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CS7GV2: Mathematics of Light and Sound

Lecture #4: Simulation

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October 30, 2020

Analytical versus numerical methods

- ▶ For a quadratic polynomial $f(x) = a x^2 + b x + c$, the roots (zero-crossings) are found with the well-known formula,

$$x_1, x_2 = \frac{-b \pm \sqrt{b^2 - 4 a c}}{2 a}.$$

- ▶ In science and engineering it's more-often-than-not the case that problems do not have neat *closed-form* or *analytical* solutions except in very specific circumstances.
- ▶ What can we do about it? Approximation, iteration. For example, the “method of bisection” for root finding: guess where a root might be; keep halving the length of an interval around it such that $f(x)$ has different signs at the start and the end.
- ▶ Such solutions often described as *numerical methods* because they use numbers (and computers) versus *analytical methods* which use symbols (and thinking.)

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Analytic or
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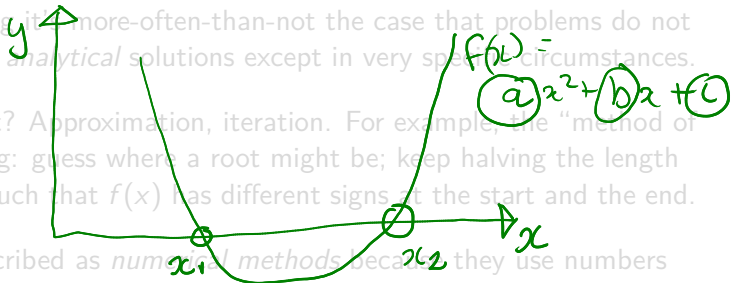
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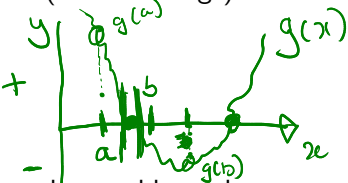
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HIDDEN FIGURES!

Wave Motion

- ▶ We've seen that wave motion is described by the second order PDE known as the wave equation,

$$\frac{\partial^2 u(x, t)}{\partial t^2} = c^2 \frac{\partial^2 u(x, t)}{\partial x^2}.$$

- ▶ We've seen a *closed-form* solution for wave propagation,

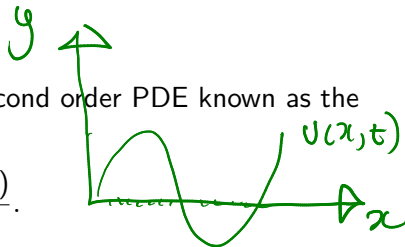
$$u(x, t) = R \cos(kx - \omega t) + (1 - R) \cos(kx + \omega t).$$

- ▶ This is perfect when there are no constraints. For example, light in a homogeneous medium, a wave on an infinitely long string (no end points,) or a sound in a huge volume of air.
- ▶ But the closed-form solution doesn't tell us, for example, how a string plucked in a particular way is going to move: <https://tinyurl.com/y4ncymx7>.

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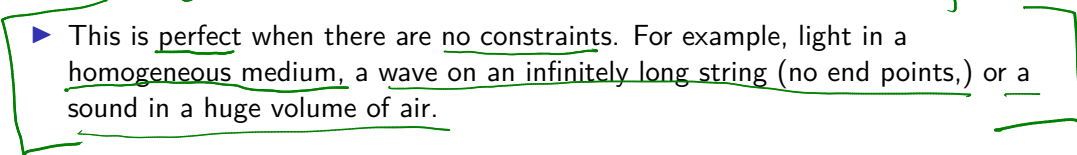
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Simulation! numerical methods

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Wave Simulation



- ▶ When there are specific constraints (also known as conditions,) there is usually no alternative but to simulate wave motion in an iterative way.
- ▶ Iterative means doing more-or-less the same sequence of calculations again and again.
- ▶ Usually the current iteration's calculations use results calculated in the previous iteration(s.)
- ▶ An iterative simulation can never be perfect. Error is inevitable, for example, because discretization is required.
- ▶ Error is typically cumulative so the results become less correct at each iteration.
- ▶ There are lots of nice interactive simulations of wave motions available, for example: <https://tinyurl.com/2xrsrz> and <https://tinyurl.com/mtwczmj>.

