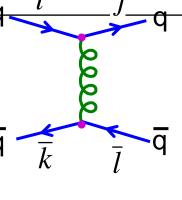
粒子物理学

习题课3



1. 考虑 $q\bar{q} \rightarrow q\bar{q}$ 的散射过程,画出费曼图并根据QCD费曼规则写出矩阵元(将颜色部分放在一起)。



$$-iM = \left[\bar{u}_q(p_3) \left(-\frac{1}{2} i g_s \lambda^a_{ji} \gamma^\mu \right) u_q(p_1) \right] \frac{-i g_{\mu\nu}}{q^2} \delta^{ab} \left[\bar{v}_{\bar{q}}(p_2) \left(-\frac{1}{2} i g_s \lambda^b_{kl} \gamma^\nu \right) v_{\bar{q}}(p_4) \right]$$

$$M = -\frac{g_s^2}{4} \lambda_{ji}^a \lambda_{kl}^a \frac{1}{q^2} g_{\mu\nu} \left[\bar{u}_q(p_3) \gamma^{\mu} u_q(p_1) \right] \left[\bar{v}_{\bar{q}}(p_2) \gamma^{\nu} v_{\bar{q}}(p_4) \right]$$

2. 计算以下四种色因子

a)
$$C(r\bar{r} \rightarrow r\bar{r})$$

b)
$$C(b\bar{g} \rightarrow b\bar{g})$$

c)
$$C(r\bar{r} \rightarrow b\bar{b})$$

d)
$$C(r\bar{g} \rightarrow g\bar{b})$$

$$M = -\frac{g_s^2}{4} \lambda_{ji}^a \lambda_{kl}^a \frac{1}{q^2} g_{\mu\nu} \left[\bar{u}_q(p_3) \gamma^{\mu} u_q(p_1) \right] \left[\bar{v}_{\bar{q}}(p_2) \gamma^{\nu} v_{\bar{q}}(p_4) \right]$$

$$C(i\overline{k} \rightarrow j\overline{l}) \equiv \frac{1}{4} \sum_{a=1}^{8} \lambda_{ji}^{a} \lambda_{kl}^{a}$$

$$\lambda^{1} = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \qquad \lambda^{4} = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{pmatrix} \qquad \lambda^{6} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} \qquad \lambda^{3} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$\lambda^{2} = \begin{pmatrix} 0 & -i & 0 \\ i & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \qquad \lambda^{5} = \begin{pmatrix} 0 & 0 & -i \\ 0 & 0 & 0 \\ i & 0 & 0 \end{pmatrix} \qquad \lambda^{7} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -i \\ 0 & i & 0 \end{pmatrix} \qquad \lambda^{8} = \frac{1}{\sqrt{3}} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -2 \end{pmatrix}$$

$$C(r\bar{r} \to r\bar{r}) = \frac{1}{4} \sum_{a=1}^{8} \lambda_{11}^{a} \lambda_{11}^{a} = \frac{1}{4} (\lambda_{11}^{3} \lambda_{11}^{3} + \lambda_{11}^{8} \lambda_{11}^{8}) = \frac{1}{3}$$

$$C(b\bar{g} \to b\bar{g}) = \frac{1}{4} \sum_{a=1}^{6} \lambda_{22}^{a} \lambda_{33}^{a} = \frac{1}{4} (\lambda_{22}^{8} \lambda_{33}^{8}) = -\frac{1}{6}$$

$$C(r\bar{r} \to b\bar{b}) = \frac{1}{4} \sum_{1}^{8} \frac{\lambda_{21}^{a} \lambda_{12}^{a}}{\lambda_{21}^{a} \lambda_{12}^{a}} = \frac{1}{4} (\lambda_{21}^{1} \lambda_{12}^{1} + \lambda_{21}^{2} \lambda_{12}^{2}) = \frac{1}{2}$$

$$C(r\bar{g} \to g\bar{b}) = \frac{1}{4} \sum_{a=1}^{8} \lambda_{31}^{a} \lambda_{32}^{a} = 0$$



3. 计算出 $q\bar{q} \rightarrow q\bar{q}$ 的散射过程的平均色因子 $\langle |C|^2 \rangle$

$$\langle |C|^2 \rangle = \frac{1}{9} \sum_{i,j,k,l=1}^{3} |C(ij \to kl)|^2$$

$$\overline{q}$$
 \overline{k} \overline{q}

$$\langle |C|^2 \rangle = \frac{1}{9} \left[3 \times \left(\frac{1}{3} \right)^2 + 6 \times \left(-\frac{1}{6} \right)^2 + 6 \times \left(\frac{1}{2} \right)^2 \right] = \frac{2}{9}$$



4. 用Appendix D的方法计算以下四种色因子 $C(ik \rightarrow jl) = \frac{1}{2}c_1c_2$

a)
$$C(r\bar{r} \rightarrow r\bar{r})$$

b)
$$C(b\bar{g} \rightarrow b\bar{g})$$

c)
$$C(r\bar{r} \rightarrow b\bar{b})$$

d)
$$C(r\bar{g} \to g\bar{b})$$

a)
$$C(r\bar{r} \to r\bar{r})$$
 $\lambda^1 = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$
 $\lambda^4 = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{pmatrix}$
 $\lambda^6 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}$
 $\lambda^3 = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{pmatrix}$

 c) $C(r\bar{r} \to b\bar{b})$
 $\lambda^2 = \begin{pmatrix} 0 & -i & 0 \\ i & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$
 $\lambda^5 = \begin{pmatrix} 0 & 0 & -i \\ 0 & 0 & 0 \\ i & 0 & 0 \end{pmatrix}$
 $\lambda^7 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -i \\ 0 & i & 0 \end{pmatrix}$
 $\lambda^8 = \frac{1}{\sqrt{3}} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -2 \end{pmatrix}$

 胶子: $r\bar{g}, g\bar{r}$
 $r\bar{b}, b\bar{r}$
 $g\bar{b}, b\bar{g}$
 $\frac{1}{\sqrt{2}}(r\bar{r} - g\bar{g})$
 $\frac{1}{\sqrt{6}}(r\bar{r} + g\bar{g} - 2b\bar{b})$

$$C(r\bar{r} \to r\bar{r}) = \frac{1}{2} \left(\frac{1}{2} + \frac{1}{6} \right) = \frac{1}{3} \qquad \frac{\frac{1}{\sqrt{2}}(r\bar{r} - g\bar{g})}{\frac{1}{\sqrt{6}}(r\bar{r} + g\bar{g} - 2b\bar{b})}$$

$$C(b\bar{g} \to b\bar{g}) = \frac{1}{2} \left(\frac{1}{6} (1 \times -2) \right) = -\frac{1}{6} \qquad \frac{\frac{1}{\sqrt{6}}(r\bar{r} + g\bar{g} - 2b\bar{b})}{\frac{1}{\sqrt{6}}(r\bar{r} + g\bar{g} - 2b\bar{b})}$$

$$C(r\bar{r} \to b\bar{b}) = \frac{1}{2} (1) = \frac{1}{2} \qquad r\bar{b}$$

$$C(r\bar{g} \to g\bar{b}) = 0$$

1.试着解释为什么存在 $\rho^0 \to \pi^-\pi^+$ 的衰变过程,但是不存在 $\rho^0 \to \pi^0\pi^0$ 的衰变过程。 (提示:可以从角动量,宇称,波函数的对称性等角度思考)

