

Volumetric Lattice Boltzmann Models in General Curvilinear Coordinates: Theoretical Formulation

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Front. Appl. Math. Stat. 7:691582 doi: 10.3389/fams.2021.691582

Published: 16 June 2021

A volumetric formulation for lattice Boltzmann models on general curvilinear coordinates is presented, preserving one-to-one advection and recovering the Navier–Stokes equations in the hydrodynamic limit.

Presenter: Chen Peng Chung

Core Problem

In **general curvilinear coordinates**, macroscopic values $(\rho, \rho \vec{u})$ experience **momentum loss** due to:

- Curvature effects
- Nonuniform grid properties

Key Question

How to maintain **momentum conservation** in the streaming step of LBM when using curvilinear coordinates?

Solution Approach

Define a **velocity-space discretized force field** to compensate for the momentum loss (interpreted as an **inertial force**).

The standard LBGK equation in curvilinear coordinates:

$$N_{\alpha}(\vec{q} + \vec{e}_{\alpha}\delta t, t + 1) = N_{\alpha}(\vec{q}, t) + \Omega_{\alpha} + \delta N_{\alpha}(\vec{q}, t)$$

where the **correction term** is defined as:

$$\delta N_{\alpha}(\vec{q}, t) \equiv N_{\alpha}(\vec{q}, t) - N_{\alpha}^{\star}(\vec{q} - \vec{e}_{\alpha}\delta t, t - 1)$$

- \vec{q} : nondimensional position (q_1, q_2, q_3)
- N^{\star} : **post-collision** distribution function
- Ω_{α} : collision operator
- \vec{e}_{α} : discrete velocity directions

Mass Conservation (✓)

$$\sum_{\alpha=0}^{q-1} \delta N_{\alpha}(\vec{q}, t) = 0$$

Reason: Sum of distributions over velocity space is conserved.

Momentum Loss (✗)

$$\sum_{\alpha=1}^q \vec{e}_{\alpha} \delta N_{\alpha} \neq 0$$

Reason: Discrete velocity vectors \vec{e}_{α} vary with position in curvilinear coords!

Key Insight: $N_{\alpha}(\vec{q}, t) \neq N_{\alpha}^*(\vec{q} - \vec{e}_{\alpha}\delta t, t - 1)$

Pre-streaming \neq Post-collision at different positions

Definition: Nondimensional Curvature Change

Definition

The **nondimensional change of curvature tangent vector**:

$$\vec{\Theta}_i^j(\vec{q} + \vec{e}_\alpha \delta t, \vec{q}) = [\vec{g}_i(\vec{q} + \vec{e}_\alpha \delta t) - \vec{g}_i(\vec{q})] \cdot \vec{g}^j(\vec{q})$$

where the **contravariant basis vector** \vec{g}^j is:

$$\vec{g}^1 = \frac{\vec{g}_2 \times \vec{g}_3}{(\vec{g}_2 \times \vec{g}_3) \cdot \vec{g}_1}$$

- \vec{g}_i : covariant basis vectors (tangent to coordinate curves)
- \vec{g}^j : contravariant basis vectors (normal to coordinate surfaces)
- Θ_i^j measures how the tangent vector changes along \vec{e}_α

Definition

Total Momentum Loss at point (\vec{q}, t) :

$$\vec{M}_{\text{Loss}}(\vec{q}, t) \equiv \sum_{\alpha=1}^q [N_{\alpha}(\vec{q}, t) \vec{e}_{\alpha}(\vec{q}) - N_{\alpha}^{\star}(\vec{q} - \vec{e}_{\alpha}\delta t, t-1) \vec{e}_{\alpha}(\vec{q} - \vec{e}_{\alpha}\delta t)]$$

Physical Meaning:

- Momentum carried by particles **arriving** at \vec{q}
- Minus momentum carried by particles **leaving** from $\vec{q} - \vec{e}_{\alpha}\delta t$
- Difference arises because $\vec{e}_{\alpha}(\vec{q}) \neq \vec{e}_{\alpha}(\vec{q} - \vec{e}_{\alpha}\delta t)$

Inner & Outer Momentum Loss

The paper separates momentum loss into **two contributions**:

Definition

Inner Momentum Loss (particles entering):

$$\mathcal{J}(\vec{q}) \bar{\chi}^{\text{in}}(\vec{q}, t) \equiv - \sum_{\alpha=1}^q N_{\alpha}(\vec{q}, t) \cdot (\vec{e}_{\alpha}(\vec{q}) - \vec{e}_{\alpha}(\vec{q} - \vec{e}_{\alpha} \delta t))$$

Definition

Outer Momentum Loss (particles leaving):

$$\mathcal{J}(\vec{q}) \bar{\chi}^{\text{out}}(\vec{q}, t) \equiv - \sum_{\alpha=1}^q (\vec{e}_{\alpha}(\vec{q} + \vec{e}_{\alpha} \delta t) - \vec{e}_{\alpha}(\vec{q})) \cdot N_{\alpha}^{\star}(\vec{q}, t)$$

where \mathcal{J} is the Jacobian of the coordinate transformation.

