\rangle = |\tilde{\psi}_i\rangle, \quad \forall i,$$
(8)

且 U 是幺正的:

$$U^{\dagger}U = UU^{\dagger} = \left(\sum_{i} |\psi_{i}\rangle\langle\tilde{\psi}_{i}|\right) \left(\sum_{j} |\tilde{\psi}_{j}\rangle\langle\psi_{j}|\right) = \sum_{ij} \delta_{ij} |\psi_{i}\rangle\langle\psi_{j}| = \sum_{i} |\psi_{i}\rangle\langle\psi_{i}| = I.$$
 (9)

第 3 题 得分: _______. 空间 H 中存在两组归一化态 $\{|\psi_i\rangle\}$ 、 $\{|\tilde{\psi}_i\rangle\}$,它们满足: $\forall i,j,$ 有 $\langle\psi_i|\psi_j\rangle=\langle\tilde{\psi}_i|\tilde{\psi}_j\rangle$. 请 证明, 则存在 U, 使得 $U|\psi_i\rangle = |\tilde{\psi}_i\rangle$, 并构造出该 U 变换.

解:设

$$U = \sum_{jk} a_{jk} |\tilde{\psi}_j\rangle\langle\psi_k|. \tag{10}$$

要使 $U|\psi_i\rangle = |\tilde{\psi}_i\rangle$, 即需

$$U|\psi_i\rangle = \sum_{jk} a_{jk} |\tilde{\psi}_j\rangle \langle \psi_k | \psi_i\rangle = \sum_j \left(\sum_k a_{jk} \langle \psi_k | \psi_i\rangle\right) |\tilde{\psi}_j\rangle = |\tilde{\psi}_i\rangle, \quad \forall i,$$
(11)

$$\Longrightarrow \sum_{k} a_{jk} \langle \psi_k | \psi_i \rangle = \delta_{ij}, \quad \forall i, j, \tag{12}$$

$$\Longrightarrow AB = I \tag{13}$$

其中矩阵

$$A = \begin{bmatrix} a_{11} & a_{12} & \cdots & a_{1n} \\ a_{21} & a_{22} & \cdots & a_{2n} \\ \vdots & \vdots & \ddots & \vdots \\ a_{n1} & a_{n2} & \cdots & a_{nn} \end{bmatrix},$$

$$B = \begin{bmatrix} \langle \psi_1 | \psi_1 \rangle & \langle \psi_1 | \psi_2 \rangle & \cdots & \langle \psi_1 | \psi_n \rangle \\ \langle \psi_2 | \psi_1 \rangle & \langle \psi_2 | \psi_2 \rangle & \cdots & \langle \psi_2 | \psi_n \rangle \\ \vdots & \vdots & \ddots & \vdots \\ \langle \psi_1 | \psi_1 \rangle & \langle \psi_2 | \psi_2 \rangle & \cdots & \langle \psi_2 | \psi_n \rangle \end{bmatrix}.$$

$$(14)$$

$$B = \begin{bmatrix} \langle \psi_1 | \psi_1 \rangle & \langle \psi_1 | \psi_2 \rangle & \cdots & \langle \psi_1 | \psi_n \rangle \\ \langle \psi_2 | \psi_1 \rangle & \langle \psi_2 | \psi_2 \rangle & \cdots & \langle \psi_2 | \psi_n \rangle \\ \vdots & \vdots & \ddots & \vdots \\ \langle \psi_n | \psi_1 \rangle & \langle \psi_n | \psi_2 \rangle & \cdots & \langle \psi_n | \psi_n \rangle \end{bmatrix}.$$

$$(15)$$

因此我们只需由上式构造矩阵 B, 再计算

$$A = B^{-1},\tag{16}$$

然后构造

$$U = \sum_{jk} a_{jk} |\tilde{\psi}_j\rangle\langle\psi_k| \tag{17}$$

即可满足 $U|\psi_i\rangle = |\tilde{\psi}_i\rangle$.

第 4 题 得分: _______. 对两比特态 $|\phi\rangle = \frac{1}{\sqrt{2}}|0\rangle_A \left(\frac{1}{2}|0\rangle_B + \frac{\sqrt{3}}{2}|1\rangle_B\right) + \frac{1}{\sqrt{2}}|1\rangle_A \left(\frac{\sqrt{3}}{2}|0\rangle_B + \frac{1}{2}|1\rangle_B\right)$,

- i) 求约化密度矩阵 ρ_A , ρ_B ;
- ii) 求 $|\phi\rangle$ 的 Schmidt 分解形式.

