

The Tensor Renormalization Group

- A efficient numerical method

The Tensor Renormalization Group

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numerical methods

1

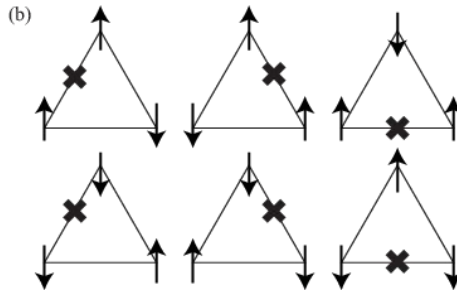
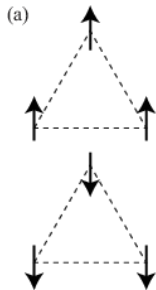
Exact diagonalization

Deal with small size systems

3

Quantum monte carlo

Sign problem [fermion systems and frustrated spin systems]



2

density matrix renormalization group

1-D system

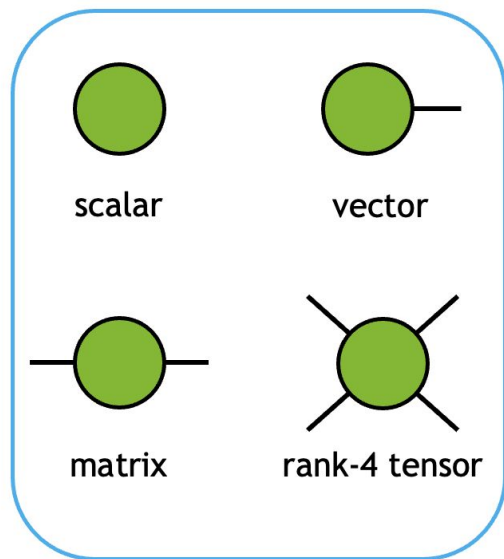


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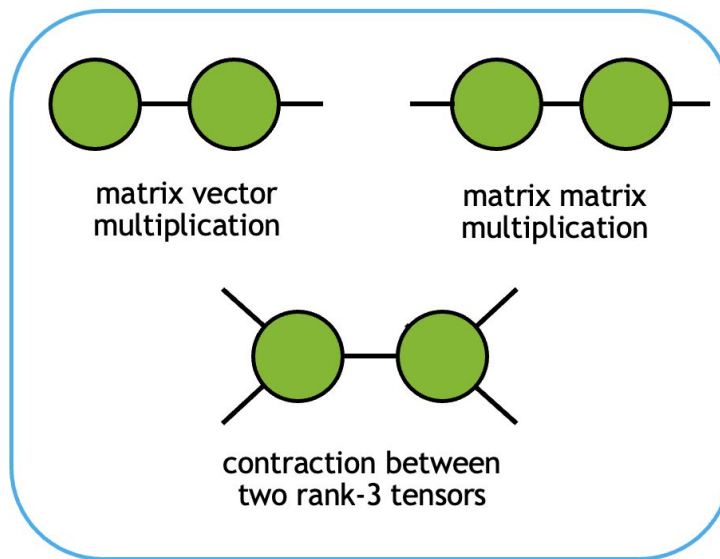
Tensor renormalization group

M. Levin and C. P. Nave. Tensor renormalization group approach to two-dimensional classical lattice models. Phys. Rev. Lett., 99:120601, Sep 2007.

What is tensor ?



tensor diagrams

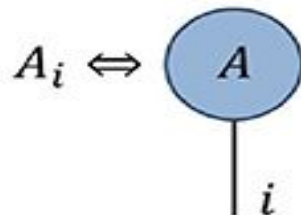


tensor contraction diagrams

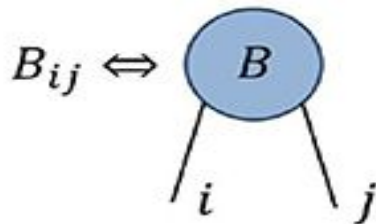
What is tensor ?

A

$$A = \begin{bmatrix} A_1 \\ A_2 \\ \vdots \\ A_m \end{bmatrix}$$

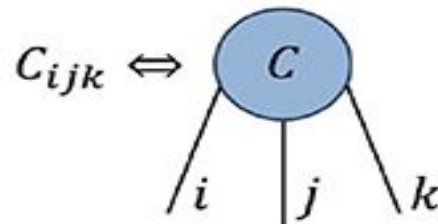
**B**

$$B = \begin{bmatrix} B_{11} & \cdots & B_{1n} \\ \vdots & \ddots & \vdots \\ B_{m1} & \cdots & B_{mn} \end{bmatrix}$$



C

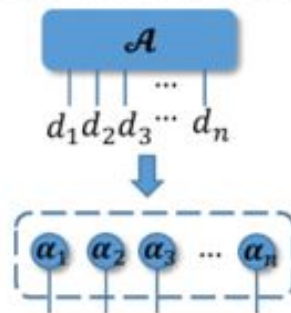
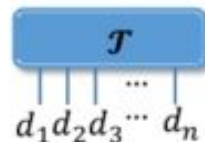
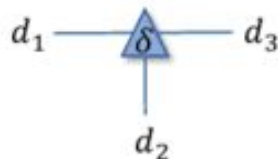
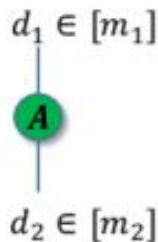
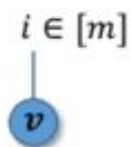
$$C = \left[\begin{matrix} C_{111} & \dots & C_{1n1} \\ \vdots & \ddots & \vdots \\ C_{m11} & \dots & C_{mn1} \end{matrix} \right]^1 \Bigg]^3$$