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after a certain amount
of computation

Wave Simulation

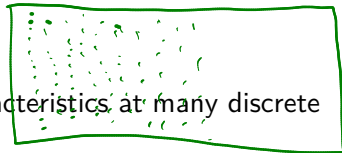
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EM Wave simulation

- ▶ Solve Maxwell's equations to find local wave characteristics at many discrete volumes of space at successive steps in time.
- ▶ The results for one discrete volume are used in the calculation of the characteristics of its neighbors.
- ▶ One of the most used techniques (e.g. in MEEP) is called *finite difference time domain* (FDTD.)
- ▶ Approaches like this in general are called *finite element methods* for the approximate solution of *boundary value problems* with *partial differential equations*.
- ▶ Advantages: can deal with complex geometries and different materials.
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cf. String Simulation



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Simulate weather < 3 days

Initial and Boundary Conditions

- ▶ To simulate a specific solution for $u(x, t)$ described by the wave equation,



$$\frac{\partial^2 u(x, t)}{\partial t^2} = c^2 \frac{\partial^2 u(x, t)}{\partial x^2} \quad x \in [0, L], \quad t \in [0, T],$$

for a string of length L over a time period T , we need:

- ▶ two *initial conditions* at time $t = 0$,

$$u(x, 0) = I(x), \quad x \in [0, L]$$

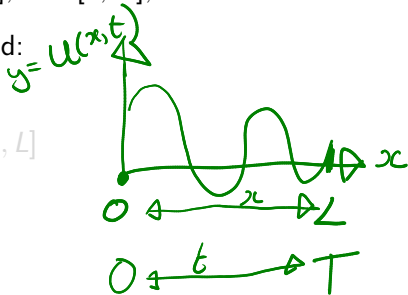
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where $I(x)$ specifies the initial shape of the string,

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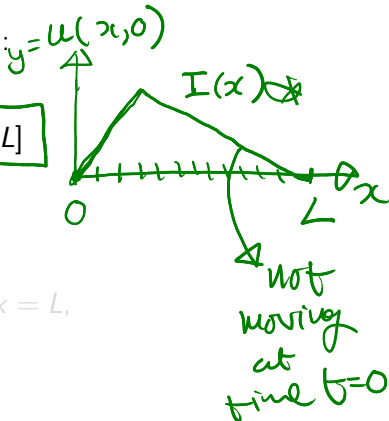
$$\begin{cases} u(x, 0) = I(x), & x \in [0, L] \\ \frac{\partial u(x, 0)}{\partial t} = 0, & x \in [0, L] \end{cases}$$

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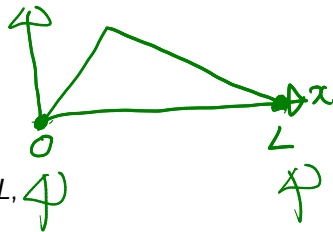
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$y = u(x, t)$



Discretization of domain

- ▶ Computer operations take a finite amount of time to complete so there can't be infinitely many time steps in the simulation.

The time period $[0, T]$ has to be discretized, e.g. into intervals of equal duration Δt ,

$$t_i = i \Delta t, \quad i = 0, \dots, N_t \text{ (where } N_t = T/\Delta t.)$$

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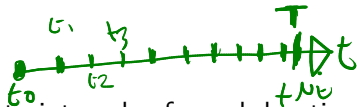
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$$T = 5s$$



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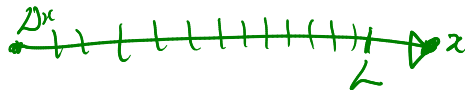
to keep calculation + memory requirements finite

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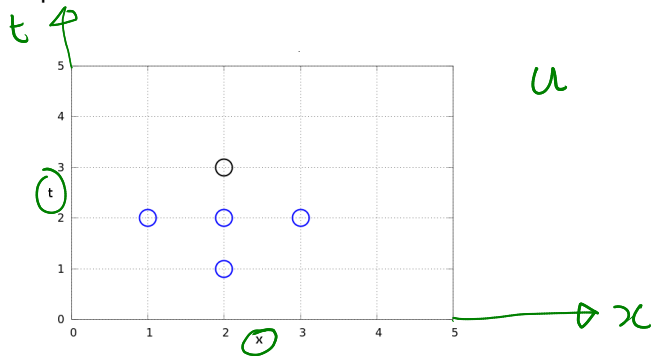


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Solution mesh

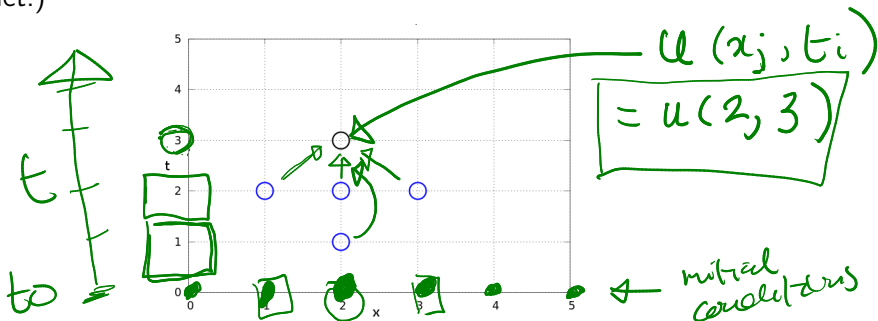
- The discrete points in space and time can be visualized as a two-dimensional *mesh* (or net.)



- The solution for wave height $u(x_j, t_i)$ at each mesh point is found using already-calculated solutions at neighbouring mesh points ...
- ... except for certain exterior mesh points whose values have been specified through the initial conditions, i.e. $I(x)$.

