

FYS 4411/Jan

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Reminder from last week

$$w_i(t) = \sum_j w(j \rightarrow i) w_j(t')$$

$$\lim_{t \rightarrow \infty} w_i(t) = w_i$$

Transition probability

$$\sum_j \underbrace{w(j \rightarrow i)}_{\text{unknown}} = 1$$

$$w(j \rightarrow i) = \frac{\overline{T}(j \rightarrow i)}{\text{Prob for a specific transition}} \frac{A(j \rightarrow i)}{\text{acceptance prob}}$$

Detailed balance

$$\underline{A(j \rightarrow i)}$$

$$A(i \rightarrow j)$$

$$= \begin{bmatrix} w_i \\ w_j \end{bmatrix} \begin{bmatrix} \tau(i \rightarrow j) \\ \tau(j \rightarrow i) \end{bmatrix}$$

Known

$$\tau(j \rightarrow i) = \tau(i \rightarrow j)$$

Metropolis

$$w_i \rightarrow \frac{|\psi_j(\vec{r}_i; \vec{\alpha})|^2}{\int_{\vec{R} \in D} d\vec{R} |\psi_j(\vec{r}_j; \vec{\alpha})|^2}$$

aside: computational techniques.

(i) move one particle at the time

Ex no in d-dim and 3 particles-

$$\frac{w_i}{w_j} = \frac{e^{-\alpha^2(r_1^2 + r_2^2 + r_3^2)}}{e^{-\alpha^2(r_1^2 + r_2^2 + r_3^2)} e^{-\alpha^2(r_1'^2 - r_1^2)}} \\ = e^{\frac{-\alpha^2(r_1'^2 + r_2^2 + r_3^2 - r_1^2 - r_2^2 - r_3^2)}{}} \\ (ii) move all$$

Importance Sampling =
Model for $\tau(i \Rightarrow j)$, based
on $\psi_i(\vec{r}_i; \vec{\alpha})$
Metropolis Test

$$A(j \Rightarrow i) = \min\left(1, \frac{w_i \tau(i \Rightarrow j)}{w_j \tau(j \Rightarrow i)}\right)$$

Basics of Monte Carlo (Stochastic)
Integration.

$$\bar{I} = \int_0^1 f(x) dx$$

uniform distribution - U

$$X \sim U$$

$$P(X)dx = \begin{cases} \frac{1}{b-a} & 0 \leq x \leq 1 \\ 0 & \text{else} \end{cases}$$



$M = \# \text{ of samples}$

$M = 1000$, with uniform PDF

$0 \rightarrow 0.1$; $0.1 \rightarrow 0.2$, ...
100 100 - - -

$$I = \int f(x) dx = \int \rho(x) f(x) dx$$

$$\left(\int_0^1 \rho(x) dx = 1 \right) \quad = E[f(x)]$$

$$\int_{-\infty}^{\infty} \rho(x) dx = 1 \quad 0 \leq \rho(x) \leq 1$$

$$\int_{-\infty}^a \rho(x) dx \leq \int_{-\infty}^b \rho(x) dx$$

$$a \leq b$$

(i) change of variable
[RNG are based on uniform
PDF]

(ii) importance sampling

change of variables

Ex 1 $P(y)dy = e^{-y}dy \quad y \in [0, \infty)$

conservation of probability

$$P(y)dy = P(x)dx$$

$$\underbrace{p'(x) dx}_{\text{uniform}} = dx \quad x \in [c_1, 1]$$

$$p(y) dy = e^{-y} dy = dx \\ y \in [c_1, \infty)$$

$$x(y) = \int_0^y e^{-y'} dy' \\ = 1 - e^{-y} \Rightarrow$$

$$y(x) = -\ln(1-x)$$

$y=0$ when $x=0 \wedge y=\infty$ when $x=1$

Ex 2

$$P(x) = \frac{1}{3}(4-2x) \quad x \in [0, 1]$$

$$\int_0^1 P(x) dx = 1$$

$P(y) dy$ is a uniform PDF
 $y \sim U \quad y \in [0, 1]$

$$y(x) = \frac{1}{3}(4-x)x \Rightarrow x(y)$$

$y=0$ when $x=0$

$y=1$ when $x=1$

why ?