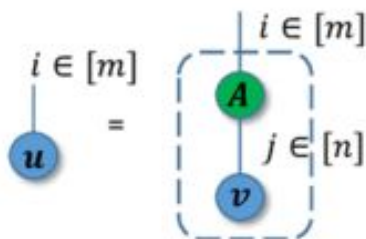
D

$$F_{ik} = \sum_j D_{ij} E_{jk} \Leftrightarrow \text{---}_i \bigcirc^F \text{---}_k = \text{---}_i \bigcirc^D \text{---}_j \bigcirc^E \text{---}_k$$

- 1) Vector v : 2) Matrix A : 3) 3-order δ tensor: 4) n -order tensor \mathcal{T} : 5) n -order rank-one tensor \mathcal{A} :



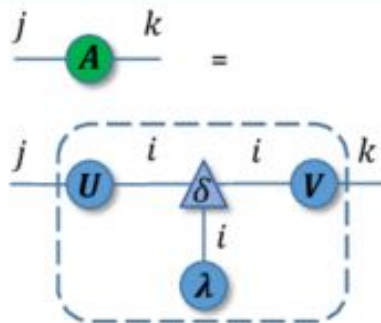
(a)



$$u = Av$$

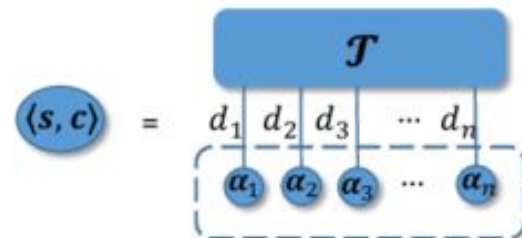
$$u_i = \sum_{j=1}^n A_{ij} v_j$$

(b)



$$A = \sum_{i=1}^r \lambda_i u_i \otimes v_i$$

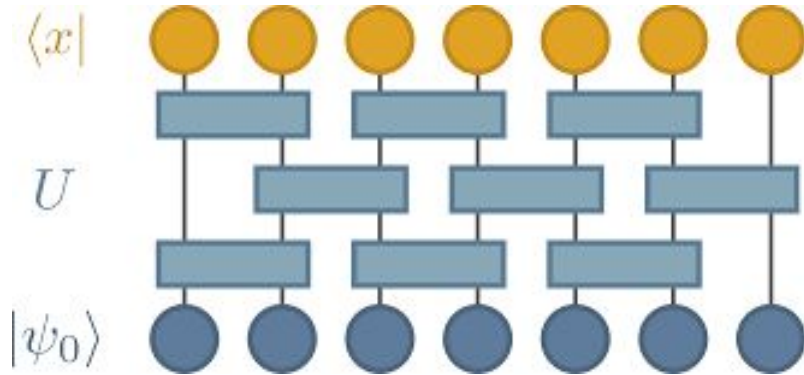
(c)



$$\langle s, c \rangle = \sum_{d_1, \dots, d_n=1}^m \mathcal{T}_{d_1 \dots d_n} \mathcal{A}_{d_1 \dots d_n}$$

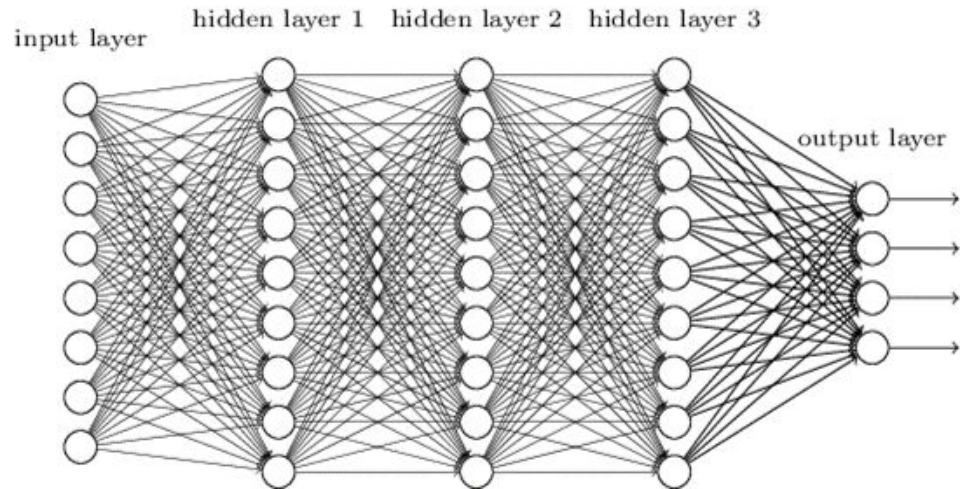
(d)

