

数学物理方法

Methods in Mathematical Physics

第十章 格林函数法 Method of Green's Function

武汉大学物理科学与技术学院





- -、 δ 函数及其在物理上的应用
- 二、用电像法求格林函数
- 三、用格林函数法求泊松方程的狄氏问题



、 δ 函数及其在物理上的应用

1.
$$\begin{cases} \delta(x - x_0) = \begin{cases} 0 & x \neq x_0 \\ \infty & x = x_0 \end{cases} \\ \int_{-\infty}^{\infty} \delta(x - x_0) dx = 1 \end{cases}$$
 2.
$$\int_{-\infty}^{\infty} f(x) \delta(x - x_0) dx = f(x_0)$$

$$2. \int_{-\infty}^{\infty} f(x)\delta(x-x_0)dx = f(x_0)$$

3.
$$\int_{-\infty}^{\infty} f(x)\delta^{(n)}(x-x_0)dx = (-1)^n f^{(n)}(x_0)$$

4.
$$\delta[\varphi(x)] = \sum_{i=1}^{n} \frac{\delta(x - x_i)}{|\varphi'(x_i)|}$$
,其中 $\varphi(x_i) = 0$



一、 δ 函数及其在物理上的应用

1.
$$H(x) = \begin{cases} 0, x < 0 \\ 1, x > 0 \end{cases}$$
 证明: $\delta(x) = \frac{dH(x)}{dx}$

设 $\varphi(x)$ 为任意一无穷次可微且在无穷区域等于零的函数,则

$$\int_{-\infty}^{\infty} \frac{dH}{dx} \varphi(x) dx = \varphi(x) H(x) \Big|_{-\infty}^{\infty} - \int_{-\infty}^{\infty} H(x) \varphi'(x) dx$$
$$= -\int_{0}^{\infty} \varphi'(x) dx = \varphi(0)$$

$$\int_{-\infty}^{\infty} \delta(x) \varphi(x) dx = \varphi(0)$$

$$\delta(x) = \frac{dH(x)}{dx}$$



一、 δ 函数及其在物理上的应用

2. iE:
$$\begin{cases} \delta(x^2 - a^2) = \frac{1}{2|a|} [\delta(x - a) + \delta(x - a)] \\ \delta(x^2) = \frac{\delta(x)}{|x|} \end{cases}$$

(1) ::
$$\delta[\varphi(x)] = \sum_{k=1}^{n} \frac{\delta(x - x_k)}{|\varphi'(x_k)|}$$
 其中 $\varphi(x_k) = 0, k = 1, 2, \dots, n$ 令 $\varphi(x) = x^2 - a^2$ 則 $\varphi'(x) = 2x$,
且由 $x^2 - a^2 = 0$ 有 $x_1 = a, x_2 = -a$
故有 $\delta(x^2 - a^2) = \delta[\varphi(x)] = \sum_{k=1}^{2} \frac{\delta(x - x_k)}{|2x_k|}$

$$\delta(x^2 - a^2) = \frac{1}{2|a|} [\delta(x - a) + \delta(x + a)]$$



$oldsymbol{----}$ 、 δ 函数及其在物理上的应用

$$(2) :: \int_{-\infty}^{\infty} f(x) \frac{1}{2|a|} \left[\delta(x-a) + \delta(x+a) \right] dx = \frac{1}{2|a|} \left[f(a) + f(-a) \right]$$

$$\int_{-\infty}^{\infty} f(x) \frac{1}{2|x|} \left[\delta(x-a) + \delta(x+a) \right] dx$$

$$= f(a) \frac{1}{2|a|} + f(-a) \cdot \frac{1}{2|-a|} = \frac{1}{2|a|} \left[f(a) + f(-a) \right]$$

$$\therefore \delta(x^2 - a^2) = \frac{1}{2|a|} \left[\delta(x - a) + \delta(x + a) \right] = \frac{1}{2|x|} \left[\delta(x - a) + \delta(x + a) \right]$$

