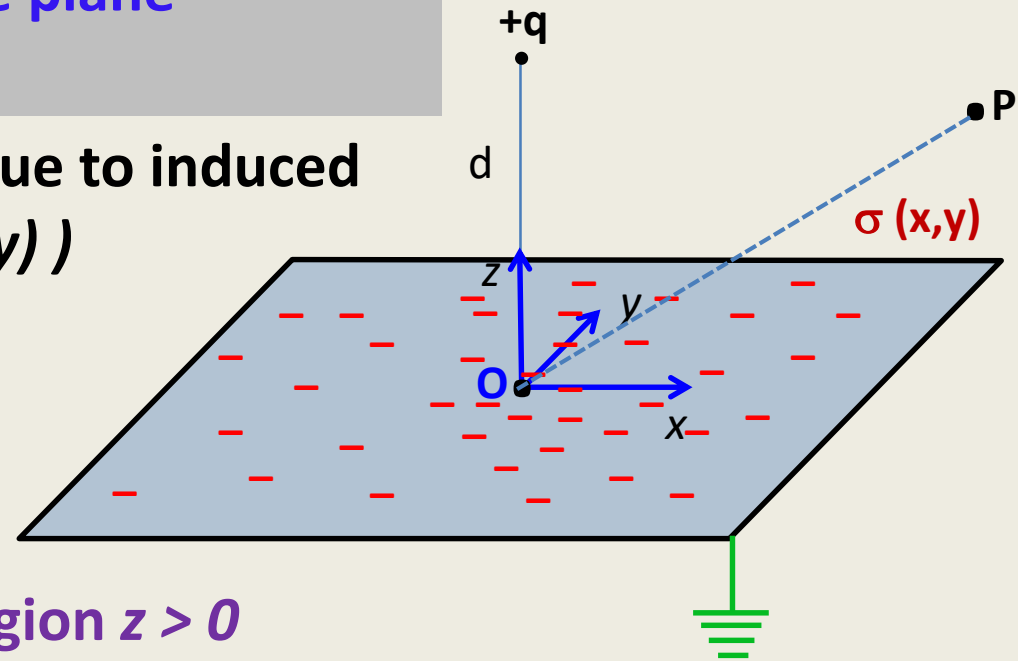


Method of Images

Point charge $+q$ near an infinite plane CONDUCTOR

$V(P) = (V(P) \text{ due to } q) + (V(P) \text{ due to induced charges on the conductor } \sigma(x,y))$

$\sigma(x,y)$ is NOT known



Solve Poisson's eqn. in a region $z > 0$
with a single point charge at $(0,0,d)$
and conductor at origin in the x - y
plane

Boundary conditions:

(1) $V = 0$ at $z = 0$ (conductor is grounded)

(2) $V \rightarrow 0$ as $(x^2 + y^2 + z^2) \gg d^2$

Solve a different problem (Analogue problem)

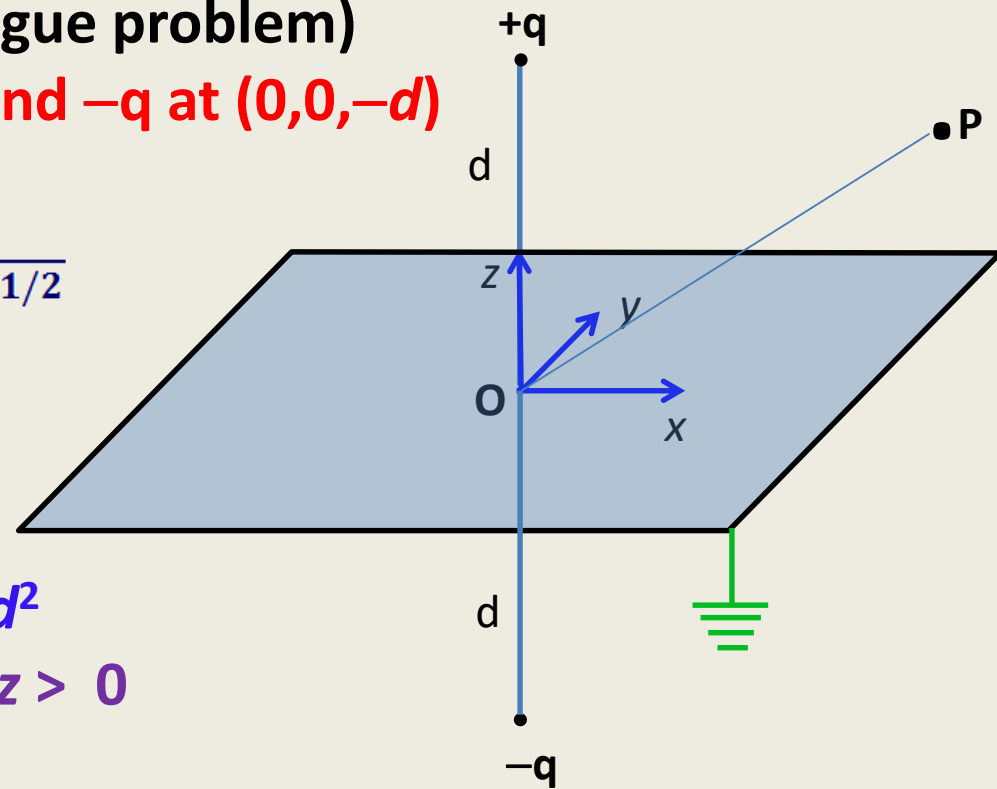
Two point charges $+q$ at $(0,0,d)$ and $-q$ at $(0,0,-d)$

