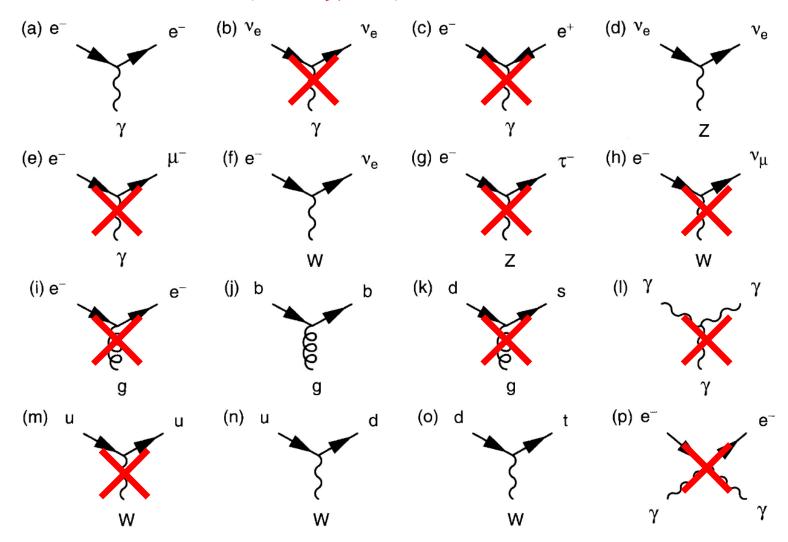
# 粒子物理学

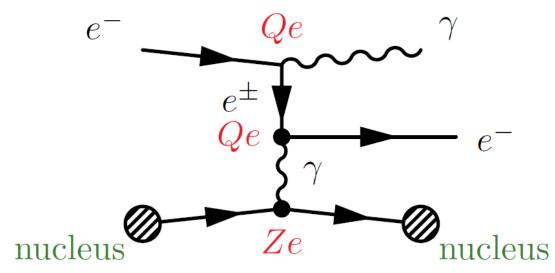
### 习题课2

### 1. 下面这些过程是否是标准模型费曼图?



2. 画出电子与原子核(核电荷数为Ze)发生电磁相互作用, 韧致辐射出光子的最低阶费曼图, 并且估计其反应截面

### Bremsstrahlung ( $e^- \rightarrow e^- \gamma$ )



$$M \propto Ze^3$$

$$\sigma \propto |M|^2 \propto Z^2e^6$$

$$\propto (4\pi)^3 Z^2 \alpha^3$$

#### 3. γ矩阵相关证明题:

a) 利用
$$\{\gamma^{\mu}, \gamma^{\nu}\} = 2g^{\mu\nu}$$
,尝试证明: 
$$\gamma^{\mu}\gamma_{\mu} = 4, \gamma^{\mu}\not{a}\gamma_{\mu} = -2\not{a}, \gamma^{\mu}\not{a}\not{b}\gamma_{\mu} = 4a \cdot b$$

 $= 4a \cdot b - 2ab - 2ab + 4ab = 4a \cdot b$ 

$$\begin{aligned} \textbf{i)} \ \gamma^{\mu} \gamma_{\mu} &= g_{\mu\nu} \gamma^{\mu} \gamma^{\nu} = g_{\mu\nu} (2g^{\mu\nu} - \gamma^{\nu} \gamma^{\mu}) = 8 - g_{\mu\nu} \gamma^{\nu} \gamma^{\mu} = 8 - \gamma_{\mu} \gamma^{\mu} \\ \gamma^{\mu} \gamma_{\mu} &= 4 \end{aligned} \\ \textbf{ii)} \ \gamma^{\mu} \alpha \gamma_{\mu} &= \gamma^{\mu} \gamma^{\nu} a_{\nu} \gamma_{\mu} = (2g^{\mu\nu} - \gamma^{\nu} \gamma^{\mu}) a_{\nu} \gamma_{\mu} = 2g^{\mu\nu} a_{\nu} \gamma_{\mu} - \gamma^{\nu} \gamma^{\mu} a_{\nu} \gamma_{\mu} = 2\alpha - 4\alpha = -2\alpha \end{aligned} \\ \textbf{iii)} \ \gamma^{\mu} \alpha \beta \gamma_{\mu} &= \gamma^{\mu} \gamma^{\nu} a_{\nu} \gamma^{\sigma} b_{\sigma} \gamma_{\mu} = \gamma^{\mu} \gamma^{\nu} a_{\nu} b^{\sigma} \gamma_{\sigma} \gamma_{\mu} = (2g^{\mu\nu} - \gamma^{\nu} \gamma^{\mu}) a_{\nu} b^{\sigma} (2g_{\sigma\mu} - \gamma_{\mu} \gamma_{\sigma}) \\ &= 4g^{\mu\nu} a_{\nu} b^{\sigma} g_{\sigma\mu} - 2g_{\sigma\mu} \gamma^{\nu} \gamma^{\mu} a_{\nu} b^{\sigma} - 2g^{\mu\nu} a_{\nu} b^{\sigma} \gamma_{\mu} \gamma_{\sigma} + \gamma^{\nu} \gamma^{\mu} a_{\nu} b^{\sigma} \gamma_{\mu} \gamma_{\sigma} \end{aligned}$$

