



## Tutorial Advanced Quantum Mechanics

Winter semester 2021/2022

Tübingen, 17th January 2022

### Problem Set 12

#### Problem 37 (Separable potential)

(8 points)

The Lippmann–Schwinger equation for a scattering state with wavenumber  $\mathbf{k}$  and energy  $E_{\mathbf{k}} = \hbar^2 \mathbf{k}^2 / (2m)$  reads

$$|\psi_{\mathbf{k}}^+\rangle = |\mathbf{k}\rangle + \frac{1}{E_{\mathbf{k}} - H_0 + i\varepsilon} V |\psi_{\mathbf{k}}^+\rangle.$$

- a. (2 points) Show that the momentum representation  $\langle \mathbf{p} | V | \mathbf{q} \rangle$  of a *local* potential  $V$  depends only on the difference of the two momenta.
- b. (4 points) Consider a non-local but *separable* potential, i.e.

$$\langle \mathbf{p} | V | \mathbf{q} \rangle = v(\mathbf{p})v(\mathbf{q}).$$

Start from the Lippmann–Schwinger equation and derive an equation for the quantity

$$A(\mathbf{k}) \equiv \int \frac{d^3 q}{(2\pi)^3} v(\mathbf{q}) \langle \mathbf{q} | \psi_{\mathbf{k}}^+ \rangle.$$

Find the exact solution for the Lippmann–Schwinger equation.

$$\text{Solution: } \langle \mathbf{p} | \psi_{\mathbf{k}}^+ \rangle = (2\pi)^3 \delta(\mathbf{p} - \mathbf{k}) + \frac{v(\mathbf{p})v(\mathbf{k})}{E_{\mathbf{k}} - E_{\mathbf{p}} + i\varepsilon} \left[ 1 - \int \frac{d^3 q}{(2\pi)^3} \frac{v^2(\mathbf{q})}{E_{\mathbf{k}} - E_{\mathbf{q}} + i\varepsilon} \right]^{-1}.$$

- c. (2 points) Give an exact expression for the scattering amplitude

$$f_{\mathbf{k}}(\theta, \phi) = -\frac{m}{2\pi\hbar^2} \langle \mathbf{k}' = k\hat{\mathbf{x}} | V | \psi_{\mathbf{k}}^+ \rangle$$

for a separable potential.

#### Problem 38 (Born approximation and Fermi's golden rule)

(6 points)

A spinless particle with mass  $m$  is elastically scattered with energy  $E$  by a short-ranged potential  $V(\mathbf{x})$ . In order to avoid non-normalizable states we work in a finite box with volume  $L^3$ . The scalar product of states reads thus

$$\langle \psi_1 | \psi_2 \rangle = \int_{L^3} d^3 x \psi_1^*(\mathbf{x}) \psi_2(\mathbf{x}).$$

As initial and final state we consider the normalized plane waves

$$\begin{aligned} t \rightarrow -\infty : \quad \psi_{\text{in}}(\mathbf{x}) &= \langle \mathbf{x} | \mathbf{k} \rangle = L^{-3/2} e^{i\mathbf{k} \cdot \mathbf{x}}, \\ t \rightarrow +\infty : \quad \psi_{\text{out}}(\mathbf{x}) &= \langle \mathbf{x} | \mathbf{k}' \rangle = L^{-3/2} e^{i\mathbf{k}' \cdot \mathbf{x}}, \end{aligned}$$

where  $|\mathbf{k}| = |\mathbf{k}'| = \sqrt{2mE}/\hbar$ .

- a. (2 points)** Show that the density of states (number of states per energy and solid angle) of a free particle ( $V(\mathbf{x}) = 0$ ) with momentum in a solid angle  $d\Omega$  is given by

$$\rho(E) = \left( \frac{L}{2\pi\hbar} \right)^3 \sqrt{2m^3 E}.$$

- b. (2 points)** Use Fermi's golden rule

$$\Gamma_{i \rightarrow f} = \frac{2\pi}{\hbar} |\langle f | V | i \rangle|^2 \rho(E_f)$$

to calculate the transition rate  $d\Gamma$  into states with energy  $E$  scattered into a solid angle  $d\Omega$ .

- c. (2 points)** The differential cross section is the ratio of the transition rate  $d\Gamma$  and the incoming particle current

$$d\sigma = |\mathbf{j}_{\text{in}}|^{-1} d\Gamma \quad \text{with } \mathbf{j} = \frac{\hbar}{m} \Im(\psi^* \nabla \psi).$$

Calculate  $\frac{d\sigma}{d\Omega}$ . Is the limit  $L \rightarrow \infty$  well defined?

### Problem 39 (Properties of the cross section in Born approximation)

(5 points)

Consider the differential cross section for elastic scattering in Born approximation

$$\frac{d\sigma}{d\Omega} = \left( \frac{m}{2\pi\hbar^2} \right)^2 \left| \int d^3x e^{-i\mathbf{q} \cdot \mathbf{x}} V(\mathbf{x}) \right|^2,$$

where  $\mathbf{q} \equiv \mathbf{k}' - \mathbf{k}$  is as usual the momentum transfer. Prove the following features of the cross section:

- a. (2 points)** If the scattering potential  $V(\mathbf{x})$  never changes its sign, then the differential cross section  $d\sigma/d\Omega$  takes the maximum value for  $E \rightarrow 0$  (if finite).
- b. (3 points)** For a *central* potential  $V(\mathbf{x}) = V(r)$  the product  $E \cdot \sigma(E)$ , (where  $\sigma(E) = \int d\Omega \frac{d\sigma}{d\Omega}$  is the *total* cross section) is a monotonic increasing function of the energy  $E$ .

*Hint:* Work out the angular integrations and show that  $E \cdot \sigma(E)$  can be cast into the form  $E \cdot \sigma(E) = \int_0^{aE} dx |f(x)|^2$  with a positive constant  $a > 0$  and an appropriate function  $f(x)$ .