tensor network

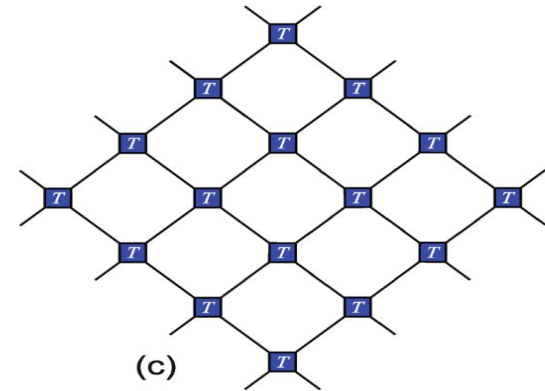
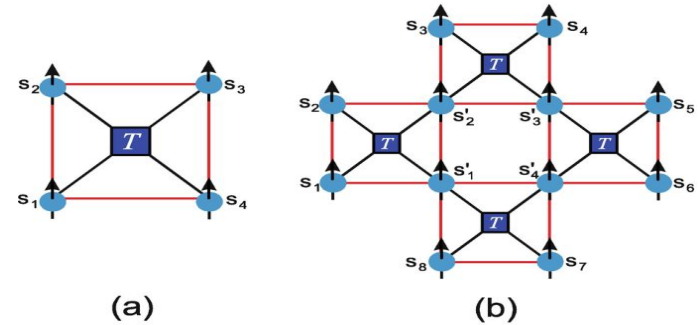
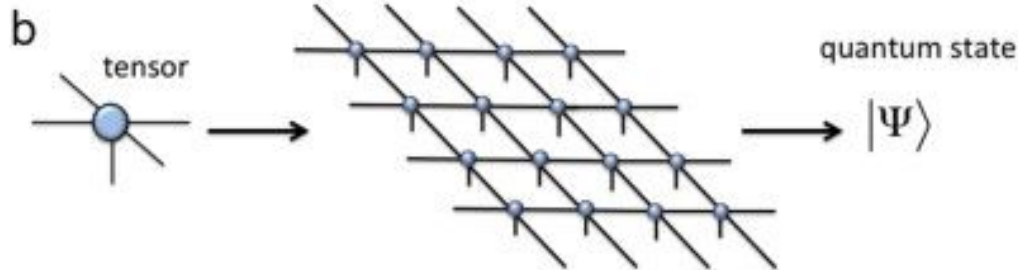
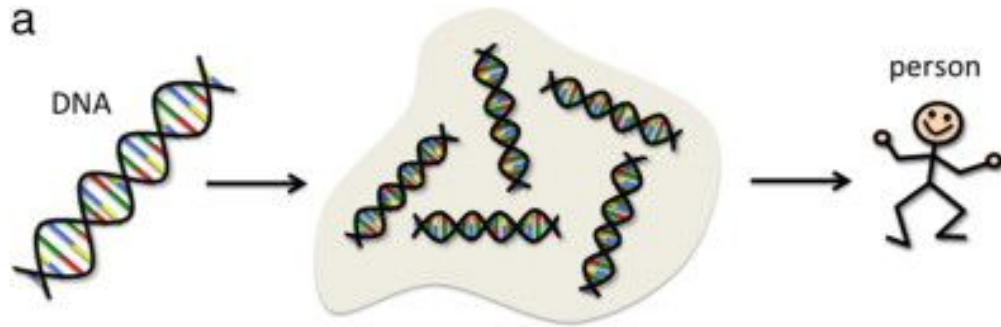


Quantum circuit

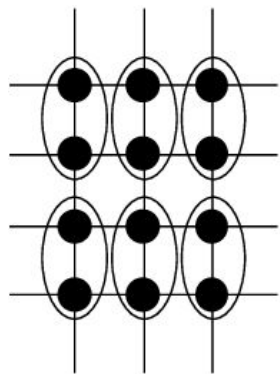
Deep neural network



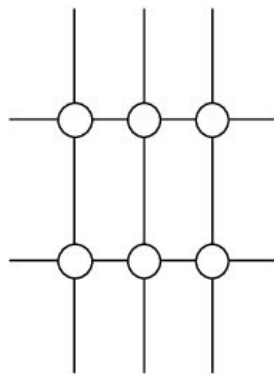
Tensor network represent quantum state and hamiltonian



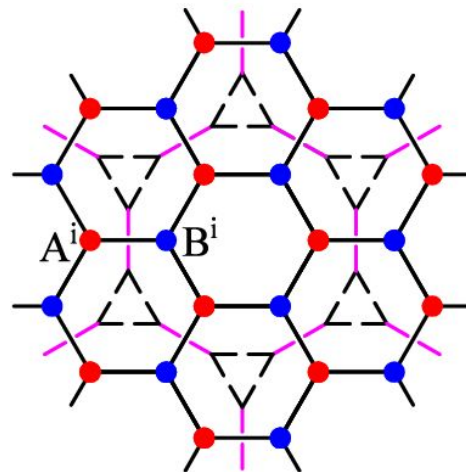
Examples of TRG



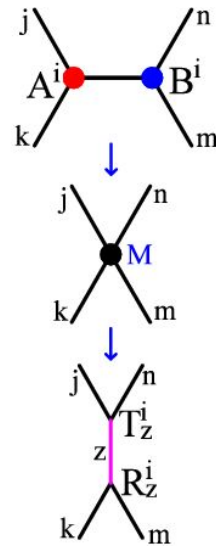
$\bullet T$



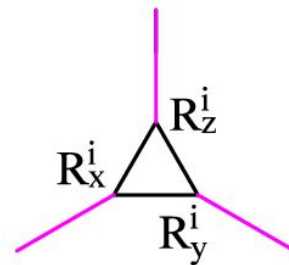
$\circ \mathcal{T}^{(\text{new})}$



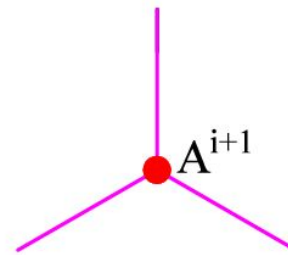
(a)



(b)



(c)



(d)