### 3. γ矩阵相关证明题:

b) 由手征算符
$$P_R = \frac{1}{2} (1 + \gamma^5), P_L = \frac{1}{2} (1 - \gamma^5),$$
 证明:  
 $P_L + P_R = 1, \quad P_R P_R = P_R, \quad P_L P_L = P_L, \quad P_L P_R = 0$ 

i) 
$$P_L + P_R = \frac{1}{2} (1 + \gamma^5) + \frac{1}{2} (1 - \gamma^5) = 1$$

ii)
$$P_R P_R = \frac{1}{4} (1 + \gamma^5)(1 + \gamma^5) = \frac{1}{4} (1 + 2\gamma^5 + 1) = P_R$$

iii) 
$$P_L P_L = \frac{1}{4} (1 - \gamma^5) (1 - \gamma^5) = \frac{1}{4} (1 - 2\gamma^5 + 1) = P_L$$

iv) 
$$P_L P_R = \frac{1}{4}(1 - \gamma^5)(1 + \gamma^5) = \frac{1}{4}(1 + \gamma^5 - \gamma^5 - 1) = 0$$

c) 
$$\left[\overline{\psi}\gamma^{\mu}\gamma^{5}\phi\right]^{\dagger} = \overline{\phi}\gamma^{\mu}\gamma^{5}\psi$$

$$\begin{split} [\bar{\psi}\gamma^{\mu}\gamma^{5}\phi]^{\dagger} &= \left[\psi^{\dagger}\gamma^{0}\gamma^{\mu}\gamma^{5}\phi\right]^{\dagger} = \phi^{\dagger}\gamma^{5}(\gamma^{\mu})^{\dagger}\gamma^{0}\psi = \phi^{\dagger}\gamma^{0}\gamma^{0}\gamma^{5}\gamma^{0}\gamma^{\mu}\gamma^{0}\gamma^{0}\psi \\ &= \bar{\phi}\gamma^{0}\gamma^{5}\gamma^{0}\gamma^{\mu}\psi = -\bar{\phi}\gamma^{0}\gamma^{0}\gamma^{5}\gamma^{\mu}\psi = \bar{\phi}\gamma^{\mu}\gamma^{5}\psi \end{split}$$

#### 4. 证明在极端相对论近似下:

$$\frac{(p_1 \cdot p_3)^2 + (p_1 \cdot p_4)^2}{(p_1 \cdot p_2)^2} = \frac{t^2 + u^2}{s^2}$$

$$t = (p_1 - p_3)^2 \approx -2p_1p_3$$
,  $s \approx 2p_1p_2 = 2p_3p_4$ ,  $u \approx 2p_1p_4 = 2p_2p_3$ 

#### 5. 模仿课件中计算muon current的做法, 计算electron current。

$$u_{\uparrow} = \sqrt{E} \begin{pmatrix} 1 \\ 0 \\ 1 \\ 0 \end{pmatrix}, u_{\downarrow} = \sqrt{E} \begin{pmatrix} 0 \\ 1 \\ 0 \\ -1 \end{pmatrix}, v_{\uparrow} = \sqrt{E} \begin{pmatrix} 1 \\ 0 \\ -1 \\ 0 \end{pmatrix}, v_{\downarrow} = -\sqrt{E} \begin{pmatrix} 1 \\ 0 \\ 1 \\ 0 \end{pmatrix}$$

$$\overline{\psi} \gamma^{0} \phi = \psi^{\dagger} \gamma^{0} \gamma^{0} \phi = \psi_{1}^{*} \phi_{1} + \psi_{2}^{*} \phi_{2} + \psi_{3}^{*} \phi_{3} + \psi_{4}^{*} \phi_{4}$$

$$\overline{\psi} \gamma^{1} \phi = \psi^{\dagger} \gamma^{0} \gamma^{1} \phi = \psi_{1}^{*} \phi_{4} + \psi_{2}^{*} \phi_{3} + \psi_{3}^{*} \phi_{2} + \psi_{4}^{*} \phi_{1}$$

$$\overline{\psi} \gamma^{2} \phi = \psi^{\dagger} \gamma^{0} \gamma^{2} \phi = -i(\psi_{1}^{*} \phi_{4} - \psi_{2}^{*} \phi_{3} + \psi_{3}^{*} \phi_{2} - \psi_{4}^{*} \phi_{1})$$

$$\overline{\psi} \gamma^{3} \phi = \psi^{\dagger} \gamma^{0} \gamma^{3} \phi = \psi_{1}^{*} \phi_{3} - \psi_{2}^{*} \phi_{4} + \psi_{3}^{*} \phi_{1} - \psi_{4}^{*} \phi_{2}$$

$$j_{e}^{\mu} = \overline{v}(p_{2}) \gamma^{\mu} u(p_{1})$$

$$\overline{v}_{\uparrow}(p_{2}) \gamma^{\mu} u_{\uparrow}(p_{1}) = (0,0,0,0)$$

$$\overline{v}_{\uparrow}(p_{2}) \gamma^{\mu} u_{\downarrow}(p_{1}) = 2E(0,-1,i,0)$$

$$\overline{v}_{\downarrow}(p_{2}) \gamma^{\mu} u_{\uparrow}(p_{1}) = 2E(0,-1,-i,0)$$

$$\overline{v}_{\downarrow}(p_{2}) \gamma^{\mu} u_{\downarrow}(p_{1}) = (0,0,0,0)$$

6. 阅读Thomson书本6.5节,尝试使用trace的技巧,计 $fe^-\mu^- \rightarrow e^-\mu^-$ 的矩阵元的平方:

$$-iM = \bar{u}(p_{3}s_{3})ie\gamma^{\mu}u(p_{1}s_{1})\frac{-ig_{\mu\nu}}{q^{2}}\bar{u}(p_{4}s_{4})ie\gamma^{\mu}u(p_{2}s_{2})$$

$$\langle \left|M_{fi}\right|^{2}\rangle = 2e^{4}\frac{s^{2}+u^{2}}{t^{2}}$$

第五章PPT——3-8页 Thomson书——6.5节

- 1. 证明一些γ矩阵迹的性质:
- a)  $\operatorname{Tr}(\gamma^{\mu}\gamma^{\nu}) = 4g^{\mu\nu}$

