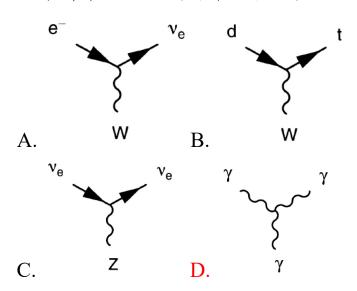
一、单选

1. 以下四组国际单位制-自然单位制的换算, 正确的有:

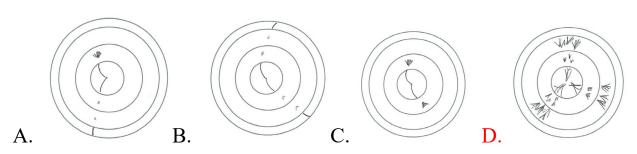
 $c={\rm speed}$ of light = $2.9979\times 10^8~{\rm m/s}$ $\hbar={\rm reduced~Planck~constant}=1.0546\times 10^{-34}\,{\rm J~s}$

- ①1GeV = 1.8×10^{-27} kg
- $21 \text{GeV}^2 = 8.19 \times 10^{-5} \text{N}$
- 31GeV = 5.39×10^{-19} kg · m/s
- $401 \text{GeV}^{-1} = 6.58 \times 10^{-28} \text{s}$
- A.14
- B.13
- C.24
- D.1123

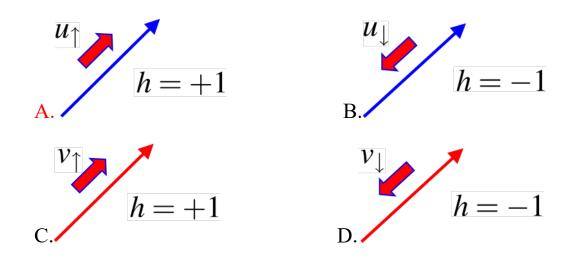
2. 以下哪个选项不是标准模型顶点:



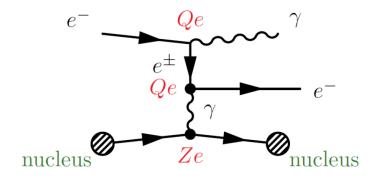
3. 以下选项中, 哪项最可能是 $e^+e^- \rightarrow Z \rightarrow qq$ 留下的径迹:



4. 选出正粒子的右手螺旋度本征态:



5. 根据图像估算并选出韧致辐射的矩阵元量级 (原子核核电荷数为Z):



 $A.M \propto e^3$ $B.M \propto Z^2 e^4$ $C.M \propto Z e^3$ $D.M \propto Z^2 e^2$

6. 以下手征算符作用结果错误的是:

$$A.P_R u_R = u_R$$

$$A.P_R u_R = u_R \qquad B. P_L u_R = 0$$

$$C.P_L v_R = v_L \qquad D. P_R v_L = 0$$

$$D. P_R v_L = 0$$

7. $\rho^0 \rightarrow e^+e^-$ 的衰变宽度与 $\omega^0 \rightarrow e^+e^-$ 的衰变宽度之比大约等于? 其中, π^0 , ρ^0 组分为 $(u\bar{u}-d\bar{d})/\sqrt{2}$, ω^0 为 $(u\bar{u}+d\bar{d})/\sqrt{2}$

C.3 $D.\frac{1}{3}$ A.9 B.81

二、判断

- 考虑一个二体散射过程,若为弹性碰撞,质心系中算得的衰变截面与 质心能量的平方成反比。√
- 2. 相较电磁簇射,强子簇射虽然过程更加复杂,但是相互作用强度反而更短。×
- 3. 强子簇射过程中也会发生电磁簇射。√
- 4. 因为反粒子沿着时间反向传播,所以费曼图上反粒子的箭头总是与时间相反。×
- 5. t-channel 的传播子总是离壳的。√
- 6. 深度非弹性碰撞中, 微分散射截面公式中的电磁结构函数和纯磁结构 函数互相独立。×
- 7. 如果夸克自旋为 0, 深度非弹性碰撞微分散射截面公式中的纯磁结构 函数也为 0.√
- 8. 强作用和电磁作用过程中同位旋守恒。×
- 9. 自然界存在色波函数为 $\frac{1}{\sqrt{3}}(r\bar{r}+g\bar{g}+b\bar{b})$ 的胶子。×
- 10.与 QED 的跑动耦合常数不同,QCD 的跑动耦合常数随着 Q^2 下降。✓