$$V(P) = \frac{1}{4\pi\epsilon_0} \left\{ \frac{q}{[x^2 + y^2 + (z - d)^2]^{1/2}} + \frac{-q}{[x^2 + y^2 + (z + d)^2]^{1/2}} \right\}$$

(i) $V = 0$ when $z = 0$

(ii) $V \rightarrow 0$ as $(x^2 + y^2 + z^2) \gg d^2$

(iii) There is only $+q$ at $(0,0,d)$, $z > 0$



Now bring the infinite conducting sheet and keep in the x-y plane.
Adjust potential to be zero (grounding)

Top configuration has the same potential as the original problem.

Then this is the only solution (Uniqueness theorem)

Point charge +q near an infinite plane CONDUCTOR

$$V(P) = \frac{1}{4\pi\epsilon_0} \left\{ \frac{q}{[x^2 + y^2 + (z - d)^2]^{1/2}} + \frac{-q}{[x^2 + y^2 + (z + d)^2]^{1/2}} \right\}$$

Induced charge density $\sigma(x, y)$

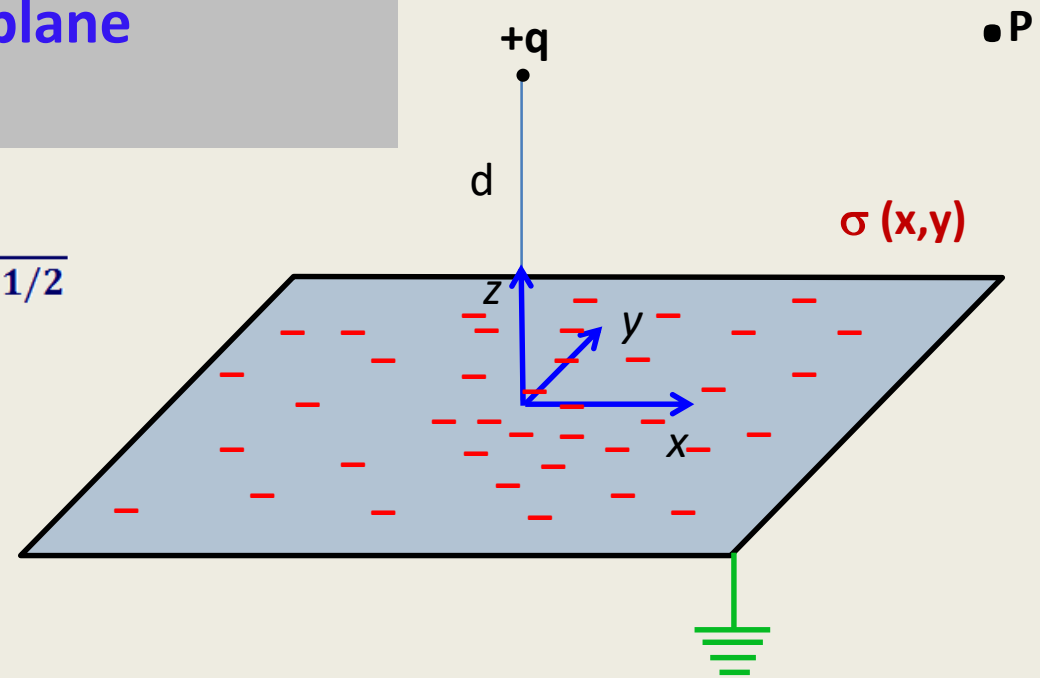
$$\sigma(x, y) = -\epsilon_0 \frac{\partial V}{\partial n} = -\epsilon_0 \frac{\partial V}{\partial z} \Big|_{z=0} = -\frac{qd}{2\pi} \frac{1}{(d^2 + x^2 + y^2)^{3/2}}$$