在上式中取
$$a=0$$
 $\delta(x^2) = \frac{\delta(x)}{|x|}$



、 δ 函数及其在物理上的应用

$$3.F[\delta(x)] = ?, \delta(x)$$
的积分表式?

$$F[\cos ax] = ?$$
, $F[\sin ax] = ?$

$$F[\delta(x)] = 1 \quad \therefore \delta(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{i\omega x} d\omega \quad \longrightarrow \int_{-\infty}^{\infty} e^{i\omega x} d\omega = 2\pi \delta(x)$$

$$\to \int_{-\infty}^{\infty} e^{i\omega x} d\omega = 2\pi \delta(x)$$

$$F[\cos ax] = \int_{-\infty}^{\infty} \frac{e^{iax} + e^{-iax}}{2} e^{-i\omega x} dx = \frac{1}{2} \left[\int_{-\infty}^{\infty} e^{i(a-\omega)x} dx + \int_{-\infty}^{\infty} e^{-i(a+\omega)x} dx \right]$$

$$\to \int_{-\infty}^{\infty} e^{i(a-\omega)x} dx = 2\pi \delta(a-\omega), \int_{-\infty}^{\infty} e^{-i(a+\omega)x} dx = 2\pi \delta(a+\omega),$$

$$\therefore F[\cos ax] = \pi[\delta(a-\omega) + \delta(a+\omega)]$$

$$F[\sin ax] = i\pi[\delta(a+\omega) - \delta(a-\omega)]$$

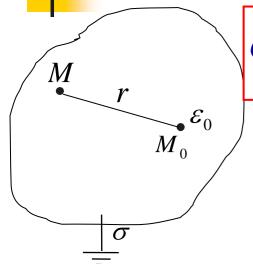
思考:
$$F[e^{iax}] = \int_{-\infty}^{\infty} e^{iax} e^{-i\omega x} dx = \int_{-\infty}^{\infty} e^{i(a-\omega)x} dx = 2\pi\delta(a-\omega) \rightarrow$$

$$F[\cos x] + iF[\sin x] = 2\pi\delta(a-\omega) \rightarrow F[\cos x] = 2\pi\delta(a-\omega), F[\sin x] = 0$$

是否正确?



二、用电像法求狄氏格林函数



$$G(M, M_0) = \frac{1}{4\pi r} + g$$

$$\begin{cases} \Delta g = 0, M \in \tau \\ g \mid_{\sigma} = -\frac{1}{4\pi r} \mid_{\sigma} \end{cases}$$

$$G = \frac{1}{2\pi} \ln \frac{1}{r} + g$$

$$\begin{cases} \Delta g = 0, M \in \sigma \\ g|_{l} = -\frac{1}{2\pi} \ln \frac{1}{r}|_{l} \end{cases}$$

 $\bar{x}G \rightarrow \bar{x}M$ 点电位 $\rightarrow \bar{x}$ 感应电荷产生的电位

用电像法求:在像点放一虚构的点电荷,来等效的代替边界面上的感应电荷所产生的电位。

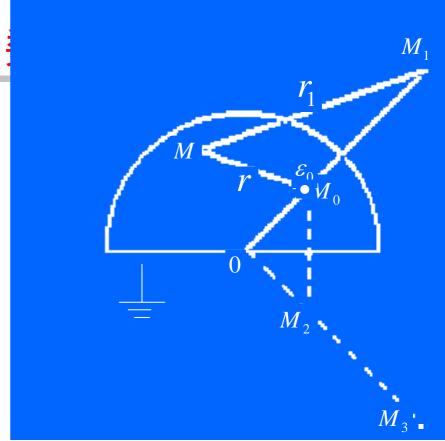
二、用电像法求格林函

1. 求上半球域的G

$$G = \frac{1}{4\pi r} + g$$

$$\begin{cases} \Delta g = 0 , M \in \tau \\ g|_{\sigma} = -\frac{1}{4\pi r}|_{\sigma} \end{cases}$$