Solution mesh

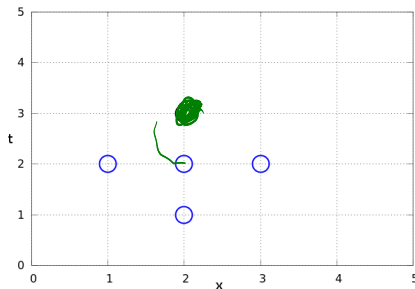
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Discretization of equations

- Wave equation. Use the symmetric second difference approximation of the second derivative,

$$\frac{\partial^2 u(x,t)}{\partial t^2}$$

$$\frac{u(x_j, t_{i+1}) - 2u(x_j, t_i) + u(x_j, t_{i-1}))}{\Delta t^2} \approx c^2 \frac{u(x_{j+1}, t_i) - 2u(x_j, t_i) + u(x_{j-1}, t_i)}{\Delta x^2}$$

$y = u(x, t)$

$\frac{\partial^2 u(x, t)}{\partial^2 x}$

Alternative notation can be used to make the parameters more obvious,

$$\frac{u_j^{i+1} - 2u_j^i + u_j^{i-1}}{\Delta t^2} \approx c^2 \frac{u_{j+1}^i - 2u_j^i + u_{j-1}^i}{\Delta x^2}, \quad (1)$$

- Initial condition. Use the centered first difference approximation of the first derivative,

$$\frac{\partial}{\partial t} u(x_j, t_i) \approx \frac{u_j^{i+1} - u_j^{i-1}}{2\Delta t} \quad (2)$$

Note division by $2\Delta t$ because the difference is between values of $u(x, t)$ separated by two time intervals.

Discretization of equations

- *Wave equation.* Use the symmetric second difference approximation of the second derivative,

$$\frac{u(x_j, t_{i+1}) - 2u(x_j, t_i) + u(x_j, t_{i-1}))}{\Delta t^2} \approx c^2 \frac{u(x_{j+1}, t_i) - 2u(x_j, t_i) + u(x_{j-1}, t_i)}{\Delta x^2}.$$

Handwritten notes: A green checkmark is next to the equation. To the right, a green bracket groups the terms $u(x_j, t_{i+1})$ and $u(x_j, t_{i-1})$, with a note $= u_j^{i+1}$ pointing to the first term.

Alternative notation can be used to make the parameters more obvious,

$$\frac{u_j^{i+1} - 2u_j^i + u_j^{i-1}}{\Delta t^2} \approx c^2 \frac{u_{j+1}^i - 2u_j^i + u_{j-1}^i}{\Delta x^2}, \quad (1)$$

Handwritten notes: A green checkmark is next to the equation. A green bracket is on the right side of the equation.

- *Initial condition.* Use the centered first difference approximation of the first derivative,

$$\frac{\partial}{\partial t} u(x_j, t_i) \approx \frac{u_j^{i+1} - u_j^{i-1}}{2\Delta t} \quad (2)$$

Note division by $2\Delta t$ because the difference is between values of $u(x, t)$ separated by two time intervals.

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Initial Conditions

- ▶ Using approximation (2), initial condition $\frac{\partial}{\partial t} u(x_j, 0) = 0$ means,

$$u_j^{i-1} = u_j^{i+1}, \quad j = 0, \dots, N_x. \quad i = 0.$$

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Formulae

$$u(x_j, t_{i+1})$$

LHS = RHS

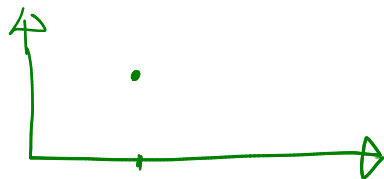
$$u_j^{i+1} = -u_j^{i-1} + 2u_j^i + C^2 (u_{j+1}^i - 2u_j^i + u_{j-1}^i)$$

$$C = c \frac{\Delta t}{\Delta x}.$$

$$u_j^1 = u_j^0 - \frac{1}{2} C^2 (u_{j+1}^0 - 2u_j^0 + u_{j-1}^0)$$

RHS terms
are
blue

look at
solution
mesh



where did
this
come from?

Rearranged
wave
equation

t_{i-1}
 t_i
 t_{i+1}

Iterative Simulation Algorithm

1. Initialize $u_j^0 = I(x_j)$ for $j = 0, \dots, N_x$.
2. Compute u_j^1 and set $u_j^1 = 0$ for the boundary points $i = 0$ and $i = N_x$, for $i = 1, \dots, N - 1$
3. For each time level $i = 1, \dots, N_t - 1$
 - 3.1 find u_j^{i+1} for $j = 1, \dots, N_x - 1$.
 - 3.2 set $u_j^{i+1} = 0$ for the boundary points $j = 0, j = N_x$.

Write a program in SciPy
to simulate string
motion
using this
algorithm!