先考虑几个粒子的自旋与宇称:

$$J^{P}(\rho^{0}) = 1^{-}, J^{P}(\pi^{0}) = J^{P}(\pi^{+}) = J^{P}(\pi^{-}) = 0^{-}$$

对于 $\rho^0 \to \pi\pi$ 的过程, 宇称守恒:

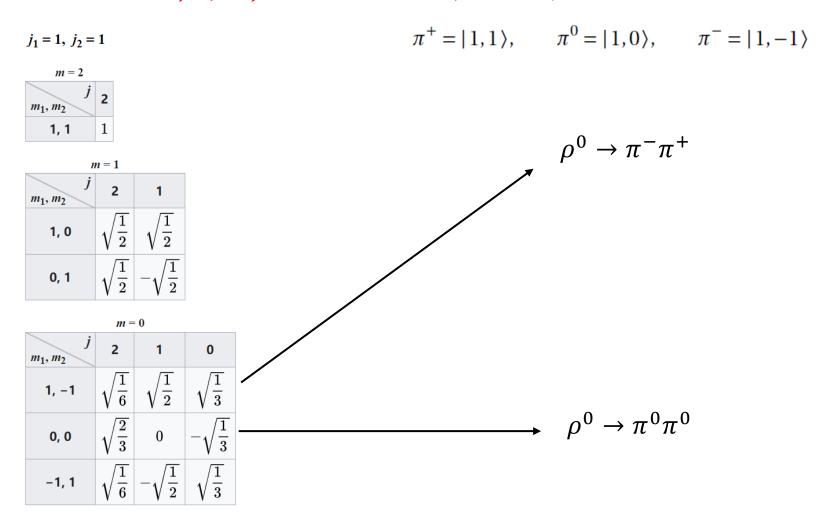
$$P_{\rho^0} = P_{\pi} P_{\pi} (-1)^L$$

(-1) = (-1)(-1)(-1)^L
L必须为奇数

若衰变到 $\pi^-\pi^+$, L=1并无不妥;

但若衰变到 $\pi^0\pi^0$,末态为全同的玻色子,波函数交换对称,要求L为偶数,矛盾。

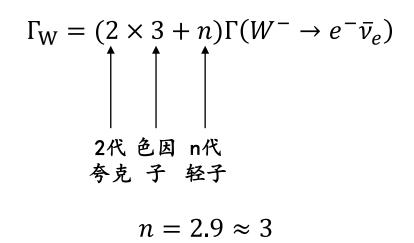
1.试着解释为什么存在 $\rho^0 \to \pi^- \pi^+$ 的衰变过程,但是不存在 $\rho^0 \to \pi^0 \pi^0$ 的衰变过程。 (提示:可以从角动量,宇称,波函数的对称性等角度思考)



2.轻子的代数可以通过W玻色子的总衰变宽度估计。标准模型预测 $W^- \to e^- \bar{\nu}_e$ 的衰变宽度写作:

$$\Gamma(W^- \to e^- \bar{\nu}_e) = \frac{G_F}{\sqrt{2}} \frac{m_W^3}{6\pi}$$

已知W玻色子的质量 $m_W = 80.385 \text{GeV}$,总衰变宽度 $\Gamma_W = 2.085 \text{GeV}$,试着由此估计轻子的代数。($G_F = 1.2 \times 10^{-5} \text{GeV}^{-2}$)



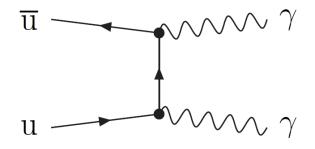
3. 考虑以下几组衰变。在每一组中,在标准模型范围内考虑每一个衰变是否可以发生,如果可以发生,画出对应的费曼图,如果不能,给出理由;在后面两组中,如果有多个过程可以发生,按照衰变率排名。

a)
$$\pi^0 \rightarrow \gamma \gamma$$
,

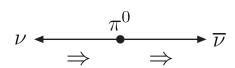
$$\pi^0 \rightarrow \pi^- e^+ \nu_e$$
,

$$\pi^0 \rightarrow \nu \bar{\nu}$$

$$J(\pi^0) = 0$$



$$m_{\pi^\pm} > m_{\pi^0}$$

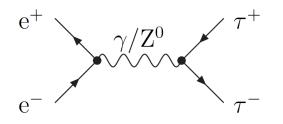


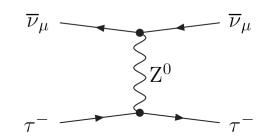
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b)
$$e^+e^- \to \tau^+\tau^-$$
,

$$\bar{\nu}_{\mu} + \tau^- \rightarrow \bar{\nu}_{\mu} + \tau^-, \qquad \qquad \nu_{\tau} + p \rightarrow \tau^+ + n$$

$$\nu_{\tau} + p \rightarrow \tau^{+} + n$$





轻子数不守恒

是散射, 题意不明

$$\underline{ V_{\rm CKM}} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \sim \begin{pmatrix} \cos\theta_{\mathcal{C}} & \sin\theta_{\mathcal{C}} & \sin^3\theta_{\mathcal{C}} \\ -\sin\theta_{\mathcal{C}} & \cos\theta_{\mathcal{C}} & \sin^2\theta_{\mathcal{C}} \\ \sin^3\theta_{\mathcal{C}} & -\sin^2\theta_{\mathcal{C}} & 1 \end{pmatrix} \underline{ }$$

3. 考虑以下几组衰变。在每一组中, 在标准模型范围内考虑每一个衰变是否可以发生 ,如果可以发生,画出对应的费曼图,如果不能,给出理由;在后面两组中,如果 有多个过程可以发生,按照衰变率排名。

c)
$$B^0(\bar{b}d) \rightarrow D^-(\bar{c}d)\pi^+$$
, $B^0 \rightarrow \pi^+\pi^-$,

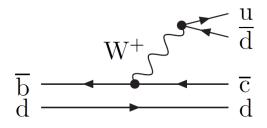
$$B^0 \rightarrow \pi^+\pi^-$$

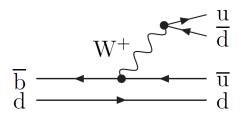
$$B^0 \to J/\psi K^0$$

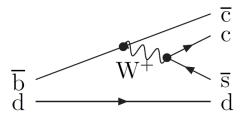
$$2.52 \times 10^{-3}$$

$$5.12 \times 10^{-6}$$

$$8.68 \times 10^{-4}$$







 $V_{bc}V_{ud}$

$$V_{bu}V_{ud}$$

$$V_{bc}V_{cs}$$

$$V_{
m CKM} = \left(egin{array}{ccc} V_{ud} & V_{us} & V_{ub} \ V_{cd} & V_{cs} & V_{cb} \ V_{td} & V_{ts} & V_{tb} \end{array}
ight) \sim \left(egin{array}{ccc} \cos heta_{C} & \sin heta_{C} & \sin^{3} heta_{C} \ -\sin heta_{C} & \cos heta_{C} & \sin^{2} heta_{C} \ \sin^{3} heta_{C} & -\sin^{2} heta_{C} \end{array}
ight) - \left(egin{array}{ccc} \cos heta_{C} & \sin^{2} heta_{C} \ \sin^{3} heta_{C} & -\sin^{2} heta_{C} \end{array}
ight)$$

