

1 Derivations from Maxwell Equations

$$\begin{aligned} \mathbf{B} &= \nabla \times \mathbf{A} \\ \mathbf{E} &= -\nabla \cdot \mathbf{v} - \frac{\partial \mathbf{A}}{\partial t} \end{aligned}$$

1.1 Gauss' Law

$$\begin{aligned} \nabla \cdot \mathbf{E} &= \frac{\rho}{\epsilon} \\ \nabla \cdot \left(-\nabla \cdot \mathbf{v} - \frac{\partial \mathbf{A}}{\partial t} \right) &= \frac{\rho}{\epsilon} \end{aligned}$$

The divergence is a linear operator :

$$-\nabla^2 \cdot \mathbf{v} - \nabla \cdot \frac{\partial \mathbf{A}}{\partial t} = \frac{\rho}{\epsilon}$$

We now multiply both terms by a test function X , and integrate it over a domain Ω :

$$\begin{aligned} - \int_{\Omega} \epsilon \cdot (\nabla^2 \cdot \mathbf{v}) \cdot X \cdot d\Omega - \int_{\Omega} \epsilon \nabla \cdot \frac{\partial \mathbf{A}}{\partial t} \cdot X \cdot d\Omega &= \int_{\Omega} \rho \cdot X \cdot d\Omega \\ - \int_{\Omega} \epsilon \cdot \nabla \mathbf{v} \cdot \nabla X \cdot d\Omega - \int_{\Omega} \epsilon \nabla \cdot \frac{\partial \mathbf{A}}{\partial t} \cdot X \cdot d\Omega &= \int_{\Omega} \rho \cdot X \cdot d\Omega \end{aligned}$$

1.1.1 In a conductor

In a conductor, a common approximation is the current density being proportional to the electric field, $J = \sigma \cdot E$. σ being the conductivity of the conductor.

The law equivalent to Gauss' law in conductors becomes :

$$\begin{aligned} \nabla \cdot \mathbf{J} &= 0 \\ \nabla \cdot \sigma \mathbf{E} &= 0 \\ \sigma \nabla \cdot \left(-\nabla \cdot \mathbf{v} - \frac{\partial \mathbf{A}}{\partial t} \right) &= 0 \end{aligned}$$

Thus,

$$- \int_{\Omega} \sigma \cdot \nabla \mathbf{v} \cdot \nabla X \cdot d\Omega - \int_{\Omega} \sigma \nabla \cdot \frac{\partial \mathbf{A}}{\partial t} \cdot X \cdot d\Omega = 0$$

1.1.2 In a conductor -DC

$$- \int_{\Omega} \sigma \cdot \nabla \mathbf{v} \cdot \nabla X \cdot d\Omega = 0$$

1.1.3 In a dielectric

There is no net charge in a dielectric. (That is, in a PCB)

$$-\int_{\Omega} \epsilon \cdot \nabla v \cdot \nabla X \cdot d\Omega - \int_{\Omega} \epsilon \nabla \cdot \frac{\partial A}{\partial t} \cdot X \cdot d\Omega = 0$$

1.1.4 In a dielectric - DC

$$\int_{\Omega} \epsilon \cdot \nabla v \cdot \nabla X \cdot d\Omega = 0$$

1.2 Gauss' Law for Magnetism

$$\nabla \cdot B = 0$$

$$\nabla \cdot \nabla \times A = 0$$

The divergence of a curl is 0, hence :

$$0 = 0$$

Therefore, we cannot get anything useful from this law

1.3 Lenz-Faraday's Law

$$\nabla \times E = -\frac{\partial B}{\partial t}$$

$$\nabla \times \left(-\nabla \cdot v - \frac{\partial A}{\partial t} \right) = -\frac{\partial (\nabla \times A)}{\partial t}$$

$$\nabla \times \nabla \cdot v - \frac{\partial (\nabla \times A)}{\partial t} = -\frac{\partial (\nabla \times A)}{\partial t}$$

$$\nabla \times \nabla \cdot v = 0$$

The curl of a gradient is 0, hence :

$$0 = 0$$

Therefore, we cannot get anything useful from this law

1.4 Ampere's Law

$$\nabla \times B = \mu \left(J + \epsilon \frac{\partial E}{\partial t} \right)$$

$$\nabla \times (\nabla \times A) = \mu \left(J + \epsilon \frac{\partial (-\nabla \cdot v - \frac{\partial A}{\partial t})}{\partial t} \right)$$

$$\frac{1}{\mu} \nabla \times (\nabla \times A) = J - \epsilon \frac{\partial (\frac{\partial A}{\partial t})}{\partial t} - \epsilon \frac{\partial (\nabla \cdot v)}{\partial t}$$

$$\frac{1}{\mu} \nabla \times (\nabla \times \mathbf{A}) = \mathbf{J} - \epsilon \frac{\partial^2 \mathbf{A}}{\partial t^2} - \epsilon \nabla \cdot \frac{\partial(\mathbf{v})}{\partial t}$$

