

# Circuits and Signals

## Periodic current circuits

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Marek Rupniewski

2022 spring semester



**Faculty of Electronics  
and Information  
Technology**

WARSAW UNIVERSITY OF TECHNOLOGY

# Periodic signal

## Definition

**Periodic signal** is a function  $x: \mathbb{R} \rightarrow \mathbb{R}$ , for which there exists a constant  $T > 0$ , called **the period of the signal**  $x(t)$ , such that

$$x(t) = x(t + T).$$

For convenience, we restrict periodic signals to signals that are piecewise smooth and such that their derivatives have one-side limits everywhere.

**Fundamental period  $T$** : is the smallest non-zero period,

**Fundamental frequency  $f$** :  $f = \frac{1}{T}$  [Hz],

**Fundamental pulsation  $\omega$** :  $\omega = 2\pi f = \frac{2\pi}{T}$  [rad/s].

# Steady state

## Definition

**Periodic current circuit** is a circuit admitting solutions that consist of periodic signals (of common period). Every such solution is called a **steady state**.

# Superposition rule

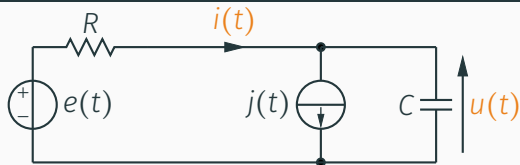
## Superposition rule for per. current circuits

Each  $T$ -periodic signal  $x$  (a voltage or a current) that is a part of a steady state in a periodic current circuit, can be expressed as the sum of a constant signal and signals of pulsations that are (integer) multiples of the fundamental pulsation  $\omega = \frac{2\pi}{T}$ :

$$x(t) = X_0 + x_1(t) + x_2(t) + \dots$$

The constant term  $X_0$  can be obtained by the means of DC analysis. Each  $k$ -th harmonics  $x_k$  (of pulsation  $k\omega$ ) can be computed by the means of AC analysis of the circuit that results from the original if all other harmonics (i.e. all but the  $k$ -th) are reduced to zero.

## Superposition rule — example



$$e(t) = E_0 + E_{2m} \cos(2\omega t), \quad j(t) = J_{1m} \cos(\omega t) + J_{2m} \cos(2\omega t + \frac{\pi}{2}),$$

$$u(t) = ?, \quad i(t) = ?.$$

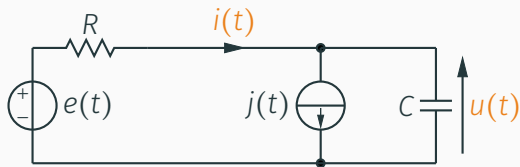
$$u(t) = U_0 + \underbrace{U_{1m} \cos(\omega t + \phi_1)}_{u_1(t)} + \underbrace{U_{2m} \cos(2\omega t + \phi_2)}_{u_2(t)}.$$

$$i(t) = I_0 + \underbrace{I_{1m} \cos(\omega t + \psi_1)}_{i_1(t)} + \underbrace{I_{2m} \cos(2\omega t + \psi_2)}_{i_2(t)}.$$

## Superposition rule — example — DC term

$$e(t) = E_0 + E_{2m} \cos(2\omega t),$$

$$j(t) = J_{1m} \cos(\omega t) + J_{2m} \cos(2\omega t + \frac{\pi}{2}),$$

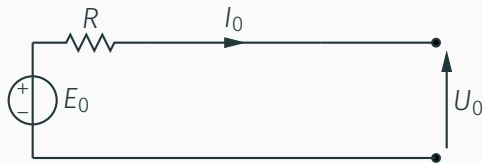


$$I_0 = 0, \quad U_0 = E_0.$$

## Superposition rule — example — DC term

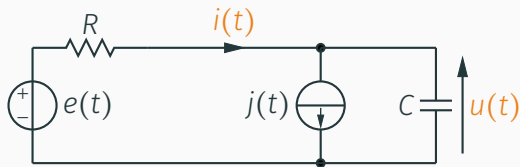
$$e(t) = E_0 + E_{2m} \cos(2\omega t),$$