Total Induced charge q_{ind}

$$q_{\text{ind}} = \int \sigma dS = -\frac{qd}{2\pi} \int \frac{1}{(d^2 + x^2 + y^2)^{3/2}} dx dy = -q$$

Force of attraction

$$F = \frac{q(-q)}{4\pi\epsilon_0(2d)^2} \text{ (analogue problem)}$$

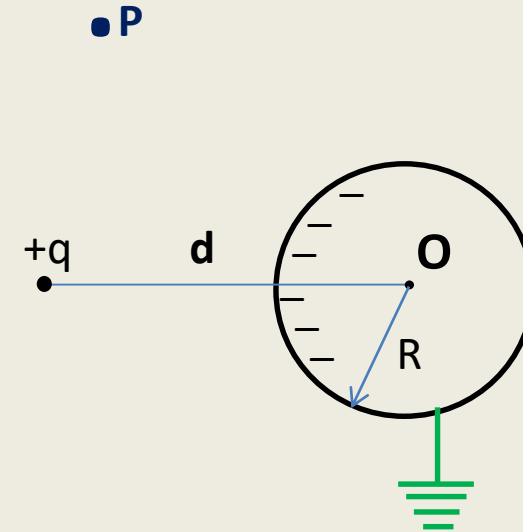


Point charge +q near an infinite plane CONDUCTING SPHERE

$V(P) = V(P)$ due to q + $V(P)$ due to induced charges on the conductor $\sigma(\theta, \phi)$

$\sigma(\theta, \phi)$ is NOT known

Solve Poisson's eqn. in a region $r > R$ with a single point charge at ($r = d$) and conductor with origin at the centre of the coordinate system

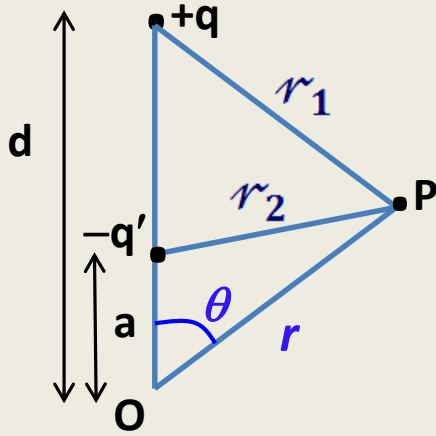


Boundary conditions:

(1) $V = 0$ at $r = R$ (conductor is grounded)

(2) $V \rightarrow 0$ as $r \rightarrow \infty$

Analogue problem



$$\begin{aligned} V(P) &= \frac{1}{4\pi\epsilon_0} \left(\frac{q}{r_1} - \frac{q'}{r_2} \right) \\ &= \frac{1}{4\pi\epsilon_0} \left(\frac{q}{(r^2 + d^2 - 2rd \cos \theta)^{1/2}} - \frac{q'}{(r^2 + a^2 - 2ra \cos \theta)^{1/2}} \right) \\ &= \frac{1}{4\pi\epsilon_0} \left(\frac{q}{d \left(1 + \frac{r^2}{d^2} - 2 \frac{r}{d} \cos \theta \right)^{1/2}} - \frac{q'}{r \left(1 + \frac{a^2}{r^2} - 2 \frac{a}{r} \cos \theta \right)^{1/2}} \right) \end{aligned}$$

Let us make $V = 0$ at some point $r = R$

$$\frac{q}{d \left(1 + \frac{R^2}{d^2} - 2 \frac{R}{d} \cos \theta \right)^{1/2}} = \frac{q'}{R \left(1 + \frac{a^2}{R^2} - 2 \frac{a}{R} \cos \theta \right)^{1/2}}$$

Solve for q'

$$q' = \left\{ q \frac{R}{d} \right\} \frac{\left(1 + \frac{a^2}{R^2} - 2 \frac{a}{R} \cos \theta \right)^{1/2}}{\left(1 + \frac{R^2}{d^2} - 2 \frac{R}{d} \cos \theta \right)^{1/2}}$$

For this to be true for any θ ,
 $q' = qR/d$ and $a = R^2/d$

$$V(r, \theta, \phi) = \frac{1}{4\pi\epsilon_0} \left(\frac{q}{(r^2 + d^2 - 2rd \cos \theta)^{1/2}} - \frac{q'}{(r^2 + a^2 - 2ra \cos \theta)^{1/2}} \right)$$