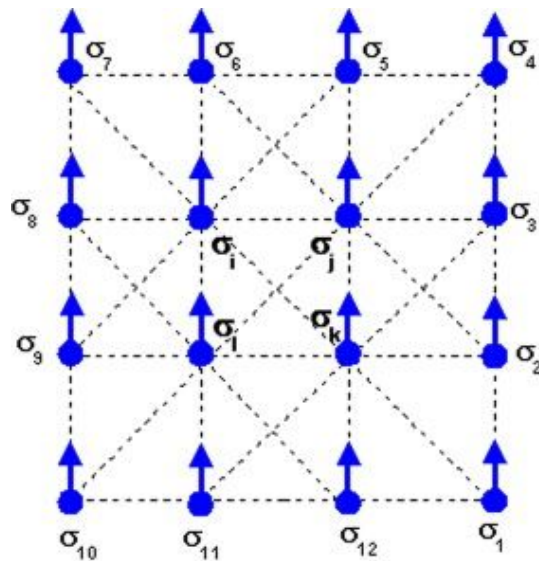


TRG on a square lattice Ising model

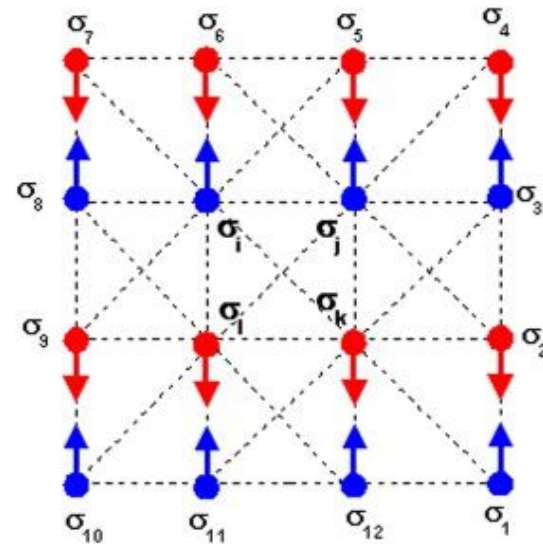


Ising model

$$H = -J \sum_{\langle ij \rangle} s_i s_j$$



(a) ferromagnetic state

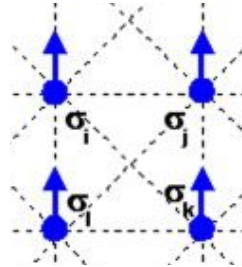


(b) superantiferromagnetic state

$$Z = \text{Tr} e^{-\beta H} = \text{Tr} \prod_{\square_{ijkl}} e^{\beta J (s_i s_j + s_j s_k + s_k s_l + s_l s_i) / 2}$$

Switch spin operator to bond variables

$$\sigma_{ij} = s_i s_j = \pm 1$$



$$Z = \text{Tr} \prod_{\langle ij \rangle} \delta(\sigma_{ij} - s_i s_j) \prod_{\square_{ijkl}} e^{\beta J(\sigma_{ij} + \sigma_{jk} + \sigma_{kl} + \sigma_{li})/2}$$

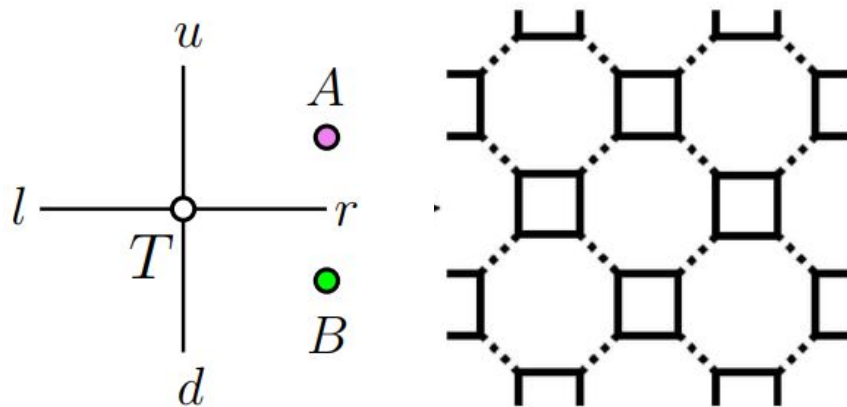
The product of bonds around a loop has value -1 is impossible.

$$Z = \text{Tr} \prod_{\square_{ijkl}} \frac{1 + \sigma_{ij} \sigma_{jk} \sigma_{kl} \sigma_{li}}{2} e^{\beta J(\sigma_{ij} + \sigma_{jk} + \sigma_{kl} + \sigma_{li})/2}$$

How to represent partition function as tensor network

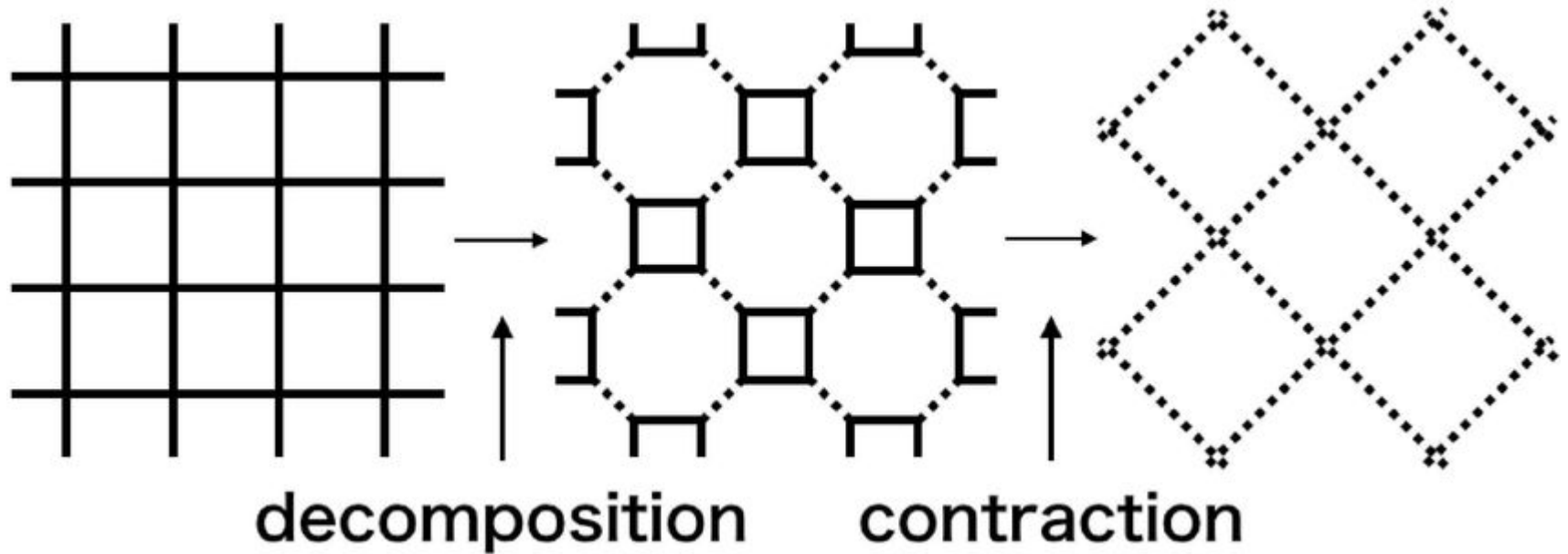
$$Z = \text{Tr} \prod_{\square_{ijkl}} \frac{1 + \sigma_{ij}\sigma_{jk}\sigma_{kl}\sigma_{li}}{2} e^{\beta J(\sigma_{ij} + \sigma_{jk} + \sigma_{kl} + \sigma_{li})/2}$$

