

# 数学物理方法

Methods in Mathematical Physics

第十章 格林函数法 Method of Green's Function

武汉大学物理科学与技术学院





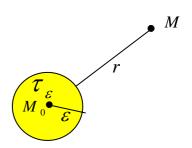
# 第十章 格林函数法 § 10.3- § 10.4 格林函数 Green's Functions

# 泊松方程的格林函数



1、三维: 
$$\Delta G = -\delta(M - M_0)$$

$$\Delta G = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 \frac{\partial G}{\partial r}) = -\delta(r)$$



(1) 若
$$r \neq 0$$
:  $G = -C_1 \frac{1}{r}$ 

(2) 若
$$r = 0$$
: 考虑: 
$$\iiint_{\tau_{\varepsilon}} \Delta G dv = -\iiint_{\tau_{\varepsilon}} \delta(r) dv = -1$$

$$\mathbf{X} \iiint_{\tau_{\varepsilon}} \Delta G dv = \iiint_{\tau_{\varepsilon}} \nabla \cdot \nabla G dv = \iint_{\sigma_{\varepsilon}} \nabla G \cdot d\vec{\sigma} = C_{1} 4\pi$$

$$\rightarrow C_1 = -\frac{1}{4\pi},$$

$$(1),(2) \rightarrow G = \frac{1}{4\pi r}$$

一泊松方程格林函数。

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# 泊松方程的格林函数



2、二维: 
$$\Delta G = -\delta(M - M_0)$$

$$\Delta G = \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial G}{\partial r} \right) = -\delta(r), \quad r = \sqrt{(x - x_0)^2 + (y - y_0)^2}$$

$$\iint_{\sigma} \nabla \cdot \nabla u d\sigma = \int_{l} \nabla u \cdot d\vec{l}$$

$$G(M, M_0) = \frac{1}{2\pi} \ln \frac{1}{r}$$

# 二、狄氏格林函数



1、 三维:  $\begin{cases} \Delta G = -\delta(x - x_0, y - y_0, z - z_0), M \in \tau \\ G|_{\sigma} = 0 \end{cases}$ 

$$\Leftrightarrow G(M, M_0) = F(M, M_0) + g(M, M_0)$$

使 
$$\Delta F(M, M_0) = -\delta(M - M_0)$$
  $M \in \tau$ 

$$G(M, M_0) = \frac{1}{4\pi r} + g$$

则

$$\begin{cases} \Delta g = 0, M \in \tau \\ g \mid_{\sigma} = -\frac{1}{4\pi r} \mid_{\sigma} \end{cases}$$

一狄氏格林函数



#### 狄氏格林函数



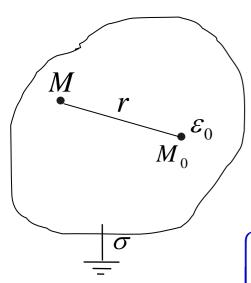
类似 
$$G = \frac{1}{2\pi} \ln \frac{1}{r} + g$$
 — 狄氏格林函数 
$$\begin{cases} \Delta g = 0, M \in \sigma \\ g|_{l} = -\frac{1}{2\pi} \ln \frac{1}{r}|_{l} \end{cases}$$







### 3、狄氏格林函数的物理意义:



$$\Delta G = -\delta(M - M_0,), M \in \tau$$

$$G|_{\sigma} = 0$$

$$G(M,M_0) = \frac{1}{4\pi r} + g$$

$$\begin{cases} \Delta g = 0, M \in \tau \\ g \mid_{\sigma} = -\frac{1}{4\pi r} \mid_{\sigma} \end{cases}$$

$$\varepsilon_0$$
 提供:  $\frac{1}{4\pi\varepsilon_0}\frac{\varepsilon_0}{r} = \frac{1}{4\pi r}$ 

$$G-M$$
点电位 $<$ 

感应电荷提供
$$v$$
: 
$$\begin{cases} \Delta v = 0, M \in \tau & \therefore v = g \\ v|_{\sigma} = -\frac{1}{4\pi r}|_{\sigma} \end{cases}$$

$$\therefore v = g$$





# 3、狄氏格林函数的物理意义:

求 $G \rightarrow$ 求M点电位  $\rightarrow$  求感应电荷产生的电位

即求:  $\begin{cases} \Delta g = 0, M \in \tau \\ g \mid_{\sigma} = -\frac{1}{4\pi r} \mid_{\sigma} \end{cases}$ 