$$\gamma^{\mu}\gamma^{\nu} + \gamma^{\nu}\gamma^{\mu} \equiv 2g^{\mu\nu}I,$$

$$\operatorname{Tr}(\gamma^{\mu}\gamma^{\nu}) + \operatorname{Tr}(\gamma^{\nu}\gamma^{\mu}) = 2g^{\mu\nu}\operatorname{Tr}(I),$$

$$\operatorname{Tr}(\gamma^{\mu}\gamma^{\nu}) = 4g^{\mu\nu}.$$

**b)**  $Tr(奇数个 \gamma矩阵) = 0$ 

$$\begin{split} \operatorname{Tr}\left(\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}\right) &= \operatorname{Tr}\left(\gamma^{5}\gamma^{5}\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}\right) \\ &= \operatorname{Tr}\left(\gamma^{5}\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}\gamma^{5}\right) \\ &= -\operatorname{Tr}\left(\gamma^{5}\gamma^{5}\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}\right) \\ &= -\operatorname{Tr}\left(\gamma^{5}\gamma^{5}\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}\right) \\ \operatorname{Tr}\left(\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}\right) &= -\operatorname{Tr}\left(\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}\right), \end{split} \tag{traces are cyclical}$$

$$\operatorname{Tr}(\gamma^{\mu}\gamma^{\nu}\gamma^{\rho})=0.$$

#### 1. 证明一些γ矩阵迹的性质:

c) 
$$\operatorname{Tr}(\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}\gamma^{\sigma}) = 4g^{\mu\nu}g^{\rho\sigma} - 4g^{\mu\rho}g^{\nu\sigma} + 4g^{\mu\sigma}g^{\nu\rho}$$

$$\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}\gamma^{\sigma} = 2g^{\mu\nu}\gamma^{\rho}\gamma^{\sigma} - \gamma^{\nu}\gamma^{\mu}\gamma^{\rho}\gamma^{\sigma}$$

$$= 2g^{\mu\nu}\gamma^{\rho}\gamma^{\sigma} - 2g^{\mu\rho}\gamma^{\nu}\gamma^{\sigma} + \gamma^{\nu}\gamma^{\rho}\gamma^{\mu}\gamma^{\sigma}$$

$$= 2g^{\mu\nu}\gamma^{\rho}\gamma^{\sigma} - 2g^{\mu\rho}\gamma^{\nu}\gamma^{\sigma} + 2g^{\mu\sigma}\gamma^{\nu}\gamma^{\rho} - \gamma^{\nu}\gamma^{\rho}\gamma^{\sigma}\gamma^{\mu}$$

$$\Rightarrow \gamma^{\mu}\gamma^{\nu}\gamma^{\rho}\gamma^{\sigma} + \gamma^{\nu}\gamma^{\rho}\gamma^{\sigma}\gamma^{\mu} = 2g^{\mu\nu}\gamma^{\rho}\gamma^{\sigma} - 2g^{\mu\rho}\gamma^{\nu}\gamma^{\sigma} + 2g^{\mu\sigma}\gamma^{\nu}\gamma^{\rho}. \qquad (6.57)$$

$$2\text{Tr}\left(\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}\gamma^{\sigma}\right) = 2g^{\mu\nu}\text{Tr}\left(\gamma^{\rho}\gamma^{\sigma}\right) - 2g^{\mu\rho}\text{Tr}\left(\gamma^{\nu}\gamma^{\sigma}\right) + 2g^{\mu\sigma}\text{Tr}\left(\gamma^{\nu}\gamma^{\rho}\right),$$

$$\operatorname{Tr}\left(\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}\gamma^{\sigma}\right)=4g^{\mu\nu}g^{\rho\sigma}-4g^{\mu\rho}g^{\nu\sigma}+4g^{\mu\sigma}g^{\nu\rho}.$$

#### 2. 第一课中得到, 微分截面满足公式:

$$\frac{d\sigma}{dt} = \frac{1}{64\pi s \left|\vec{p}_i^*\right|^2} \left|M_{fi}\right|^2$$

第四节课中又得到,对 $e^-\mu^- \rightarrow e^-\mu^-$ 散射过程,有:

$$\left| M_{fi} \right|^2 = 2e^4 \frac{s^2 + u^2}{t^2}$$

从这两个式子证明:

$$\begin{split} \frac{d\sigma}{dq^2} &= \frac{2\pi\alpha^2}{q^4} \left[ 1 + \left( 1 + \frac{q^2}{s} \right)^2 \right] \\ |\vec{p}_i^*|^2 &= \frac{1}{4s} \left[ s - (m_1 + m_2)^2 \right] \left[ s - (m_1 - m_2)^2 \right] \approx \frac{s}{4} \\ \frac{d\sigma}{dt} &= \frac{1}{16\pi s^2} 2e^4 \frac{s^2 + u^2}{t^2} = \frac{e^4}{8\pi t^2} \frac{s^2 + u^2}{s^2} & t = q^2 \\ &= \frac{2\pi\alpha^2}{q^4} \left[ 1 + \frac{(s+t)^2}{s^2} \right] & \alpha &= \frac{e^2}{4\pi} \\ &= \frac{2\pi\alpha^2}{q^4} \left[ 1 + \left( 1 + \frac{q^2}{s} \right)^2 \right] & s + t + u = \sum m^2 = 0 \end{split}$$