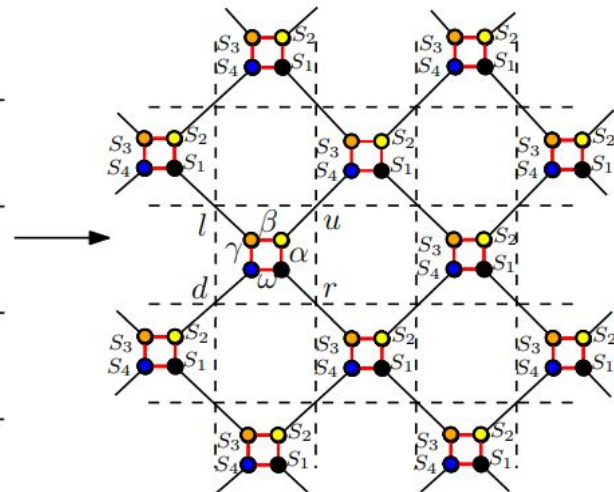
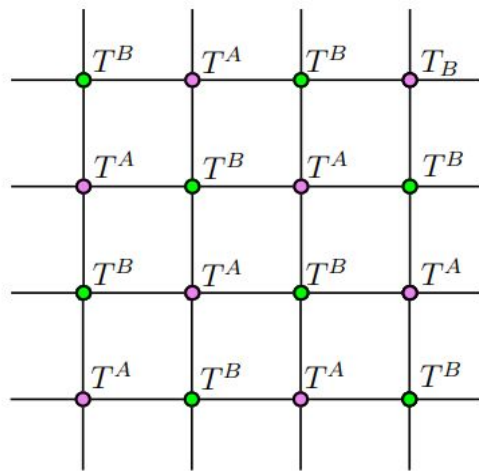
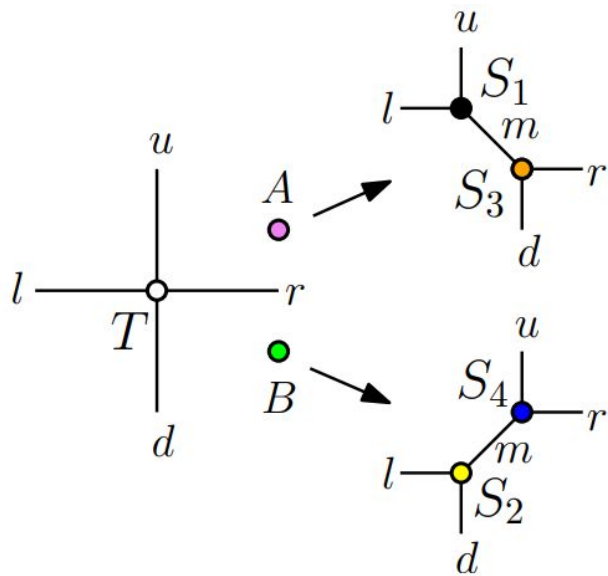
$$Z = \text{tTr} \left[\bigotimes_{p=1}^N T \right]$$



$$T_{r_p u_p l_p d_p} = \frac{1 + \sigma_r^p \sigma_u^p \sigma_l^p \sigma_d^p}{2} e^{\beta J(\sigma_r^p + \sigma_u^p + \sigma_l^p + \sigma_d^p)/2}$$

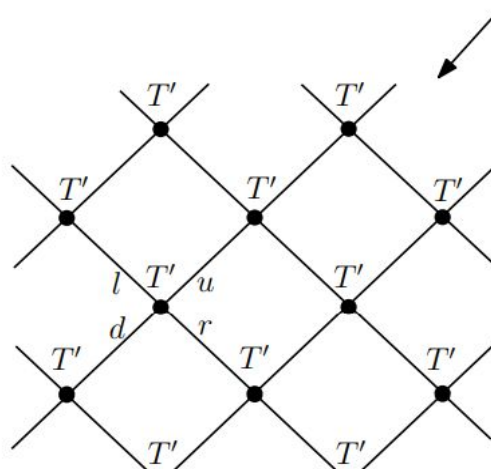
Tensor renormalization group of square lattice