解: 1) 复合系统的密度矩阵为

$$\begin{split} &\rho_{AB} = |\phi\rangle\langle\phi| \\ &= \left[\frac{1}{\sqrt{2}}|0\rangle_{A}\left(\frac{1}{2}|0\rangle_{B} + \frac{\sqrt{3}}{2}|1\rangle_{B}\right) + \frac{1}{\sqrt{2}}|1\rangle_{A}\left(\frac{\sqrt{3}}{2}|0\rangle_{B} + \frac{1}{2}|1\rangle_{B}\right)\right] \left[\frac{1}{\sqrt{2}}\langle0|_{A}\left(\frac{1}{2}\langle0|_{B} + \frac{\sqrt{3}}{2}\langle1|_{B}\right) + \frac{1}{\sqrt{2}}\langle1|_{A}\left(\frac{\sqrt{3}}{2}\langle0|_{B} + \frac{1}{2}\langle1|_{B}\right)\right)\right] \\ &= \frac{1}{8}|0\rangle_{A}|0\rangle_{B}\langle0|_{A}\langle0|_{B} + \frac{\sqrt{3}}{8}|0\rangle_{A}|0\rangle_{B}\langle0|_{A}\langle1|_{B} + \frac{\sqrt{3}}{8}|0\rangle_{A}|0\rangle_{B}\langle1|_{A}\langle0|_{B} + \frac{1}{8}|0\rangle_{A}|0\rangle_{B}\langle1|_{A}\langle1|_{B} \\ &+ \frac{\sqrt{3}}{8}|0\rangle_{A}|1\rangle_{B}\langle0|_{A}\langle0|_{B} + \frac{3}{8}|0\rangle_{A}|1\rangle_{B}\langle0|_{A}\langle1|_{B} + \frac{3}{8}|0\rangle_{A}|1\rangle_{B}\langle1|_{A}\langle0|_{B} + \frac{\sqrt{3}}{8}|0\rangle_{A}|1\rangle_{B}\langle1|_{A}\langle1|_{B} \\ &+ \frac{\sqrt{3}}{8}|1\rangle_{A}|0\rangle_{B}\langle0|_{A}\langle0|_{B} + \frac{3}{8}|1\rangle_{A}|0\rangle_{B}\langle0|_{A}\langle1|_{B} + \frac{3}{8}|1\rangle_{A}|0\rangle_{B}\langle1|_{A}\langle0|_{B} + \frac{\sqrt{3}}{8}|1\rangle_{A}|0\rangle_{B}\langle1|_{A}\langle1|_{B} \\ &+ \frac{1}{8}|1\rangle_{A}|1\rangle_{B}\langle0|_{A}\langle0|_{B} + \frac{\sqrt{3}}{8}|1\rangle_{A}|1\rangle_{B}\langle0|_{A}\langle1|_{B} + \frac{\sqrt{3}}{8}|1\rangle_{A}|1\rangle_{B}\langle1|_{A}\langle0|_{B} + \frac{1}{8}|1\rangle_{A}|1\rangle_{B}\langle1|_{A}\langle1|_{B}. \end{split} \tag{18}$$

约化密度矩阵

$$\rho_A = \text{Tr}_B[\rho_{AB}] = \sum_{i=0,1} \langle i|_B(|\phi\rangle\langle\phi|)|i\rangle_B = \frac{1}{2}|0\rangle_A\langle 0|_A + \frac{\sqrt{3}}{4}|0\rangle_A\langle 1|_A + \frac{\sqrt{3}}{4}|1\rangle_A\langle 0|_A + \frac{1}{2}|1\rangle_A\langle 1|_A, \tag{19}$$

$$\rho_B = \text{Tr}_A[\rho_{AB}] = \sum_{i=0,1} \langle i|_A(|\phi\rangle\langle\phi|)|i\rangle_A = \frac{1}{2}|0\rangle_B\langle 0|_B + \frac{\sqrt{3}}{4}|0\rangle_B\langle 1|_B + \frac{\sqrt{3}}{4}|1\rangle_B\langle 0|_B + \frac{1}{2}|1\rangle_B\langle 1|_B.$$
 (20)

(2) $|\phi\rangle$ 可表为

$$|\phi\rangle = \sum_{m,n\in\{0,1\}} a_{mn} |m\rangle_A |n\rangle_B, \tag{21}$$

其中矩阵

$$a = \begin{bmatrix} \frac{\sqrt{2}}{4} & \frac{\sqrt{6}}{4} \\ \frac{\sqrt{6}}{4} & \frac{\sqrt{2}}{4} \end{bmatrix}. \tag{22}$$