3. 模仿课上对卢瑟福散射的推导,尝试推导莫特散射公式(提示:莫特散射中靶的反弹可忽略且忽略电子质量)

$$u_{\uparrow}(p_1) = N_e \begin{pmatrix} 1 \\ 0 \\ \alpha \\ 0 \end{pmatrix} \quad u_{\downarrow}(p_1) = N_e \begin{pmatrix} 0 \\ 1 \\ 0 \\ -\alpha \end{pmatrix} \qquad u_{\uparrow}(p_3) = N_e \begin{pmatrix} c \\ s \\ \alpha c \\ \alpha s \end{pmatrix} \qquad u_{\downarrow}(p_3) = N_e \begin{pmatrix} -s \\ c \\ \alpha s \\ -\alpha c \end{pmatrix} . \mathbf{e}^{-} \underbrace{\begin{array}{c} p_3 \\ -s \\ \alpha s \\ -\alpha c \end{array}}_{\text{(neglect proton recoil)}}$$

$$\overline{u}_{\uparrow}(p_3)\gamma^{\mu}u_{\uparrow}(p_1) = (E + p_e)\left[(\alpha^2 + 1)c, 2\alpha s, -2i\alpha s, 2\alpha c\right]$$

$$\overline{u}_{\downarrow}(p_3)\gamma^{\mu}u_{\downarrow}(p_1) = (E + p_e)\left[(\alpha^2 + 1)c, 2\alpha s, -2i\alpha s, 2\alpha c\right]$$

$$\overline{u}_{\uparrow}(p_3)\gamma^{\mu}u_{\downarrow}(p_1) = (E + m_e)\left[(1 - \alpha^2)s, 0, 0, 0\right]$$

$$\overline{u}_{\downarrow}(p_3)\gamma^{\mu}u_{\uparrow}(p_1) = (E + m_e)\left[(\alpha^2 - 1)s, 0, 0, 0\right]$$

极端相对论极限:  $\alpha \to 1$ 

$$u_{\uparrow}(0) = \sqrt{2M_p} \begin{pmatrix} 1\\0\\0\\0 \end{pmatrix} \quad u_{\downarrow}(0) = \sqrt{2M_p} \begin{pmatrix} 0\\1\\0\\0 \end{pmatrix} \quad j_{p\uparrow\uparrow} = j_{p\downarrow\downarrow} = 2M_p (1,0,0,0) \quad j_{p\uparrow\downarrow} = j_{p\downarrow\uparrow} = 0$$

$$\mathcal{M}_{fi} = \frac{e^2}{q^2} j_{e} \cdot j_{p} \qquad \left\langle \left| M_{fi} \right|^2 \right\rangle = \frac{1}{4} \frac{e^4}{q^4} \left( 16E^2 M_p^2 \right) \left( 4\cos^2\frac{\theta}{2} \right) \qquad q^2 = -4E^2 \sin^2\frac{\theta}{2}$$

$$\frac{\mathrm{d}\sigma}{\mathrm{d}\Omega} = \frac{1}{64\pi^2 M_p^2} |M_{fi}|^2 \qquad \left(\frac{d\sigma}{d\Omega}\right)_{Mott} = \frac{\alpha^2}{4E^2 \sin^4 \frac{\theta}{2}} \cos^2 \frac{\theta}{2}$$

4. 在一个 $e^-p$ 散射的实验中,入射电子的能量 $E_1 = 529.5 MeV$ ,出射电子在相对于入射 电子角度等于75°的地方被探测,试求出射电子的能量 $E_3$ 和四动量转移的平方 $Q^2$ 

$$E_3 = \frac{ME_1}{M + E_1(1 - \cos\theta)}$$

$$E_3 = \frac{938 \times 529}{938 + 529(1 - \cos 75^\circ)} = 373 \text{ MeV}$$

$$q^{2} = -\frac{2ME_{1}^{2}(1 - \cos\theta)}{M + E_{1}(1 - \cos\theta)}$$

$$q^{2} = -\frac{2ME_{1}^{2}(1-\cos\theta)}{M+E_{1}(1-\cos\theta)} \qquad |q^{2}| = \frac{2\times938\times529^{2}(1-\cos75^{\circ})}{938+529(1-\cos75^{\circ})} = 29400 \text{MeV}^{2}$$

- 5. 在将弹性碰撞过渡到非弹性碰撞时,介绍了四个新的运动学变量:  $x, y, v, Q^2$ 
  - a) 写出四个变量的定义式并给出文字描述
  - b) 将x, y表示为v的形式,将 $Q^2$ 表示为x, y的形式

$$x \equiv \frac{Q^2}{2p_2.q}$$
  $Q^2 \equiv -q^2$   $y \equiv \frac{p_2.q}{p_2.p_1}$   $v \equiv \frac{p_2.q}{M}$ 

$$x = \frac{Q^2}{2Mv}$$
  $y = \frac{2M}{s - M^2}v$   $Q^2 = (s - M^2)xy$ 

### $6.在一个固定靶的e^-p弹性碰撞中,已知:$

$$y = \frac{M(E_1 - E_3)}{ME_1} = 1 - \frac{E_3}{E_1}$$

$$Q^{-} = 2M(E_1 - E_3) = 2ME_1y,$$

$$Q^2 = 2M(E_1 - E_3) = 2ME_1y$$
,  $Q^2 = 4E_1E_3\sin^2\frac{\theta}{2}$ 

a) 证明:

$$\sin^2 \frac{\theta}{2} = \frac{E_1}{E_3} \frac{M^2}{Q^2} y^2, \qquad \frac{E_3}{E_1} \cos^2 \frac{\theta}{2} = 1 - y - \frac{M^2 y^2}{Q^2}$$

i) 
$$Q^4 = 4M^2E_1^2y^2$$
,  $Q^2 = 4E_1E_3\sin^2\frac{\theta}{2}$ 

两式相除: 
$$Q^2 = \frac{4M^2E_1y^2}{4E_3\sin^2\frac{\theta}{2}}$$
, 即  $\sin^2\frac{\theta}{2} = \frac{E_1}{E_3}\frac{M^2}{Q^2}y^2$ 

ii) 
$$\frac{E_3}{E_1} \cos^2 \frac{\theta}{2} = \frac{E_3}{E_1} \left( 1 - \sin^2 \frac{\theta}{2} \right) = \frac{E_3}{E_1} - \frac{M^2}{Q^2} y^2 = 1 - y - \frac{M^2}{Q^2} y^2$$