对于二维:

即求: 
$$\begin{cases} \Delta g = 0, M \in \sigma \\ g|_{l} = -\frac{1}{2\pi} \ln \frac{1}{r}|_{l} \end{cases}$$



10.3-10.4 格林函数



#### 1、问题的引入:

求解球内的狄氏问题: 
$$\begin{cases} \Delta u = 0, \ \rho < a \\ u|_{\rho=a} = f(M) \end{cases}$$

解: 
$$u(M) = -\iint_{\sigma} f(M_0) \frac{\partial G}{\partial n_0} d\sigma_0$$
$$G(M, M_0) = \frac{1}{4\pi r} + g \quad \begin{cases} \Delta g = 0, M \in \rho < a \\ g \mid_{\rho = a} = -\frac{1}{4\pi r} \mid_{\rho = a} \end{cases}$$

 $\ddot{x}u \to \ddot{x}G \to \ddot{x}M$ 点电位  $\to \ddot{x}$ 感应电荷产生的电位 g



10.3-10.4 格林函数

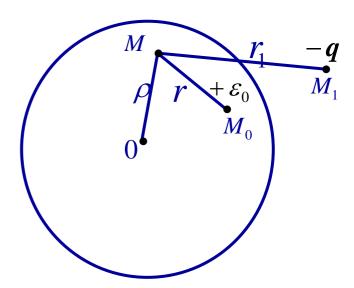
#### 2、用电像法求g:

(1) 分析: 若能在 $\sigma$ 外的某点 $M_1$ 放一适当的负q,则

$$\Delta(\frac{-q}{4\pi\varepsilon_0 r_1}) = 0, \ M \in \rho < a$$

使: 
$$-\frac{q}{4\pi\varepsilon_0 r_1}\Big|_{\rho=a} = -\frac{1}{4\pi r}\Big|_{\rho=a}$$

则 
$$g = -\frac{q}{4\pi\varepsilon_0 r_1}$$



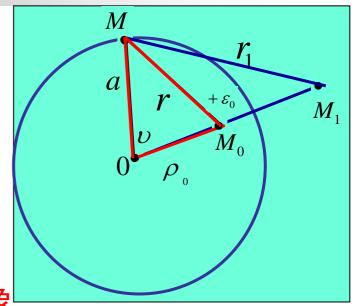
::  $\bar{x}g \rightarrow a$ )确定 $M_1$ 的位置; b)确定q大小问题





#### 2、用电像法求g:

- (2) 求球域的G
- a)  $r_1 = ?$ 记  $|OM_0| = \rho_0$ ,  $|OM_1| = \rho_1$  使  $\rho_0 \cdot \rho_1 = a^2$  即  $\frac{\rho_0}{a} = \frac{a}{\rho_1}$



#### 则称 $M_1$ 为 $M_0$ 关于球面 $\rho = a$ 的像

b) 
$$q = ?$$
  $\frac{1}{r}|_{\sigma} = ?$   $\therefore \Delta OM_0 M \hookrightarrow \Delta OM_1 M$ 

$$\therefore \frac{\rho_0}{a} = \frac{a}{\rho_1} = \frac{r}{r_1}, \quad \text{RD} \left[ \frac{1}{r} \big|_{\rho=a} = \frac{a/\rho_0}{r_1} \big|_{\rho=a} \right]$$