3. 考虑以下几组衰变。在每一组中,在标准模型范围内考虑每一个衰变是否可以发生,如果可以发生,画出对应的费曼图,如果不能,给出理由;在后面两组中,如果有多个过程可以发生,按照衰变率排名。

d)
$$D^0(c\bar{u}) \rightarrow K^-\pi^+$$
,

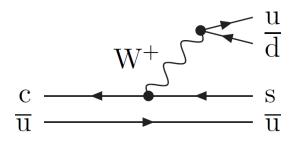
$$D^0 \rightarrow \pi^+\pi^-$$

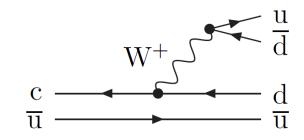
$$D^0 \rightarrow K^+\pi^-$$

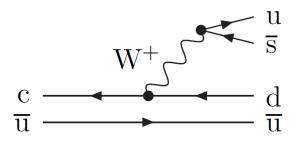
$$3.95 \times 10^{-2}$$

$$1.455 \times 10^{-3}$$

$$1.5\times10^{-4}$$







$$V_{ud}V_{cs}$$

$$V_{ud}V_{cd}$$

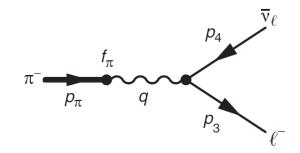
$$V_{us}V_{cd}$$

4. π 介子衰变计算:课上我们学到,实验测量得到的 $\frac{\Gamma(\pi^- \to e^- \overline{\nu}_e)}{\Gamma(\pi^- \to \mu^- \overline{\nu}_e)}$ 值与预期并不相同,其原因是衰变到电子的过程被螺旋度压低了。之后课上进行了一部分的计算,但没有最终完成,在这里将其补完:

a) 考虑 π 介子静止系,先写出轻子流(注意手征算符);至于夸克流,写成正比于 $f_{\pi}p_{\pi}^{\mu}$ 的形式(参考书本11.6)

$$j_{\ell}^{\nu} = \frac{g_W}{\sqrt{2}} \bar{u}(p_3) \frac{1}{2} \gamma^{\nu} (1 - \gamma^5) \nu(p_4)$$

$$j_{\pi}^{\mu} = \frac{g_W}{\sqrt{2}} \frac{1}{2} f_{\pi} p_{\pi}^{\mu}$$

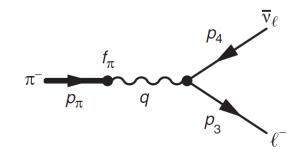


4. (b) 写出矩阵元 M_{fi} 的具体形式(考虑 $q^2=m_\pi^2\ll m_W^2$)

$$M_{fi} = \left[\frac{g_W}{\sqrt{2}} \frac{1}{2} f_\pi p_\pi^\mu \right] \times \left[\frac{g_{\mu\nu}}{m_W^2} \right] \times \left[\frac{g_W}{\sqrt{2}} \bar{u}(p_3) \frac{1}{2} \gamma^\nu (1 - \gamma^5) \nu(p_4) \right]$$
$$= \frac{g_W^2}{4m_W^2} g_{\mu\nu} f_\pi p_\pi^\mu \bar{u}(p_3) \frac{1}{2} \gamma^\nu (1 - \gamma^5) \nu(p_4)$$

在 π^- 的静止系中,只有零分量不为零, $\bar{u}\gamma^0 = u^\dagger$

$$M_{fi} = \frac{g_W^2}{4m_W^2} f_{\pi} m_{\pi} u^{\dagger}(p_3) \frac{1}{2} (1 - \gamma^5) \nu(p_4)$$
$$= \frac{g_W^2}{4m_W^2} f_{\pi} m_{\pi} u^{\dagger}(p_3) \nu_{\uparrow}(p_4)$$



4. (c)写出狄拉克方程的螺旋度解,并代入矩阵元具体计算,得到矩阵元平方 $|M_{fi}|^2$ 的具

体形式 $M_{fi} = \frac{g_W^2}{4m_W^2} f_\pi m_\pi u^\dagger(p_3) \nu_\uparrow(p_4)$

$$u_{\uparrow}(p_{3}) = \sqrt{E_{\ell} + m_{\ell}} \begin{pmatrix} 1 \\ 0 \\ p \\ \hline E_{\ell} + m_{\ell} \end{pmatrix}, u_{\downarrow}(p_{3}) = \sqrt{E_{\ell} + m_{\ell}} \begin{pmatrix} 1 \\ 0 \\ 0 \\ -\frac{p}{E_{\ell} + m_{\ell}} \end{pmatrix} \qquad \nu_{\uparrow}(p_{3}) = \sqrt{p} \begin{pmatrix} 1 \\ 0 \\ -1 \\ 0 \end{pmatrix}$$

$$\nu_{\uparrow}(p_3) = \sqrt{p} \begin{pmatrix} 1\\0\\-1\\0 \end{pmatrix}$$

根据四动量守恒

$$\begin{split} M_{fi} &= \frac{g_W^2}{4m_W^2} f_\pi m_\pi \sqrt{E_\ell + m_\ell} \sqrt{p} \left(1 - \frac{p}{E_\ell + m_\ell} \right) \\ &= \frac{g_W^2}{4m_W^2} f_\pi m_\ell \sqrt{m_\pi^2 - m_\ell^2} \end{split} \qquad E_\ell = \frac{m_\pi^2 + m_e^2}{2m_\pi}, |\vec{p}| = \frac{m_\pi^2 - m_e^2}{2m_\pi} \end{split}$$

$$\left| M_{fi} \right|^2 = 2G_F^2 f_{\pi}^2 m_{\ell}^2 \left(m_{\pi}^2 - m_{\ell}^2 \right)$$

$$\frac{G_F}{\sqrt{2}} = \frac{g_W^2}{8m_W^2}$$

4. (d)计算出
$$\frac{\Gamma(\pi^- \to e^- \overline{\nu}_e)}{\Gamma(\pi^- \to \mu^- \overline{\nu}_e)}$$
的值

$$\left| M_{fi} \right|^2 = 2G_F^2 f_\pi^2 m_\ell^2 \left(m_\pi^2 - m_\ell^2 \right)$$

$$\Gamma = \frac{4\pi}{32\pi^2 m_{\pi}^2} p \left| M_{fi} \right|^2 = \frac{G_F^2}{8\pi m_{\pi}^3} f_{\pi}^2 \left[m_{\ell} \left(m_{\pi}^2 - m_{\ell}^2 \right) \right]^2$$

$$\frac{\Gamma(\pi^- \to e^- \bar{\nu}_e)}{\Gamma(\pi^- \to \mu^- \bar{\nu}_e)} = \left[\frac{m_e (m_\pi^2 - m_e^2)}{m_\mu (m_\pi^2 - m_\mu^2)} \right]^2 = 1.26 \times 10^{-4}$$

