Letter to the Editor

Another integrable case in the Lorenz model

Tat-Leung Yee* and Robert Conte Service de physique de l'état condensé (URA 2464) CEA-Saclay, F-91191 Gif-sur-Yvette Cedex, France

E-mail: TonYee@ust.hk and Conte@drecam.saclay.cea.fr

To appear in Journal of Physics A. 13 February 2004

Abstract

A scaling invariance in the Lorenz model allows one to consider the usually discarded case $\sigma = 0$. We integrate it with the third Painlevé function.

Keywords: Lorenz model, first integral, third Painlevé equation.

PACS 1995: 02.30.-f, 05.45.+b, 47.27.-i,

1 Introduction

The Lorenz model [1]

$$\frac{\mathrm{d}x}{\mathrm{d}t} = \sigma(y - x), \quad \frac{\mathrm{d}y}{\mathrm{d}t} = rx - y - xz, \quad \frac{\mathrm{d}z}{\mathrm{d}t} = xy - bz, \tag{1}$$

in which (b, σ, r) are real constants, is a prototype of chaotic behaviour [3]. In particular, it fails the Painlevé test unless the parameters obey the constraints [4]

$$Q_{2} \equiv (b - 2\sigma)(b + 3\sigma - 1) = 0,$$

$$\forall x_{2}: Q_{4} \equiv -4i(b - \sigma - 1)(b - 6\sigma + 2)x_{2} - (4/3)(b - 3\sigma + 5)b\sigma r$$

$$+ (-4 + 10b + 30b^{2} - 20b^{3} - 16b^{4})/27$$

$$+ (-38b - 56b^{2} - (28/3)b^{3} + 88\sigma + 86b^{2}\sigma)\sigma/3$$

$$-32\sigma/9 + 70b\sigma^{2} - 64\sigma^{3} - 58b\sigma^{3} + 36\sigma^{4} = 0.$$
(2)

This system (2)–(3) depends on r only through the product $b\sigma r$, as a consequence of an obvious scaling invariance in the model, and it admits four solutions,

$$(b, \sigma, b\sigma r) = (1, 1/2, 0), (2, 1, 2/9), (1, 1/3, 0), (1, 0, 0). \tag{4}$$

In the first three cases, i.e. when the system (1) is nonlinear, which excludes $\sigma = 0$, the system can be explicitly integrated [4], and the general solution (x, y, z) is a singlevalued function of time expressed with, respectively, an elliptic function, the second and the third Painlevé functions.

In this letter, we consider the fourth case

$$(b, \sigma, r) = (1, 0, r).$$
 (5)

The apparently linear nature of the dynamical system can be removed by eliminating y and z and considering the third order differential equation for x(t) [5],

$$y = x + x'/\sigma, \ z = r - 1 - [(\sigma + 1)x' + x'']/(\sigma x),$$

$$xx''' - x'x'' + x^3x' + \sigma x^4 + (b + \sigma + 1)xx'' + (\sigma + 1)(bxx' - x'^2)$$

$$+b(1-r)\sigma x^2 = 0,$$
(6)

^{*}Permanent address: Department of Mathematics, The Hong-Kong university of science and technology, Clear Water Bay, Kowloon, Hong Kong. S2004/003.

which also depends on r only through the product $b\sigma r$, and thus implements the above mentioned scaling invariance. The necessary conditions for (7) to pass the Painlevé test are the same ($Q_2 = 0, Q_4 = 0$) as for the dynamical system (1), the restriction $\sigma \neq 0$ being now removed.

2 Integration for $b = 1, \sigma = 0$

Because of the scaling invariance, the following first integral [4] of the dynamical system (1),

$$(b, \sigma, r) = (1, \sigma, 0) : K_3 = (y^2 + z^2)e^{2t},$$
 (8)

is also a first integral of the third order equation for $(b, \sigma, b\sigma r) = (1, \sigma, 0)$, which includes the particular case of interest to us $(b, \sigma, b\sigma r) = (1, 0, 0)$,

$$(b, \sigma, b\sigma r) = (1, 0, 0): K^2 = \lim_{\sigma \to 0} \sigma^2 K_3 = \left[\left(\frac{x'' + x'}{x} \right)^2 + x'^2 \right] e^{2t}.$$
 (9)

For K = 0, the general solution is

$$x = ik \tanh \frac{k}{2}(t - t_0) - i, \ i^2 = -1, \ (k, t_0) \text{ arbitrary.}$$
 (10)

For $K \neq 0$, after taking the usual parametric representation

$$\frac{x'' + x'}{x} = Ke^{-t}\cos\lambda, \ x' = Ke^{-t}\sin\lambda,\tag{11}$$

the second order ODE for $\lambda(t)$ is found to be

$$\lambda'' - Ke^{-t}\sin\lambda = 0, (12)$$

with the link

$$x(t) = \lambda'(t). \tag{13}$$

In the variable $\cos \lambda$, the differential equation (12) becomes algebraic and belongs to an already integrated class [2]. The overall result is the general solution

$$x = i + 2i\frac{\mathrm{d}}{\mathrm{d}t} \log w(\xi(t)), \ i^2 = -1, \ \xi = ae^{-t},$$
 (14)

in which $w(\xi)$ is the particular third Painlevé function defined by

$$\frac{\mathrm{d}^2 w}{\mathrm{d}\xi^2} = \frac{1}{w} \left(\frac{\mathrm{d}w}{\mathrm{d}\xi}\right)^2 - \frac{\mathrm{d}w}{\xi \mathrm{d}\xi} + \frac{\alpha w^2 + \gamma w^3}{4\xi^2} + \frac{\beta}{4\xi} + \frac{\delta}{4w},\tag{15}$$

$$\alpha = 0, \ \beta = 0, \ \gamma \delta = -(K/a)^2. \tag{16}$$

3 Conclusion

Out of the two cases selected by the condition $Q_2 = 0$, one admits a first integral [4],

$$b = 2\sigma: K_1 = (x^2 - 2\sigma z)e^{2\sigma t},$$
 (17)

but, in the second case $b=1-3\sigma$, the first integral whose existence has been conjectured [6] is not yet known. The present result, which belongs to this unsettled case $b=1-3\sigma$, should help to solve this open question.

Acknowledgments

T.-L. Yee thanks the Croucher Foundation for a postdoc grant at CEA. R. Conte acknowledges the support of the France-Hong Kong PROCORE grant 04807SF, which allowed starting this work.

References

- [1] E. N. Lorenz, Deterministic nonperiodic flow, J. Atmos. Sci. 20 (1963) 130–141.
- [2] B. Gambier, Sur les équations différentielles du second ordre et du premier degré dont l'intégrale générale est à points critiques fixes, Acta Math. **33** (1910) 1–55.
- [3] J. Guckenheimer and P. Holmes, Nonlinear oscillations, dynamical systems and bifurcations of vector fields, Applied mathematical sciences 42 (Springer-Verlag, Berlin, 1983).
- [4] H. Segur, Solitons and the inverse scattering transform, *Topics in ocean physics*, 235–277, eds. A. R. Osborne and P. Malanotte Rizzoli (North-Holland publishing co., Amsterdam, 1982).
- [5] T. Sen and M. Tabor, Lie symmetries of the Lorenz model, Physica D 44 (1990) 313–339.
- [6] M. Tabor and J. Weiss, Analytic structure of the Lorenz system, Physical Review A 24 (1981) 2157–2167.