



# Simulating electromagnetic cascades in AGN-magnetospheres

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# Conventional jet-scenarios are threatened by IC 310 and Mrk 501

Recently, the peculiar radiogalaxy IC 310 has shown extreme flux variations (Fig. 1) with time-scales of about 5min at TeV-energies. This is evidence of emission-regions of length-scales below  $r_{\rm g}$  and thus cannot be explained by shock-in-jet models [1].

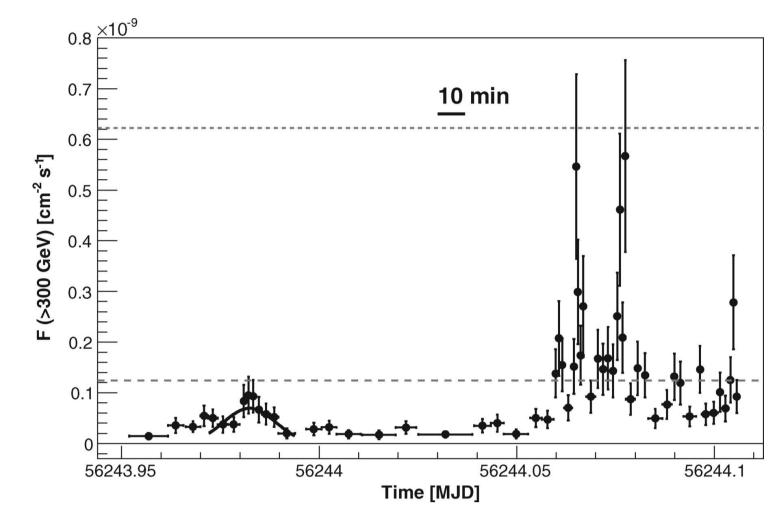


Fig. 1: IC 310's light-curve above 300 GeV, detected by the MAGIC-telescopes on 12./13. 11. 2012. [1]

Moreover, a very narrow, TeV-bump (Fig. 3b) observed in the SED of the HBL Mrk 501 by the MAGIC-telescopes in June 2012 cannot be explained by conventional models [2].

### Are vacuum gaps the enigmatic machines?

Alternative models for understanding the origin of gamma-ray emission are magnetospheric models. According to these, charge-depleted regions (vacuum gaps) in AGN-magnetospheres are the birthplace of the highest-energetic particles (both UHE-cosmic rays and GeV- to TeV-photons). These gaps are expected to be located in the magnetosphere's polar regions and permeated by electric and magnetic fields [3], giving rise to particle acceleration up to ultrahigh energies and photon emission [4, 5]. Subsequently, an electromagnetic cascade develops.

#### Electromagnetic cascades

Particles, that intrude into the gap, are accelerated to high energies and act as seed particles for the development of an electromagnetic cascade (Fig. 2) via inverse-Compton-upscattering soft background-photons (which may stem from the accretion disk or from photo-ionised clouds). By this, HE-photons are created, which, in addition to injected HE-photons (which may stem e. g. from pion-decay), can again pair-produce new particles via interaction with the soft background-photons and thus sustain the cascade.

