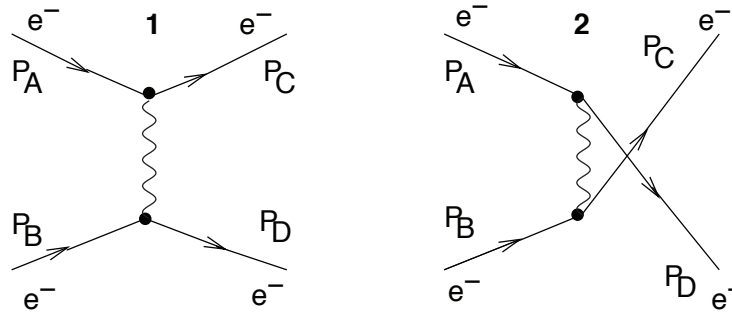


Quantum Field Theory

Exercise 40:

We consider Møller scattering which is the QED process $e^- e^- \rightarrow e^- e^-$. In lowest order of perturbative QED this scattering process is described by following Feynman diagrams with the 4-momenta p_i (i=A,B,C,D):



- (a) Write down an expression for the invariant Feynman amplitude $\mathcal{M} = \mathcal{M}_1 + \mathcal{M}_2$ following the Feynman rules which includes the contributions $\mathcal{M}_{i=1,2}$ of both diagrams. (4P)
- (b) Show that the spin-summed modulus squared of the Feynman amplitude in the high-energy limit (the mass m_e of the electron can be neglected) is given by

$$\overline{|\mathcal{M}|^2} = \frac{1}{4} \sum_{\text{spin projections}} |\mathcal{M}_1 + \mathcal{M}_2|^2 = 2e^4 \left\{ \frac{s^2 + u^2}{t^2} + \frac{2s^2}{tu} + \frac{s^2 + t^2}{u^2} \right\}.$$

The quantities s, u and t are the invariant Mandelstam variables which are defined with respect to the p_i (i=A,B,C,D) as

$$\begin{aligned} s &\equiv (p_A + p_B)^2 \approx 2p_A \cdot p_B \approx 2p_C \cdot p_D \\ t &\equiv (p_A - p_C)^2 \approx -2p_A \cdot p_C \approx -2p_B \cdot p_D \\ u &\equiv (p_A - p_D)^2 \approx -2p_A \cdot p_D \approx -2p_C \cdot p_B \end{aligned}$$

and where the last identities correspond to the high-energy limit. (8P)

Hint: to evaluate the spin-summed modulus squared of \mathcal{M}_1 and \mathcal{M}_2 you can use results already derived in the context of $e^+ e^- \rightarrow \mu^+ \mu^-$.

- (c) Show that the differential cross-section in the center-of-momentum (CM) system for electron-electron scattering in the high-energy limit ($E \gg m_e$) is given by (8P)

$$\left(\frac{d\sigma}{d\Omega}\right)_{CM} = \frac{\alpha^2}{8E^2} \left[\frac{1 + \cos^4(\theta/2)}{\sin^4(\theta/2)} + \frac{2}{\sin^2(\theta/2) \cos^2(\theta/2)} + \frac{1 + \sin^4(\theta/2)}{\cos^4(\theta/2)} \right]$$

where θ is the scattering angle and E is the energy of either electron in the CM frame. The result for the differential cross section is the high-energy limit of the Møller formula already derived in C. Møller, *Ann. Phys.*, **14**, 531 (1932).

Exercise 41 (bonus points):

In the lecture we introduced the polarization vector $\varepsilon_r^\mu(k)$ in the plane wave solutions of the photon. Those polarization vectors are gauge dependent. For example, for a free photon, described in a Lorentz gauge by a plane wave $A^\mu(x) = \text{const } \varepsilon_r^\mu(k) e^{\pm i k \cdot x}$, the gauge transformation $A^\mu(x) \rightarrow A^\mu(x) + \partial^\mu f(x)$, with $f(x) = \tilde{f}(k) e^{\pm i k \cdot x}$, implies

$$\varepsilon_r^\mu(k) e^{\pm i k \cdot x} \rightarrow [\varepsilon_r^\mu(k) \pm i \tilde{f}(k) k^\mu] e^{\pm i k \cdot x}.$$

Consider the invariant amplitude \mathcal{M} for the Compton scattering process of a photon and an electron, $\gamma(\mathbf{k}) + e^-(\mathbf{p}) \rightarrow \gamma(\mathbf{k}') + e^-(\mathbf{p}')$.

- (a) Draw all of the lowest order QED Feynman diagrams for this process and give explicit expressions for them using Feynman rules. (2P)
- (b) Show that the sum of the lowest order Compton scattering Feynman diagrams is gauge invariant, although the individual contributions of the diagrams are not, by considering the gauge transformations (8P)

$$\varepsilon_r^\mu(k) \rightarrow \varepsilon_r^\mu(k) + \lambda k^\mu \quad ; \quad \varepsilon_s'^\mu(k') \rightarrow \varepsilon_s'^\mu(k') + \lambda' k'^\mu.$$

Worked-out solutions to the homework problems should be handed in latest at the beginning of the lecture of January 31, 2023.