

## Quantum dots

- electrons that are in trapped in 2-dim systems.
- Trapped in harmonic oscillator-like potentials.
- $N = 2, 6, 12, 20$

$$\hat{H} = \sum_{i=1}^N \left( -\frac{1}{2} \nabla_i^2 + \frac{1}{2} \omega_i^2 r_i^2 \right) \quad \text{one-body operator}$$

$$+ H_I$$

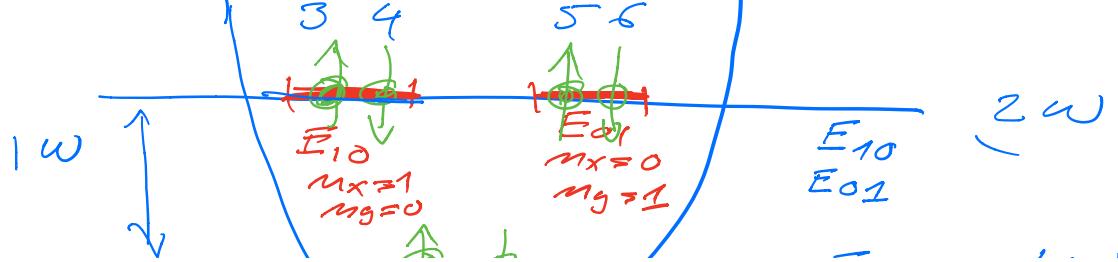
$$H_I = \sum_{\substack{i,j \\ i < j}} \frac{1}{|\vec{r}_i - \vec{r}_j|} = \sum_{i < j} \frac{1}{r_{ij}}$$