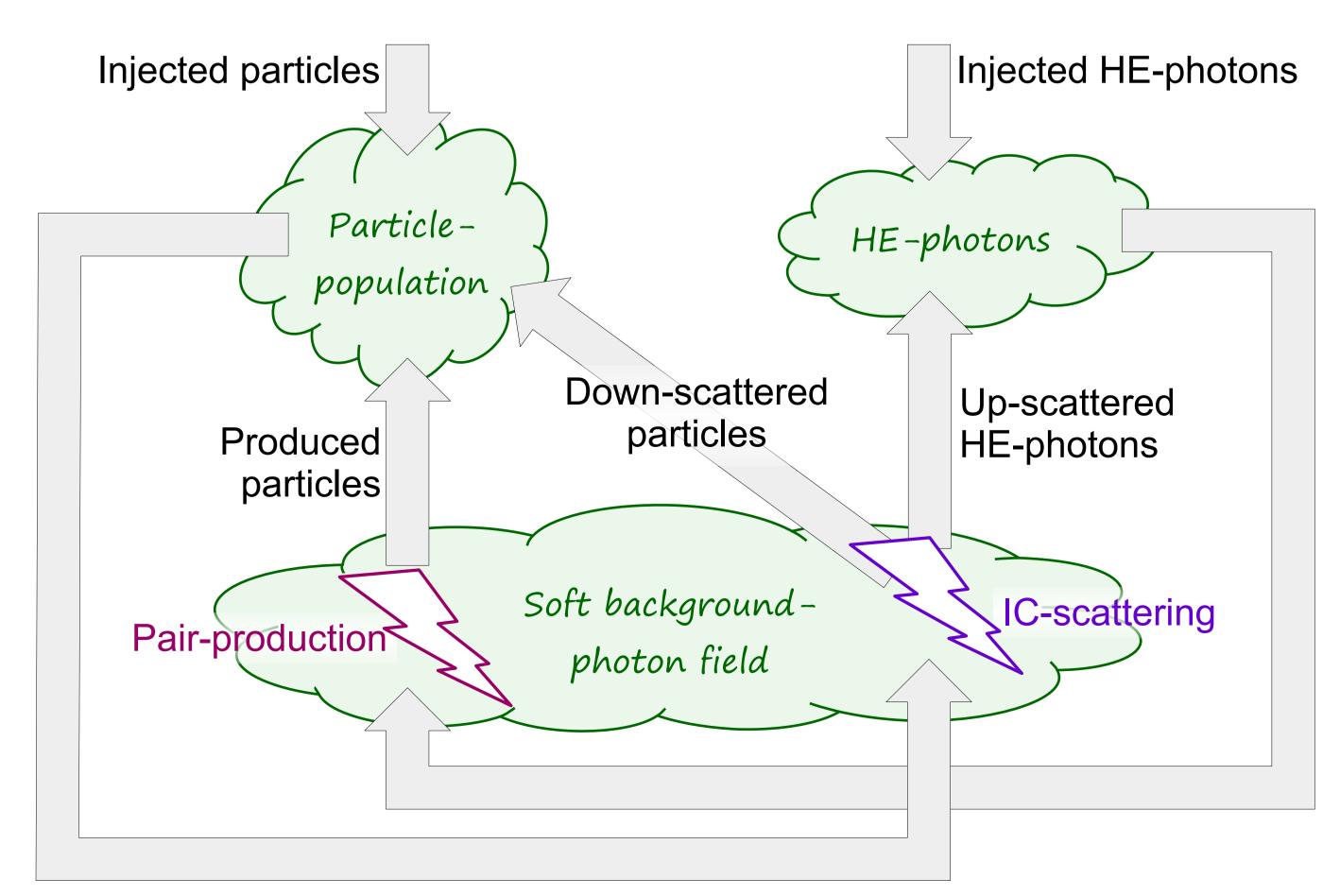
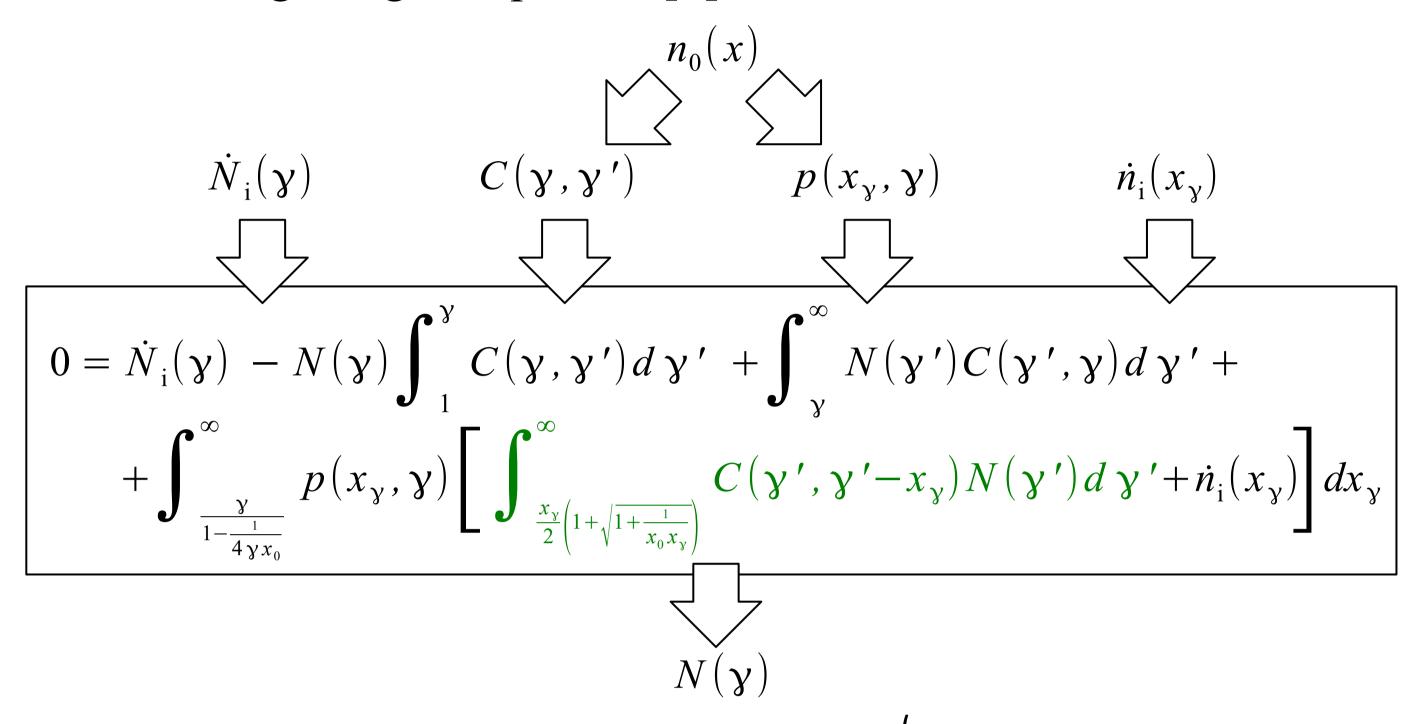