$$j(t) = I_{1m} \cos(\omega t) + I_{2m} \cos(2\omega t + \frac{\pi}{2}),$$



$$I_0 = 0, \quad U_0 = E_0.$$

## Superposition rule — example — first harmonics

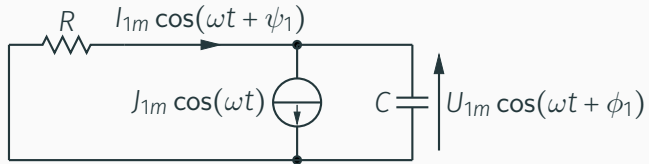


$$e(t) = E_0 + E_{2m} \cos(2\omega t),$$

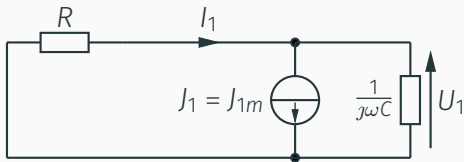
$$j(t) = J_{1m} \cos(\omega t) + J_{2m} \cos\left(2\omega t + \frac{\pi}{2}\right),$$



## Superposition rule — example — first harmonics



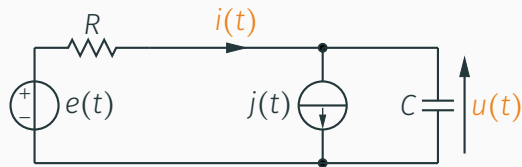
## Superposition rule — example — first harmonics



$$I_1 = J_1 \frac{1/j\omega C}{R + 1/j\omega C} = J_1 \frac{1}{1 + j\omega RC}, \quad U_1 = -J_1 \frac{R}{1 + j\omega RC}.$$

$$\begin{cases} I_{1m} = \frac{J_{1m}}{\sqrt{1 + (\omega RC)^2}}, & \psi_1 = -\arctan \omega RC, \\ U_{1m} = \frac{J_{1m} R}{\sqrt{1 + (\omega RC)^2}}, & \phi_1 = \pi - \arctan \omega RC, \end{cases}$$

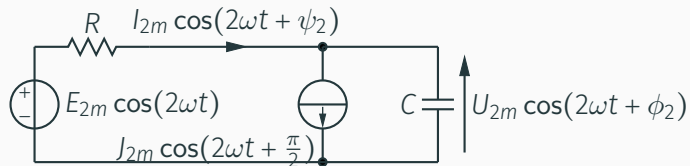
## Superposition rule — example — second harmonics



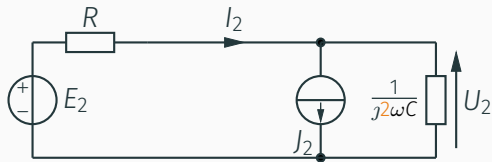
$$e(t) = E_0 + E_{2m} \cos(2\omega t),$$

$$j(t) = J_{1m} \cos(\omega t) + J_{2m} \cos\left(2\omega t + \frac{\pi}{2}\right),$$

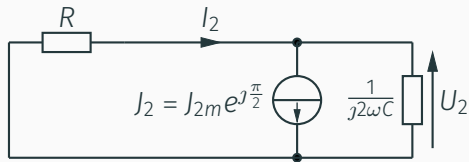
## Superposition rule — example — second harmonics



## Superposition rule — example — second harmonics



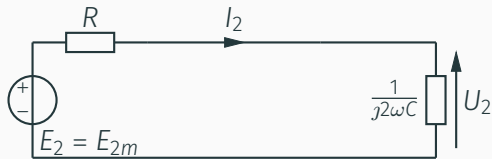
## Superposition rule — example — second harmonics



$$I_2 = J_2 \frac{1}{1 + j2\omega RC} +$$

$$U_2 = -J_2 \frac{R}{1 + j2\omega RC} +$$

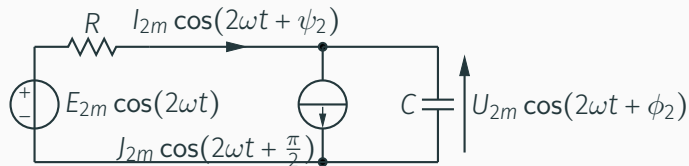
## Superposition rule — example — second harmonics



$$I_2 = J_2 \frac{1}{1 + j2\omega RC} + E_2 \frac{j2\omega C}{1 + j2\omega RC},$$

