

Gauss's Law:-

Gauss's law states that the total electric flux ψ through any closed surface is equal to the total charge enclosed by the surface.

$$\psi = Q_{enc}$$

i.e.
$$\psi = \oint_S d\psi = \oint_S \mathbf{D} \cdot d\mathbf{s}$$

= total charge enclosed $Q = \int_V \rho_v dv$

(or)

$$Q = \oint_S \mathbf{D} \cdot d\mathbf{s} = \int_V \rho_v dv \quad \text{--- (1)}$$

By applying divergence theorem to the middle term

$$\oint_S \mathbf{D} \cdot d\mathbf{s} = \int_V \nabla \cdot \mathbf{D} dv \quad \text{--- (2)}$$

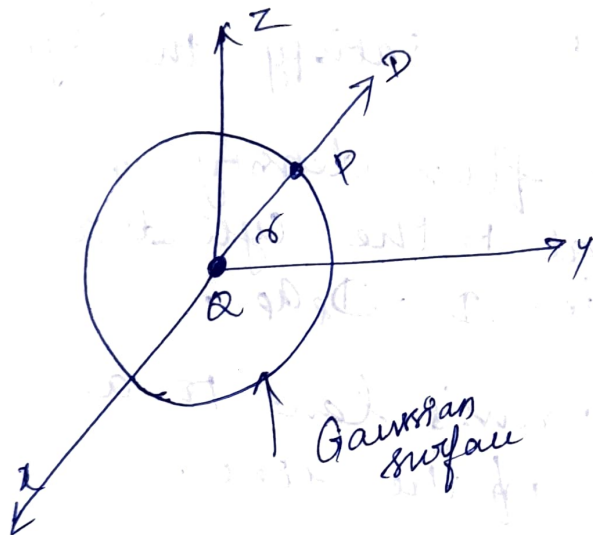
Comparing the two volume integrals

$$\boxed{\rho_v = \nabla \cdot \mathbf{D}}$$

This is the 1st Maxwell's eqn.

Point charge:-

- * Suppose a point charge Q is located at the origin.
- * Determine \vec{D} at a point P , choosing a spherical surface containing P will satisfy symmetry conditions.
- * Spherical surface centered at the origin is the Gaussian surface.



Since \vec{D} is everywhere normal to the Gaussian surface, $\vec{D} = D_r \hat{a}_r$ applying Gauss's law ($\psi = Q_{enclosed}$) gives

$$Q = \oint_S \vec{D} \cdot d\vec{s} = D_r \oint_S ds = D_r 4\pi r^2$$

where $\oint ds = \int_0^{2\pi} \int_0^{\pi} r^2 \sin\theta \, d\theta \, d\phi$

$$\begin{aligned} &= r^2 [-\cos\theta]_0^{\pi} (\phi)_0^{2\pi} \\ &= r^2 (2)(2\pi) = 4\pi r^2 \end{aligned}$$

$$D_r = \frac{Q}{4\pi r^2}$$

$$D = \frac{Q}{4\pi r^2} \cdot a_r$$

Infinite line charge:-

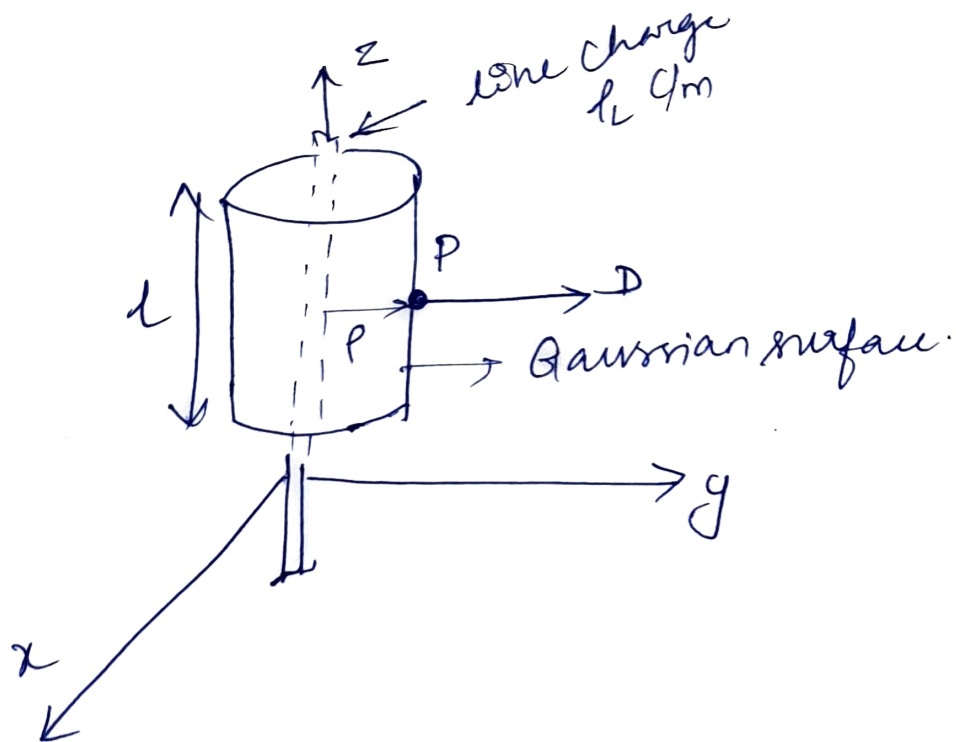
- Suppose the infinite line of uniform charge ρ_L C/m lies along the z -axis. To determine D at a point P , we choose a cylindrical surface containing P to satisfy the symmetry condition
- The electric flux density D is constant on & normal to the cylindrical Gaussian surface. i.e. $D = D_r a_r$.
- Apply Gauss's law to an arbitrary length l of the line.
- $$Q = \rho_L l = \oint_S D \cdot ds$$
- $$= D_r \oint_S ds \rightarrow \int_0^l \int_0^{2\pi} \rho_L r dr d\phi dz$$
- $$\rho_L l = D_r 2\pi r l \quad D_r = \frac{\rho_L}{2\pi r}$$