(e)计算
$$\frac{\Gamma(K^- \to e^- \overline{\nu}_e)}{\Gamma(K^- \to \mu^- \overline{\nu}_e)}$$
的值

$$2.55 \times 10^{-5}$$

 $E_{\ell} = \frac{m_{\pi}^2 + m_e^2}{2m_{-}}, |\vec{p}| = \frac{m_{\pi}^2 - m_e^2}{2m_{-}}$

 $rac{1}{ au}=\Gamma=rac{|ec{p}^*|}{32\pi^2m^2}\int |M_{fi}|^2\mathrm{d}\Omega$



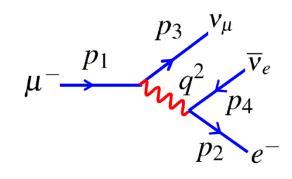
4. μ子衰变计算:

(a)写出图示过程的矩阵元(不用化简)

$$-i\mathcal{M}_{fi} = \left[\overline{u}(p_3) \frac{-ig_W}{\sqrt{2}} \gamma^{\mu} \frac{1-\gamma^5}{2} u(p_1) \right]$$

$$\frac{ig_{\mu\nu}}{m_W^2}$$

$$\left[\overline{u}(p_2) \frac{-ig_W}{\sqrt{2}} \gamma^{\nu} \frac{1-\gamma^5}{2} v(p_4) \right]$$



$$\mathcal{M}_{fi} = \frac{g_W^2}{8m_W^2} g_{\mu\nu} \left[\bar{u}(p_3) \gamma^{\mu} (1 - \gamma^5) u(p_1) \right] \left[\bar{u}(p_2) \gamma^{\nu} (1 - \gamma^5) v(p_4) \right]$$

4. μ子衰变计算:

$$\frac{G_F}{\sqrt{2}} = \frac{g_W^2}{8m_W^2}$$

(b)使用之前使用过的求迹技巧进行计算

$$|M_{fi}|^2 = \frac{G_F^2}{2} \frac{1}{2} L_{\mu\nu} W_{\mu\nu}$$

$$L_{\mu\nu} = \frac{1}{2} \text{Tr} \left[(p_1 + m) \gamma^{\mu} \frac{1 - \gamma^5}{2} (p_3) \gamma^{\nu} \frac{1 - \gamma^5}{2} \right] \frac{1 - \gamma^5}{\epsilon^{abcd} \epsilon_{ebfd}} = -2 \delta_{ae} \delta_{cf} + 2 \delta_{af} \delta_{ce}$$
$$= 2 \left[p_1^{\mu} p_3^{\nu} + p_1^{\nu} p_3^{\mu} - g^{\mu\nu} p_1 p_3 + i \epsilon^{\mu\nu\nu\sigma} p_{3\rho} p_{1\sigma} \right]$$

$$W_{\mu\nu} = 2 \left[p_{2\mu} p_{4\nu} + p_{2\nu} p_{4\mu} - g_{\mu\nu} p_2 p_4 + i \epsilon_{\mu\alpha\nu\beta} p_2^{\alpha} p_4^{\beta} \right]$$

对称项与反对称项混合的乘法,可以分别相乘

$$|M_{fi}|^2 = 64G_F^2(p_1 \cdot p_4)(p_2 \cdot p_3)$$

$$P_L P_L = P_L$$

$$Tr(\text{odd number of } \gamma s) = 0$$

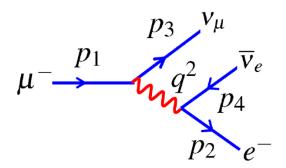
$$Tr(\gamma^{\mu} \gamma^{\nu} \gamma^{\rho} \gamma^{\sigma}) = 4g^{\mu\nu} g^{\rho\sigma} - 4g^{\mu\rho} g^{\nu\sigma} + 4g^{\mu\sigma} g^{\nu\rho}$$

$$Tr[\gamma^{\mu} \gamma^{\nu} \gamma^{5}] = 0$$

$$Tr[\gamma^{\mu} \gamma^{\nu} \gamma^{\rho} \gamma^{\sigma} \gamma^{5}] = 4i \epsilon^{\mu\nu\rho\sigma}$$

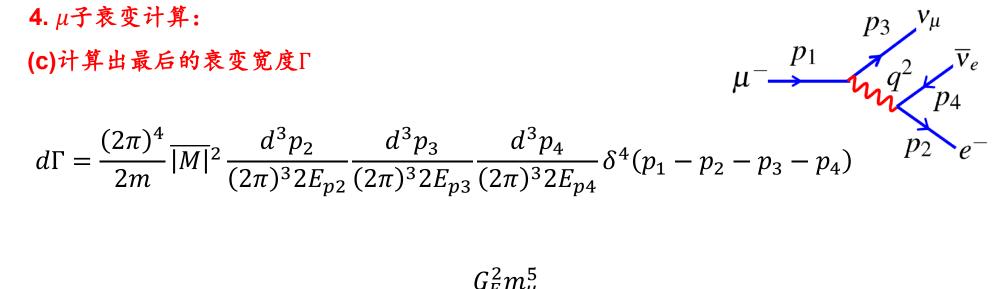
$$\epsilon^{abcd} \epsilon_{ebfd} = -2\delta_{ae} \delta_{cf} + 2\delta_{af} \delta_{ce}$$

 $Tr(\gamma^{\mu}\gamma^{\nu}) = 4g^{\mu\nu}$





4. μ子衰变计算:



$$\Gamma = \frac{G_F^2 m_\mu^5}{192\pi^3}$$

1. 填表

粒子 (场)	I_{W}	I_{W}^3	Q	Y
$L_{iL} = \begin{pmatrix} \nu_{iL} \\ \ell_{iL} \end{pmatrix} : \begin{pmatrix} \nu_{eL} \\ e_L \end{pmatrix} \begin{pmatrix} \nu_{\mu L} \\ \mu_L \end{pmatrix} \begin{pmatrix} \nu_{\tau L} \\ \tau_L \end{pmatrix}$				
没有右手中微子 ℓ _{iR} : e _R , μ _R , τ _R				
$Q_{iL} = \begin{pmatrix} u_{iL} \\ d_{iL} \end{pmatrix} : \begin{pmatrix} u_L \\ d_L \end{pmatrix} \begin{pmatrix} c_L \\ s_L \end{pmatrix} \begin{pmatrix} t_L \\ b_L \end{pmatrix}$				
u_{iR} : u_R , c_R , t_R d_{iR} : d_R , s_R , b_R				