Fig. 2: Sketch of a saturated electromagnetic cascade

## Numerical solution of the cascade equation

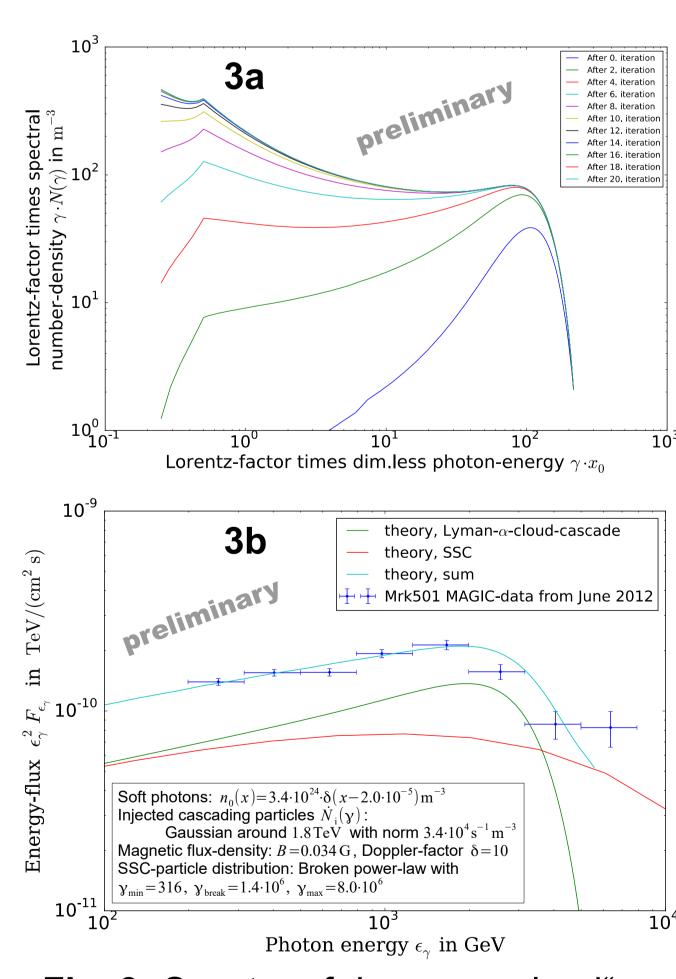
Involving pair-production and IC-scattering in the Klein-Nishina-regime, the steady-state particle-distribution  $N(\gamma)$  is described by the following integral equation [6]:



Firstly, after specifying the input functions (distributions of soft photons  $n_0(x)$ , injected HE-photons  $\dot{n}_i(x_y)$  and injected particles  $\dot{N}_i(\gamma)$ , the interaction-probabilities (kernels  $C(\gamma,\gamma')$  and  $p(x_y,\gamma)$  of the integral equation) for IC-scattering and for pair-production have to be yielded via numerical integration of the double-spectral interaction-probabilities. Secondly, the integral equation can be solved iteratively. In this way, the steady-state particle-distribution  $N(\gamma)$  in the Klein-Nishina-regime and the HE-photon-distribution  $n(x_y)$  are yielded. Next, to determine the particle-distribution in the Thomson-regime, a continuity-equation has to be solved. As a byproduct the IC-upscattered HE-photons (green term above) are yielded.

#### **Dissecting Mrk 501**

We have developed a python-code, that is able to execute the described algorithm. In the case of Mrk 501 one can assume 1.8TeV-electrons (produced in a gap) as  $N_i(\gamma)$ and Ly-α-photons (from an ionised cloud) as  $n_0(x)$ . These species interact in the jet of Mrk 501 via the described cascade. The solution  $N(\gamma)$  converges after approx. 20 iterations (Fig. 3a). The emerging HE-photon distribution is producing the bump on top of the SSC-background (which explains X-ray and γ-ray data as well) in the Mrk 501-spectrum (Fig. 3b).



**Fig. 3:** Spectra of "Lyman-α-cloud"-cascade in Mrk 501:

**a:** Iteration of particle-distribution **b:** HE-photon-distribution (data points taken from [2]) as superposition of cascade- and SSC-emission

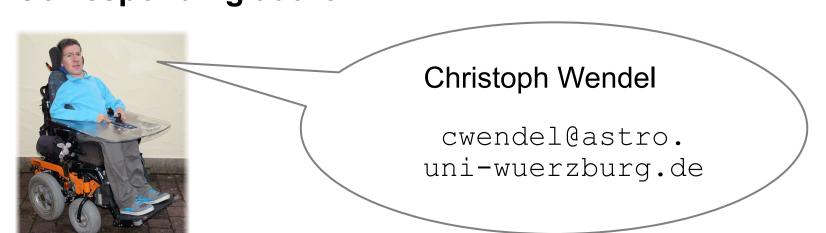
#### References:

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