We now multiply both terms by a test function \mathbf{X} , and integrate it over a domain Ω :

$$\int_{\Omega} \frac{1}{\mu} \nabla \times (\nabla \times \mathbf{A}) \cdot \mathbf{X} \cdot d\Omega = \int_{\Omega} \mathbf{J} \cdot \mathbf{X} \cdot d\Omega - \int_{\Omega} \epsilon \frac{\partial^2 \mathbf{A}}{\partial t^2} \cdot \mathbf{X} \cdot d\Omega - \int_{\Omega} \epsilon \nabla \cdot \frac{\partial(\mathbf{v})}{\partial t} \cdot \mathbf{X} \cdot d\Omega$$

$$\int_{\Omega} \frac{1}{\mu} \nabla \times (\nabla \times \mathbf{A}) \cdot \mathbf{X} \cdot d\Omega = \int_{\Omega} \mathbf{J} \cdot \mathbf{X} \cdot d\Omega - \int_{\Omega} \epsilon \frac{\partial^2 \mathbf{A}}{\partial t^2} \cdot \mathbf{X} \cdot d\Omega - \int_{\Omega} \epsilon \nabla \cdot \frac{\partial(\mathbf{v})}{\partial t} \cdot \mathbf{X} \cdot d\Omega$$

1.4.1 In a conductor - General case

In a conductor, a common approximation is the current density being proportional to the electric field, $\mathbf{J} = \sigma \cdot \mathbf{E}$. σ being the conductivity of the conductor.

$$\int_{\Omega} \frac{1}{\mu} \nabla \times (\nabla \times \mathbf{A}) \cdot \mathbf{X} \cdot d\Omega = - \int_{\Omega} \sigma \nabla v \cdot \mathbf{X} \cdot d\Omega - \int_{\Omega} \sigma \frac{\partial \mathbf{A}}{\partial t} \cdot \mathbf{X} \cdot d\Omega - \int_{\Omega} \epsilon \frac{\partial^2 \mathbf{A}}{\partial t^2} \cdot \mathbf{X} \cdot d\Omega - \int_{\Omega} \epsilon \nabla \cdot \frac{\partial(\mathbf{v})}{\partial t} \cdot \mathbf{X} \cdot d\Omega$$

We could even add that in a good conductor, there is little to no displacement current:

$$\int_{\Omega} \frac{1}{\mu} \nabla \times (\nabla \times \mathbf{A}) \cdot \mathbf{X} \cdot d\Omega = - \int_{\Omega} \sigma \nabla v \cdot \mathbf{X} \cdot d\Omega - \int_{\Omega} \sigma \frac{\partial \mathbf{A}}{\partial t} \cdot \mathbf{X} \cdot d\Omega$$

1.4.2 In a conductor - DC

$$\int_{\Omega} \frac{1}{\mu} \nabla \times (\nabla \times \mathbf{A}) \cdot \mathbf{X} \cdot d\Omega = - \int_{\Omega} \sigma \nabla v \cdot \mathbf{X} \cdot d\Omega$$

1.4.3 In a conductor - DC - Ignoring magnetic field

$$\int_{\Omega} \sigma \nabla v \cdot \mathbf{X} \cdot d\Omega = 0$$

1.4.4 In a dielectric

There cannot be any current flowing in a good dielectric, hence $\mathbf{J} = 0$

$$\int_{\Omega} \frac{1}{\mu} \nabla \times (\nabla \times \mathbf{A}) \cdot \mathbf{X} \cdot d\Omega = \epsilon \frac{\partial^2 \mathbf{A}}{\partial t^2} \cdot \mathbf{X} \cdot d\Omega - \int_{\Omega} \epsilon \nabla \cdot \frac{\partial(\mathbf{v})}{\partial t} \cdot \mathbf{X} \cdot d\Omega$$

1.4.5 In a dielectric - DC

$$\int_{\Omega} \frac{1}{\mu} \nabla \times (\nabla \times \mathbf{A}) \cdot \mathbf{X} \cdot d\Omega = 0$$