三、简答与计算

1. 考虑一个二体衰变过程 $X \to 1 + 2$, 计算末态的不变质量。

$$M_X^2 = m_1^2 + m_2^2 + 2(E_1E_2 - |\vec{p_1}||\vec{p_2}|\cos\theta)$$

2. 考虑 $1+2 \rightarrow 3+4$ 的过程, 画出 s-channel 和 t-channel 的图像, 并写出 s 和 t 的定义式。

$$s = (p_1 + p_2)^2$$
, $t = (p_1 - p_3)^2$

3. 推导对撞机和固定靶的质心能量。计算要达到质子-质子 200GeV 质心

Fixed Target Collision:

$$s = m_1^2 + m_2^2 + 2E_1 m_2$$
 $p_1^{\mu}(\mathbf{E}, \mathbf{\bar{p}})$ $p_2^{\mu}(\mathbf{m}_2, \mathbf{0})$

For $E_1\gg m_1,m_2$, we have $s=2E_1m_2$ The energy in the CM frame $\sqrt{s}=\sqrt{2E_1m_2}$ e.g. $100\,\mathrm{GeV}$ proton hitting a proton ar rest

$$\sqrt{s} = \sqrt{2E_p m_p} \approx \sqrt{2 \times 100 \times 1} \approx 14 \,\mathrm{GeV}$$

Collider Experiment (pp colliding head on):

$$s = m_1^2 + m_2^2 + 2(E_1 E_2 - |\mathbf{p}_1||\mathbf{p}_2|\cos\theta) \qquad \qquad \overrightarrow{\mathbf{p}_1^{\mu}(\mathbf{E},\overline{\mathbf{p}})} \qquad \qquad \overleftarrow{\mathbf{p}_2^{\mu}(\mathbf{E},\overline{\mathbf{p}})}$$

If $E\gg m_1,m_2$, then $s=2(E^2-E^2\cos\theta)=4E^2$ The energy in the CM frame $\sqrt{s}=2E$ e.g. $100\,{\rm GeV}$ proton colliding with a $100\,{\rm GeV}$ proton

$$\sqrt{s} = 2E = 2 \times 100 = 200 \,\text{GeV}$$

4. γ矩阵相关计算:

a) 定义:
$$\gamma^5 = i\gamma^0\gamma^1\gamma^2\gamma^3$$
, 证明:

i.
$$(\gamma^5)^2 = I$$
.

ii.
$$\{\gamma^{\mu}, \gamma^{5}\} = 0$$
.

iii.
$$(\gamma^5)^{\dagger} = \gamma^5$$
.

(i)
$$(\gamma^5)^2 = -1 \times \gamma^0 \gamma^1 \gamma^2 \gamma^3 \gamma^0 \gamma^1 \gamma^2 \gamma^3 = -1 \times 1 \times -1 \times -1 \times -1 = 1$$

(ii)
$$\{\gamma^{0}, \gamma^{5}\} = i\gamma^{0}\gamma^{0}\gamma^{1}\gamma^{2}\gamma^{3} + i\gamma^{0}\gamma^{1}\gamma^{2}\gamma^{3}\gamma^{0} = i(\gamma^{1}\gamma^{2}\gamma^{3} - \gamma^{1}\gamma^{2}\gamma^{3}) = 0$$

 $\{\gamma^{1}, \gamma^{5}\} = i\gamma^{1}\gamma^{0}\gamma^{1}\gamma^{2}\gamma^{3} + i\gamma^{0}\gamma^{1}\gamma^{2}\gamma^{3}\gamma^{1} = i(\gamma^{0}\gamma^{2}\gamma^{3} - \gamma^{0}\gamma^{2}\gamma^{3}) = 0$
 $\{\gamma^{2}, \gamma^{5}\} = i\gamma^{2}\gamma^{0}\gamma^{1}\gamma^{2}\gamma^{3} + i\gamma^{0}\gamma^{1}\gamma^{2}\gamma^{3}\gamma^{2} = i(-\gamma^{0}\gamma^{1}\gamma^{3} + \gamma^{0}\gamma^{1}\gamma^{3}) = 0$
 $\{\gamma^{3}, \gamma^{5}\} = i\gamma^{1}\gamma^{0}\gamma^{1}\gamma^{2}\gamma^{3} + i\gamma^{0}\gamma^{1}\gamma^{2}\gamma^{3}\gamma^{3} = i(\gamma^{0}\gamma^{1}\gamma^{2} - \gamma^{0}\gamma^{1}\gamma^{2}) = 0$
(iii) $(\gamma^{5})^{\dagger} = -i(\gamma^{3})^{\dagger}(\gamma^{2})^{\dagger}(\gamma^{1})^{\dagger}(\gamma^{0})^{\dagger} = i\gamma^{3}\gamma^{2}\gamma^{1}\gamma^{0} = i\gamma^{0}\gamma^{1}\gamma^{2}\gamma^{3} = \gamma^{5}$