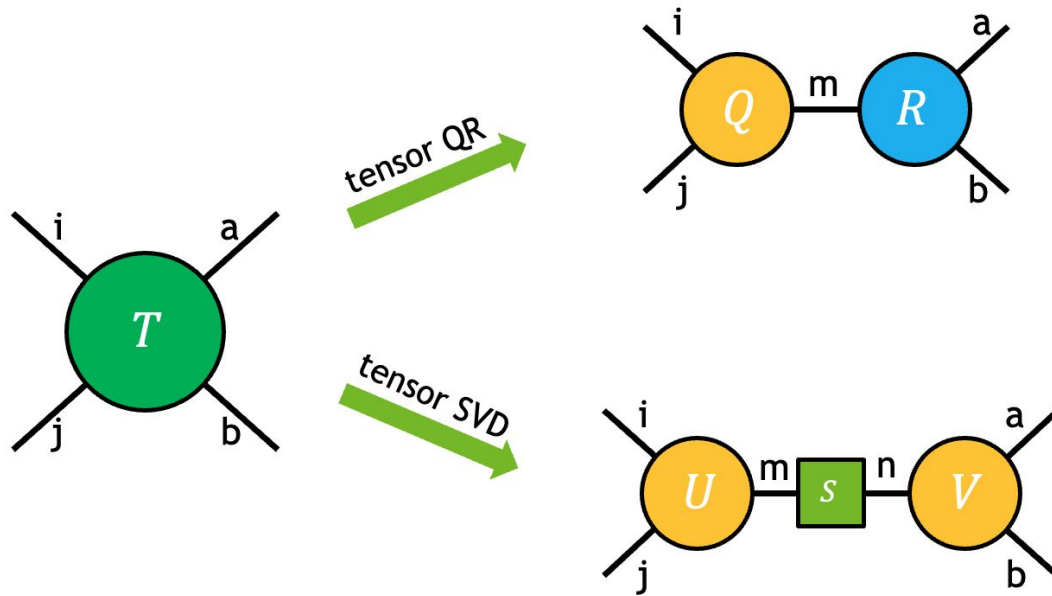


$$T_{ruld}^A = \sum_{m=1}^{D^2} S_{1,ulm} S_{3,drm}$$

$$T_{ruld}^B = \sum_{m=1}^{D^2} S_{2,l dm} S_{4,rum}$$

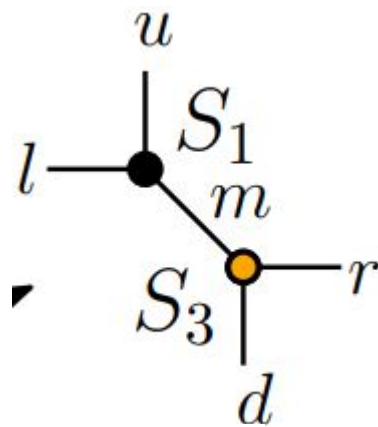
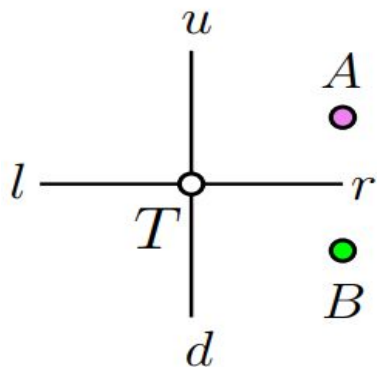


SVD decomposition

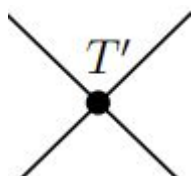


$$M = U\Sigma V^\dagger$$

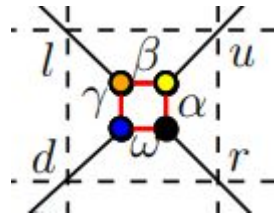
$$T_{ruld}^A = M_{lu,rd} \quad D^2 \times D^2 \text{ matrix} \quad T_{ruld}^A = \sum_{m=1}^{D^2} S_{1,ulm} S_{3,drm}$$



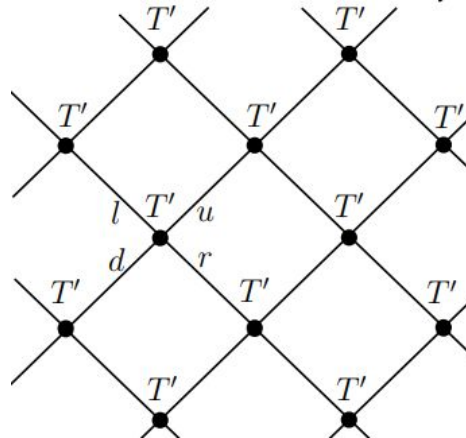
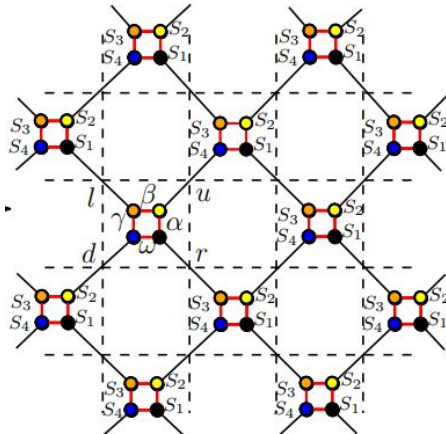
$$T' \quad T_{ruld}^A \approx \sum_{m=1}^{D'} \underbrace{\left(\sqrt{\lambda_m} U_{lu,m} \right)}_{S_{1,ulm}} \underbrace{\left(\sqrt{\lambda_m} V_{m,rd}^\dagger \right)}_{S_{3,drm}}$$