10.3-10.4 格林函数



#### 2、用电像法求g:

(2) 求球域的G

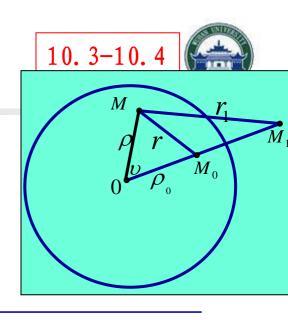
$$g = \frac{-\varepsilon_0 a/\rho_0}{4\pi\varepsilon_0 r_1} = \frac{-a/\rho_0}{4\pi r_1} \qquad q = \frac{\varepsilon_0 a}{\rho_0}$$

$$G = \frac{1}{4\pi r} - \frac{a/\rho_0}{4\pi r_1} \qquad -q = -\frac{\varepsilon_0 a}{\rho_0}$$
 是 $\varepsilon_0$ 的电像

(3) 电像法:这种在像点放一虚构的点电荷,来等效代替边界面上的感应电荷所产生的电位的方法称之为电像法

3、求u(M)

$$\frac{\partial G}{\partial n} = \frac{1}{4\pi} \left[ \frac{\partial}{\partial \rho} \left( \frac{1}{r} \right) - \frac{a}{\rho_0} \frac{\partial}{\partial \rho} \left( \frac{1}{r_1} \right) \right]$$



$$r = \sqrt{\rho^2 + \rho_0^2 - 2\rho\rho_0 \cos \gamma} \qquad r_1 = \sqrt{\rho^2 + \rho_1^2 - 2\rho\rho_1 \cos \gamma}$$

$$\frac{\partial}{\partial \rho} \frac{1}{r} = \frac{\rho_0^2 - \rho^2 - r^2}{2\rho r^3} \qquad \Rightarrow \frac{\partial}{\partial \rho} \left(\frac{1}{r}\right)_{\rho=a} = \frac{\rho_0^2 - a^2 - r^2}{2ar^3}$$

类似有: 
$$\frac{a}{\rho_0} \frac{\partial}{\partial \rho} \left( \frac{1}{r_1} \right)_{\rho=a} = \frac{a}{\rho_0} \frac{\rho_1^2 - a^2 - r_1^2}{2ar_1^3} = \frac{a^2 - r^2 - \rho_0^2}{2ar^3}$$

$$\frac{\partial G}{\partial n}\Big|_{\rho=a} = \frac{1}{4\pi a} \frac{\rho_0^2 - a^2}{r^3}$$

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3、求<math>u(M)

$$u(M) = -\iint_{\sigma} f(M_0) \frac{\partial G}{\partial n_0} d\sigma_0$$

$$= \frac{1}{4\pi a} \int_0^{2\pi} \int_0^{\pi} f(\theta_0, \varphi_0) \frac{a^2 - \rho^2}{\left(a^2 + \rho^2 - 2a\rho\cos\gamma\right)^{\frac{3}{2}}} a^2 \sin\theta_0 d\theta_0 d\varphi_0$$

$$u(M) = \frac{a}{4\pi} \int_0^{2\pi} \int_0^{\pi} f(\theta_0, \varphi_0) \frac{a^2 - \rho^2}{\left(a^2 + \rho^2 - 2a\rho\cos\gamma\right)^{\frac{3}{2}}} \sin\theta_0 d\theta_0 d\varphi_0$$

#### 一球的泊松积分公式

其中, 
$$\cos \gamma = \sin \theta \sin \theta_0 \cos(\varphi - \varphi_0) + \cos \theta \cos \theta_0$$



#### 四、注释



$$1 \cos \gamma = ?$$

设 $\overrightarrow{I}$ 为 $\overrightarrow{OM}$ 方向单位向量, $\overrightarrow{I_0}$ 为 $\overrightarrow{OM_0}$ 方向单位向量

$$\iint_{0} \vec{I} = x \vec{i} + y \vec{j} + z \vec{k} = \sin \theta \cos \varphi \vec{i} + \sin \theta \sin \varphi \vec{j} + \cos \theta \vec{k}$$