* Where $\oint ds = 2\pi r l$ is the surface area of the Gaussian surface.

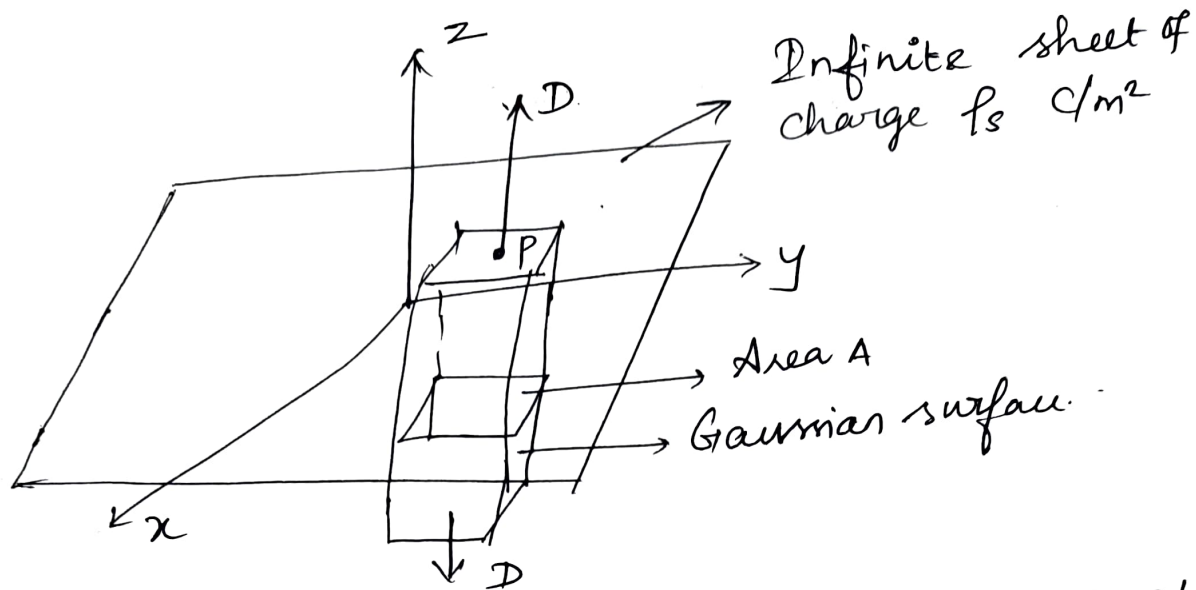
▽ Note: that $\int D \cdot ds$ evaluated on the top & bottom surfaces of the cylinder is zero.

→ D has no z -component that means D is tangential to those surfaces.

$$D = \frac{\rho_L}{2\pi r} a_r$$



Infinite sheet of charge:-



Consider an infinite sheet of charge ρ_s C/m² lying on the $z=0$ plane.