The **discrete particle velocity set** in curvilinear coordinates:

$$\vec{e}_\alpha(\vec{q}) = \sum_{i=1}^3 (c_\alpha^i) \vec{g}_i \frac{\Delta x}{\Delta t}$$

- c_α^i : lattice velocity components (integers like 0, ± 1)
- \vec{g}_i : local covariant basis vectors
- $\Delta x / \Delta t$: lattice speed ratio

Important

Unlike Cartesian coordinates, \vec{e}_α depends on position \vec{q} through the basis vectors $\vec{g}_i(\vec{q})$.

The author imposes the **momentum constraint** for the streaming step:

$$\sum_{\alpha=1}^q \vec{e}_{\alpha}(\vec{q}) \delta N_{\alpha}(\vec{q}) = \mathcal{J} \cdot \frac{\vec{\chi}^{\text{in}}(\vec{q}, t) + \vec{\chi}^{\text{out}}(\vec{q}, t)}{2}$$

Expanding:

$$= \mathcal{J} \cdot \frac{-\sum_{\alpha} N_{\alpha}(\vec{q}) (\vec{e}_{\alpha}(\vec{q}) - \vec{e}_{\alpha}(\vec{q} - \vec{e}_{\alpha} \delta t)) - \sum_{\alpha} N_{\alpha}^*(\vec{q}) (\vec{e}_{\alpha}(\vec{q} + \vec{e}_{\alpha} \delta t) - \vec{e}_{\alpha}(\vec{q}))}{2}$$

- **Blue**: contribution from pre-collision distribution
- **Red**: contribution from post-collision distribution

Definition

The **velocity-space discretized force field** in curvilinear coordinates:

$$F^i(\vec{q}, t) = \frac{\vec{\chi}^{\text{in}}(\vec{q}, t) + \vec{\chi}^{\text{out}}(\vec{q}, t)}{2} \cdot \vec{g}^i(\vec{q})$$

Physical Interpretation:

- This force represents the **inertial force** arising from curvilinear coordinates
- Analogous to centrifugal/Coriolis forces in rotating frames
- Compensates for the “loss” in momentum during streaming

General Force Definition

$$\mathcal{J}(\vec{q})F^i(\vec{q}, t) \equiv \frac{\text{Momentum Loss}}{\Delta t} \cdot \vec{e}_j \frac{1}{|\vec{g}_j|}$$

After derivation, the force field becomes:

$$\mathcal{J}(\vec{q})F^j(\vec{q}, t) = \frac{-1}{2} \left(\sum_{\alpha} N_{\alpha} c_{\alpha}^i \Theta_i^j(\vec{q} - \vec{e}_{\alpha} \delta t, \vec{q}) + \sum_{\alpha} N_{\alpha}^{\star} c_{\alpha}^i \Theta_i^j(\vec{q} + \vec{e}_{\alpha} \delta t, \vec{q}) \right)$$

Key Components:

Symbol	Meaning
\mathcal{J}	Jacobian of transformation
c_{α}^i	Lattice velocity component
Θ_i^j	Curvature change tensor
N_{α}	Pre-collision distribution
N_{α}^{\star}	Post-collision distribution

Three-Step Derivation:

1 Define Momentum Loss:

$$\sum_{\alpha=1}^q \vec{e}_{\alpha}(\vec{q}) \delta N_{\alpha}(\vec{q}, t) \equiv \text{Momentum Loss}$$

2 Express via Inner/Outer contributions:

$$\sum_{\alpha=1}^q \vec{e}_{\alpha}(\vec{q}) \delta N_{\alpha} \equiv \mathcal{J}(\vec{q}) \cdot \frac{\vec{\chi}^{\text{in}} + \vec{\chi}^{\text{out}}}{2}$$

3 Project onto coordinate directions:

$$\mathcal{J}(\vec{q}) F^i \equiv \frac{\text{Momentum Loss}}{\Delta t} \cdot \vec{e}_j \frac{1}{|\vec{g}_j|}$$

Correction term:

$$\delta N_{\alpha}(\vec{q}, t) = w_{\alpha} \left(\frac{c_{\alpha}^i F^i}{c_s^2} + \frac{c_{\alpha}^i c_{\alpha}^j - c_s^2 \delta^{ij}}{2c_s^4} \right)$$

Moment constraints (for Chapman-Enskog \rightarrow Navier-Stokes):

$$\sum_{\alpha} f^{\alpha q} = \rho$$

$$\sum_{\alpha} f^{\alpha q} c_{\alpha}^i = \rho u^i$$

$$\sum_{\alpha} f^{\alpha q} c_{\alpha}^i c_{\alpha}^j = \rho c_s^2 g^{ij} + \rho \left(u^i + \frac{F^i}{2\rho} \right) \left(u^j + \frac{F^j}{2\rho} \right)$$

Note: The force F^i enters the velocity shift: $u^i \rightarrow u^i + \frac{F^i}{2\rho}$

Summary & Key Takeaways

Main Contributions

- 1 Identified **momentum loss** in streaming step for curvilinear LBM
- 2 Introduced **curvature change tensor** Θ_i^j to quantify geometric effects
- 3 Derived **velocity-space discretized force field** to restore conservation

Physical Insight

The force field F^i acts as an **inertial force** (like centrifugal force), arising purely from the curvilinear coordinate system.

Practical Use

Add correction term δN_α to LBGK equation:

$$N_\alpha(\vec{q} + \vec{e}_\alpha \delta t, t + 1) = N_\alpha(\vec{q}, t) + \Omega_\alpha + \delta N_\alpha$$