对 a 做奇异值分解, 得

$$a = udv, (23)$$

其中

$$u = \begin{bmatrix} \frac{\sqrt{2}}{2} & \frac{\sqrt{2}}{2} \\ -\frac{\sqrt{2}}{2} & \frac{\sqrt{2}}{2} \end{bmatrix},\tag{24}$$

$$d = \begin{bmatrix} \frac{\sqrt{2} - \sqrt{6}}{4} & 0\\ 0 & \frac{\sqrt{2} + \sqrt{6}}{4} \end{bmatrix}, \tag{25}$$

$$v = \begin{bmatrix} \frac{\sqrt{2}}{2} & -\frac{\sqrt{2}}{2} \\ \frac{\sqrt{2}}{2} & \frac{\sqrt{2}}{2} \end{bmatrix}. \tag{26}$$

令子系统 A 和子系统 B 的正交基

$$|\psi_i\rangle_A = \sum_{j=0,1} u_{ji}|j\rangle,\tag{27}$$

$$|\psi_i\rangle_B = \sum_{k=0,1} v_{ik}|k\rangle,\tag{28}$$

以及

$$\lambda_i = d_{ii}, \tag{29}$$

即

$$|\psi_1\rangle_A = \frac{\sqrt{2}}{2}|0\rangle_A - \frac{\sqrt{2}}{2}|1\rangle_A, \qquad |\psi_2\rangle_A = \frac{\sqrt{2}}{2}|0\rangle_A + \frac{\sqrt{2}}{2}|1\rangle_A, \tag{30}$$

$$|\psi_1\rangle_B = \frac{\sqrt{2}}{2}|0\rangle_B - \frac{\sqrt{2}}{2}|1\rangle_B,$$
 $|\psi_2\rangle_B = \frac{\sqrt{2}}{2}|0\rangle_B + \frac{\sqrt{2}}{2}|1\rangle_B,$ (31)

$$\lambda_1 = \frac{\sqrt{2} - \sqrt{6}}{4}, \qquad \qquad \lambda_2 = \frac{\sqrt{2} + \sqrt{6}}{4}, \qquad (32)$$

从而得到 $|\phi\rangle$ 的 Schmidt 分解式为

$$|\phi\rangle = \sum_{i=1}^{2} \lambda_{i} |\psi_{i}\rangle_{A} |\psi_{i}\rangle_{B}$$

$$= \frac{\sqrt{2} - \sqrt{6}}{4} \left(\frac{\sqrt{2}}{2}|0\rangle_{A} - \frac{\sqrt{2}}{2}|1\rangle_{A}\right) \left(\frac{\sqrt{2}}{2}|0\rangle_{B} - \frac{\sqrt{2}}{2}|1\rangle_{B}\right) + \frac{\sqrt{2} + \sqrt{6}}{4} \left(\frac{\sqrt{2}}{2}|0\rangle_{A} + \frac{\sqrt{2}}{2}|1\rangle_{A}\right) \left(\frac{\sqrt{2}}{2}|0\rangle_{B} + \frac{\sqrt{2}}{2}|1\rangle_{B}\right). \tag{33}$$

第 5 题 得分: ______. 对三粒子系统纯态 $|\phi_{ABC}\rangle$, 在空间 $H_A\otimes H_B\otimes H_C$ 中是否存在 H_A , H_B , H_C 中的正交基, 使得 $|\phi_{ABC}\rangle=\sum_i\sqrt{p_i}|i_A\rangle\otimes|i_B\rangle\otimes|i_C\rangle$ 一定成立? 给出理由.

解: 不一定. 理由如下:

对三粒子系统纯态 $|\phi_{ABC}\rangle$, 利用 Schmidt 分解总是可以得到

$$|\phi_{ABC}\rangle = \sum_{i} \sqrt{p_i} |i_A\rangle \otimes |i_{BC}\rangle.$$
 (34)

对于 $|i_{BC}\rangle$, 利用 Schmidt 分解得

$$|i_{BC}\rangle = \sum_{j} \sqrt{p_{ij}} |j_{i,B}\rangle \otimes |j_{i,C}\rangle,$$
 (35)

从而

$$|\phi_{ABC}\rangle = \sum_{i} \sqrt{p_i} |i_A\rangle \otimes \sum_{j} \sqrt{p_{ij}} |j_{i,B}\rangle \otimes |j_{i,C}\rangle$$
 (36)

因此, $|\phi_{ABC}\rangle$ 无法写成 $|\phi_{ABC}\rangle = \sum_{i} \sqrt{p_i} |i_A\rangle \otimes |i_B\rangle \otimes |i_C\rangle$ 的形式.

只有当对任意 $i, |i_{BC}\rangle$ 的分解式 (式 35) 中的求和都只有一项, 才能保证 $|\phi_{ABC}\rangle = \sum_i \sqrt{p_i} |i_A\rangle \otimes |i_B\rangle \otimes |i_C\rangle$ 的分解形式成立.