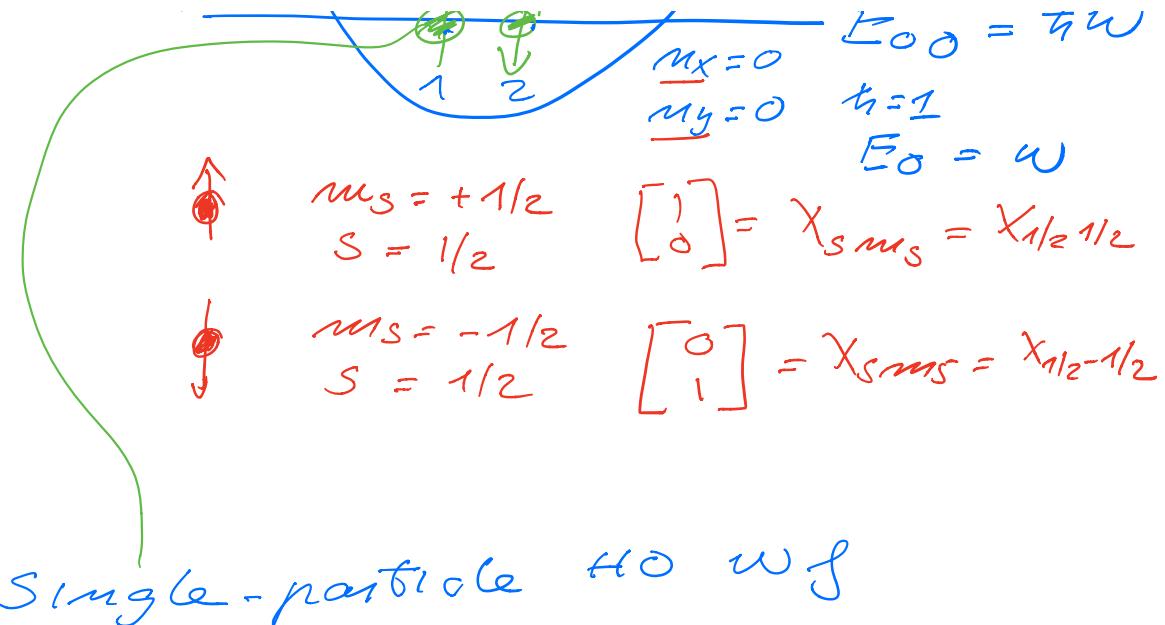
$$\hbar = m = c = \epsilon = 1$$

N = 6 1D system

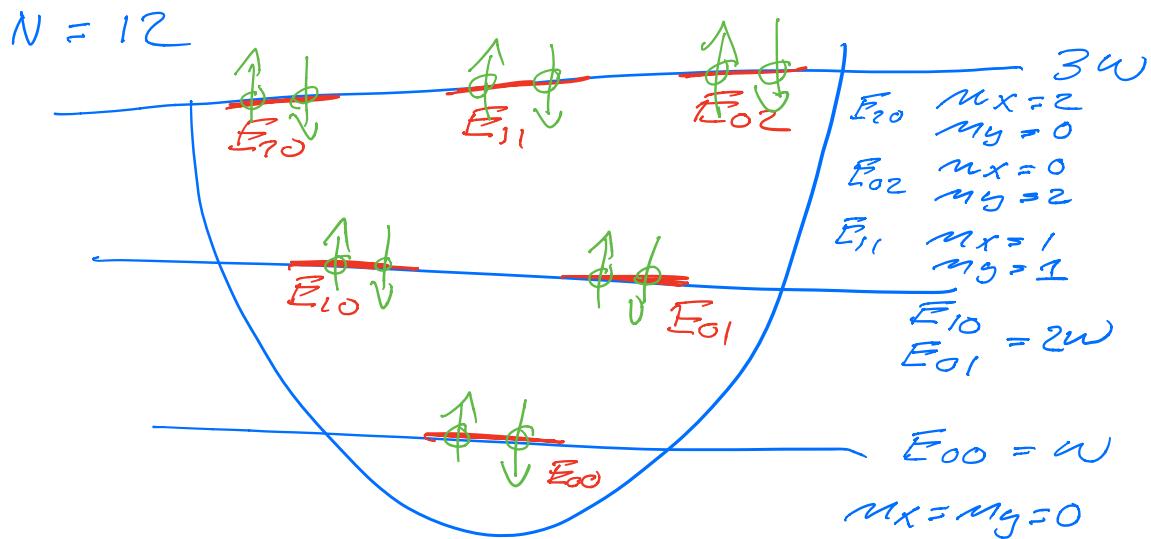
1-particle  $E_{mx,mg} = \hbar \omega (m_x + m_g + 1)$   
2-dim

1D-potential





$$\Psi_i = \underbrace{\varphi_{m_x, m_y}}_{\text{HO wf}} \underbrace{\chi_{S, m_s}}_{1/2}$$



$E_0$  total energy for  $N$  particles

$$\underline{N = 6}$$

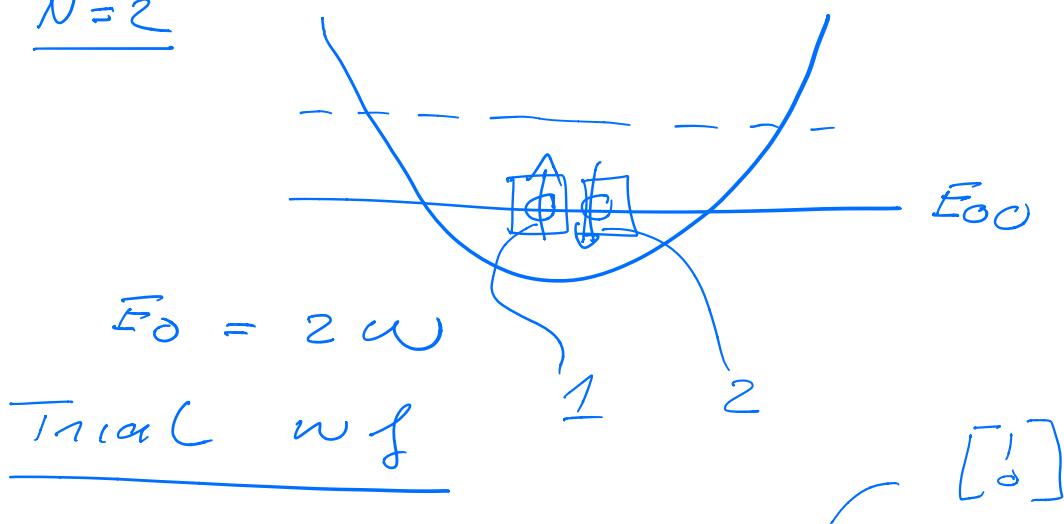
$$E_0 = \frac{2\omega}{(E_{00} \times 2)} + 4 \times 2\omega = 10\omega$$

useful test for our code  
with fermions.

