# Lindblad

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# 1 Lie Geometry of Density Matrix

Rather than parameterize  $\rho$  through matrix entries, parameterize it via an exponential family. Use the known importance of generalized Gibbs ensembles. Also reveals the Poisson geometry underlying Fisher metric.

# 2 Lindblad

$$\dot{\rho}(t) = -i[H, \rho(t)] + \sum_{k=1}^{M} L_k \rho(t) L_k^{\dagger} - \frac{1}{2} \{ L_k^{\dagger} L_k, \rho(t) \}$$

The action of H on  $V \otimes V^{\dagger}$  is given by coproduct and a sign flip for the  $V^{\dagger}$  (More generally the antipode, but in this case just -1).

$$\Delta H = H \otimes 1 - 1 \otimes H$$

$$\dot{\rho} \in V \otimes V^{\dagger} 
\dot{\rho} = (-i\Delta H + \mathcal{G})\rho 
\mathcal{G} = \sum_{m=0}^{M} \bar{L}_{m} \otimes L_{m} - \frac{1}{2}Id \otimes L_{m}^{\dagger}L_{m} - \frac{1}{2}\bar{L}_{m}^{\dagger}\bar{L}_{m} \otimes Id$$

# 3 Effective Hamiltonian

# 3.1 Definition $(H_{eff})$

$$H_{eff} = H - i \sum_{k} \frac{\gamma_k}{2} L_k^{\dagger} L_k$$

The antiHermitian part is negative for  $\gamma_k \geq 0$  in the sense of  $L_k^{\dagger} L_k$  being a positive operator. In Lie algebraic terms  $iH_{eff} \in \mathfrak{u}(n) \bigoplus \mathfrak{u}(n)_{>0}^*$ 

**3.2 Definition (Iwasawa Poisson-Lie Group)** The one coming from KAN decomposition. So for  $GL(n, \mathbb{C})$  regarded as a real Lie group, we get Hermitian and anti-Hermitian split. TODO: Write down this parameterization.

$$\frac{d\rho}{dt} = -i(H_{eff}\rho - \rho H_{eff}^{\dagger}) + \sum_{k} \gamma_{k} L_{k} \rho L_{k}^{\dagger}$$

### 4 Lie Parameterization

Let  $\rho =$ 

### 5 Examples

#### 5.1 Example

$$H = 0$$

$$\frac{\partial \rho}{\partial t} = -i[0, \rho] + \frac{\gamma}{2} (\sigma_z \rho \sigma_z - \frac{1}{2} {\{\sigma_z^{\dagger} \sigma_z, \rho\}})$$

#### 5.2 Example

$$H = \hbar \omega a^{\dagger} a$$

$$L_{1} = a$$

$$\gamma_{1} = \frac{\gamma}{2} (\bar{n} + 1)$$

$$L_{2} = a^{\dagger}$$

$$\gamma_{2} = \frac{\gamma}{2} (\bar{n})$$

$$H_{eff} = \hbar \omega a^{\dagger} a - i \frac{\gamma}{4} ((\bar{n} + 1) a^{\dagger} a + (\bar{n}) a a^{\dagger})$$

$$= \hbar \omega a^{\dagger} a - i \frac{\gamma}{4} ((2\bar{n} + 1) a^{\dagger} a + \bar{n})$$

### 6 Fisher Information Norm of Lindblad Vector Field

The ordinary differential equation described above determines a vector field on the space of density matrices (M,g) where g is the Fisher information metric tensor. So one can evaluate  $v(\rho)$  which is the direction of flow when starting at  $\rho$ . One can use g to give  $g(v(\rho), v(\rho))$  which measures the length thereof. The actual dynamics is a trajectory for this vector field by solving for the flow  $\phi_t$ . We want to calculate the length of the curve.

$$L = \int_0^T \sqrt{g(v(\rho(t)), v(\rho(t)))} dt$$

where  $\rho(t) = \phi_t \rho(0)$  is the solution to the ordinary differential equation on M.

# 7 Hochschild Homology

### 8 Quantum Chemistry

https://dash.harvard.edu/bitstream/handle/1/4646014/0807.0929v2.pdf?sequence=1

### 8.1 Unitary Evolution

https://arxiv.org/pdf/0905.0887.pdf

### 8.2 Open

For a particular small open system with known jump operators, can we come up with an auxiliary system to approximate the infinite dimensional bath Hilbert space in such a way that following the gates of "Towards quantum chemisty" on this auxiliary closed system problem implements the Lindblad evolution for the qubits  $H_S = (\mathbb{C}^2)^d$ ?

Start with well studied  $H_2$  molecule example. Decide what environment we should put it in... Closest reference could find on this subject: https://www.pks.mpg.de/~eisfeld/wwww/pdf/Mostame\_et\_al\_Emulation\_of\_open\_quant\_scqb.pdf