

B. ENSO

- explain
- model results

Box model

$$\dot{h}_1 = -\Gamma \left(h_1 + \frac{bL}{2} \frac{\mu}{\beta} (T_2 - T_1) \right)$$

$$\dot{T}_1 = -\alpha (T_1 - T_r) - \varepsilon \mu (T_2 - T_1)^2$$

$$\dot{T}_2 = -\alpha (T_2 - T_r) + \gamma \mu (T_2 - T_1) \\ * \left(T_2 - \frac{1}{T_{sub}} (h_2) \right)$$

$$T_{sub}(h) = T_r - \frac{T_r - T_0}{2} *$$

$$* \left(1 - \tanh \left(\frac{h + H - z_0}{h_*} \right) \right)$$

$$h_2 = h_1 + bL \frac{\mu}{\beta} (T_2 - T_1)$$

$$\text{Define } S = T_2 - \bar{T}_1$$

$$T = \bar{T}_1 - \bar{T}_r$$

$$h = h_1 + H - z_0$$

$$\rightarrow \dot{S} = -\alpha S + \mu \sum S^* *$$

$$* \left(S + T + \frac{\bar{T}_1 - \bar{T}_r}{2} \right) -$$

$$+ \tanh \left(\frac{h}{h_*} + \frac{b L \mu}{2 B h_*} S \right)$$

$$+ \epsilon \mu S^2$$

$$\dot{T} = -\alpha T - \epsilon \mu S^2$$

$$\dot{h} = -\alpha \left(h + z_0 - H + \frac{b \mu L}{2 B} S \right)$$

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$$S = \chi S_0, T = \gamma T_0$$

$$h = z h_0, t = t_0 \tau$$

$$\rightarrow \frac{\partial}{\partial t} \dot{x} = \varepsilon \mu S_0^2 x^2 - \alpha S_0 x +$$

$$\mu \left(S_0 x \left(S_0 x + T_0 y + \frac{T - T_0}{2} \right) \right)$$

$$\left(1 - \tanh \left(z \frac{h_0}{h_*} + \frac{BL\mu}{Bh_*} S_0 x \right) \right)$$

$$\frac{\partial}{\partial t} \dot{y} = -\alpha T_0 y - \varepsilon \mu S_0^2 x^2$$

$$\frac{\partial}{\partial t} \dot{z} = -r \left(z h_0 + z_0 - T + \frac{BL\mu}{2B} S_0 x \right)$$

$$S_0 = T_0 = \frac{h_* B}{BL\mu} \quad h_0 = b_* \\ t_0 = \frac{BL}{B S_0 h_*}$$

$$\dot{x} = \epsilon \mu t_0 S_0 x^2 - a t_0 x + S_0 t_0 \mu \zeta x \left(x + y + \frac{T_r - T_b}{2S_0} \right) *$$

$$* \left(1 - \tanh(z + x) \right)$$

$$S_0 t_0 \mu \zeta = \frac{h_* \beta}{b L \mu} \cdot \frac{b L}{\beta \zeta h_*} \cdot \mu \zeta = 1$$

$$\boxed{C} = \frac{T_r - T_b}{2S_0}$$

$$\boxed{S} = \frac{\epsilon h_* \beta}{r b L} \quad \boxed{\bar{o}} = \frac{r b L}{S h_* \beta}$$

$$S^\delta = \epsilon / S$$

$$\boxed{a} = \frac{a}{\epsilon \mu} \cdot \frac{b L \mu}{h_* \beta}$$

$$S_0 t_0 = \frac{1}{\mu \zeta}$$

$$a t_0 = \epsilon \mu t_0 S_0 \cdot \frac{a}{\epsilon \mu S_0}$$

$$= S^\delta \cdot a$$

Finally

$$\dot{x} = \delta \left(x^2 - ax \right) + x *$$

$$(x + y + c - c \tanh(x + z))$$

$$\dot{y} = -\alpha t_0 y - \varepsilon e^{\delta t_0} t_0 x^2$$

$$\rightarrow = - (a y + x^2) \delta$$

$$\dot{z} = -\tau t_0 \left(z + \frac{z_0 - h}{\tau_*} + \right.$$

$$\left. + x/2 \right)$$

$$\bar{\delta} = \tau t_0$$

$$= \delta \left(k - z - x/2 \right)$$

$$\boxed{k} = \frac{h - z_0}{\tau_*}$$

5. Bifurcations of fixed points

Scalar equation, autonomous
($n=1$)

$$\dot{x} = f(x, \lambda)$$

Fixed points \bar{x}

$$f(\bar{x}, \lambda) = 0$$

Small perturbations \tilde{x} ,

with $|\tilde{x}| \ll \bar{x} \rightarrow$

$$\dot{\tilde{x}} + \frac{\dot{\bar{x}}}{\bar{x}} = f(\bar{x} + \tilde{x}, \lambda)$$