$$U_2 = -J_2 \frac{R}{1 + j2\omega RC} + E_2 \frac{1}{1 + j2\omega RC}.$$

## Superposition rule — example — second harmonics



$$I_{2m} = \frac{J_{2m} + 2\omega C E_{2m}}{\sqrt{1 + (2\omega RC)^2}},$$

$$\psi_2 = \frac{\pi}{2} - \arctan 2\omega RC,$$

$$U_{2m} = \sqrt{\frac{E_{2m}^2 + (J_{2m} R)^2}{1 + (2\omega RC)^2}},$$

$$\phi_2 = -\arctan \frac{R J_{2m}}{E_{2m}} - \arctan 2\omega RC.$$



What if a  $T$ -periodic signal  $x$  (electromotive force, current generated by a current source and so on) is not a finite sum of alternating signals (of pulsations that are multiples of  $\frac{2\pi}{T}$ )?

We can approximate the signal by such a finite sum!

How can we get such an approximation?

# Fourier series — definition

## Definition

Fourier series of a  $T$ -periodic signal  $x(t)$  is the series

$$\sum_{k=-\infty}^{+\infty} X_k e^{jk\omega t}, \quad \omega = \frac{2\pi}{T},$$

where

$$X_k = \frac{1}{T} \int_{t_0}^{t_0+T} x(t) e^{-jk\omega t} dt, \quad T = \frac{2\pi}{\omega}.$$

$X_0$  is the mean value of signal  $x$ !

$$X_{-k} = \overline{X_k}.$$

## Fourier series — point-wise convergence

Theorem on point-wise convergence of FS

The partial sums of the Fourier series of a  $T$ -periodic signal  $x$ :

$$\sum_{k=-N}^N x_k e^{jk\omega t},$$

converge point-wise to the mean of one-sided limits of the signal, i.e. for a fixed  $t$ , the above sums converge to

$$\frac{x(t^-) + x(t^+)}{2}.$$

That's how we should interpret the equality:

$$x(t) = \sum_{k=-\infty}^{+\infty} x_k e^{jk\omega t}.$$

## Fourier series — uniqueness

### Theorem on uniqueness of FS

Two  $T$ -periodic signals  $x_1$  and  $x_2$  have equal Fourier series if, and only if, they are equal up to their values at non-continuity points (i.e., equality  $x_1(t) = x_2(t)$  holds at every point at which both signals are continuous).

## Fourier series — expansion into harmonics

$$x(t) = X_0 + \sum_{k=1}^{\infty} (X_k e^{jk\omega t} + X_{-k} e^{-jk\omega t}) =$$

$$\overset{X_{-k} = \overline{X_k}, \quad \overset{z + \bar{z} = 2 \operatorname{Re} z}{=} X_0 + \sum_{k=1}^{\infty} 2 \operatorname{Re}(X_k e^{jk\omega t})$$

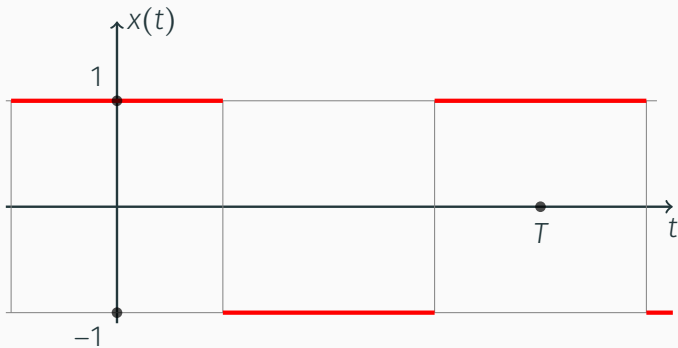
$$\overset{X_k = |X_k| e^{j\phi_k}}{=} X_0 + \sum_{k=1}^{\infty} 2 \operatorname{Re}(|X_k| e^{jk\omega t + j\phi_k}) = X_0 + \underbrace{\sum_{k=1}^{\infty} 2|X_k| \cos(k\omega t + \phi_k)}_{\text{trigonometric Fourier series}}.$$