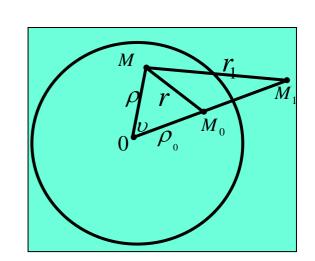
$$\vec{I}_{0} = x_{0} \vec{i} + y_{0} \vec{j} + z_{0} \vec{k} = \sin \theta_{0} \cos \varphi_{0} \vec{i} + \sin \theta_{0} \sin \varphi_{0} \vec{j} + \cos \theta_{0} \vec{k}$$

$$\vec{I} \cdot \vec{I}_0 = |\vec{I} \cdot \vec{I}_0| \cos \gamma = \cos \gamma$$

$$= \sin \theta \cos \varphi \sin \theta_0 \cos \varphi_0$$

$$+ \sin \theta \sin \varphi \sin \theta_0 \sin \varphi_0$$

$$+ \cos \theta \cos \theta_0$$



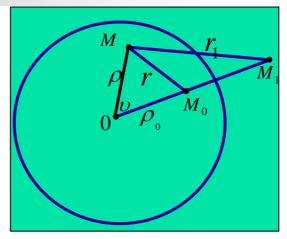




#### 2、对于球外的狄氏问题:

$$\begin{cases} \Delta u_1 = 0 , \rho > a \\ u_1 \mid_{\rho = a} = f(M) \end{cases}$$

$$u_1(M) = -u(M)$$



$$u_1(M) = \frac{-a}{4\pi} \int_0^{2\pi} \int_0^{\pi} f(\theta_0, \varphi_0) \frac{a^2 - \rho^2}{\left(a^2 + \rho^2 - 2a\rho\cos\gamma\right)^{\frac{3}{2}}} \sin\theta_0 d\theta_0 d\varphi_0$$

#### 四、注释





#### 3、对于圆内的狄氏问题:

$$\begin{cases} \Delta u = 0, & \rho < a \\ u|_{\rho = a} = f(\varphi) \end{cases}$$

$$G = \frac{1}{2\pi} \ln \frac{1}{r} + g$$

$$\int \Delta a = 0 \quad M \in \sigma$$

$$\begin{cases} \Delta g = 0, M \in \sigma \\ g|_{\rho=a} = -\frac{1}{2\pi} \ln \frac{1}{r}|_{\rho=a} \end{cases}$$

$$u(M) = \frac{1}{2\pi} \int_0^{2\pi} f(\varphi_0) \frac{a^2 - \rho^2}{a^2 + \rho^2 - 2a\rho\cos(\varphi - \varphi_0)} d\varphi_0$$

#### 4、对于圆外的狄氏问题:

$$\begin{cases} \Delta u_1 = 0, & \rho > a \\ u_1 \mid_{\rho = a} = f(\varphi) \end{cases} \qquad u_1(M) = -u(M)$$

$$u(M) = \frac{1}{2\pi} \int_0^{2\pi} f(\varphi_0) \frac{\rho^2 - a^2}{a^2 + \rho^2 - 2a\rho\cos(\varphi - \varphi_0)} d\varphi_0$$

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# 五、小结



$$\Delta G = -\delta(M - M_0) \rightarrow G = \frac{1}{4\pi r}$$

$$G = \frac{1}{4\pi r}$$

$$\Delta G = -\delta(M - M_0) \rightarrow G = \frac{1}{4\pi r}$$

$$\begin{cases} \Delta G = -\delta(M - M_0), M \in \tau \\ G|_{\sigma} = 0 \end{cases} \rightarrow G(M, M_0) = \frac{1}{4\pi r} + G(M, M_0) = \frac$$

$$G(M, M_0) = \frac{1}{4\pi r} + g$$

$$\begin{cases} \Delta g = 0 \\ g|_{\sigma} = -\frac{1}{4\pi r}|_{\sigma} \end{cases}$$

$$u(M) = \frac{a}{4\pi} \int_0^{2\pi} \int_0^{\pi} f(\theta_0, \varphi_0) \frac{a^2 - \rho^2}{\left(a^2 + \rho^2 - 2a\rho\cos\gamma\right)^{\frac{3}{2}}} \sin\theta_0 d\theta_0 d\varphi_0$$





习题 10.3:1;

习题 10.4: 1;2;3

# Good-by!