$$N=12$$

$$\bar{E}_0 = 2\omega + 8\omega + 18\omega = 28\omega$$

$$\underline{N=2}$$



$$\psi_1(\vec{r}_i, s_i) = \varphi_{00}(\vec{r}_i) \chi_{1/2, 1/2}(s_i)$$

$$r'_i = \sqrt{x_i^2 + y_i^2} \quad \vec{r}'_i = x_i \hat{\vec{e}}_x + y_i \hat{\vec{e}}_y$$

$$\psi_2(\vec{r}_i, s_i) = \varphi_{00}(\vec{r}_i) \chi_{1/2, -1/2}(s_i)$$

$$\Psi(\vec{r}_1, \vec{r}_2; s_1, s_2) = \frac{1}{\sqrt{2!}} \begin{vmatrix} \psi_1(1) & \psi_1(2) \\ \psi_2(1) & \psi_2(2) \end{vmatrix}$$

Slater determinant

$$\frac{1}{\sqrt{2}} (\psi_1(1)\psi_2(2) - \psi_1(2)\psi_2(1))$$

$$\psi(\vec{r}_1 \vec{r}_2; s_1 s_2) = \boxed{\psi(r_1, z) = -\psi(z, r)}$$

our ansatz

$$= \frac{1}{\sqrt{2}} \left( \frac{\varphi_{00}(r_1) \varphi_{00}(z) \chi_{\uparrow}(r_1) \chi_{\downarrow}(z) - \varphi_{00}(z) \varphi_{00}(r_1)}{+0} \times \chi_{\uparrow}(z) \chi_{\downarrow}(r_1) \right)$$

$$S=0$$

$$\frac{1}{\sqrt{2}} \left[ \begin{array}{c} \chi_{\uparrow}(r_1) \chi_{\downarrow}(z) \\ - \chi_{\uparrow}(z) \chi_{\downarrow}(r_1) \end{array} \right]$$

6 electrons  $6 \times 6$  det,

12 —  $\rightarrow$   $12 \times 12$  det

$$\text{trial function } \varphi_{00}(r) \propto e^{-\frac{1}{2} \alpha^2 (x_r^2 + y_r^2)}$$

Hermite polynomial

$$H_m(x_i) H_{m'}(y_i)$$

$$m_x = m_y = 0 \Rightarrow f_0(i) = 1$$

- spatial part  $(\varphi_{00}(r) \varphi_{00}(z))$  is symmetric

- spin part is antisymmetric

$$S=0 \quad M_S = m_{S_1} + m_{S_2} = 0$$

- Hamiltonian is spin independent,

$$\langle \Psi' | H | \Psi \rangle = 0$$

$\downarrow S' \neq S'$

$$H = H_0 + H_1$$

$$H_0 = \sum_{i=1}^2 \left( -\frac{1}{2} \partial_i^2 + \frac{1}{2} \omega_i^2 \alpha_i^2 \right) \quad x_1 = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$

$$\frac{1}{2} \left[ \underbrace{(x_1^*(1)x_0^*(2) - x_1^*(2)x_0^*(1))}_{\times H_0} \right] \quad x_0 = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

$$\times H_0 \left[ \underbrace{(x_1(1)x_0(2) - x_1(2)x_0(1))}_{\times H_0} \right]$$

$$\frac{1}{2} \left[ \underbrace{\begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}}_{1} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} - \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}}_{2} \right] = 0$$

$$- \frac{1}{2} \cdot 2 + \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} = 0$$

$$\frac{1}{2} \cdot 2 = 1$$

How to deal with a  
determinant efficiently  
in our VMC code?

$$\Psi_{...} = \Psi_n \times \Psi_l$$

more

$\begin{smallmatrix} \downarrow \\ \text{anti} \\ \text{symmetry} \end{smallmatrix}$

