## Quantum Physics II Spring 2019 Midterm Exam

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You must show your work. No credits will be given if you don't show how you get your answers.

You may use the following formula:

• One can determine how a state  $|\psi\rangle$  evolves with time

$$\psi(x',t') = \int_{-\infty}^{\infty} dx_0 \left\langle x',t' \middle| x_0,t_0 \right\rangle \psi(x_0,t_0)$$

if one knows the amplitude  $\langle x', t' | x_0, t_0 \rangle$ :

$$\langle x', t' | x_0, t_0 \rangle = \int_{x_0}^{x'} D[x(t)] \exp\left(i\frac{1}{\hbar}S[x(t)]\right),$$

where D[x(t)] means summing over all possible path x(t) between  $(x_0, t_0)$  and (x', t'). S[x(t)] is the action for a path x(t):

$$S[x(t)] = \int_{t_0}^{t'} dt L(x, \dot{x}),$$

and  $L(x, \dot{x})$  being the Lagrangian,

$$L(x, \dot{x}) = \frac{1}{2}m\dot{x}^2 - V(x).$$

The Planck constant  $\hbar = 1.055 \times 10^{-34}$  Joule · sec.

• The Schrödinger equation:

$$i\hbar \frac{d|\psi(t)\rangle}{dt} = \hat{H}|\psi(t)\rangle$$

where  $\hat{H}$  is the Hamiltonian  $\frac{\hat{\mathbf{p}}^2}{2\mu} + V(\hat{\mathbf{r}})$  and  $\mu$  is the mass of the particle.

For energy eigenstates, this reduces to the time-independent Schrödinger equation (in spherical coordinates)

$$-\frac{\hbar^2}{2\mu}(\frac{\partial^2}{\partial r^2} + \frac{2}{r}\frac{\partial}{\partial r})\psi(\mathbf{r}) + \langle \mathbf{r} | \frac{\hat{\mathbf{L}}^2}{2\mu r^2} | \psi \rangle + V(\mathbf{r})\psi(\mathbf{r}) = E\psi(\mathbf{r})$$

 $\hat{\mathbf{L}}$  is the orbital angular momentum operator,  $\hat{\mathbf{L}} = \hat{\mathbf{r}} \times \hat{\mathbf{p}}$ , which satisfies the usual commutation relations for angular momenta, e.g.  $[\hat{L}_x, \hat{L}_y] = i\hbar \hat{L}_z$ . In the spherical coordinate,  $\hat{L}_z \to -i\hbar \frac{\partial}{\partial \phi}$ .

• For spherical symmetric potential V(r), there exist simultaneous eigenstates of  $\hat{H}$ ,  $\hat{\mathbf{L}}^2$ , and  $\hat{L}_z$ :

$$\hat{H}\left|E,l,m\right\rangle = E\left|E,l,m\right\rangle,\; \hat{\mathbf{L}}^{2}\left|E,l,m\right\rangle = l(l+1)\hbar^{2}\left|E,l,m\right\rangle,\; \hat{L}_{z}\left|E,l,m\right\rangle = m\hbar\left|E,l,m\right\rangle.$$

• In terms of wavefunctions,  $|E,l,m\rangle$  is of the form  $R(r)Y_{l,m}(\theta,\phi)$ , where the spherical harmonics  $Y_{l,m}(\theta,\phi)=\langle \theta,\phi|l,m\rangle$ , and

$$\int d\Omega |Y_{l,m}(\theta,\phi)|^2 = 1,$$

with  $d\Omega$  being the solid angle.

• Given one  $|l,m\rangle$ , one can obtain other eigenkets with the same l by the ladder operators  $\hat{L}_{\pm} = \hat{L}_x \pm i\hat{L}_y$ :

$$\hat{L_{\pm}}|l,m\rangle = \sqrt{l(l+1) - m(m\pm 1)}\hbar |l,m\pm 1\rangle.$$

 $\bullet$  The parity operator  $\hat{\pi}$  acts on a ket  $|\psi\rangle$  and flips the sign of the position expectation value:

$$|\psi'\rangle = \hat{\pi} |\psi\rangle, \langle x\rangle_{\psi'} = -\langle x\rangle_{\psi}$$

From this definition one can show that  $\hat{\pi}\hat{x}\hat{\pi} = -\hat{x}$ ,  $\hat{\pi}|x\rangle = |-x\rangle$ , and  $\hat{\pi} = \hat{\pi}^{-1} = \hat{\pi}^{\dagger}$ . The eigenvalues of  $\hat{\pi}$  are  $\pm 1$ . The wavefunctions of its eigenkets have the property  $\psi(-x) = \pm \psi(x)$ , and thus they are called parity-even or parity-odd states.