$k$ -th harmonics of signal  $x$  ( $k \geq 1$ ) is the signal

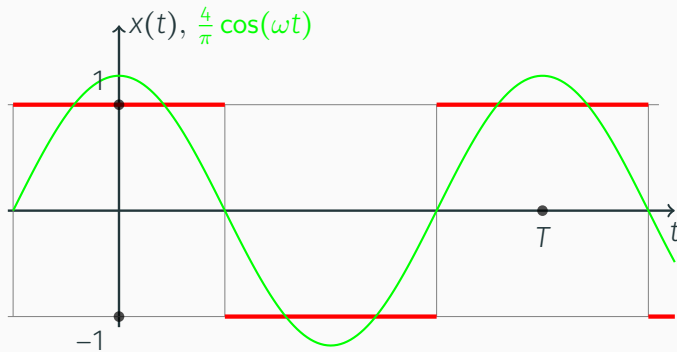
$$x_k(t) = X_{km} \cos(k\omega t + \phi_k), \quad X_{km} = 2|X_k|, \quad \phi_k = \arg X_k.$$

$$x(t) = X_0 + x_1(t) + x_2(t) + \dots$$

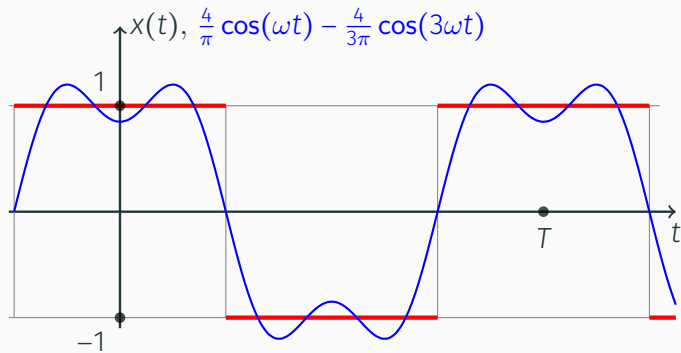
# Example



# Example

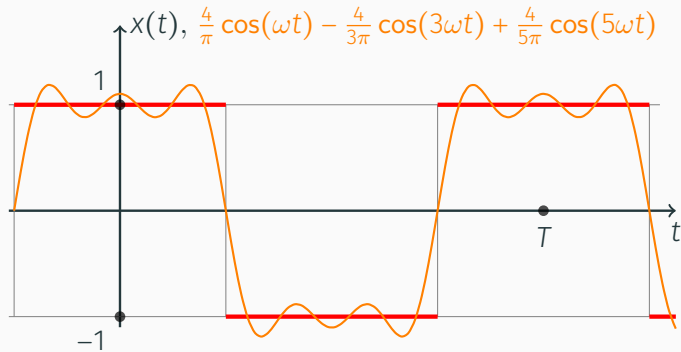


# Example

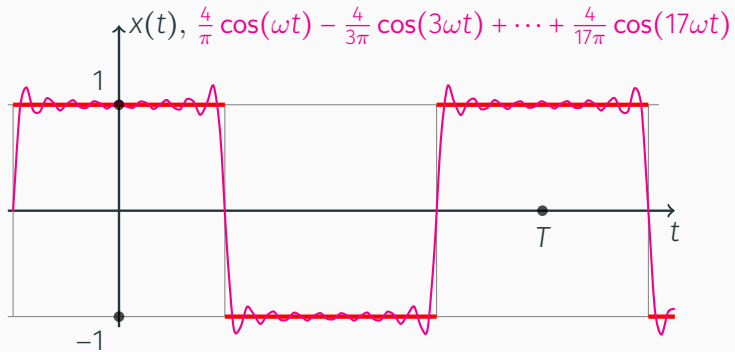




## Example



# Example



## Parseval's theorem

$$p(t) = i(t)u(t), P = \frac{1}{T} \int_{t_0}^{t_0+T} p(t)dt = \frac{1}{T} \int_{t_0}^{t_0+T} i(t)u(t)dt.$$

### Parseval's theorem

For  $T$ -periodic signals  $x, y$

$$\frac{1}{T} \int_{t_0}^{t_0+T} x(t)y(t) dt = \sum_{k=-\infty}^{+\infty} X_k \overline{Y_k},$$

where  $t_0$  is an arbitrary time instant. In particular:

$$\frac{1}{T} \int_{t_0}^{t_0+T} x^2(t) dt = \sum_{k=-\infty}^{+\infty} |X_k|^2, \quad X_{\text{RMS}} = \sqrt{X_0^2 + \frac{1}{2} \sum_{k=1}^{+\infty} X_{km}^2}.$$