(1)确定象点: 设 $M_0(\rho_0,\theta_0,\varphi_0)$



則
$$M_1(\frac{a^2}{\rho_0},\theta_0,\varphi_0), M_2(\rho_0,-\theta_0,\varphi_0), M_3(\frac{a^2}{\rho_0},-\theta_0,\varphi_0),$$

$$\frac{\rho_0}{\zeta} = \frac{a}{\rho_1} = \frac{r}{r_1}$$
 $\frac{\rho_0}{a} = \frac{a}{\rho_1} = \frac{r}{r_1}$
 $\frac{\rho_0}{a} = \frac{r}{\rho_1} = \frac{r}{r_1}$

$$(2)$$
确定 g 及 q_i :在 M_i 点放置 $-q_i$ 则

$$(\mathbf{Q}_i : \mathbf{E} \mathbf{M}_i \mathbf{A} \mathbf{X} \mathbf{E} - \mathbf{Q}_i \mathbf{M}) = 0, M \in \tau \quad g = \sum_{i=1}^{3} \frac{1}{2} \frac{1}{2}$$

$$\begin{array}{c|c}
1 & 5 \\
 & 5 \\
 & 7 \\
\hline
 & 7_i
\end{array}$$

$$\Delta(\sum_{i=1}^{3} \left(\frac{1}{4\pi\varepsilon_{0}} \frac{q_{i}}{r_{i}}\right)) = 0, M \in \tau \quad g = \sum_{i=1}^{3} \frac{1}{4\pi\varepsilon_{0}} \frac{q_{i}}{r_{i}}$$

而由
$$\frac{1}{4\pi\varepsilon_0} \frac{q_i}{r_i} |_{\sigma} = -\frac{1}{4\pi r} |_{\sigma}$$

$$\bar{\mathbf{J}} \stackrel{\mathbf{I}}{=} \frac{1}{4\pi\varepsilon_0} \frac{q_i}{r_i} \Big|_{\sigma} = -\frac{1}{4\pi r} \Big|_{\sigma}$$

$$\begin{array}{c|c}
& \pm \frac{1}{4\pi\varepsilon_0} \frac{q_i}{r_i} \Big|_{\sigma} = -\frac{1}{4\pi r} \Big|_{\sigma} \\
& q_1 \\
& q_1
\end{array}$$

$$\frac{1}{4\pi\varepsilon_0} \frac{r_i}{r_i} \Big|_{\sigma} = -\frac{1}{4\pi r} \Big|_{\sigma}$$

$$q_1 \qquad q_1 \qquad q_1 \qquad q_2 \qquad q_3 \qquad q_4 \qquad q_4 \qquad q_4 \qquad q_5 \qquad q_6 \qquad q$$

有:
$$\frac{1}{4\pi\varepsilon_0} \frac{q_1}{r_1} \Big|_{\rho=a} = \frac{q_1}{4\pi\varepsilon_0} \frac{1}{a} \Big|_{\rho=a} = -\frac{1}{4\pi r} \Big|_{\rho=a} \therefore q_1 = -\varepsilon_0 \frac{a}{\rho_0}$$

$$\frac{1}{4\pi\varepsilon_0} \frac{q_2}{r_2} \Big|_{\sigma} = -\frac{1}{4\pi r} \Big|_{\sigma} (\because \sigma : r_2 = r) \therefore q_2 = -\varepsilon_0$$
 显然, $q_3 = \varepsilon_0 \frac{a}{\rho_0}$

$$G = \frac{1}{4\pi} \left[\left(\frac{1}{r} - \frac{a/\rho_0}{r_1} \right) - \left(\frac{1}{r_2} - \frac{a/\rho_0}{r_3} \right) \right]$$