费米子	Q	I_w^3	c_L	c_R	c_V	c_A
$ u_{e,\mu, au}$						
e, μ, τ						
u, c, t						
d, s, b						

1. 填表

粒子 (场)	I_{W}	I_{W}^{3}	Q	Y
$L_{iL} = \begin{pmatrix} \nu_{iL} \\ \ell_{iL} \end{pmatrix} : \begin{pmatrix} \nu_{eL} \\ e_L \end{pmatrix} \begin{pmatrix} \nu_{\mu L} \\ \mu_L \end{pmatrix} \begin{pmatrix} \nu_{\tau L} \\ \tau_L \end{pmatrix}$	$\frac{1}{2}$	$+\frac{1}{2}$ $-\frac{1}{2}$	0 -1	-1
没有右手中微子 ℓ_{iR} : e_R , μ_R , $ au_R$	_ 0	_ 0	- -1	- -2
$Q_{iL} = \begin{pmatrix} u_{iL} \\ d_{iL} \end{pmatrix} : \begin{pmatrix} u_L \\ d_L \end{pmatrix} \begin{pmatrix} c_L \\ s_L \end{pmatrix} \begin{pmatrix} t_L \\ b_L \end{pmatrix}$	$\frac{1}{2}$	$+\frac{1}{2}$ $-\frac{1}{2}$	$+\frac{2}{3}$ $-\frac{1}{3}$	$\frac{1}{3}$
u_{iR} : u_R , c_R , t_R d_{iR} : d_R , s_R , b_R	0	0 0	$+\frac{2}{3}$ $-\frac{1}{3}$	$\frac{4}{3}$ $-\frac{2}{3}$

费米子	Q	I_w^3	c_L	c_R	c_V	c_A
$ u_{e,\mu, au}$	0	$+\frac{1}{2}/-$	$\frac{1}{2}$	0	$\frac{1}{2}$	$\frac{1}{2}$
e, μ, τ	-1	$-\frac{1}{2}/0$	$-\frac{1}{2} + \sin^2 \theta_w$	$\sin^2 \theta_w$	$-\frac{1}{2} + 2\sin^2\theta_w$	$-\frac{1}{2}$
u, c, t	$+\frac{2}{3}$	$+\frac{1}{2}/0$	$\frac{1}{2} - \frac{2}{3}\sin^2\theta_w$	$-\frac{2}{3}\sin^2\theta_w$	$\frac{1}{2} - \frac{4}{3}\sin^2\theta_w$	$\frac{1}{2}$
d, s, b	$-\frac{1}{3}$	$-\frac{1}{2}/0$	$-\frac{1}{2} + \frac{1}{3}\sin^2\theta_w$	$\frac{1}{3}\sin^2\theta_w$	$-\frac{1}{2} + \frac{2}{3}\sin^2\theta_w$	$-\frac{1}{2}$

2. 仿照课件,从 j_{μ}^{em} , $j_{\mu}^{W^3}$, j_{μ}^{Y} 推导出 j_{μ}^{Z} 的具体形式

$$j_{\mu}^{em} = e\overline{\psi}Q_{e}\gamma_{\mu}\psi = e\overline{e}_{L}Q_{e}\gamma_{\mu}e_{L} + e\overline{e}_{R}Q_{e}\gamma_{\mu}e_{R}$$

$$j_{\mu}^{W^{3}} = -\frac{g_{W}}{2}\overline{e}_{L}\gamma_{\mu}e_{L}$$

$$j_{\mu}^{Y} = \frac{g'}{2}\overline{\psi}Y_{e}\gamma_{\mu}\psi = \frac{g'}{2}\overline{e}_{L}Y_{e_{L}}\gamma_{\mu}e_{L} + \frac{g'}{2}\overline{e}_{R}Y_{e_{R}}\gamma_{\mu}e_{R}$$

$$j_{\mu}^{em} = j_{\mu}^{Y}\cos\theta_{W} + j_{\mu}^{W^{3}}\sin\theta_{W}$$

$$A_{\mu} = B_{\mu} \cos \theta_W + W_{\mu}^3 \sin \theta_W$$

 $Z_{\mu} = -B_{\mu} \sin \theta_W + W_{\mu}^3 \cos \theta_W$

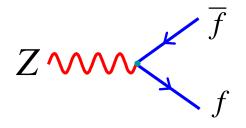
$$e = g_W \sin \theta_W = g' \cos \theta_W$$

$$j_{\mu}^{Z} = g_{Z}(I_{W}^{3} - Q\sin^{2}\theta_{W})[\bar{e}_{L}\gamma_{\mu}e_{L}] - g_{Z}Q\sin^{2}\theta_{W}[e_{R}\gamma_{\mu}e_{R}] \quad g_{Z} = \frac{g_{W}}{\cos\theta_{W}}$$

$$= g_{Z}c_{L}[\bar{e}_{L}\gamma_{\mu}e_{L}] + g_{Z}c_{R}[e_{R}\gamma_{\mu}e_{R}]$$

$$c_{L} = I_{W}^{3} - Q\sin^{2}\theta_{W} \qquad c_{R} = -Q\sin^{2}\theta_{W}$$

3. 考虑 $Z^0 \to f\bar{f}$ 的过程,其中费米子与 Z^0 耦合常数为 c_V , c_A (a) 使用费曼规则写出 $Z^0 \to f\bar{f}$ 的矩阵元,并整理成左手部分和右手部分



$$-iM = \epsilon_{\mu}(p_1)\bar{u}(p_3) \left(-ig_Z \gamma^{\mu} \frac{1}{2} (c_V - c_A \gamma^5)\right) v(p_4)$$

$$\frac{1}{2}(c_V - c_A \gamma^5) = c_L \frac{1}{2}(1 - \gamma^5) + c_R \frac{1}{2}(1 + \gamma^5)$$

$$\begin{split} M &= g_Z \epsilon_\mu(p_1) \bar{u}(p_3) \gamma^\mu c_L \frac{1}{2} (1 - \gamma^5) v(p_4) + g_Z \epsilon_\mu(p_1) \bar{u}(p_3) \gamma^\mu c_R \frac{1}{2} (1 + \gamma^5) v(p_4) \\ &= M_L + M_R \end{split}$$

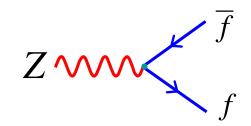
$$\mathcal{L}M_{fi} = \frac{g_W}{\sqrt{2}} \varepsilon_{\mu}(p_1) \overline{u}(p_3) \gamma^{\mu} \frac{1}{2} (1 - \gamma^5) v(p_4)$$

