

# Quantum Physics II Spring 2018

## Midterm Exam

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**You must show your work. No credits will be given if you don't show how you get your answers.**

**You may use the following formula:**

- One can determine how a state  $|\psi\rangle$  evolves with time

$$\psi(x', t') = \int_{-\infty}^{\infty} dx_0 \langle x', t' | x_0, t_0 \rangle \psi(x_0, t_0)$$

if one knows the amplitude  $\langle x', t' | x_0, t_0 \rangle$ :

$$\langle x', t' | x_0, t_0 \rangle = \int_{x_0}^{x'} D[x(t)] \exp\left(i \frac{1}{\hbar} S[x(t)]\right),$$

where  $D[x(t)]$  means summing over all possible path  $x(t)$  between  $(x_0, t_0)$  and  $(x', t')$ .  $S[x(t)]$  is the action for a path  $x(t)$ :

$$S[x(t)] = \int_{t_0}^{t'} dt L(x, \dot{x}),$$

and  $L(x, \dot{x})$  being the Lagrangian,

$$L(x, \dot{x}) = \frac{1}{2} m \dot{x}^2 - V(x).$$

The Planck constant  $\hbar = 1.055 \times 10^{-34}$  Joule  $\cdot$  sec.

- The Schrödinger equation:

$$i\hbar \frac{d|\psi(t)\rangle}{dt} = \hat{H} |\psi(t)\rangle$$

where  $\hat{H}$  is the Hamiltonian  $\frac{\hat{\mathbf{p}}^2}{2\mu} + V(\hat{\mathbf{r}})$  and  $\mu$  is the mass of the particle.

For energy eigenstates, this reduces to the time-independent Schrödinger equation (in spherical coordinates)

$$-\frac{\hbar^2}{2\mu}\left(\frac{\partial^2}{\partial r^2} + \frac{2}{r}\frac{\partial}{\partial r}\right)\psi(\mathbf{r}) + \langle \mathbf{r} | \frac{\hat{\mathbf{L}}^2}{2\mu r^2} | \psi \rangle + V(\mathbf{r})\psi(\mathbf{r}) = E\psi(\mathbf{r})$$

$\hat{\mathbf{L}}$  is the orbital angular momentum operator,  $\hat{\mathbf{L}} = \hat{\mathbf{r}} \times \hat{\mathbf{p}}$ , which satisfies the usual commutation relations for angular momenta, e.g.  $[\hat{L}_x, \hat{L}_y] = i\hbar\hat{L}_z$ . In the spherical coordinate,  $\hat{L}_z \rightarrow -i\hbar\frac{\partial}{\partial\phi}$ .

- For spherical symmetric potential  $V(r)$ , there exist simultaneous eigenstates of  $\hat{H}$ ,  $\hat{\mathbf{L}}^2$ , and  $\hat{L}_z$ :

$$\hat{H} |E, l, m\rangle = E |E, l, m\rangle, \quad \hat{\mathbf{L}}^2 |E, l, m\rangle = l(l+1)\hbar^2 |E, l, m\rangle, \quad \hat{L}_z |E, l, m\rangle = m\hbar |E, l, m\rangle.$$

- In terms of wavefunctions,  $|E, l, m\rangle$  is of the form  $R(r)Y_{l,m}(\theta, \phi)$ , where the **spherical harmonics**  $Y_{l,m}(\theta, \phi) = \langle \theta, \phi | l, m \rangle$ , and

$$\int d\Omega |Y_{l,m}(\theta, \phi)|^2 = 1,$$

with  $d\Omega$  being the solid angle.

- Given one  $|l, m\rangle$ , one can obtain other eigenkets with the same  $l$  by the ladder operators  $\hat{L}_{\pm} = \hat{L}_x \pm i\hat{L}_y$ :

$$\hat{L}_{\pm} |l, m\rangle = \sqrt{l(l+1) - m(m \pm 1)}\hbar |l, m \pm 1\rangle.$$

- The parity operator  $\hat{\pi}$  acts on a ket  $|\psi\rangle$  and flips the sign of the position expectation value:

$$|\psi'\rangle = \hat{\pi} |\psi\rangle, \quad \langle x \rangle_{\psi'} = -\langle x \rangle_{\psi}$$

From this definition one can show that  $\hat{\pi}\hat{x}\hat{\pi} = -\hat{x}$ ,  $\hat{\pi}|x\rangle = |-x\rangle$ , and  $\hat{\pi} = \hat{\pi}^{-1} = \hat{\pi}^{\dagger}$ . The eigenvalues of  $\hat{\pi}$  are  $\pm 1$ . The wavefunctions of its eigenkets have the property  $\psi(-x) = \pm\psi(x)$ , and thus they are called parity-even or parity-odd states.