To determine D at point P , we choose a rectangular box that is cut symmetrically by the sheet of charge and has two ~~sides~~ of its faces parallel to the sheet as shown in fig. As D is normal to the sheet $D = D_z a_z$. apply Gauss's law gives

$$\rho_s \int_S ds = Q = \oint_S D \cdot ds$$

$$= D_z \left[\int_{\text{top}} ds + \int_{\text{bottom}} ds \right]$$

$D \cdot ds \rightarrow$ on the sides of the box is zero because D has no component along a_x & a_y

If the top and bottom area of the box each has area A .

$$\rho_s A = D_z (A + A)$$

$$\& \quad \mathcal{Q}_z = \frac{\rho_s A}{2A} = \frac{\rho_s}{2}$$

$$\mathcal{Q} = \frac{\rho_s}{2} a_z$$

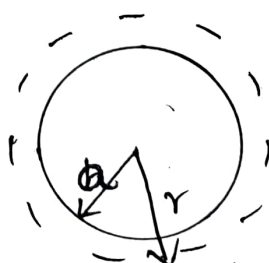
$$(or) \quad E = \frac{D}{\epsilon_0} = \frac{\rho_s}{2\epsilon_0} a_z$$

Uniformly charged sphere

Consider a sphere of radius ' a ' with a uniform charge ρ C/m³.

To determine ' \mathcal{Q} ' everywhere, construct Gaussian surfaces for cases $r \leq a$ & $r > a$ separately.

spherical surface \rightarrow appropriate Gaussian surface.



Gaussian surface for a uniformly charged sphere when $r \geq a$ & $r \leq a$.

For $r \leq a$, The total charge enclosed by the spherical surface of radius 'r'.

$$\begin{aligned}
 Q_{enc} &= \int_V \rho_v dv \\
 &= \rho_0 \int dv \\
 &= \rho_0 \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} \int_{r=0}^r r^2 \sin\theta dr d\theta d\phi \\
 &= \rho_0 \left[\left(\frac{r^3}{3} \right)_0^r (\cos\theta)_0^{\pi} [\phi]_0^{2\pi} \right] \\
 &= \rho_0 \left[\frac{r^3}{3} - 0 \right] [(-1) - 1] [2\pi - 0] \\
 &= \rho_0 \left(\frac{r^3}{3} \right) (2) (2\pi) \\
 &= \frac{\rho_0 4\pi r^3}{3}
 \end{aligned}$$

$$\begin{aligned}
 \& \quad \psi = \oint_S \mathbf{D} \cdot d\mathbf{s} = D_r \oint_S ds = D_r \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} r^2 \sin\theta d\theta d\phi \\
 & \quad D_r \left[r^2 (-\cos\theta)_0^{\pi} [\phi]_0^{2\pi} \right] \\
 & \quad = D_r r^2 (2) (2\pi) \\
 & \quad = D_r 4\pi r^2
 \end{aligned}$$

$\psi = Q_{enc}$ gives

$$\begin{aligned}
 D_r 4\pi r^2 &= \frac{4\pi r^3}{3} \rho_0 \\
 D_r &= \frac{r}{3} \rho_0
 \end{aligned}$$

$$\boxed{D = \frac{r}{3} \rho_0 \text{ or } 0 \leq r \leq a}$$

For $r \geq a$, the charge enclosed by the surface is the entire charge in this case,

$$\begin{aligned} Q_{enc} &= \int_V \rho_v dv = \rho_0 \int_V dv \\ &= \rho_0 \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} \int_{r=0}^a r^2 \sin\theta dr d\theta d\phi \\ &= \rho_0 \cdot \frac{4}{3} \pi a^3 \end{aligned}$$

$$\begin{aligned} \psi &= \oint_S \mathbf{D} \cdot d\mathbf{s} = D_r \oint d\mathbf{s} \\ &= D_r \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} r^2 \sin\theta d\theta d\phi \\ &= D_r 4\pi r^2 \end{aligned}$$

$$\psi = Q_{enc}$$

$$\rho_0 \frac{4}{3} \pi a^3 = D_r 4\pi r^2$$

$$D_r = \frac{\rho_0 a^3}{3r^2}$$

$$\boxed{D = \frac{a^3 \rho_0}{3r^2} \mathbf{a}_r} \quad r \geq a$$

\mathbf{D} everywhere is given by,

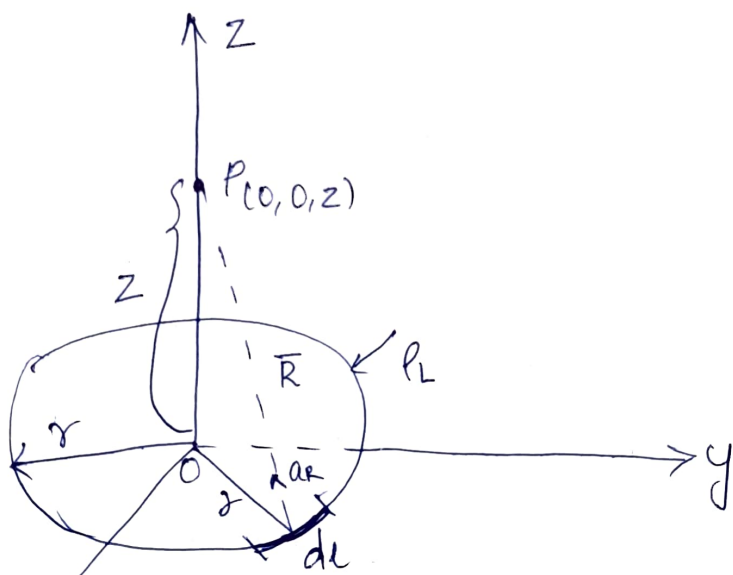
$$\mathbf{D} = \begin{cases} \frac{r}{3} \rho_0 \mathbf{a}_r & 0 \leq r \leq a \\ \frac{a^3}{3r^2} \rho_0 \mathbf{a}_r & r \geq a \end{cases}$$