3. 考虑 $Z^0 \to f\bar{f}$ 的过程,其中费米子与 Z^0 耦合常数为 c_V, c_A

(b) 仿照课件中计算W玻色子衰变宽度的方法, 计算出 $Z^0 \to f\bar{f}$ 的衰变宽度的表达式

$$M_{L} = g_{Z} \epsilon_{\mu}(p_{1}) \bar{u}(p_{3}) \gamma^{\mu} c_{L} \frac{1}{2} (1 - \gamma^{5}) v(p_{4})$$

$$M_{R} = g_{Z} \epsilon_{\mu}(p_{1}) \bar{u}(p_{3}) \gamma^{\mu} c_{R} \frac{1}{2} (1 + \gamma^{5}) v(p_{4})$$



$$M_R = g_Z c_R \epsilon_\mu(p_1) \bar{u}_\uparrow(p_3) \gamma^\mu v_\downarrow(p_4)$$

$$\begin{split} M_{R-} &= \frac{1}{\sqrt{2}} g_Z c_R m_Z (\cos\theta - 1) \\ M_{RL} &= g_Z c_R m_Z \sin\theta \\ M_{R+} &= \frac{1}{\sqrt{2}} g_Z c_R m_Z (-\cos\theta - 1) \end{split}$$

$$\overline{u}_{\uparrow}(p_3)\gamma^{\nu}v_{\downarrow}(p_4) = 2E(0, -\cos\theta, i, \sin\theta)$$

$$\varepsilon_{-}^{\mu} = \frac{1}{\sqrt{2}}(0, 1, -i, 0); \quad \varepsilon_{L} = \frac{1}{m}(p_{z}, 0, 0, E) \quad \varepsilon_{+}^{\mu} = -\frac{1}{\sqrt{2}}(0, 1, i, 0)$$

$$\langle |M_R|^2 \rangle = \frac{2}{3} g_Z^2 m_Z^2 c_R^2$$

$$\langle |M|^2 \rangle = \frac{2}{3} g_Z^2 m_Z^2 (c_L^2 + c_R^2)$$

$$c_V^2 + c_A^2 = 2(c_L^2 + c_R^2)$$
$$\frac{d\Gamma}{d\Omega} = \frac{|p^*|}{32\pi^2 m_W^2} |M|^2$$

$$c_V^2 + c_A^2 = 2(c_L^2 + c_R^2)$$

$$\frac{d\Gamma}{d\Omega} = \frac{|p^*|}{32\pi^2 m_W^2} |M|^2 \qquad \Gamma(Z \to f\bar{f}) = \frac{g_Z^2 m_Z}{48\pi} (c_V^2 + c_A^2)$$

3. 考虑 $Z^0 \to f\bar{f}$ 的过程, 其中费米子与 Z^0 耦合常数为 c_V, c_A

(c)
$$R_{\mu} = \frac{\Gamma(Z \to \mu^{+}\mu^{-})}{\Gamma(Z \to hadrons)}$$

$$\Gamma(Z \to f\bar{f}) = \frac{g_Z^2 m_Z}{48\pi} (c_V^2 + c_A^2)$$

费米子	Q	I_W^3	c_L	c_R	c_V	c_A
$ u_e, u_\mu, u_ au$	0	$+\frac{1}{2}$	$+\frac{1}{2}$	0	$+\frac{1}{2}$	$+\frac{1}{2}$
e^-, μ^-, au^-	-1	$-\frac{1}{2}$	-0.27	0.23	-0.04	$-\frac{1}{2}$
u, c, t	$+\frac{2}{3}$	$+\frac{1}{2}$	0.35	-0.15	+0.19	$+\frac{1}{2}$
d,s,b	$-\frac{1}{3}$	$-\frac{1}{2}$	-0.42	0.08	-0.35	$-\frac{1}{2}$

$$R_{\mu} = \frac{\Gamma(Z \to \mu^{+}\mu^{-})}{\Gamma(Z \to hadrons)} = \frac{0.04^{2} + 0.5^{2}}{9 \times (0.35^{2} + 0.5^{2}) + 6 \times (0.19^{2} + 0.5^{2})} = 0.049$$

- 4. 在考虑了QED的效应后, $e^+e^- \rightarrow \mu^+\mu^-$ 和 $e^+e^- \rightarrow hadrons$ 经过Z共振态的截面为 $\sigma^0(e^+e^- \rightarrow Z \rightarrow \mu^+\mu^-) = 1.9993~nb$ $\sigma^0(e^+e^- \rightarrow Z \rightarrow hadrons) = 41.476~nb$
- (a) 假设轻子普适性,计算 $\frac{\Gamma_{ee}}{\Gamma_Z}$ 和 $\frac{\Gamma_{hadrons}}{\Gamma_Z}$ (使用公式 $\sigma^0 = \frac{12\pi}{m_Z^2} \frac{\Gamma_{ee}\Gamma_{ff}}{\Gamma_Z^2}$,记得单位换算)

$$\sigma_{\mu}^0 = \frac{12\pi}{m_Z^2} \frac{\Gamma_{ee}^2}{\Gamma_Z^2}$$

$$\frac{\Gamma_{ee}}{\Gamma_Z} = \sqrt{\frac{\sigma_\mu^0 m_Z^2}{12\pi}} = 0.0337$$

$$\sigma_{had}^{0} = \frac{12\pi}{m_Z^2} \frac{\Gamma_{ee} \Gamma_{hadrons}}{\Gamma_Z^2}$$

$$\frac{\Gamma_{hadrons}}{\Gamma_{Z}} = \frac{\sigma_{had}^{0} m_{Z}^{2}}{12\pi} / \frac{\Gamma_{ee}}{\Gamma_{Z}} = 0.6992$$

- **4.** 在考虑了**QED**的效应后, $e^+e^- \to \mu^+\mu^-$ 和 $e^+e^- \to hadrons$ 经过**Z**共振态的截面为 $\sigma^0(e^+e^- \to Z \to \mu^+\mu^-) = 1.9993 \ nb$ $\sigma^0(e^+e^- \to Z \to hadrons) = 41.476 \ nb$
- (b) 使用测量值 $\Gamma_Z = 2.4952 \text{GeV}$ 和理论值 $\Gamma_{\nu\nu} = 167 \text{MeV}$,试着估算中微子有几代

$$\Gamma_{\rm Z} = 3\Gamma_{\ell\ell} + \Gamma_{\rm hadrons} + N_{\rm v}\Gamma_{\rm vv}$$

$$N_{\nu} = \frac{\Gamma_{\rm Z} - 3\Gamma_{\ell\ell} - \Gamma_{\rm hadrons}}{\Gamma_{\nu\nu}} = \frac{2.4952 \times (1 - 3 \times 0.0337 - 0.6992)}{0.167} \approx 2.98$$

1. 缪子前后向不对称性参数的测量值 $A_{\mu}=0.1456$,试计算 $\sin^2\theta$ 的值。

$$A_f \equiv \frac{(c_L^e)^2 - (c_R^e)^2}{(c_L^e)^2 + (c_R^e)^2} = \frac{2c_V^f c_A^f}{(c_V^f)^2 + (c_A^f)^2}$$