≈ 0

f smooth \rightarrow

$$\dot{\tilde{x}} = f(\bar{x}, \lambda) + \tilde{x} f_x(\bar{x}, \lambda) + \mathcal{O}(\tilde{x}^2)$$

$$\text{Let } \delta = f_x(\bar{x}, \lambda)$$

then $\dot{x} = \sigma \bar{x}$

$$\rightarrow \bar{x}(t) = \bar{x}_0 e^{\sigma t}$$

$\sigma < 0$: \bar{x} stable ; $\sigma > 0$: \bar{x} unstable

Changes in perturbation

behavior occur at $\sigma = 0$,

i.e. at values of λ_c for
which $f_x(\bar{x}, \lambda_c) = 0$;

λ_c are called bifurcation
points.

Classification of
bifurcation points based
on Taylor series of f .

For a single parameter λ ,
there are three cases :

(i) saddle-node

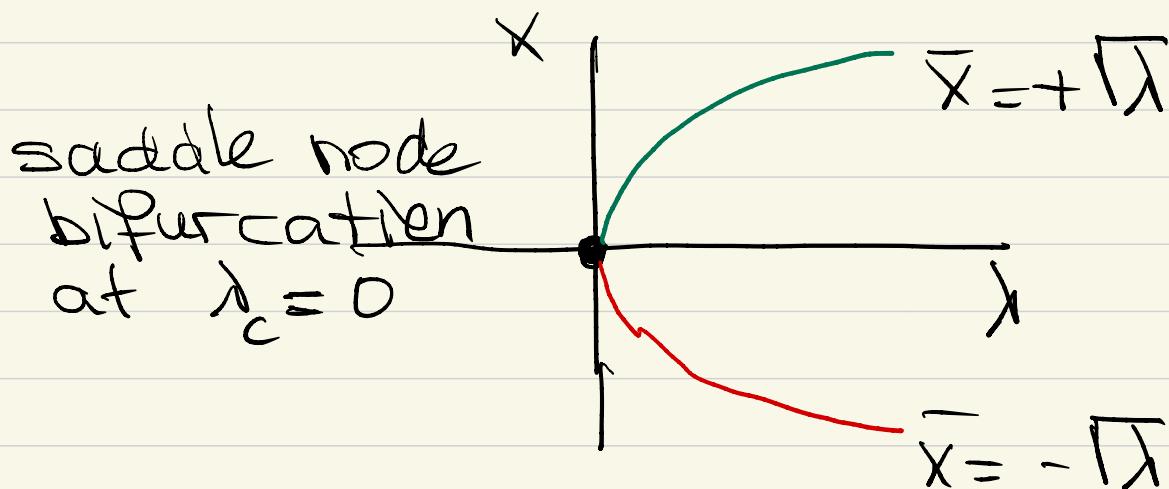
normal form: $\dot{x} = \lambda - x^2$

fixed points: $\bar{x} = \pm\sqrt{\lambda}, \lambda \geq 0$

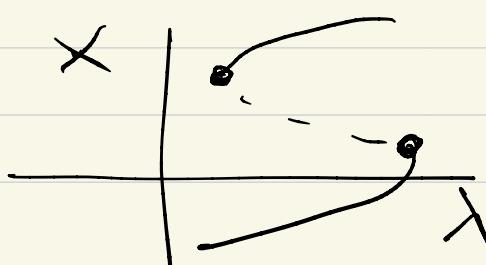
stability: $f_x(\bar{x}, \lambda) = -2\bar{x} = \mp 2\sqrt{\lambda}$

hence: $\bar{x} = +\sqrt{\lambda} \rightarrow \sigma < 0$

$\bar{x} = -\sqrt{\lambda} \rightarrow \sigma > 0$



- physical systems
(bounded solutions)