- The energy eigenvalues of a hydrogen-like atom (with Coulomb potential  $-\frac{Ze^2}{r}$ ) can be written as  $E_n = -\frac{\mu c^2 Z^2 \alpha^2}{2n^2}$ , where  $\mu$  is the reduced mass (can be take as the mass of an electron  $m_e$  for most cases), and  $\alpha$  is the fine structure constant  $\alpha = e^2/\hbar c \approx 1/137$ .
- For a 1-D simple harmonic oscillator (SHO),  $\hat{H} = \frac{\hat{p}_x^2}{2m} + \frac{1}{2}m\omega^2\hat{x}^2$ :

The raising and lowering operators are

$$\hat{a}^{\dagger} = \sqrt{\frac{m\omega}{2\hbar}}(\hat{x} - \frac{i}{m\omega}\hat{p_x}), \quad \hat{a} = \sqrt{\frac{m\omega}{2\hbar}}(\hat{x} + \frac{i}{m\omega}\hat{p_x})$$

and  $[\hat{a}, \hat{a}^{\dagger}] = 1$ . The operators get their names from the facts that

$$\hat{a}^{\dagger}\left|n\right\rangle = \sqrt{n+1}\left|n+1\right\rangle, \quad \hat{a}\left|n\right\rangle = \sqrt{n}\left|n-1\right\rangle.$$

One can rewrite  $\hat{x}$  and  $\hat{p_x}$  as

$$\hat{x} = \sqrt{\frac{\hbar}{2m\omega}}(\hat{a} + \hat{a}^{\dagger}), \quad \hat{p_x} = -i\sqrt{\frac{m\omega\hbar}{2}}(\hat{a} - \hat{a}^{\dagger})$$

and the Hamiltonian as

$$\hat{H} = \hbar\omega(\hat{a}^{\dagger}\hat{a} + \frac{1}{2}) = \hbar\omega(\hat{N} + \frac{1}{2}),$$

where the number operator  $\hat{N} \equiv \hat{a}^{\dagger} \hat{a}$ .

- The energy eigenvalues of an SHO are  $E_n=(n+\frac{1}{2})\hbar\omega, \quad n=0,1,2,...$
- Expectation value of an operator  $\hat{A}$  for a state  $|\psi\rangle$  is  $\langle A \rangle = \langle \psi | \hat{A} | \psi \rangle$
- Uncertainty of an operator  $\hat{A}$  is defined as  $\Delta A = \sqrt{\langle A^2 \rangle \langle A \rangle^2}$
- Heisenberg Equation:

$$\frac{d}{dt}\langle A\rangle = \langle \frac{\partial A}{\partial t}\rangle + \frac{1}{i\hbar}\langle [\hat{A},\hat{H}]\rangle$$

where  $\hat{H}$  is the Hamiltonian.

- 1. Quantum v.s. Classical descriptions:
  - (a) Use the Heisenberg equation prove that (in 1-D)

$$\frac{d}{dt}\langle xp\rangle = 2\langle T\rangle - \langle x\frac{dV}{dx}\rangle,$$

where  $\hat{T} = \hat{H} - \hat{V}$  is the kinetic energy. (2 points)

[Hint:  $[\hat{A}\hat{B}, \hat{C}] = \hat{A}[\hat{B}, \hat{C}] + [\hat{A}, \hat{C}]\hat{B}$ , and  $[\hat{p}_x, f(x)] = -i\hbar \frac{df}{dx}$ .]

(b) Explain why the result reduces to

$$2\langle T\rangle = \langle x\frac{dV}{dx}\rangle$$

for energy eigenstates. This is the Virial Theorem you've seen in Classical Mechanics. (2 points)

- (c) Express the Virial Theorem in 3-D, and show that  $\langle T \rangle = -\frac{1}{2} \langle V \rangle$  for a hydrogen atom (Coulomb potential). (2 points)
- (d) Argue that the motion of a ground-state electron in a hydrogen atom has to be described quantum-mechanically. (2 points) [Hint: Use the results from part (c) to estimate the size of v/c (v is the typical speed) for the electron, and take the typical size of a hydrogen atom as the traveling distance, and estimate the size of its classical action.  $m_e \approx 9 \times 10^{-31}$  kg. An order-of-magnitude estimation is all you need.]
- 2. Electron in a Coulomb potential: Let  $\psi_{nlm}(\mathbf{r})$  denote the normalized energy eigenfunctions of a Coulomb potential, with principle quantum number n and angular momentum quantum numbers l and m. Consider an electron in the state

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at time t = 0.