费米子	Q	I_w^3	c_L	c_R	c_V	c_A
$ u_{e,\mu, au}$	0	$+\frac{1}{2}/-$	$\frac{1}{2}$	0	$\frac{1}{2}$	$\frac{1}{2}$
e, μ, τ	-1	$-\frac{1}{2}/0$	$-\frac{1}{2} + \sin^2 \theta_w$	$\sin^2 \theta_w$	$-\frac{1}{2} + 2\sin^2\theta_w$	$-\frac{1}{2}$
u, c, t	$+\frac{2}{3}$	$+\frac{1}{2}/0$	$\frac{1}{2} - \frac{2}{3}\sin^2\theta_w$	$-\frac{2}{3}\sin^2\theta_w$	$\frac{1}{2} - \frac{4}{3}\sin^2\theta_w$	$\frac{1}{2}$
d, s, b	$-\frac{1}{3}$	$-\frac{1}{2}/0$	$-\frac{1}{2} + \frac{1}{3}\sin^2\theta_w$	$\frac{1}{3}\sin^2\theta_w$	$-\frac{1}{2} + \frac{2}{3}\sin^2\theta_w$	$-\frac{1}{2}$

$$\sin^2 \theta \approx 0.2317$$

2. 使用幺正规范下的 $\phi(x)$,写出自由复标量场最终的拉氏量,并解释每一项的意义。

$$\mathcal{L} = (\partial_{\mu}\phi)^{*}(\partial^{\mu}\phi) + \mu^{2}\phi^{*}\phi - \lambda(\phi^{*}\phi)^{2} - \frac{1}{4}F_{\mu\nu}F^{\mu\nu} - qi(\phi^{*}\partial^{\mu}\phi - \phi\partial^{\mu}\phi^{*})A_{\mu} + q^{2}A_{\mu}A^{\mu}\phi^{*}\phi$$

$$\phi(x) = \frac{1}{\sqrt{2}}[\eta(x) + \nu]$$

$$\begin{split} \mathcal{L}_{\eta} &= \frac{1}{2} \partial_{\mu} \eta \partial^{\mu} \eta - \lambda \nu^{2} \eta^{2} - \lambda \nu \eta^{3} - \frac{\lambda}{4} \eta^{4} - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ &+ \frac{1}{2} q^{2} A_{\mu} A^{\mu} \eta^{2} + q^{2} \nu A_{\mu} A^{\mu} \eta + \frac{1}{2} q^{2} \nu^{2} A_{\mu} A^{\mu} \end{split}$$

3. 尝试化简 $\mathcal{L}_S = (\mathcal{D}_\mu \Phi)^\dagger (\mathcal{D}_\mu \Phi) + \mu^2 \Phi^\dagger \Phi - \lambda (\Phi^\dagger \Phi)^2$,整理成包含物理场的形式,并

指出H、W、Z的质量项。

$$\mathcal{D}_{\mu}(x)\Phi(x) = \left(\partial_{\mu} + ig\frac{\sigma^{a}}{2}W_{\mu}^{a} + ig'\frac{1}{2}B_{\mu}\right)\Phi(x)$$

$$= \frac{1}{\sqrt{2}} \begin{pmatrix} \partial_{\mu} + \frac{ig_{W}}{2} W_{\mu}^{3} + \frac{ig'}{2} B_{\mu} & \frac{ig_{W}}{2} [W_{\mu}^{1} - iW_{\mu}^{2}] \\ \frac{ig_{W}}{2} [W_{\mu}^{1} + iW_{\mu}^{2}] & \partial_{\mu} - \frac{ig_{W}}{2} W_{\mu}^{3} + \frac{ig'}{2} B_{\mu} \end{pmatrix} \begin{pmatrix} 0 \\ v + H \end{pmatrix}$$

$$=\frac{1}{\sqrt{2}}\begin{pmatrix} \partial_{\mu}+\frac{ig_{W}}{2\cos\theta_{W}}Z_{\mu} & \frac{ig_{W}}{\sqrt{2}}W_{\mu}^{+}\\ \frac{ig_{W}}{\sqrt{2}}W_{\mu}^{-} & \partial_{\mu}-\frac{ig_{W}}{2\cos\theta_{W}}Z_{\mu} \end{pmatrix}\begin{pmatrix} 0\\ \nu+H \end{pmatrix}$$

$$= \frac{1}{\sqrt{2}} \begin{pmatrix} i \frac{g_W}{\sqrt{2}} W_{\mu}^+(x) [\nu + H(x)] \\ \partial_{\mu} H(x) - i \frac{g_W}{2 \cos \theta_W} Z_{\mu}(x) [\nu + H(x)] \end{pmatrix}$$

$$\Phi(\mathbf{x}) = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ \nu + H(\mathbf{x}) \end{pmatrix}$$

$$Z_{\mu} = -B_{\mu} \sin \theta_{W} + W_{\mu}^{3} \cos \theta_{W}$$
$$e = g_{W} \sin \theta_{W} = g' \cos \theta_{W}$$

3. 尝试化简 $\mathcal{L}_S = (\mathcal{D}_\mu \Phi)^\dagger (\mathcal{D}_\mu \Phi) + \mu^2 \Phi^\dagger \Phi - \lambda (\Phi^\dagger \Phi)^2$,整理成包含物理场的形式,并指出H、W、Z的质量项。

$$\mathcal{D}_{\mu}(x)\Phi(x) = \frac{1}{\sqrt{2}} \begin{pmatrix} i\frac{g_W}{\sqrt{2}}W_{\mu}^+(x)[\nu + H(x)]\\ \partial_{\mu}H(x) - i\frac{g_W}{2\cos\theta_W}Z_{\mu}(x)[\nu + H(x)] \end{pmatrix}$$

$$\Phi(\mathbf{x}) = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ \nu + H(\mathbf{x}) \end{pmatrix}$$

$$\begin{split} \mathcal{L}_{S} &= \frac{1}{2} \left(\partial_{\mu} H \right) (\partial^{\mu} H) - \lambda v^{2} H^{2} - \lambda v H^{3} - \frac{1}{4} \lambda H^{4} \\ &+ \frac{v^{2} g^{2}}{4} W_{\mu}^{+} W^{-\mu} + \frac{v^{2} g^{2}}{8 \cos^{2} \theta_{W}} Z_{\mu} Z^{\mu} \\ &+ \frac{v g^{2}}{2} W_{\mu}^{+} W^{-\mu} H + \frac{v g^{2}}{4 \cos^{2} \theta_{W}} Z_{\mu} Z^{\mu} H + \frac{g^{2}}{4} W_{\mu}^{+} W^{-\mu} H^{2} + \frac{g^{2}}{8 \cos^{2} \theta_{W}} Z_{\mu} Z^{\mu} H^{2} \end{split}$$