back to back saddle nodes

(ii) pitchfork bifurcation

normal form : $\dot{x} = x(\lambda - x^2)$

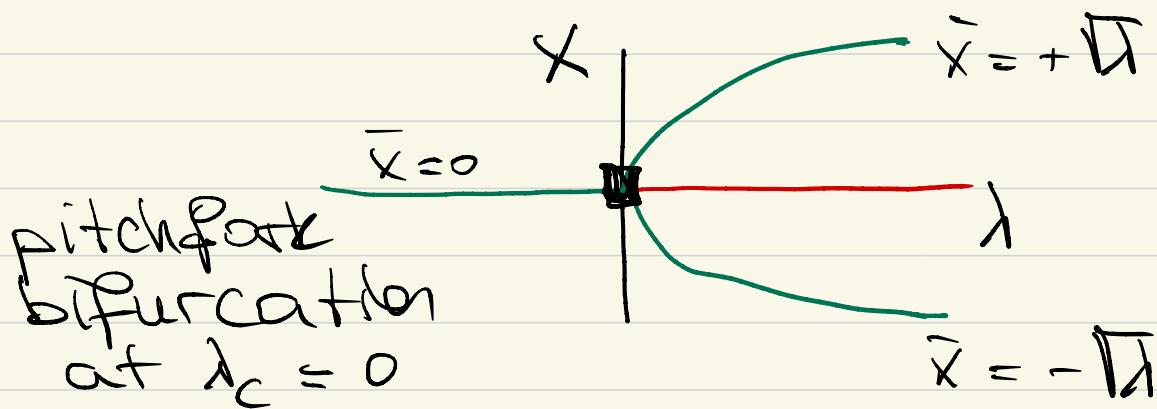
fixed points : $\bar{x} = 0 \quad \bar{x} = \pm\sqrt{\lambda}, \lambda > 0$

stability : $\sigma = \frac{d}{dx}(\bar{x})\lambda = \lambda - 3\bar{x}^2$

$$\bar{x} = 0 \rightarrow \sigma = \lambda \quad \begin{array}{c} \sigma < 0 \\ 0 \\ \sigma > 0 \end{array}$$

$$\bar{x} = +\sqrt{\lambda} \rightarrow \sigma = -2\lambda < 0$$

$$\bar{x} = -\sqrt{\lambda} \rightarrow \sigma = -2\lambda < 0$$



- physical systems

symmetry breaking

(iii) transcritical bifurcation

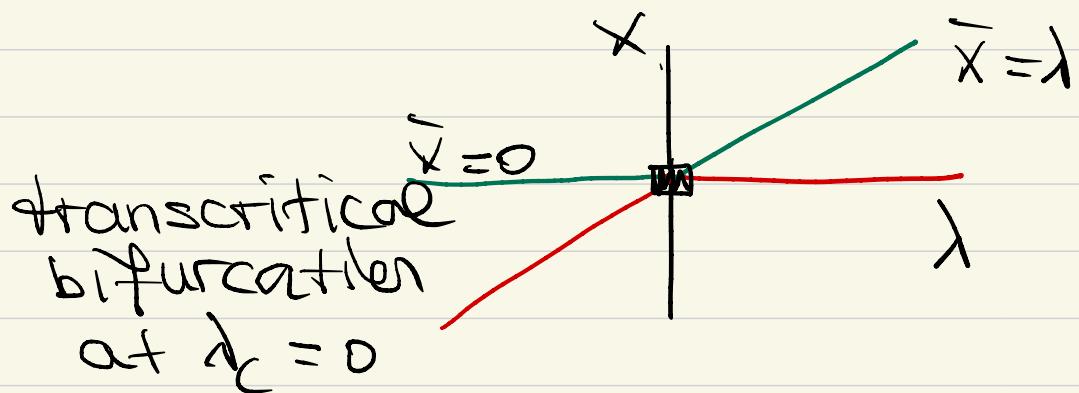
normal form : $\dot{x} = \lambda x - x^2$

fixed points : $\bar{x} = 0, \bar{x} = \lambda$

stability : $\sigma = f_x(\bar{x}, \lambda) = \lambda - 2\bar{x}$

$$\bar{x} = 0 \rightarrow \sigma = \lambda \quad \begin{array}{c} \sigma < 0 \\ \hline 0 \\ \sigma > 0 \end{array}$$

$$\bar{x} = \lambda \rightarrow \sigma = -\lambda \quad \begin{array}{c} \sigma > 0, \sigma < 0 \\ \hline 0 \\ \sigma < 0 \end{array}$$



- Physics

- no symmetry,
(trivial) solution for all
values of λ .