b) 利用
$$\{\gamma^{\mu}, \gamma^{\nu}\} = 2g^{\mu\nu}$$
, 尝试证明:

$$\gamma^{\mu}\gamma_{\mu} = 4$$
, $\gamma^{\mu}\alpha\gamma_{\mu} = -2\alpha$, $\gamma^{\mu}\alpha\beta\gamma_{\mu} = 4\alpha \cdot b$

$$\mathbf{i)} \, \gamma^\mu \gamma_\mu = g_{\mu\nu} \gamma^\mu \gamma^\nu = g_{\mu\nu} (2g^{\mu\nu} - \gamma^\nu \gamma^\mu) = 8 - g_{\mu\nu} \gamma^\nu \gamma^\mu = 8 - \gamma_\mu \gamma^\mu \\ \gamma^\mu \gamma_\mu = 4$$

$$\text{ii) } \gamma^{\mu} \not a \gamma_{\mu} = \gamma^{\mu} \gamma^{\nu} a_{\nu} \gamma_{\mu} = (2g^{\mu\nu} - \gamma^{\nu} \gamma^{\mu}) a_{\nu} \gamma_{\mu} = 2g^{\mu\nu} a_{\nu} \gamma_{\mu} - \gamma^{\nu} \gamma^{\mu} a_{\nu} \gamma_{\mu} = 2\not a - 4\not a = -2\not a$$

$$\begin{aligned} \text{iii)} \ \gamma^{\mu} \mathbf{\alpha} \mathbf{b} \gamma_{\mu} &= \gamma^{\mu} \gamma^{\nu} a_{\nu} \gamma^{\sigma} b_{\sigma} \gamma_{\mu} = \gamma^{\mu} \gamma^{\nu} a_{\nu} b^{\sigma} \gamma_{\sigma} \gamma_{\mu} = (2g^{\mu\nu} - \gamma^{\nu} \gamma^{\mu}) a_{\nu} b^{\sigma} \left(2g_{\sigma\mu} - \gamma_{\mu} \gamma_{\sigma} \right) \\ &= 4g^{\mu\nu} a_{\nu} b^{\sigma} g_{\sigma\mu} - 2g_{\sigma\mu} \gamma^{\nu} \gamma^{\mu} a_{\nu} b^{\sigma} - 2g^{\mu\nu} a_{\nu} b^{\sigma} \gamma_{\mu} \gamma_{\sigma} + \gamma^{\nu} \gamma^{\mu} a_{\nu} b^{\sigma} \gamma_{\mu} \gamma_{\sigma} \\ &= 4a \cdot b - 2\mathbf{a} \mathbf{b} - 2\mathbf{a} \mathbf{b} + 4\mathbf{a} \mathbf{b} = 4a \cdot b \end{aligned}$$

c) 由手紅算符
$$P_R = \frac{1}{2}(1+\gamma^5), P_L = \frac{1}{2}(1-\gamma^5),$$
 证明:
$$P_L + P_R = 1, \qquad P_R P_R = P_R, \qquad P_L P_L = P_L, \qquad P_L P_R = 0.$$
 i) $P_L + P_R = \frac{1}{2}(1+\gamma^5) + \frac{1}{2}(1-\gamma^5) = 1$ ii) $P_R P_R = \frac{1}{4}(1+\gamma^5)(1+\gamma^5) = \frac{1}{4}(1+2\gamma^5+1) = P_R$ iii) $P_L P_L = \frac{1}{4}(1-\gamma^5)(1-\gamma^5) = \frac{1}{4}(1-2\gamma^5+1) = P_L$ iv) $P_L P_R = \frac{1}{4}(1-\gamma^5)(1+\gamma^5) = \frac{1}{4}(1+\gamma^5-\gamma^5-1) = 0$

5. 只考虑*u*, *d*夸克, 考虑它们的 SU(2)对称性, 用图文并茂的方式给出重 子所有可能的同位旋态。

$$\begin{split} |\frac{3}{2}, +\frac{3}{2}\rangle &= uuu \\ |\frac{3}{2}, +\frac{1}{2}\rangle &= \frac{1}{\sqrt{3}}(uud + udu + duu) \\ |\frac{3}{2}, -\frac{1}{2}\rangle &= \frac{1}{\sqrt{3}}(ddu + dud + udd) \\ |\frac{3}{2}, -\frac{3}{2}\rangle &= ddd \\ |\frac{1}{2}, -\frac{1}{2}\rangle &= -\frac{1}{\sqrt{6}}(2ddu - udd - dud) \\ |\frac{1}{2}, +\frac{1}{2}\rangle &= \frac{1}{\sqrt{6}}(2uud - udu - duu) \\ |\frac{1}{2}, -\frac{1}{2}\rangle &= \frac{1}{\sqrt{2}}(udd - dud) \\ |\frac{1}{2}, +\frac{1}{2}\rangle &= \frac{1}{\sqrt{2}}(udu - duu) \end{split}$$