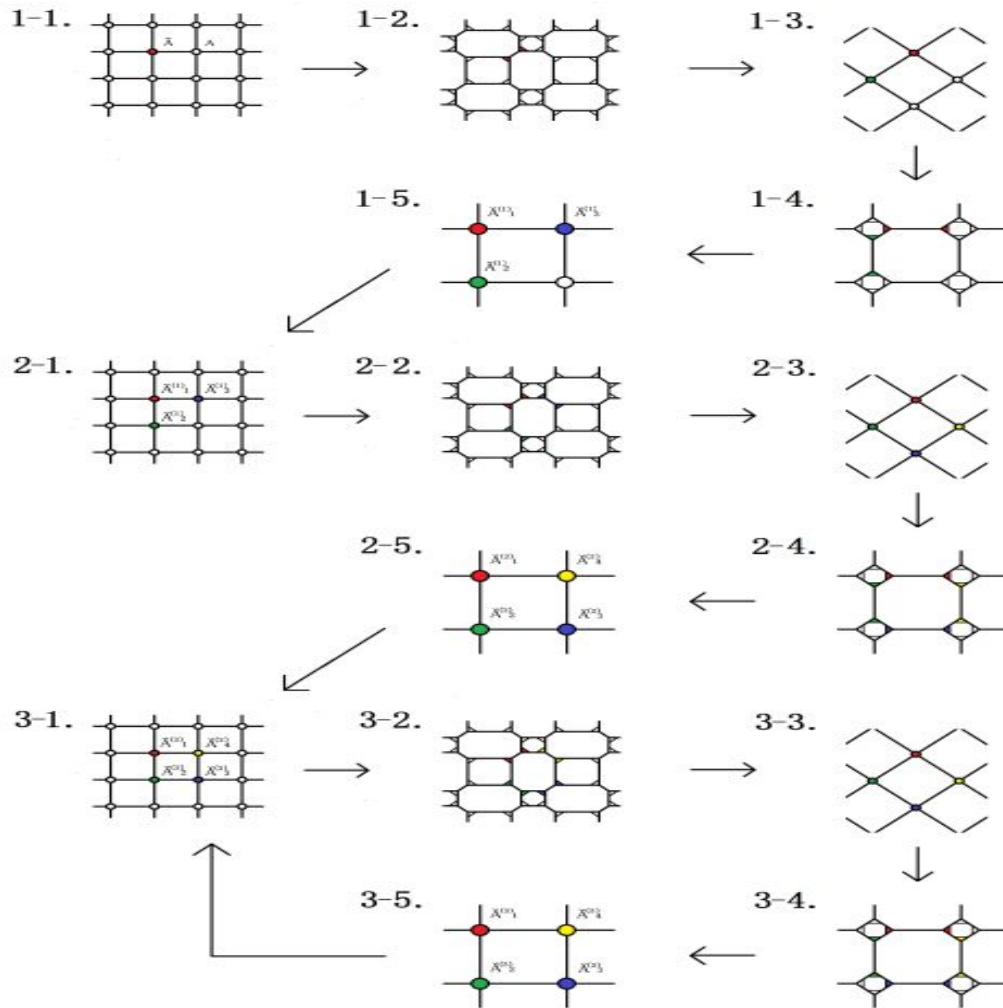


only kept a fixed number $D' = \min(D2, D_{\text{cut}})$ of the largest eigenvalues λ_m