5.在一个固定靶的 $e^-p$ 弹性碰撞中,已知:

$$Q^2 = 2M(E_1 - E_3) = 2ME_1y$$
,  $Q^2 = 4E_1E_3\sin^2\frac{\theta}{2}$ 

b) 假设角度对称,并利用
$$Q^2 = \frac{2ME_1^2(1-\cos\theta)}{M+E_1(1-\cos\theta)}$$
,证明:
$$\frac{d\sigma}{dQ^2} = \left| \frac{d\Omega}{dQ^2} \right| \frac{d\sigma}{d\Omega} = \frac{\pi}{E_3^2} \frac{d\sigma}{d\Omega}$$

$$d\Omega = d(\cos\theta)d\phi = 2\pi d(\cos\theta)$$

$$dQ^{2} = \dots = -2E_{3}^{2}d(\cos\theta)$$

$$\left|\frac{d\Omega}{dQ^{2}}\right| = \frac{\pi}{E_{3}^{2}}$$

$$\frac{d\sigma}{dQ^{2}} = \left|\frac{d\Omega}{dQ^{2}}\right|\frac{d\sigma}{d\Omega} = \frac{\pi}{E_{2}^{2}}\frac{d\sigma}{d\Omega}$$

#### 5.在一个固定靶的 $e^-p$ 弹性碰撞中,已知:

$$Q^2 = 2M(E_1 - E_3) = 2ME_1y$$
,  $Q^2 = 4E_1E_3\sin^2\frac{\theta}{2} \sin^2\frac{\theta}{2} = \frac{E_1}{E_2}\frac{M^2}{O^2}y^2$ ,

c)利用上面两小问的结果,将Rosenbluth公式

两小问的结果,将Rosenbluth公式 
$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2}{4E_1^2 \sin^4 \theta/2} \frac{E_3}{E_1} \left( \frac{G_E^2 + \tau G_M^2}{(1+\tau)} \cos^2 \frac{\theta}{2} + 2\tau G_M^2 \sin^2 \frac{\theta}{2} \right) \frac{E_3}{E_1} \cos^2 \frac{\theta}{2} = 1 - y - \frac{M^2 y^2}{Q^2}$$
 伦兹不变的形式:

改写为洛伦兹不变的形式:

$$\frac{d\sigma}{dQ^2} = \frac{4\pi\alpha^2}{Q^4} \left[ \frac{G_E^2 + \tau G_M^2}{(1+\tau)} \left( 1 - y - \frac{M^2 y^2}{Q^2} \right) + \frac{1}{2} y^2 G_M^2 \right]$$

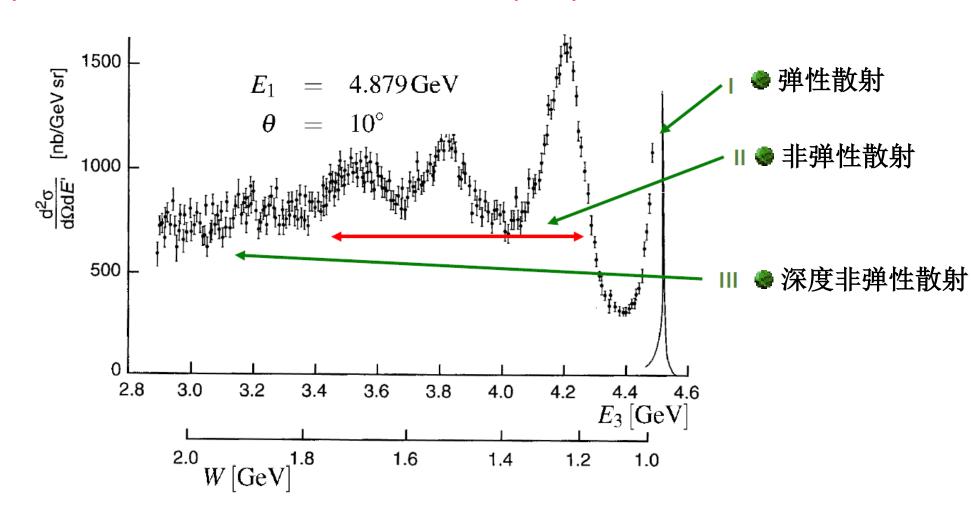
$$\frac{d\sigma}{dO^2} = \frac{\pi}{E_2^2} \frac{\alpha^2}{4E_1^2 \sin^4 \theta / 2} \frac{E_3}{E_1} \left( \frac{G_E^2 + \tau G_M^2}{(1+\tau)} \cos^2 \frac{\theta}{2} + 2\tau G_M^2 \sin^2 \frac{\theta}{2} \right)$$

$$=\frac{4\pi\alpha^2}{Q^4}\frac{E_3}{E_1}\left[\frac{G_E^2+\tau G_M^2}{(1+\tau)}\frac{E_1}{E_3}\left(1-y-\frac{M^2y^2}{Q^2}\right)+2\frac{Q^2}{4M^2}G_M^2\frac{E_1}{E_3}\frac{M^2}{Q^2}y^2\right]$$

$$= \frac{4\pi\alpha^2}{Q^4} \left[ \frac{G_E^2 + \tau G_M^2}{(1+\tau)} \left( 1 - y - \frac{M^2 y^2}{Q^2} \right) + \frac{1}{2} y^2 G_M^2 \right]$$