(ii) importance sampling

$$\bar{I} = \int_0^1 \frac{dx}{1+x^2} = \frac{\pi}{4}$$

$$f(x) dx$$

$$= \int_0^1 p(x) \frac{f(x)}{p(x)} dx$$

$$p(x) = \frac{1}{3}(4-2x)$$

$$\frac{f(0)}{p(0)} = \frac{f(1)}{p(1)} = \frac{3}{4}$$

(constant)

The ideal is that $\frac{f(x)}{p(x)} =$
constant for all x

$$I = \int_c^1 p(x) \left[\frac{f(x)}{p(x)} \right] dx$$

$F(x)$ if const

$$I = \int_0^1 p(x) F(x) dx = \text{const} \underbrace{\int_0^1 p(x) dx}_{=1}$$

$$E[F(x)] = \int_0^1 p(x) F(x) dx$$

$$\text{Var}[F(x)] = \int_0^1 p(x) [F(x) - E[F]]^2 dx$$

$$= \int_0^1 p(x) F(x) dx - (E[F])^2$$

↑ if const

with zero variance, then
STD is also zero

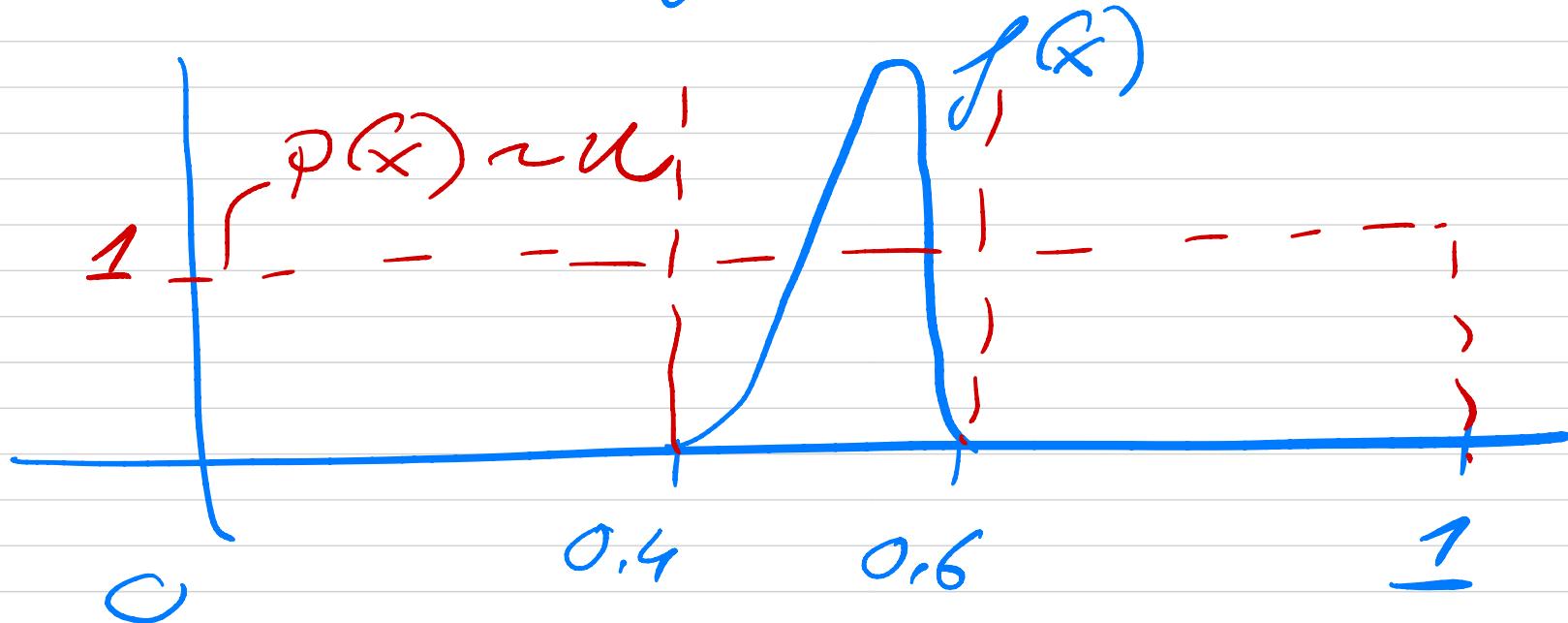
with a change of variable

$$P(Y)dy = dy = \frac{1}{3}(4-x)dx$$

$$\begin{aligned} I &= \int_c^1 P(X) \frac{f(x)}{P(X)} dx \\ &= \int_0^1 \frac{f(x(y))}{P(x(y))} dy \quad \text{RNG } \sim U \end{aligned}$$

$$P(X)dx = dy$$

Ex 3 $f(x)$ $x \in [\bar{g}_1, 1]$



$$\mathcal{I} = \int_0^1 f(x) dx \approx \frac{1}{M} \sum_{i=1}^M f(x_i)$$

$M = 1000$
only 200
are in
 $x \in [0.4, 0.6]$

ideal: find a PDF $p'(g)$

$g \in [0.4, 0.6]$ where $f(x)$
is nonzero [change of
variable post]

then have $\frac{f(g)}{p'(g)} \approx \text{const}$

in this case [importance
sampling]

$\text{var}[\hat{f}] \approx 0$ (small)

Ansatz for ψ_T

$$\tilde{E}_L(\vec{R}) = \frac{1}{\psi_T(\vec{R})} \hat{H} \psi_T(\vec{R})$$

electron in atomic hydrogen.

in spherical coordinates

$$\psi_{nl}(r, \theta, \varphi) = R_{nl}(r) Y_{lm}^{(e, \varphi)}$$