- For a 1-D simple harmonic oscillator (SHO),  $\hat{H} = \frac{\hat{p}_x^2}{2m} + \frac{1}{2}m\omega^2\hat{x}^2$ :

The raising and lowering operators are

$$\hat{a}^{\dagger} = \sqrt{\frac{m\omega}{2\hbar}}\left(\hat{x} - \frac{i}{m\omega}\hat{p}_x\right), \quad \hat{a} = \sqrt{\frac{m\omega}{2\hbar}}\left(\hat{x} + \frac{i}{m\omega}\hat{p}_x\right)$$

and  $[\hat{a}, \hat{a}^{\dagger}] = 1$ . The operators get their names from the facts that

$$\hat{a}^{\dagger} |n\rangle = \sqrt{n+1} |n+1\rangle, \quad \hat{a} |n\rangle = \sqrt{n} |n-1\rangle.$$

One can rewrite  $\hat{x}$  and  $\hat{p}_x$  as

$$\hat{x} = \sqrt{\frac{\hbar}{2m\omega}}(\hat{a} + \hat{a}^\dagger), \quad \hat{p}_x = -i\sqrt{\frac{m\omega\hbar}{2}}(\hat{a} - \hat{a}^\dagger)$$

and the Hamiltonian as

$$\hat{H} = \hbar\omega(\hat{a}^\dagger\hat{a} + \frac{1}{2}) = \hbar\omega(\hat{N} + \frac{1}{2}),$$

where the number operator  $\hat{N} \equiv \hat{a}^\dagger\hat{a}$ .

- The energy eigenvalues of an SHO are  $E_n = (n + \frac{1}{2})\hbar\omega$ ,  $n = 0, 1, 2, \dots$
- Expectation value of an operator  $\hat{A}$  for a state  $|\psi\rangle$  is  $\langle A \rangle = \langle \psi | \hat{A} | \psi \rangle$
- Uncertainty of an operator  $\hat{A}$  is defined as  $\Delta A = \sqrt{\langle A^2 \rangle - \langle A \rangle^2}$

1. **Gravity-induced Quantum interference:** (2 points each) A beam of neutrons goes into the neutron interferometer as sketched in Fig. 1, where it gets split at A into two beams, then merged again at D. The lengths of each side of the interferometer  $l_1$  and  $l_2$  are of the order of a few cm. The mass of neutrons is about  $1.67 \times 10^{-27}$  kg. The de Broglie wavelength of the neutron beam is  $\lambda = 1.42 \text{ \AA}$ .

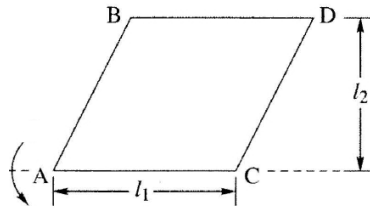


Figure 1: Problem 1

- (a) Can we approximate the paths of the two neutron beams with classical trajectories ABD and ACD? Explain why.
  - (b) Suppose the neutrons beams have no phase difference at D when the interferometer lies on a horizontal plane. Then we rotate around  $\overline{AC}$  until the interferometer becomes vertical. What will be the phase difference of the two beams at D?
2. **Properties of orbital angular momentum:** (2 points each) A particle is in an eigenstate of the orbital angular momentum  $|l, m\rangle$ .
    - (a) Evaluate  $\langle L_x \rangle$  and  $\langle L_y \rangle$  for this state.  
[Hint: Use  $\hat{L}_+$  and  $\hat{L}_-$  operators.]
    - (b) Evaluate  $\langle L_x^2 \rangle$  and  $\langle L_y^2 \rangle$ .  
[Hint: Again you can use  $\hat{L}_+$  and  $\hat{L}_-$  operators. Another approach is to use the symmetry of  $L_x$  and  $L_y$  for this state.]
    - (c) What will  $|l, m\rangle$  be if the state satisfies the equality of the uncertainty limit  $\Delta L_x \Delta L_y \geq \frac{\hbar}{2} |\langle L_z \rangle|$ ?  
[Note: Partial credits will be given if you don't get part (b) right, but argue the correct answer for part (c) with physics intuition.]
  3. **Energy of a rigid motor:** (2 points each) Consider a spherically symmetric rigid motor (like a ball) with moment of inertia  $I_x = I_y = I_z = I$ . Suppose its center-of-mass is at rest.

- (a) What is the Hamiltonian?
- (b) What are the wavefunctions of the energy eigenstates?
- (c) What are the energy levels and degeneracies (i.e. For a given energy level, how many eigenstates have the same energy eigenvalue)?
- (d) Suppose the rigid body gets stretched along the  $z$ -axis, so that the moment of inertia about  $z$  becomes  $I_z = I' \neq I$ . Show that  $[\hat{H}, \hat{L}_z] = 0$  so the energy eigenstates remain the same as in (b).
- (e) Are the eigenstates from (d) non-degenerate? Explain why or why not using the symmetry of the system.