Two-dimensional case

$$\begin{cases} \dot{x} = f(x, y, t) \\ \dot{y} = g(x, y, t) \end{cases}$$

fixed points : $f(\bar{x}, \bar{y}, \bar{t}) = 0$

$$g(\bar{x}, \bar{y}, \bar{t}) = 0$$

stability : $\begin{cases} \dot{x} = \bar{x} + x^1 \\ \dot{y} = \bar{y} + y^1 \end{cases}$

$$\begin{aligned} \dot{x}^1 &= f(\bar{x} + x^1, \bar{y} + y^1, \bar{t}) \\ \dot{y}^1 &= g(\bar{x} + x^1, \bar{y} + y^1, \bar{t}) \end{aligned}$$

$$\begin{aligned} \dot{x}^1 &= f_x \bar{x} + f_x x^1 + f_y \bar{y} + f_y y^1 + \dots \\ &= g_x \bar{x} + g_x x^1 + g_y \bar{y} + g_y y^1 + \dots \end{aligned}$$

Jacobian matrix

$$\mathbb{J} = \begin{pmatrix} \frac{\partial f}{\partial x} & \frac{\partial f}{\partial y} \\ \frac{\partial g}{\partial x} & \frac{\partial g}{\partial y} \end{pmatrix}$$

$\rightarrow \begin{pmatrix} \frac{\partial x^1}{\partial t} & \frac{\partial x^2}{\partial t} \end{pmatrix} = \mathbb{J} \begin{pmatrix} x^1 \\ x^2 \end{pmatrix}$

Solution: $\begin{pmatrix} \frac{\partial x^1}{\partial t} \\ \frac{\partial x^2}{\partial t} \end{pmatrix} = e^{t\mathbb{J}} \begin{pmatrix} x_0^1 \\ x_0^2 \end{pmatrix}$

$$\mathbb{J} = T \sum_i \frac{1}{\lambda_i} \mathbb{J}_i$$

eigenvalues of \mathbb{J}

In the 2-dimensional case,

there is one additional
bifurcation

(iv) Hopf bifurcation

normal form

$$\begin{cases} \dot{x} = \lambda x - \omega y - x(x^2 + y^2) \\ \dot{y} = \omega x + \lambda y - y(x^2 + y^2) \end{cases}$$

ω fixed; λ varying

fixed point : $\bar{x} = \bar{y} = 0$

stability :

$$J\left(\bar{x} = \bar{y} = 0\right) = \begin{pmatrix} \lambda & -\omega \\ \omega & \lambda \end{pmatrix}$$

eigenvalues : $\det(J - \lambda I) = 0$

$$\begin{vmatrix} \lambda - \sigma & -\omega \\ \omega & \lambda - \sigma \end{vmatrix} = 0$$

$$\rightarrow (\lambda - \sigma)^2 + \omega^2 = 0$$

$$\sigma_{1,2} = \lambda \pm i\omega$$

$$\lambda > 0 \quad : \operatorname{Re}(\sigma) > 0$$

$$\lambda < 0 \quad : \operatorname{Re}(\sigma) < 0$$

Solutions :

$$\begin{pmatrix} x \\ y \end{pmatrix} = e^{\lambda t} \begin{pmatrix} \tilde{x}_0 \cos \omega t + \tilde{y}_0 \sin \omega t \\ -\tilde{x}_0 \sin \omega t + \tilde{y}_0 \cos \omega t \end{pmatrix}$$

Polar coordinates

$$\left. \begin{array}{l} x = r \cos \theta \\ y = r \sin \theta \end{array} \right\}$$

$$\left. \begin{array}{l} \dot{x} = \dot{r} \cos \theta - \dot{\theta} r \sin \theta \\ \dot{y} = \dot{r} \sin \theta + \dot{\theta} r \cos \theta \\ x^2 + y^2 = r^2 \end{array} \right\}$$

→

$$\begin{aligned} \dot{r} \cos \theta - \dot{\theta} r \sin \theta &= \lambda r \cos \theta - \\ -\omega r \sin \theta - r^3 \dot{\theta} \cos \theta & \end{aligned} \quad (1)$$

$$\begin{aligned} \dot{r} \sin \theta + \dot{\theta} r \cos \theta &= \omega r \sin \theta + \\ + \lambda r \sin \theta - r^3 \dot{\theta} \sin \theta & \end{aligned} \quad (2)$$

$$(\cos \theta \ (1) + \sin \theta \ (2)) \Rightarrow$$

$$\begin{aligned} \dot{r} &= \lambda r - r^3 \quad \left(\text{'pitchfork'} \right) \\ -\sin \theta \ (1) + \cos \theta \ (2) &\Rightarrow \end{aligned}$$

$$r^2 \dot{\theta} = -\omega r^2 \rightarrow \dot{\theta} = \omega$$