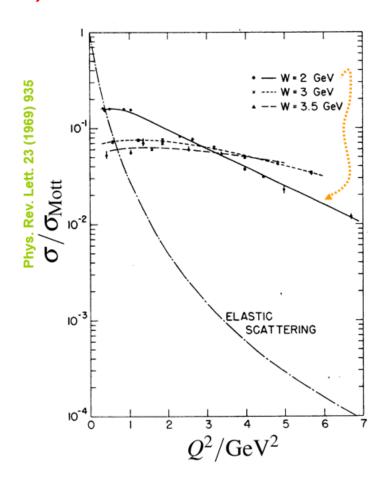
#### 1.读图答题:

a) 辨别弹性散射、非弹性散射、深度非弹性散射(I,II,III):



#### 1.读图答题:

b) 描述弹性散射、深度非弹性散射中散射截面与q²的关系

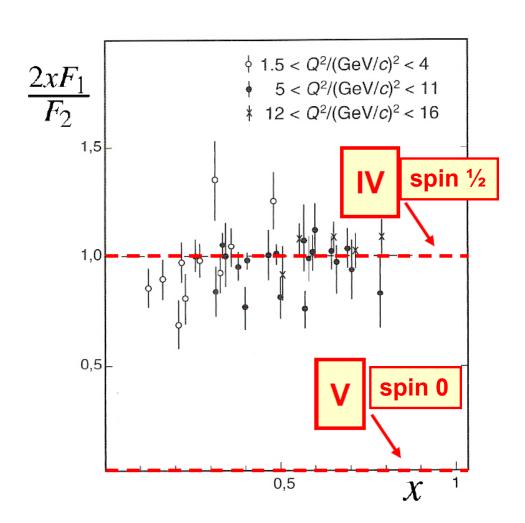


- 弹性散射截面随q<sup>2</sup>迅速下降, 由于质子不是点粒子(即 形状因子)
- · 非弹性散射截面随q2缓慢下降,
- 深度非弹性散射截面几乎与随q²无关即"形状因子(傅里叶变换)"→1
  - → 月

质子内部点状对象的散射!

#### 1.读图答题:

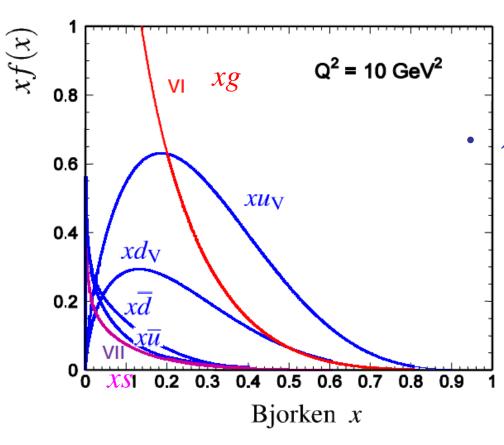
c) 写出Callan-Gross公式并辨认spin-1/2, spin-0 (IV,V)



$$F_2(x) = 2xF_1(x)$$

#### 1.读图答题:

d)辨认胶子和奇异夸克曲线(VI,VII),并解释为什么在x比较大的时候 $u_V(x) \approx 2d_V(x)$ 



• 分解价夸克和海夸克的贡献:

$$u(x) = u_{V}(x) + u_{S}(x)$$
  $d(x) = d_{V}(x) + d_{S}(x)$   
 $\overline{u}(x) = \overline{u}_{S}(x)$   $\overline{d}(x) = \overline{d}_{S}(x)$ 

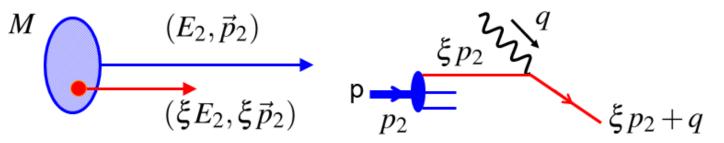
质子包含两个上型价夸克和一个下型价夸克

$$\int_0^1 u_{\rm V}(x) dx = 2 \qquad \int_0^1 d_{\rm V}(x) dx = 1$$

• 高x区域海贡献变小

2. 在夸克模型中,证明质子极高能近似下,比约肯变量x可被看作被轰击夸克携带的质子总动量的比分( $\xi = x$ )

设夸克携带的动量占质子总动量的比分为ξ



• 相互作用后被轰击后的夸克四动量为

$$(\xi p_2 + q)^2 = m_q^2 \approx 0 \implies \xi^2 p_2^2 + q^2 + 2\xi p_2 \cdot q = 0 \implies \xi = \frac{Q^2}{2p_2 \cdot q} = x$$

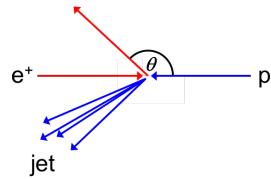
$$(\xi^2 p_2^2 = m_q^2 \approx 0)$$

3. 在HERA对撞机上,能量 $E_1=27.5$  GeV的电子与能量 $E_2=820$  GeV的质子对撞,出射的电子能量为 $E_3=31$  GeV,入射、出射电子流夹角为 $\theta\approx45^\circ$ ,考虑深度非弹性碰撞,证明:

$$x = \frac{E_3}{E_2} \left[ \frac{1 - \cos \theta}{2 - (E_3/E_1)(1 + \cos \theta)} \right]$$

#### 并计算x, $Q^2$ 的值。

$$\begin{split} E_1 \gg m_e, E_2 \gg m_p \\ p_1 &= (E_1, 0, 0, E_1) \\ p_2 &= (E_2, 0, 0, -E_2) \\ p_3 &= (E_3, E_3 \sin \theta, 0, E_3 \cos \theta) \\ q &= p_1 - p_3 = (E_1 - E_3, -E_3 \sin \theta, 0, E_1 - E_3 \cos \theta) \\ x &= \frac{Q^2}{2p_2 \cdot q} = \frac{-(E_1 - E_3)^2 + (E_3^2 \sin^2 \theta + (E_1 - E_3 \cos \theta)^2)}{2[E_2(E_1 - E_3) + E_2(E_1 - E_3 \cos \theta)]} \\ &= \frac{E_1 E_3 (1 - \cos \theta)}{2E_1 E_2 - E_2 E_3 \cos \theta} = \frac{E_3}{E_2} \left[ \frac{1 - \cos \theta}{2 - (E_3/E_1)(1 + \cos \theta)} \right] \\ x &= 0.146, Q^2 = 499.4 GeV^2 \end{split}$$