$\begin{smallmatrix} \curvearrowleft \\ \curvearrowright \end{smallmatrix}$  two  
body  
correlators

$$D = \begin{pmatrix} d_{11} & d_{12} & \dots & d_{1N} \\ d_{21} & -d_{22} & & \\ \vdots & & \ddots & \\ d_{N1} & \dots & \dots & d_{NN} \end{pmatrix}$$

$D \in \mathbb{R}^{N \times N}$  to obtain the  
determinant  $\mathcal{O}(N^3)$  FLOPs  
when moving  $N$ -particles  
 $\sim \mathcal{O}(N^4)$  FLOPs

- Metropolis ratio

$$\frac{|\psi_D(\vec{r}^{\text{new}}) \psi_J(\vec{r}^{\text{new}})|^2}{|\psi_D(\vec{r}^{\text{old}}) \psi_J(\vec{r}^{\text{old}})|^2}$$

Slater determinant = SD

$$\psi_{SD}(\vec{r}_1, \vec{r}_2, \vec{r}_3, \dots, \vec{r}_N; s_1, s_2, s_3, \dots, s_N)$$

$$= \frac{1}{N!} \begin{vmatrix} \varphi_1(1) & \dots & \varphi_1(N) \\ \varphi_2(1) & & \varphi_2(N) \\ \vdots & \varphi_i(j) & \vdots \\ \varphi_N(1) & & \varphi_N(N) \end{vmatrix}$$

columns = given particle  
 row  $s$  = given single  
 particle quantum  
 state.

$$d_{ij} = \varphi_i(\vec{r}_{ij,s})$$

$3 \times 3$  det:

$$\left| \begin{array}{ccc} d_{11} & d_{12} & d_{13} \\ d_{21} & d_{22} & d_{23} \\ d_{31} & d_{32} & d_{33} \end{array} \right|$$

$$= d_{11} \left| \begin{array}{cc} d_{22} & d_{23} \\ d_{32} & d_{33} \end{array} \right| - d_{21} \left| \begin{array}{cc} d_{12} & d_{13} \\ d_{32} & d_{33} \end{array} \right|$$

*move particle 1*

and

$$= \sum_{i=1}^n d_{ij} \cdot c_{ij}$$

$$c_{ij} = (-1)^{i+j} M_{i,j}^{(1)}$$

Minor

$$\overline{D}^{-1} D = \underline{1}$$

$$\sum_{k=1}^N d_{ik} \overline{d_{kj}} = \delta_{ij}$$

$$d_{ij}^{-1} = \frac{c_{j|i}}{|D|}$$

### Metropolis Test

$$R = \frac{\psi_{SD}(\alpha^{\text{new}})}{\psi_{SD}(\alpha^{\text{old}})}$$

$$= \frac{\sum_{j=1}^n d_{ij}(\alpha^{\text{new}}) c_{ij}(\alpha^{\text{new}})}{\sum_{j=1}^n d_{ij}(\alpha^{\text{old}}) c_{ij}(\alpha^{\text{old}})}$$

moving particle  $j'$  from  
 $\alpha_j^{\text{old}}$  to  $\alpha_j^{\text{new}}$  leaves

$c_{ij}$  unchanged.

$$c_{ij}(\alpha^{\text{old}}) = c_{ij}(\alpha^{\text{new}})$$

$$\frac{c_{ij}}{|D|} = d_{j|i}^{-1}$$

$$R = \frac{\sum_{j=1}^n d_{ij}(\alpha^{\text{new}}) d_{j|i}^{-1}(\alpha^{\text{old}})}{\sum_{j=1}^n d_{ij}(\alpha^{\text{old}}) d_{j|i}^{-1}(\alpha^{\text{old}})} = \delta_{ij}$$

$$= \sum_{j=1}^m d_{ij} (\alpha_{\text{new}}) \frac{d_{j,i}^{-1} (\alpha_{\text{old}})}{}$$

Next simplification

$$\frac{D_i |\psi_{\text{SD}}|}{|\psi_{\text{SD}}|} = \sum_{j=1}^N D_i \psi_j(\vec{\alpha}_i) \cdot d_{j,i}^{-1}(\vec{\alpha}_i)$$

Final simplification

$$\frac{D_i^2 |\psi_{\text{SD}}|}{|\psi_{\text{SD}}|}$$