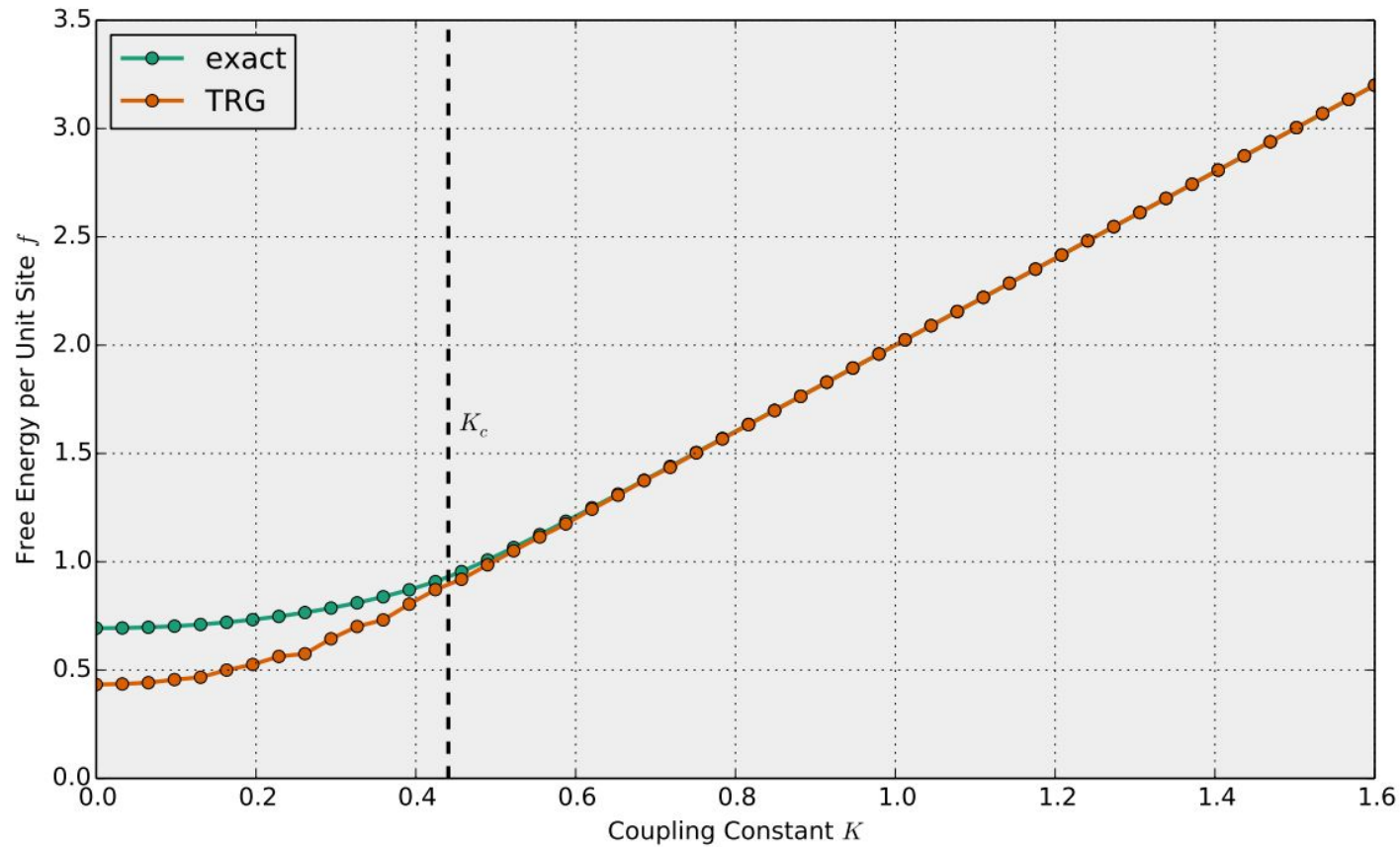


$$T'_{ruld} = \sum_{\alpha, \beta, \gamma, \omega=1}^D S_{1, \omega \alpha r} S_{2, \alpha \beta u} S_{3, \beta \gamma l} S_{4, \gamma \omega d}$$





Renormalization
Group flow



The TRG was applied 3 times with $D_{\text{cut}} = 6$. The accuracy of the TRG method is seen to decrease by several orders of magnitude as the critical point K_c is approached.

Limitation of TRG

Near criticality quantum states under the classical-quantum lattice model mapping become gapless ground states [more entangled than their gapped counterparts]

the truncated SVD decomposition failing to accurately represent the original tensor T near criticality

The 2D square lattice Ising model experiences a phase transition and therefore becomes critical at the value

$$K_c = \frac{1}{2} \ln (1 + \sqrt{2}) \approx 0.441$$

Improve TRG \rightarrow SRG

take into account the "environment" lattice when performing SVD
composition tensors node in the rewiring step

TRG : minimizes the truncation error of the local matrices M

SRG: minimizes the truncation error of entire partition function

– improve the accuracy of the TRG by several orders of magnitude

$$Z = \text{tTr} [MM_e]$$

Insight of renormalization group

Every partition function [local interaction] \rightarrow tensor network model

Exact sum of a general high-dimensional tensor network is a #P-complete problem

Transfer matrix VS coarse graining

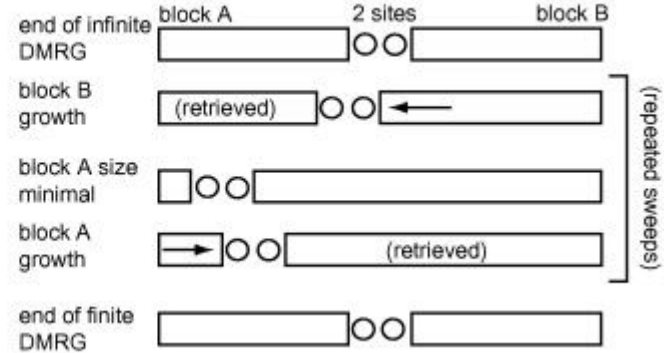
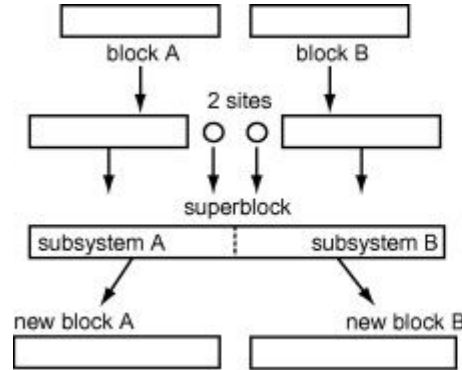
spin-blocking prescription – Kadanoff [1]

Wilson [2] the path to non-perturbative approaches based on coarse-graining a lattice

tensor renormalization group [3] a more general RG approach for classical lattice models

The tensor renormalization group may be seen as a generalization of the density matrix renormalization group (DMRG) method, introduced by White [4] to study **the ground state of Heisenberg spin chains.**

DMRG



matrix product states + Wilsonian RG

Steps: dividing the system into blocks and truncating basis states at each step

an error-minimizing way by maximizing the entanglement entropy at each step

λ are the the eigenvalues of the reduced (post-truncation) density matrix.

$$S = - \sum_l \lambda_l \ln \lambda_l$$

lack of "minus sign" problem.

Why truncation can be precise ?

Although approximate due to these truncations, the DMRG is an extraordinarily precise method in one dimensional systems.

the surface of a lattice model contains just two points and does not grow with system size.

In higher dimensions, exponential growth of the matrix dimension at each iteration is required to faithfully represent a quantum state, making the algorithm intractable.

Efficiency of dmrg

the surface of a lattice model contains just two points and does not grow with system size.

As the entanglement entropy scales with this surface area, [area law] only a small matrix dimension is necessary for a matrix product state to accurately represent a quantum state.

In higher dimensions, exponential growth of the matrix dimension at each iteration is required to faithfully represent a quantum state, making the algorithm intractable.



Thank you