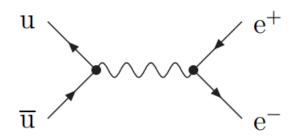


#### 4. 组分夸克模型:

a)  $\rho^0 \rightarrow e^+e^-$ 的衰变宽度为7keV,利用费曼图估算 $\omega^0 \rightarrow e^+e^-$ 的衰变宽度

b)  $\rho^0 \to \pi^0 \gamma$ 的衰变宽度为77keV,利用费曼图估算 $\omega^0 \to \pi^0 \gamma$ 的衰变宽度

 $(\pi^0, \rho^0$ 组分为 $(u\bar{u} - d\bar{d})/\sqrt{2}$ ,  $\omega^0$ 为 $(u\bar{u} + d\bar{d})/\sqrt{2}$ )

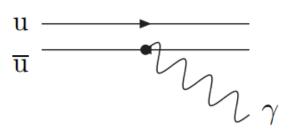


a) 
$$M(\rho^0 \to e^+ e^-) \sim \langle u\bar{u} - d\bar{d} | \gamma \rangle \sim \frac{2}{3} - \left( -\frac{1}{3} \right) = 1$$

$$M(\omega^0 \to e^+ e^-) \sim \langle u\bar{u} + d\bar{d} | \gamma \rangle \sim \frac{2}{3} + \left( -\frac{1}{3} \right) = \frac{1}{3}$$

$$\frac{\Gamma(\rho^0 \to e^+ e^-)}{\Gamma(\omega^0 \to e^+ e^-)} = \left( \frac{1}{1/3} \right)^2 = 9$$

$$\Gamma(\omega^0 \to e^+ e^-) \approx \frac{7}{9} keV$$



b) 
$$M(\rho^0 \to \pi^0 \gamma) \sim \langle u\bar{u} - d\bar{d} | u\bar{u} - d\bar{d} \rangle \sim \frac{2}{3} + \left(-\frac{1}{3}\right) = \frac{1}{3}$$

$$M(\omega^0 \to \pi^0 \gamma) \sim \langle u\bar{u} + d\bar{d} | u\bar{u} - d\bar{d} \rangle \sim \frac{2}{3} - \left(-\frac{1}{3}\right) = 1$$

$$\frac{\Gamma(\rho^0 \to e^+ e^-)}{\Gamma(\omega^0 \to e^+ e^-)} = \left(\frac{1/3}{1}\right)^2 = \frac{1}{9}$$

$$\Gamma(\omega^0 \to e^+ e^-) \approx 77 \times 9keV = 693keV$$

1. 仿照课上构造 $|p\uparrow\rangle$ 的方法,组合 $\frac{1}{\sqrt{2}}\phi(M_S)\chi(M_S)+\frac{1}{\sqrt{2}}\phi(M_A)\chi(M_A)$ ,写出自旋向下的中子波函数 $|n\downarrow\rangle$ 

$$|n\downarrow\rangle = \frac{1}{6\sqrt{2}}(2ddu - udd - dud)(2\downarrow\downarrow\uparrow-\uparrow\downarrow\downarrow-\downarrow\uparrow\downarrow) + \frac{1}{2\sqrt{2}}(udd - dud)(\uparrow\downarrow\downarrow-\downarrow\uparrow\downarrow)$$

$$= \frac{1}{\sqrt{18}} \begin{pmatrix} 2d\downarrow d\downarrow u\uparrow -d\uparrow d\downarrow u\downarrow -d\downarrow d\uparrow u\downarrow + \\ 2d\downarrow u\uparrow d\downarrow -d\downarrow u\downarrow d\uparrow -d\uparrow u\downarrow d\downarrow + \\ 2u\uparrow d\downarrow d\downarrow -u\downarrow d\uparrow d\downarrow -u\downarrow d\downarrow d\uparrow \end{pmatrix}$$

# 2.考虑只由u,d夸克及其反夸克组成的介子,尝试从 $|1,+1\rangle = -u\bar{d}$

开始, a)用下降算符得到另外两个三重态, b)用正交性给出最后一个单态, c)最后证明单态是阶梯算符的"尽头"。

a) 
$$T_{-}|1,+1\rangle = \sqrt{2}|1,0\rangle = T_{-}(-u\bar{d}) = -d\bar{d} + u\bar{u} \rightarrow |1,0\rangle = \frac{1}{\sqrt{2}}(u\bar{u} - d\bar{d})$$

$$T_{-}|1,0\rangle = \sqrt{2}|1,-1\rangle = \frac{1}{\sqrt{2}}T_{-}(u\bar{u} - d\bar{d}) = \sqrt{2}d\bar{u} \rightarrow |1,-1\rangle = d\bar{u}$$
b)  $|0,0\rangle = \alpha(u\bar{u}) + \beta(d\bar{d})$ 
 $\langle 0,0|0,0\rangle = 1$ 
 $\langle 1,0|0,0\rangle = 0$ 
 $|0,0\rangle = \frac{1}{\sqrt{2}}(u\bar{u} + d\bar{d})$ 
c)  $T_{-}|0,0\rangle = \frac{1}{\sqrt{2}}(d\bar{u} - d\bar{u}) = 0, T_{+}|0,0\rangle = \frac{1}{\sqrt{2}}(-u\bar{d} + u\bar{d}) = 0$ 