4. **Parity of SHO eigenstates:** (2 points each)

- (a) Let  $\hat{O}$  be an operator with definite parity  $\eta$ , i.e.  $\hat{\pi}\hat{O}\hat{\pi} = \eta\hat{O}$  with  $\eta = \pm 1$ . Let  $|\psi_\lambda\rangle$  be an eigenket of parity with eigenvalue  $\lambda$ ,  $\lambda = \pm 1$ . Show that  $|\psi\rangle = \hat{O}|\psi_\lambda\rangle$  is an eigenket of parity with eigenvalue  $\eta\lambda$ .
- (b) Show that the wavefunction of the ground state  $|0\rangle$  for a 1-D SHO is Gaussian (you don't need to derive the normalization constant).  
[Hint: Use the property  $\hat{a}|0\rangle = 0$ . Note that  $\hat{p}_x \rightarrow -i\hbar \frac{d}{dx}$  in 1-D.]
- (c) Given the fact  $|0\rangle$  is Gaussian, show that the parity of each eigenstate  $|n\rangle$  of a 1-D SHO is  $(-1)^n$ .  
[Hint: How do you derive  $|n\rangle$  from  $|0\rangle$ ?]

5. **Schrödinger equation in the presence of EM field:** (2 points each) We've shown in class that the canonical momentum (to be identified with  $-i\hbar\nabla$  in the coordinate space) when there exists a magnetic vector potential  $\mathbf{A}$  is

$$\mathbf{p}_c = m\mathbf{v} + \frac{q}{c}\mathbf{A}$$

And the Hamiltonian is

$$H = \frac{(\mathbf{p}_c - q\mathbf{A}/c)^2}{2m} + q\phi$$

where  $\phi$  is the electric potential.

- (a) We know that different  $\mathbf{A}$  and  $\phi$  can result in the same  $\mathbf{E}$  and  $\mathbf{B}$  fields, called the “gauge transformation” of the potentials. Please specify whether the following quantities are gauge-dependent or gauge-invariant:  $\mathbf{p}_c$ ,  $m\mathbf{v}$ , and the wavefunction  $\psi(\mathbf{r}, t)$ . Which of them is/are physical observable (something you can measure in an experiment)?

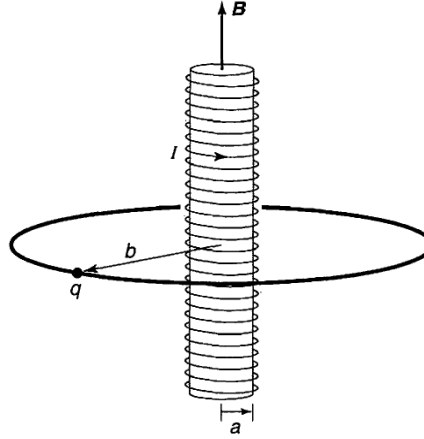


Figure 2: Problem 5

Consider a spinless particle with mass  $m$  and charge  $q$ . It is constrained to move along a ring of radius  $b$  around an ideal solenoid of radius  $a$  ( $a < b$ ) around the  $z$ -axis, as shown in Fig. 2.

The magnetic field  $\mathbf{B}$  is zero along its path, but the vector potential  $\mathbf{A}$  is not.

- (b) [Bonus] One choice of  $\mathbf{A}$  is  $\mathbf{A} = \frac{\Phi_B}{2\pi r} \hat{\phi}$ , where  $r$  is the perpendicular distance of the particle to the solenoid ( $r > a$ ), and  $\Phi_B$  is the magnetic flux through the solenoid. Hence the Hamiltonian depends only on  $\phi$ . Show that the time-independent Schrödinger equation can be written as

$$\frac{1}{2m} \left( -i\hbar \frac{1}{b} \frac{\partial}{\partial \phi} - \frac{q}{c} \frac{\Phi_B}{2\pi b} \right)^2 \psi = E\psi$$

[Hint: Recall that  $\hat{\mathbf{L}} = \hat{\mathbf{r}} \times \hat{\mathbf{p}}$ , and what you need is  $\hat{p}_\phi$ .]

- (c) Argue that the solution of the energy eigenfunction is of the form  $e^{in\phi}$  where  $n$  must be an integer (explain why). Then plug this solution into the Schrödinger equation, and show that the energy eigenvalues are

$$E_n = \frac{\hbar^2}{2mb^2} \left( n - \frac{q\Phi_B}{2\pi\hbar c} \right)^2$$

In other words, the allowed energy values depend on the field *inside* the solenoid, despite that the particle feels no magnetic force along its way!