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- (a) Find the normalization constant c. (1 point)
- (b) Let  $E_1$  be the energy eigenvalue of the state  $\psi_{100}(\mathbf{r})$ , express the expectation value of energy (E) at t=0 in terms of  $E_1$ . (1 point)

 $\Psi(\mathbf{r}) = c[\psi_{100}(\mathbf{r}) + 4i\psi_{210}(\mathbf{r}) - 2\sqrt{2}\psi_{21-1}(\mathbf{r})]$ 

(c) Evaluate the expectation values of the angular momentum  $\langle L_z \rangle$  and  $\langle L_x \rangle$  at time t=0. (2 points)

[Hint: Make use of the ladder operators  $\widehat{(L_{\pm})}$ ]

(d) Now we turn on an external magnetic field along the z-axis, and the resulting Hamiltonian is

$$\hat{H} = \frac{\hat{\mathbf{p}}^2}{2\mu} - \frac{Ze^2}{r} + \omega_0 \hat{L}_z.$$

Evaluate  $\langle L_z \rangle$  and  $\langle L_x \rangle$  as functions of time. (2 points).

[Note: You can save some efforts by arguing which is/are constant(s) of motion.]

- 3. Parity of SHO eigenstates:
  - (a) Consider a 1-D SHO, explain why its energy eigenstates have definite parity (i.e. each energy eigenstate  $|n\rangle$  is an eigenstate of the parity operator  $\hat{\pi}$ ). (1 point)
  - (b) Show that  $\hat{\pi}\hat{a}^{\dagger}\hat{\pi} = -\hat{a}^{\dagger}$ , where  $\hat{a}^{\dagger}$  is the raising operator of the 1-D SHO. (1 point)
  - (c) Show that the wavefunction of the ground state |0> for a 1-D SHO is Gaussian (you don't need to derive the normalization constant).
     (2 points)

[Hint: Use the property  $\hat{a}|0\rangle = 0$  ( $\hat{a}$  is the lowering operator). Note that  $\hat{p}_x \to -i\hbar \frac{d}{dx}$  in 1-D.]

- (d) Given the fact  $\langle x|0\rangle$  is Gaussian and the results from part (b), show that  $\hat{\pi}|n\rangle = (-1)^n|n\rangle$ . (2 points)
- 4. Schrödinger equation in the presence of EM field: We've shown in class that the canonical momentum (to be identified with  $-i\hbar\nabla$  in the coordinate space) when there exists a magnetic vector potential A is

$$\mathbf{p}_c = m\mathbf{v} + \frac{q}{c}\mathbf{A}$$

And the Hamiltonian is

$$H = \frac{(\mathbf{p}_c - q\mathbf{A}/c)^2}{2m} + q\phi$$

where  $\phi$  is the electric potential. Different **A** and  $\phi$  can result in the same **E** and **B** fields, under the following gauge transformation:

$$\mathbf{A}' = \mathbf{A} - \nabla f(\mathbf{r}, t), \quad \phi' = \phi + \frac{1}{c} \frac{\partial f}{\partial t}.$$

with some real function  $f(\mathbf{r}, t)$ .

(a) One can show that if  $\psi(\mathbf{r},t)$  is a solution of the Schrödinger equation with  $\mathbf{A}$  and  $\phi$ , then

$$\psi'(\mathbf{r},t) = \exp(-\frac{iq}{\hbar c}f(\mathbf{r},t))\psi(\mathbf{r},t)$$

is a solution of the the Schrödinger equation with A' and  $\phi'$ . Is it worrisome that  $\psi(\mathbf{r},t)$  is *not* gauge-invariant? Explain why. (1 point)

(b) Define the "velocity" operator as

$$\hat{v}_i = \frac{1}{m}(\hat{p}_{c,i} - \frac{q}{c}A_i).$$

Please rewrite the Hamiltonian in terms of  $\hat{v}_i$ . (1 point)

(c) Explicitly show whether the expectation values  $\langle \hat{p}_{c,i} \rangle$  and  $\langle \hat{v}_i \rangle$  are gauge-invariant or not (i.e. a few lines of derivations are needed). (2 points)

[Note: Partial credits will be given if you only argue by physics.]

(d) Suppose we are in an EM field with (certain gauge in which)  $\phi=0$  and

$$\mathbf{A} = \frac{1}{2}(-B_0 y, B_0 x, 0),$$

where  $B_0$  is a constant. What are the corresponding electric and magnetic fields? (2 points)

(e) Show that the commutator

$$[\hat{v}_x, \hat{v}_y] = \frac{i\hbar\omega_L}{m},$$

where  $\omega_L = \frac{qB_0}{mc}$ . (2 points)

(f) Introduce the annihilation and creation operators

$$\hat{a} = \sqrt{\frac{m}{2\hbar\omega_L}}(\hat{v}_x + i\hat{v}_y), \quad \hat{a}^{\dagger} = \sqrt{\frac{m}{2\hbar\omega_L}}(\hat{v}_x - i\hat{v}_y).$$

Rewrite the Hamiltonian in terms of  $\hat{a}$  and  $\hat{a}^{\dagger}$ . What are the energy eigenvalues of this system? (2 points)

[*Hint:* You need to show that  $[\hat{a}, \hat{a}^{\dagger}] = 1$ .)

The energy spectrum is called the Landau Levels, which is essential in understanding the Quantum Hall Effect.