6. 考虑ee → μμ的湮灭产生过程, 画出费曼图并根据 QED 费曼规则写出 矩阵元

$$\mathbf{e}^{+} \begin{array}{ccc} p_{2} & p_{4} & \mathbf{\mu}^{+} \\ \mu & & v \\ \mathbf{e}^{-} & p_{1} & p_{3} & \mathbf{\mu}^{-} \end{array} \qquad -iM = \left[\overline{v}(p_{2})ie\gamma^{\mu}u(p_{1})\right] \frac{-ig_{\mu\nu}}{q^{2}} \left[\overline{u}(p_{3})ie\gamma^{\nu}v(p_{4})\right]$$

7. 考虑 $q\bar{q} \to q\bar{q}$ 的散射过程(t道), 画出费曼图并根据QCD 费曼规则写出矩阵元(将颜色部分放在一起)。

$$\mathbf{u} = \begin{bmatrix} u & p_1 \\ i & p_3 \end{bmatrix} \mathbf{u}$$

$$\mathbf{d} = \begin{bmatrix} \overline{u}_u(p_3) \{-\frac{1}{2} i g_s \lambda_{ji}^a \gamma^{\mu}\} u_u(p_1) \end{bmatrix} \frac{-i g_{\mu\nu}}{q^2} \delta^{ab} \left[\overline{u}_d(p_4) \{-\frac{1}{2} i g_s \lambda_{lk}^b \gamma^{\nu}\} u_d(p_2) \right]$$

8. QED 中费米子和光子的基本相互作用顶角因子为 iey^{μ} ,费米子与光子的相互作用可以表示为四矢量流 $j^{\mu}=ie\bar{\psi}y^{\mu}\phi$.

$$P_{R} = \frac{1}{2}(1 + \gamma^{5}), \qquad P_{L} = \frac{1}{2}(1 - \gamma^{5})$$
 $P_{R}u_{R} = u_{R}, \qquad P_{R}u_{L} = 0, \qquad P_{L}u_{R} = 0, \qquad P_{L}u_{L} = u_{L}$
 $P_{R}v_{R} = 0, \qquad P_{R}v_{L} = v_{L}, \qquad P_{L}v_{R} = v_{R}, \qquad P_{L}v_{L} = 0$

(a) 证明任意旋量可以写成左手和右手手征分量,并将 QED 中的 四矢量流用旋量的左手和右手手征分量展开。

$$\psi = \psi_R + \psi_L = \frac{1}{2} (1 + \gamma^5) \psi + \frac{1}{2} (1 - \gamma^5) \psi$$
$$j^{\mu} = \bar{\psi} \gamma^5 \phi = \bar{\psi}_R \gamma^5 \phi_R + \bar{\psi}_R \gamma^5 \phi_L + \bar{\psi}_L \gamma^5 \phi_R + \bar{\psi}_L \gamma^5 \phi_L$$

(b) 利用简答题 4 中的性质,证明 $\bar{\mu}_L \gamma^\mu \mu_R = 0$ 和 $\bar{\nu}_L \gamma^\mu \mu_L = 0$.

$$\bar{u}_L \gamma^{\mu} u_R = \frac{1}{2} u^{\dagger} (1 - \gamma^5) \gamma^0 \gamma^{\mu} \frac{1}{2} (1 + \gamma^5) u = \frac{1}{4} \bar{u} \gamma^{\mu} (1 - \gamma^5) (1 + \gamma^5) u$$
$$= 0$$

$$\bar{v}_L \gamma^\mu u_L = \frac{1}{2} v^\dagger (1 + \gamma^5) \gamma^0 \gamma^\mu \frac{1}{2} (1 - \gamma^5) u = \frac{1}{4} \bar{v} \gamma^\mu (1 + \gamma^5) (1 - \gamma^5) u$$
$$= 0$$

(c) 实际上,弱相互作用中顶角因子不是矢量形式 (γ^{μ}) ,而是一个特殊的 $\gamma^{\mu}(1-\gamma^{5})$ 的形式,试证明:只有粒子旋量的左手手征分量 μ_{L} 和反粒子旋量的右手手征分量 ν_{R} 参与带电粒弱相互作用。

$$\frac{1}{2}\bar{\psi}\gamma^{\mu}(1-\gamma^{5})\phi = \bar{\psi}\gamma^{\mu}\phi_{L} = \bar{\psi}_{L}\gamma^{\mu}\phi_{L}$$