## Mean power in per. current circuits

$$i(t) = \sum_{k=-\infty}^{+\infty} I_k e^{jk\omega t}, \quad u(t) = \sum_{k=-\infty}^{+\infty} U_k e^{jk\omega t}.$$

$$P = \sum_{k=-\infty}^{+\infty} U_k \overline{I_k}.$$

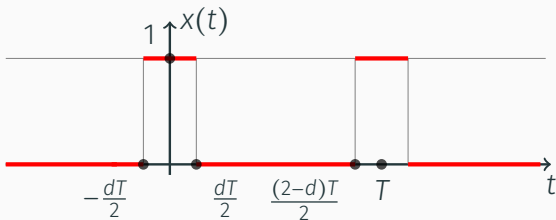
$$U_0 I_0 + 2 \sum_{k=1}^{\infty} \operatorname{Re}(U_k \overline{I_k}) = U_0 I_0 + \sum_{k=1}^{\infty} \frac{1}{2} U_{km} I_{km} \cos(\arg U_k - \arg I_k).$$

### Proposition

The mean power delivered to a device in a per. current circuit is a sum of the DC-power (power due to DC signals only) and the mean (real) powers resulting from each harmonics.

# How to compute Fourier series — pulse wave

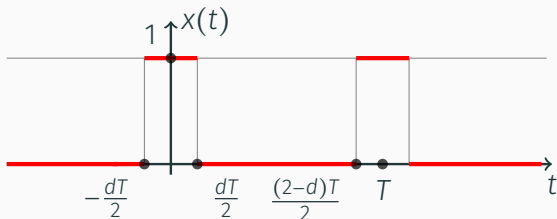
Pulse wave  $x(t)$  with duty cycle  $d$  and amplitude 1:



$$X_0 = \frac{1}{T} \int_{-\frac{dT}{2}}^{\frac{dT}{2}} 1 dt = d.$$

$$X_k = \frac{1}{T} \int_{-\frac{dT}{2}}^{\frac{dT}{2}} e^{-jk\omega t} dt = \frac{1}{T} \frac{e^{-jk\omega t}}{-jk\omega} \Big|_{-\frac{dT}{2}}^{\frac{dT}{2}} = \frac{1}{-j2\pi k} \left( e^{-jk\pi d} - e^{jk\pi d} \right) = \frac{\sin(k\pi d)}{k\pi}.$$

## How to compute Fourier series — pulse wave, cont.



Fourier series coefficients:

$$X_0 = d, \quad X_k = \frac{\sin(\pi kd)}{k\pi}, \quad k \neq 0.$$

Parameters of individual harmonics:

$$X_{km} = 2 \frac{|\sin(\pi kd)|}{k\pi}, \quad k \geq 1,$$

$$\phi_k = \begin{cases} 0 & \text{if } \sin(\pi kd) > 0 \\ \pi & \text{if } \sin(\pi kd) < 0 \end{cases}, \quad k \geq 1.$$

# Fourier series — linearity

## Linearity

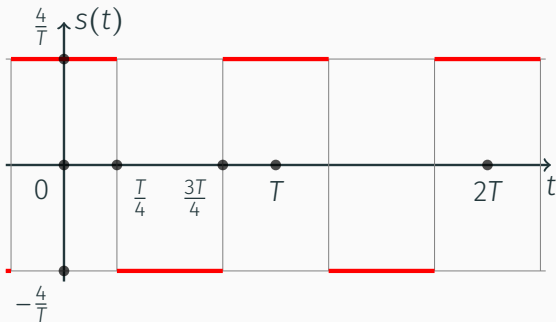
The Fourier series of a signal

$$\alpha x(t) + y(t), \quad \alpha \in \mathbb{R}$$

is the series

$$\sum_{k=-\infty}^{+\infty} (\alpha X_k + Y_k) e^{jk\omega t}.$$

## How to compute Fourier series — square wave



$$s(t) = \frac{8}{T} \left( x(t) - \frac{1}{2} \right), \quad d = \frac{1}{2}.$$

Fourier coefficients:

$$X_0 = d, \quad X_k = \frac{\sin(\pi k d)}{k\pi}, \quad d = \frac{1}{2}.$$

$$S_0 = 0, \quad S_k = \frac{8}{T} \frac{\sin(k\frac{\pi}{2})}{k\pi}, \quad k \neq 0.$$



## Fourier series — derivative of a signal

Theorem (of Fourier series of a derivative)