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二、用电像法求格林函数



2、求四分之一平面的狄氏格林函数

设
$$M_0(x_0,y_0)$$
点置有电荷 ε_0 ,则有

象点:
$$M_1(-x_0, y_0), M_2(x_0, -y_0), M_3(-x_0, -y_0)$$

电象:
$$-\varepsilon_0$$

$$\frac{1}{2\pi} \ln \frac{1}{\pi}$$
, $\frac{1}{2\pi} \ln \frac{1}{\pi}$

$$g_i$$
:

$$g_i: -\frac{1}{2\pi} \ln \frac{1}{r_1}, -\frac{1}{2\pi} \ln \frac{1}{r_2}, \frac{1}{2\pi} \ln \frac{1}{r_3}$$

$$G = \frac{1}{2\pi} \ln \frac{1}{r} + g$$

$$\begin{cases} \Delta g = 0, M \in \sigma \\ g|_{l} = -\frac{1}{2\pi} \ln \frac{1}{r}|_{l} \end{cases}$$

故
$$G = \frac{1}{2\pi} \ln \frac{1}{r} - \frac{1}{2\pi} \ln \frac{1}{r_1} - \frac{1}{2\pi} \ln \frac{1}{r_2} + \frac{1}{2\pi} \ln \frac{1}{r_3}$$

$$G = \frac{1}{2\pi} \ln \frac{r_1 r_2}{r r_3}$$



三、求泊松方程的狄氏问题

1、求上半空间的狄氏问题

$$\begin{cases} \Delta u = 0, \ z > 0 \\ u|_{z=0} = f(x,y) \end{cases} \rightarrow u(M) = -\iint_{\sigma} f(M_0) \frac{\partial G}{\partial n_0} dx_0 dy_0$$

$$\begin{cases} \Delta G = -\delta(M - M_0) \\ G|_{z=0} = 0 \end{cases} \qquad \begin{cases} G = \frac{1}{4\pi r} + g \\ \Delta g = 0, z > 0 \\ g|_{z=0} = -\frac{1}{4\pi r}|_{z=0} \end{cases}$$

$$(1) 在 M_1(x, y, -z) 放 - q, 则 \Delta(\frac{-q}{4\pi \varepsilon_0 r_1}) = 0 , \quad z > 0$$

$$(e \frac{-q}{4\pi \varepsilon_0 r_1}|_{z=0} = -\frac{\varepsilon_0}{4\pi \varepsilon_0 r}|_{z=0} \quad \text{則} \quad g = -\frac{q}{4\pi \varepsilon_0 r_1}$$

使
$$\frac{-q}{4\pi\varepsilon_0 r_1}|_{z=0} = -\frac{\varepsilon_0}{4\pi\varepsilon_0 r}|_{z=0}$$
 见 $g = -\frac{q}{4\pi\varepsilon_0 r_1}$

$$-q = -\varepsilon_0 \qquad \therefore G = \frac{1}{4\pi r} - \frac{1}{4\pi r_1}$$

问:下半空间狄氏问题?

三、求泊松方程的狄氏问G有无差别? $\frac{\partial G}{\partial G}$ 有无差别?

$$(2)\frac{\partial G}{\partial n} = \frac{\partial G}{\partial (-z)} = -\frac{1}{4\pi} \left[\frac{\partial}{\partial z} \frac{1}{r} - \frac{\partial}{\partial z} \frac{1}{r_1} \right]$$

$$\frac{\partial G}{\partial n} = \frac{\partial G}{\partial (-z)} = \frac{1}{4\pi} \frac{(z - z_0)}{\left[(x - x_0)^2 + (y - y_0)^2 + (z - z_0)^2 \right]^{\frac{3}{2}}}$$

$$-\frac{1}{4\pi} \frac{(z + z_0)}{\left[(x - x_0)^2 + (y - y_0)^2 + (z + z_0)^2 \right]^{\frac{3}{2}}}$$

$$\frac{\partial G}{\partial n} \Big|_{z=0} = \frac{1}{2\pi} \frac{-z_0}{\left[(x - x_0)^2 + (y - y_0)^2 + z_0^2 \right]^{\frac{3}{2}}}$$

$$u(M) = \frac{z}{2\pi} \int_{-\infty}^{\infty} f(x_0, y_0) \frac{1}{[(x - x_0)^2 + (y - y_0)^2 + z^2]^{\frac{3}{2}}} dx_0 dy_0$$