# 3. 写出八个盖尔曼矩阵及味道SU(3)群中用到的三组阶梯算符,之后在以下关系中任选一组进行证明(六个一组):

$$\hat{T}_{+}u = 0, \hat{T}_{-}u = d$$
  $\hat{V}_{+}u = 0, \hat{V}_{-}u = s$   $\hat{T}_{+}d = u, \hat{T}_{-}d = 0$   $\hat{V}_{+}d = 0, \hat{V}_{-}d = 0$   $\hat{T}_{+}s = 0, \hat{T}_{-}s = 0$   $\hat{V}_{+}s = u, \hat{V}_{-}s = 0$ 

$$\lambda_{1} = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \lambda_{2} = \begin{pmatrix} 0 & -i & 0 \\ i & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \lambda_{3} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$\lambda_{4} = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{pmatrix} \lambda_{5} = \begin{pmatrix} 0 & 0 & -i \\ 0 & 0 & 0 \\ i & 0 & 0 \end{pmatrix} \lambda_{5} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ i & 0 & 0 \end{pmatrix}$$

$$\lambda_{6} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} \lambda_{7} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -i \\ 0 & i & 0 \end{pmatrix}$$

$$\widehat{U}_{+}u = 0, \widehat{U}_{-}u = 0$$

$$\widehat{U}_{+}d = 0, \widehat{U}_{-}d = s$$

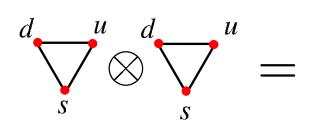
$$\widehat{U}_{+}s = d, \widehat{U}_{-}s = 0$$

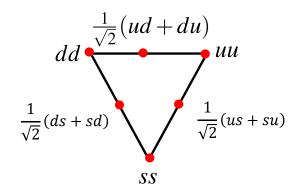
$$egin{align} T_\pm &= rac{1}{2}(\lambda_1 \pm i\lambda_2) \ V_\pm &= rac{1}{2}(\lambda_4 \pm i\lambda_5) \ U_\pm &= rac{1}{2}(\lambda_6 \pm i\lambda_7) \ \end{align}$$

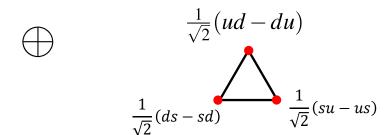
4. 模仿上课时的方法, 用画图的方式画出10重态和8重态的构造过程:

a)  $3 \otimes 3 = 6 \oplus \overline{3}$ 

计算并标出后面两个三角上各个点所对应的态



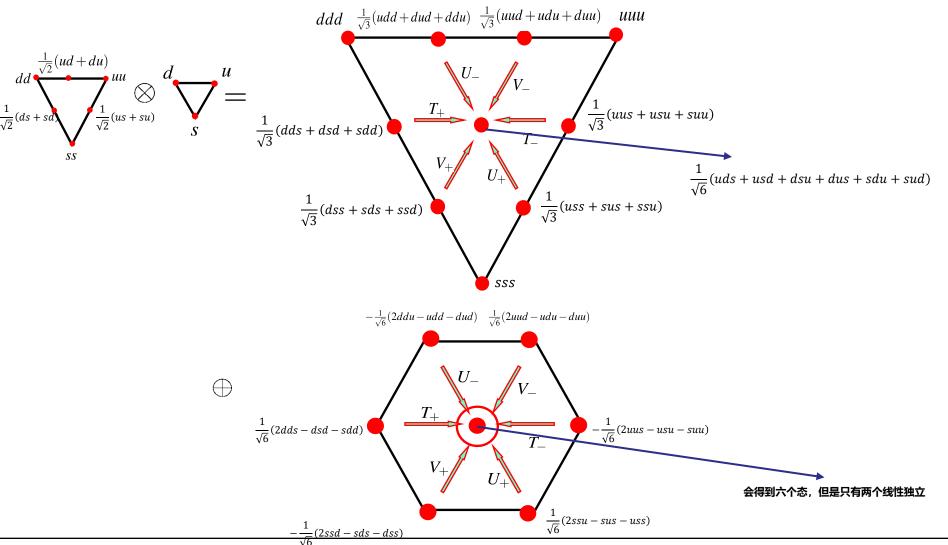




#### 4. 模仿上课时的方法, 用画图的方式画出10重态和8重态的构造过程:

b)  $3 \otimes 6 = 10 \oplus 8$ 

计算并标出最后面两个图形上各个点所对应的态

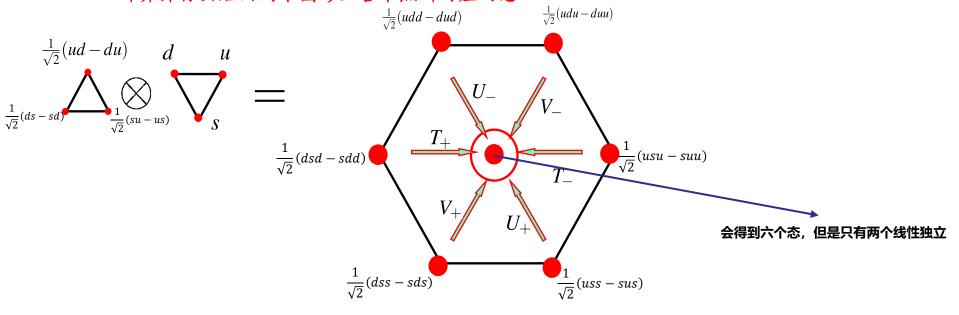




#### 4. 模仿上课时的方法, 用画图的方式画出10重态和8重态的构造过程:

c)  $\overline{3} \otimes 3 = 8 \oplus 1$ 

计算并标出后面两个图形上各个点所对应的态



$$\frac{1}{\sqrt{6}}(uds - usd + dsu - dus + sud - sdu)$$