If  $x$  is a  $T$ -periodic signal such that

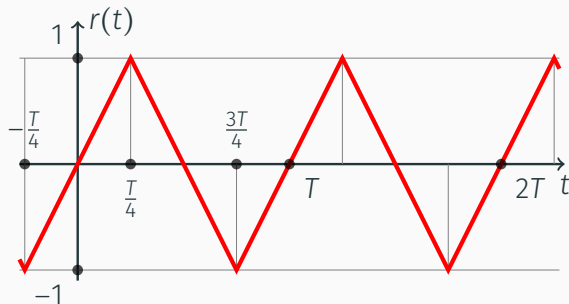
$$x(t) = \sum_{k=-\infty}^{+\infty} X_k e^{jk\omega t},$$

then

$$x'(t) = \sum_{k=-\infty}^{+\infty} (jk\omega) X_k e^{jk\omega t}.$$

In other words we may differentiate FS term-wise.

# How to compute Fourier series — sawtooth wave



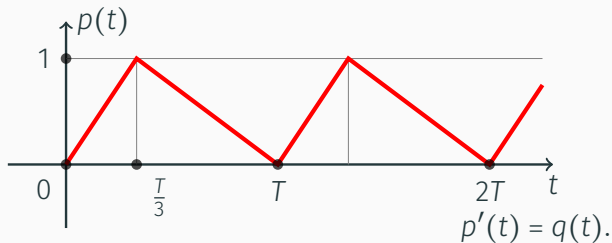
$$r'(t) = s(t).$$

Fourier series coefficients:

$$S_0 = 0, \quad S_k = \frac{8}{T} \frac{\sin(k\frac{\pi}{2})}{k\pi}, \quad k \neq 0.$$

$$R_0 = 0, \quad R_k = \frac{1}{jk\omega} \frac{8}{T} \frac{\sin(k\frac{\pi}{2})}{k\pi} = \frac{4}{jk\pi} \frac{\sin(k\frac{\pi}{2})}{k\pi}, \quad k \neq 0.$$

## How to compute Fourier series — another saw



$$q(t) = \frac{9}{2T}x\left(t - \frac{T}{6}\right) - \frac{3}{2T}, \quad d = \frac{1}{3}.$$

## Fourier series — time shifting

Theorem (on shifting in time domain)

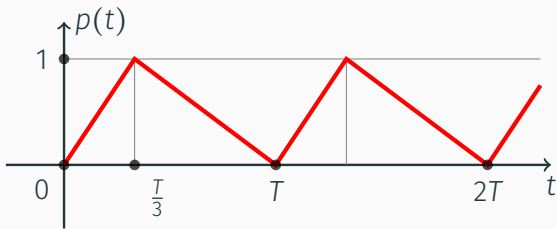
The Fourier series of a  $T$ -periodic signal  $y(t) = x(t - t_0)$  is the series

$$y(t) = \sum_{k=-\infty}^{+\infty} (X_k e^{-j\omega k t_0}) e^{jk\omega t},$$

where

$$x(t) = \sum_{k=-\infty}^{+\infty} X_k e^{jk\omega t}.$$

## How to compute Fourier series — another saw



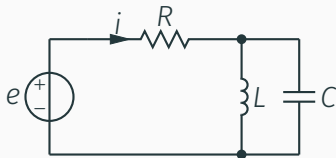
$$X_k = \frac{\sin(\pi k d)}{k\pi}, \quad d = \frac{1}{3}.$$

$$p'(t) = q(t), \quad q(t) = \frac{9}{2T}x\left(t - \frac{T}{6}\right) - \frac{3}{2T}, \quad d = \frac{1}{3}.$$

$$Q_k = \frac{9}{2T} \frac{\sin(k\frac{\pi}{3})}{k\pi} e^{-jk\frac{\pi}{3}}, \quad k \neq 0; \quad P_0 = \frac{1}{2}, \quad P_k = \frac{9}{jk4\pi} \frac{\sin(k\frac{\pi}{3})}{k\pi} e^{-jk\frac{\pi}{3}}, \quad k \neq 0.$$

## Example

Find  $i(t)$  and its RMS value  $I_{\text{RMS}}$ .



$$e(t) = 10 + 10 \cos\left(\omega t - \frac{\pi}{3}\right) + 10 \cos\left(2\omega t - \frac{\pi}{6}\right) \text{V.}$$

$$\omega = 100 \text{ krad/s}, R = 1 \text{ k}\Omega, L = 20 \text{ mH}, C = 5 \text{ nF.}$$