4. 尝试化简 $\mathcal{L} = -\sum_{i=1,2,3} \sum_{i=u,c,t} (Y'^u_{ii} \bar{Q}^i_{iL} \widetilde{\Phi} u'_{iR})$,并指出质量项和相互作用项。

$$\mathcal{L} = -\frac{\nu + H(x)}{\sqrt{2}} \sum_{j=u,c,t} \left(Y_{ij}^{\prime u} \bar{u}_{iL}^{\prime} u_{jR}^{\prime} \right) \qquad \qquad \bar{Q}_{iL}^{i} = \begin{pmatrix} \bar{u}_{iL}^{\prime} \\ \bar{d}_{iL}^{\prime} \end{pmatrix}^{T} \tilde{\Phi} = \frac{1}{\sqrt{2}} \begin{pmatrix} \nu + H(x) \\ 0 \end{pmatrix}$$

$$\mathcal{L}_{F-S,Q} = -\frac{\nu + H}{\sqrt{2}} \left[\sum_{i,j=d,s,b} (Y_{ij}^{\prime d} \bar{d}_{iL}^{\prime} d_{jR}^{\prime}) + \sum_{i,j=u,c,t} (Y_{ij}^{\prime u} \bar{u}_{iL}^{\prime} u_{jR}^{\prime}) \right] + h.c.$$

定义列矩阵

需要将
$$Y_{ij}^{\prime d}$$
和 $Y_{ij}^{\prime u}$ 对角化 $\mathbf{u}_L^\prime = \begin{pmatrix} u_L^\prime \\ c_L^\prime \\ t_L^\prime \end{pmatrix}$, $\mathbf{u}_R^\prime = \begin{pmatrix} u_R^\prime \\ c_R^\prime \\ t_R^\prime \end{pmatrix}$, $\mathbf{d}_L^\prime = \begin{pmatrix} d_L^\prime \\ s_L^\prime \\ b_L^\prime \end{pmatrix}$, $\mathbf{d}_R^\prime = \begin{pmatrix} d_R^\prime \\ s_R^\prime \\ b_R^\prime \end{pmatrix}$

$$\mathcal{L}_{F-S,Q} = -\frac{\nu + H}{\sqrt{2}} \left[\overline{\boldsymbol{d}}'_{\boldsymbol{L}} Y'^{d} \boldsymbol{d}'_{\boldsymbol{R}} + \overline{\boldsymbol{u}}'_{\boldsymbol{L}} Y'^{u} \boldsymbol{u}'_{\boldsymbol{R}} \right] + h. c.$$

$$V_L^{d\dagger} Y'^d V_R^d = Y^d, \qquad Y_{ij}^d = y_i^d \delta_{ij}$$

$$V_L^{u\dagger}Y'^uV_R^u=Y^u, \qquad Y_{ij}^u=y_i^u\delta_{ij}$$

4. 尝试化简 $\mathcal{L} = -\sum_{i=1,2,3} \sum_{i=u,c,t} (Y_{ii}^{\prime u} \bar{Q}_{iL}^i \tilde{\Phi} u_{iR}^{\prime})$,并指出质量项和相互作用项。

$$\mathcal{L}_{F-S,Q} = -\frac{v + H}{\sqrt{2}} \left[\overline{\boldsymbol{d}}_{\boldsymbol{L}} Y^{d} \boldsymbol{d}_{\boldsymbol{R}} + \overline{\boldsymbol{u}}_{\boldsymbol{L}} Y^{u} \boldsymbol{u}_{\boldsymbol{R}} \right] + h. c.$$

$$\boldsymbol{u}_{\boldsymbol{L}} = V_{L}^{u\dagger} \boldsymbol{u}_{\boldsymbol{L}}' = \begin{pmatrix} u_{L} \\ c_{L} \\ t_{L} \end{pmatrix}, \boldsymbol{u}_{\boldsymbol{R}} = V_{R}^{u\dagger} \boldsymbol{u}_{\boldsymbol{R}}' = \begin{pmatrix} u_{R} \\ c_{R} \\ t_{R} \end{pmatrix}$$

$$\boldsymbol{d}_{\boldsymbol{L}} = V_{L}^{d\dagger} \boldsymbol{d}_{\boldsymbol{L}}' = \begin{pmatrix} d_{L} \\ s_{L} \\ b_{L} \end{pmatrix}, \boldsymbol{d}_{\boldsymbol{R}} = V_{R}^{d\dagger} \boldsymbol{d}_{\boldsymbol{R}}' = \begin{pmatrix} d_{R} \\ s_{R} \\ b_{R} \end{pmatrix}$$

$$\mathcal{L}_{F-S,Q} = -\frac{\nu + H(x)}{\sqrt{2}} \sum_{i=u,c,t} y_i^u \bar{u}_i u_i = -\sum_{i=u,c,t} \frac{y_i^u \nu}{\sqrt{2}} \bar{u}_i u_i - \sum_{i=u,c,t} \frac{y_i^u}{\sqrt{2}} \bar{u}_i u_i H$$

4. 尝试化简 $\mathcal{L} = -\sum_{i=1,2,3} \sum_{j=u,c,t} (Y'^u_{ij} \bar{Q}^i_{iL} \tilde{\Phi} u'_{jR})$,并指出质量项和相互作用项。

$$j_{W,Q}^{\mu} = 2\overline{\boldsymbol{u}}_{\boldsymbol{L}}^{\prime} \boldsymbol{\gamma}^{\mu} \boldsymbol{d}_{\boldsymbol{L}}^{\prime}$$

$$j_{W,Q}^{\mu} = 2\overline{\boldsymbol{u}}_{\boldsymbol{L}} V_{L}^{u\dagger} \boldsymbol{\gamma}^{\mu} V_{L}^{d} \boldsymbol{d}_{\boldsymbol{L}} = 2\overline{\boldsymbol{u}}_{\boldsymbol{L}} \boldsymbol{\gamma}^{\mu} V_{L}^{u\dagger} V_{L}^{d} \boldsymbol{d}_{\boldsymbol{L}} = 2\overline{\boldsymbol{u}}_{\boldsymbol{L}} \boldsymbol{\gamma}^{\mu} V \boldsymbol{d}_{\boldsymbol{L}}$$

$$V \equiv V_{L}^{u\dagger} V_{L}^{d}$$

$$\begin{split} j_{Z,Q}^{\mu} &= 2g_L^u \overline{u}_L' \gamma^{\mu} u_L' + 2g_R^u \overline{u}_R' \gamma^{\mu} u_R' + 2g_L^d \overline{d}_L' \gamma^{\mu} d_L' + 2g_R^d \overline{d}_R' \gamma^{\mu} d_R' \\ &= 2g_L^u \overline{u}_L V_L^{u\dagger} \gamma^{\mu} V_L^u u_L + 2g_R^u \overline{u}_R V_R^{u\dagger} \gamma^{\mu} V_R^u u_R \\ &\quad + 2g_L^d \overline{d}_L V_L^{d\dagger} \gamma^{\mu} V_L^d d_L + 2g_R^d \overline{d}_R V_R^{d\dagger} \gamma^{\mu} V_R^d d_R \\ &= 2g_L^u \overline{u}_L \gamma^{\mu} u_L + 2g_R^u \overline{u}_R \gamma^{\mu} u_R + 2g_L^d \overline{d}_L \gamma^{\mu} d_L + 2g_R^d \overline{d}_R \gamma^{\mu} d_R \end{split}$$

5. 画出LHC上pp对撞产生Higgs四种主要模式的费曼图