$$r \in [0, \infty)$$

$$\theta \in [0, \pi]$$

$$\varphi \in [0, 2\pi]$$

\hat{H} (in atomic units)

$$\hat{H} = -\frac{1}{2} \nabla^2 - \frac{1}{r}$$

Local energy for the radial part

$$E_L(r) = \frac{1}{R_T(r)} \left(-\frac{1}{2} \frac{d^2}{dr^2} - \frac{1}{r} \frac{d}{dr} - \frac{1}{r^2} \right) R_T(r) + \text{finite terms}$$

$$\lim_{r \rightarrow 0} \frac{d^n R_T}{dr^n} = L < \infty \quad (\epsilon, e)$$

$$\lim_{r \rightarrow 0} E_L(r) = \frac{1}{R_T} \left(-\frac{1}{2} \frac{d}{dr} - \frac{1}{2} \right) R_T$$

$$= 0$$

(omitted finite terms)

$$\frac{d}{dr} R_T = -R_T \Rightarrow$$

$$R_T = e^{-r}$$

Minimal requirement

$$R_{me}(r) = e^{-\alpha r} \underbrace{L_{me}(r)}_{\text{laguerre}}$$

For potentials like

$$\frac{1}{|\vec{r}_1 - \vec{r}_2|} = \frac{1}{\sqrt{(x_1 - x_2)^2 + (y_1 - y_2)^2 + (z_1 - z_2)^2}}$$

For two particles getting close to each other

$$\lim_{|\vec{r}_1 - \vec{r}_2| \rightarrow 0} = \infty \Rightarrow \text{term in}$$

relative kinetic energy

which cancels

$$\frac{1}{|\vec{r}_1 - \vec{r}_2|} \Rightarrow R_T(g_1, g_2) \sim e^{\beta R_{12}} \quad R_{12} = |\vec{r}_1 - \vec{r}_2|$$

Ansatz for atomic helium

$$\psi_1(\vec{r}_1, \vec{r}_2) \propto e^{-\alpha(r_1 + r_2)} \frac{\beta R_{12}}{r}$$

One body correlated part

Jastrow function

H0 for two electrons in 2d.

$$-\sqrt{2} \phi - m_x = m_y = 0$$

$$\psi_T(\vec{r}_1, \vec{r}_2) = e^{-\alpha^2(r_1^2 + r_2^2)} e^{\beta r_{12}}$$

$$\vec{r}_1 = \vec{x}_1 + \vec{y}_1$$

Variational
parameters