Electric Field Due to charged circular Ring

Consider a charged circular ring of radius ' r ' placed in xy plane with centre at origin, carrying a charge uniformly along its circumference.

The charge density is ρ_L C/m

The point 'P' is at a perpendicular distance ' z ' from the ring as shown in fig below.



Consider a small differential length ' dl ' on this ring. The charge on it is dq .

$$dq = \rho_L dl \quad E = \frac{Q}{4\pi\epsilon_0 R^2}$$

$$d\vec{E} = \frac{\rho_L dl}{4\pi\epsilon_0 R^2} \vec{a}_R \quad \text{--- ①}$$

where

R = distance of point 'P' from dl .

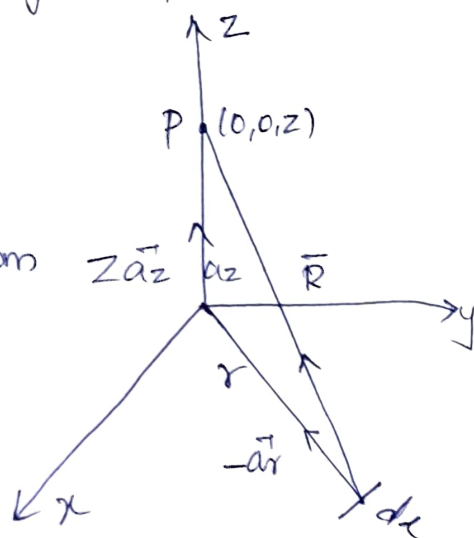
Consider the cylindrical coordinate system.

For 'dl' we are moving in ' ϕ ' direction

where $dl = r d\phi$

$$R^2 = r^2 + z^2$$

\vec{R} can be obtained from its two components in cylindrical system.



1) distance ' r ' in the direction of $-\vec{a}_r$ radially inwards
ie $-r\vec{a}_r$

2) distance ' z ' in the direction of \vec{a}_z ie $z\vec{a}_z$

$$\vec{R} = -r\vec{a}_r + z\vec{a}_z$$

$$|\vec{R}| = \sqrt{(r)^2 + (z)^2} = \sqrt{r^2 + z^2}$$

$$\vec{a}_R = \frac{\vec{R}}{|\vec{R}|} = \frac{-r\vec{a}_r + z\vec{a}_z}{\sqrt{r^2 + z^2}}$$

$$d\vec{E} = \frac{\rho_L dl}{4\pi\epsilon_0 (\sqrt{r^2 + z^2})^2} \times \frac{-r\vec{a}_r + z\vec{a}_z}{\sqrt{r^2 + z^2}}$$

$$= \frac{\rho_L (r d\phi) (-r\vec{a}_r + z\vec{a}_z)}{4\pi\epsilon_0 (r^2 + z^2)^{3/2}}$$

The radial components of \vec{E} at point 'P' will be symmetrically placed in the plane parallel to 'xy' plane and are going to cancel each other.

Hence neglecting \vec{a}_r component from $d\vec{E}$ we get

$$d\vec{E} = \frac{\rho_L (r d\phi)}{4\pi \epsilon_0 (r^2 + z^2)^{3/2}} z \vec{a}_z$$

$$\vec{E} = \int_{\phi=0}^{2\pi} \frac{\rho_L r d\phi}{4\pi \epsilon_0 (r^2 + z^2)^{3/2}} z \vec{a}_z$$

$$= \frac{\rho_L r}{4\pi \epsilon_0 (r^2 + z^2)^{3/2}} z \vec{a}_z (\phi)_0^{2\pi}$$

$$= \frac{\rho_L r z (2\pi)}{4\pi \epsilon_0 (r^2 + z^2)^{3/2}} \vec{a}_z$$

$$= \frac{\rho_L r z}{2 \epsilon_0 (r^2 + z^2)^{3/2}} \vec{a}_z$$

where

r = Radius of the ring

z = Perpendicular distance of point 'p' from the ring along the axis of the ring.