三、求泊松方程的狄氏问题

2、求上半平面的狄氏问题:

$$\begin{cases} \Delta u = 0 & , y > 0 \\ u|_{y=0} = f(x) \end{cases} \qquad u(x,y) = -\int_{-\infty}^{\infty} f(x_0) \frac{\partial G}{\partial n_0} dx_0$$

$$G = \frac{1}{2\pi} \ln \frac{1}{r} - \frac{1}{2\pi} \ln \frac{1}{r_1} = \frac{1}{2\pi} \ln \frac{r_1}{r}$$

$$r = \sqrt{(x - x_0)^2 + (y - y_0)^2},$$

$$r_1 = \sqrt{(x - x_0)^2 + (y + y_0)^2}$$

三、求泊松方程的狄氏
$$u(x,y) = -\int_{-\infty}^{\infty} f(x_0) \frac{\partial G}{\partial n_0} dx_0$$

$$G = \frac{1}{2\pi} \ln \frac{r_1}{r} = \frac{1}{4\pi} \left\{ \ln \left[(x - x_0)^2 + (y + y_0)^2 \right] - \ln \left[(x - x_0)^2 + (y - y_0)^2 \right] \right\}$$

$$\frac{\partial G}{\partial n} = \frac{\partial G}{\partial (-y)} = -\frac{\partial G}{\partial y}$$

$$\therefore \frac{\partial G}{\partial y}|_{y=0} = \frac{1}{4\pi} \left[\frac{2(y+y_0)}{(x-x_0)^2 + (y+y_0)^2} - \frac{2(y-y_0)}{(x-x_0)^2 + (y-y_0)^2} \right]_{y=0}$$

$$= \frac{y_0}{\pi} \left[\frac{1}{(x-x_0)^2 + y_0^2} \right]$$

$$u(x, y) = \frac{y}{\pi} \int_{-\infty}^{\infty} \frac{f(x_0)}{(x_0 - x)^2 + y^2} dx_0$$



三、求泊松方程的狄氏问题

3、用格林函数法重新求解

$$\begin{cases} \Delta_2 u = 0, & \rho < a \\ u \mid_{\rho = a} = A \cos \varphi \end{cases}$$

解:由圆域的狄氏积分公式有

$$u(\rho, \varphi) = \frac{1}{2\pi} \int_0^{2\pi} \frac{A\cos\varphi_0(a^2 - \rho^2)}{a^2 + \rho^2 - 2a\rho\cos(\varphi - \varphi_0)} d\varphi_0$$

将
$$\frac{分子/a^2}{分母/a^2}$$
,并记 $\varepsilon = \frac{\rho}{a}$

于是有:
$$u(\rho,\varphi) = \frac{(1-\varepsilon^2)A}{2\pi} \int_0^{2\pi} \frac{\cos\varphi_0}{1-2\varepsilon\cos(\varphi-\varphi_0)+\varepsilon^2} d\varphi_0$$

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三、求泊松方程的狄氏问题

f(z)