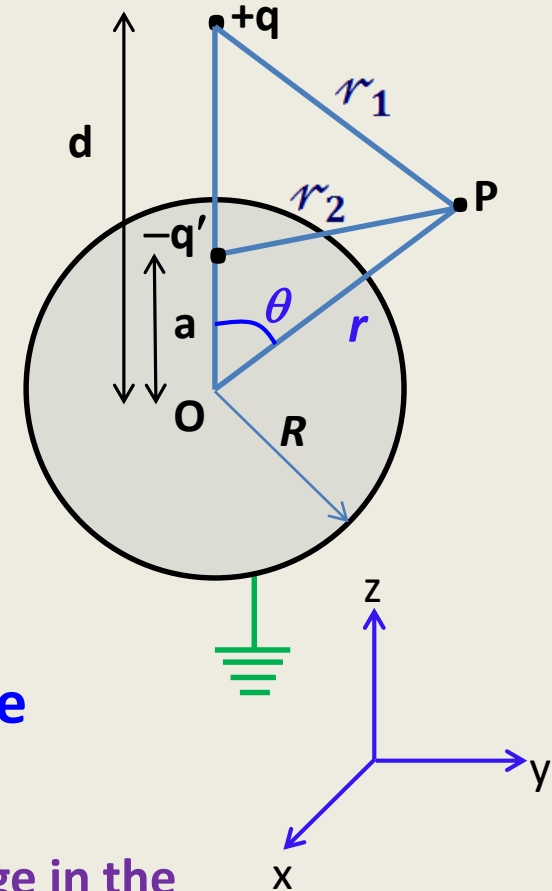
where

$$q' = \left\{ q \frac{R}{d} \right\} \quad \text{and} \quad a = \frac{R^2}{d}$$

Image charge $q' = -q R/d$ is formed at a distance $a = R^2/d$ from the origin

Since the image charge is inside the sphere ($r < R$), no change in the conditions for the solution of Poisson's equation

image charge cannot be placed in a region where V is to be calculated



Induced charge σ

$$\sigma = -\epsilon_0 \left. \frac{\partial V}{\partial n} \right|_{\text{surface}}$$

$$\begin{aligned} V(r, \theta, \phi) &= \frac{1}{4\pi\epsilon_0} \left(\frac{q}{(r^2 + d^2 - 2rd \cos \theta)^{1/2}} - \frac{q'}{(r^2 + a^2 - 2ra \cos \theta)^{1/2}} \right) \\ &= \frac{q}{4\pi\epsilon_0} \left(\frac{1}{(r^2 + d^2 - 2rd \cos \theta)^{1/2}} - \frac{1}{\left(\frac{d^2 r^2}{R^2} + R^2 - 2rd \cos \theta \right)^{1/2}} \right) \\ \left. \frac{\partial V}{\partial r} \right|_{r=R} &= \frac{-q}{4\pi\epsilon_0} \left(\frac{(R - d \cos \theta)}{(R^2 + d^2 - 2Rd \cos \theta)^{3/2}} - \frac{\left(\frac{d^2}{R} - d \cos \theta \right)}{(d^2 + R^2 - 2Rd \cos \theta)^{3/2}} \right) \end{aligned}$$

$$\sigma = \frac{-q(d^2 - R^2)}{4\pi R(R^2 + d^2 - 2Rd \cos \theta)^{3/2}} = \sigma(\theta)$$

$$q_{\text{induced}} = \int \sigma dS = \frac{-qR}{d}$$

Force of attraction

$$F = \frac{1}{4\pi\epsilon_0} \frac{q(-q')}{(d - a)^2} = \frac{-q^2 R d}{4\pi\epsilon_0 (d^2 - R^2)^2}$$

Point charge +q near a **floating** CONDUCTING SPHERE

Image charge $q' = -q R/d$ forms at distance $a = R^2/d$ from the centre, inside the sphere

Since the sphere is NOT grounded, the potential has to be raised to a value of $q/4\pi\epsilon_0 d$

This has to be achieved by the left over +ve charges since +q and $-q'$ makes the potential on the surface to zero

This can be achieved by putting the left over charges $q'' = +q R/d$ at the centre of the sphere

$$\text{Then, } V = \frac{q''}{4\pi\epsilon_0 R} = \frac{+q R/d}{4\pi\epsilon_0 R} = \frac{+q}{4\pi\epsilon_0 d}$$