三、求泊松方程的狄氏问题

奇点
$$z = \begin{cases} \frac{(1+\varepsilon^2) \pm \sqrt{(1+\varepsilon)^2 - 4\varepsilon^2}}{2e^{-i\varphi}\varepsilon} = \begin{cases} \frac{e^{-i\varphi}}{\varepsilon} > 1\\ e^{-i\varphi} < 1 \end{cases}$$

$$\therefore \sum_{k=1}^{2} resf(z_k) = resf(0) + resf(\varepsilon e^{i\varphi})$$

$$= \frac{1}{\varepsilon} e^{-i\varphi} + \frac{1 + \varepsilon^2 e^{2i\varphi}}{\varepsilon e^{i\varphi}(\varepsilon^2 - 1)} = \frac{2\varepsilon}{\varepsilon^2 - 1} \cos \varphi$$

$$\therefore I = -\frac{1}{2i} \cdot 2\pi i \cdot \frac{2\varepsilon}{\varepsilon^2 - 1} \cos \varphi = \frac{2\pi\varepsilon}{1 - \varepsilon^2} \cos \varphi$$

$$u(\rho,\varphi) = \frac{\left[1 - (\rho/a)^2\right]}{2\pi} A \frac{2\pi \rho/a}{1 - (\rho/a)^2} \cos\varphi = \frac{A}{a} \rho \cos\varphi$$

三、求泊松方程的狄氏问题

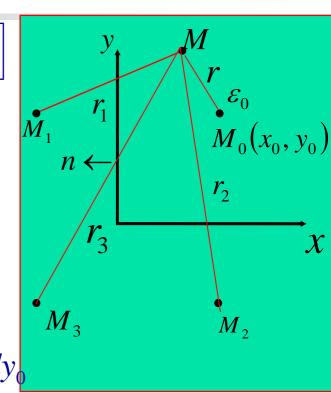
2、求四分之一平面的狄氏问题:

$$\begin{cases} u_{xx} + u_{yy} = 0 & , 0 \le x < \infty, 0 \le y < \infty \\ u(0, y) = f(y) & (0 \le y < \infty) \\ u(x, 0) = 0 \end{cases}$$

$$u(x,y) = -\int_{l} f(M_{0}) \frac{\partial G}{\partial n_{0}} dl_{0}$$

$$= -\int_{0}^{\infty} 0 \cdot \frac{\partial G}{\partial (-y_{0})} dx_{0} - \int_{0}^{\infty} f(y_{0}) \frac{\partial G}{\partial (-x_{0})} dy_{0}$$

$$u(x,y) = \int_0^\infty f(y_0) \frac{\partial G}{\partial x_0} dy_0 \qquad G = \frac{1}{2\pi} \ln \frac{r_1 r_2}{r r_3}$$





三、求泊松方程的狄氏问题

$$G = \frac{1}{2\pi} \ln \frac{r_1 r_2}{r r_3} = \frac{1}{4\pi} \left\{ \ln \left[(x + x_0)^2 + (y - y_0)^2 \right] + \left[(x - x_0)^2 + (y + y_0)^2 \right] - \left[(x + x_0)^2 + (y + y_0)^2 \right] \right\}$$

$$- \ln \left[(x - x_0)^2 + (y - y_0)^2 \right] - \left[(x + x_0)^2 + (y + y_0)^2 \right]$$

$$\therefore \frac{\partial G}{\partial x}|_{x=0} = \frac{1}{4\pi} \left[\frac{2(x+x_0)}{(x+x_0)^2 + (y-y_0)^2} + \frac{2(x-x_0)}{(x-x_0)^2 + (y+y_0)^2} - \frac{2(x-x_0)}{(x-x_0)^2 + (y-y_0)^2} - \frac{2(x+x_0)}{(x+x_0)^2 + (y+y_0)^2} \right]_{x=0}$$

$$= \frac{x_0}{\pi} \left[\frac{1}{x_0^2 + (y-y_0)^2} - \frac{1}{x_0^2 + (y+y_0)^2} \right]$$

$$u(x,y) = \frac{x}{\pi} \int_0^\infty \left[\frac{1}{x^2 + (y_0 - y)^2} - \frac{1}{x^2 + (y_0 + y)^2} \right] f(y_0) dy_0$$

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10.3-10.4 格林函数

Good-by!

