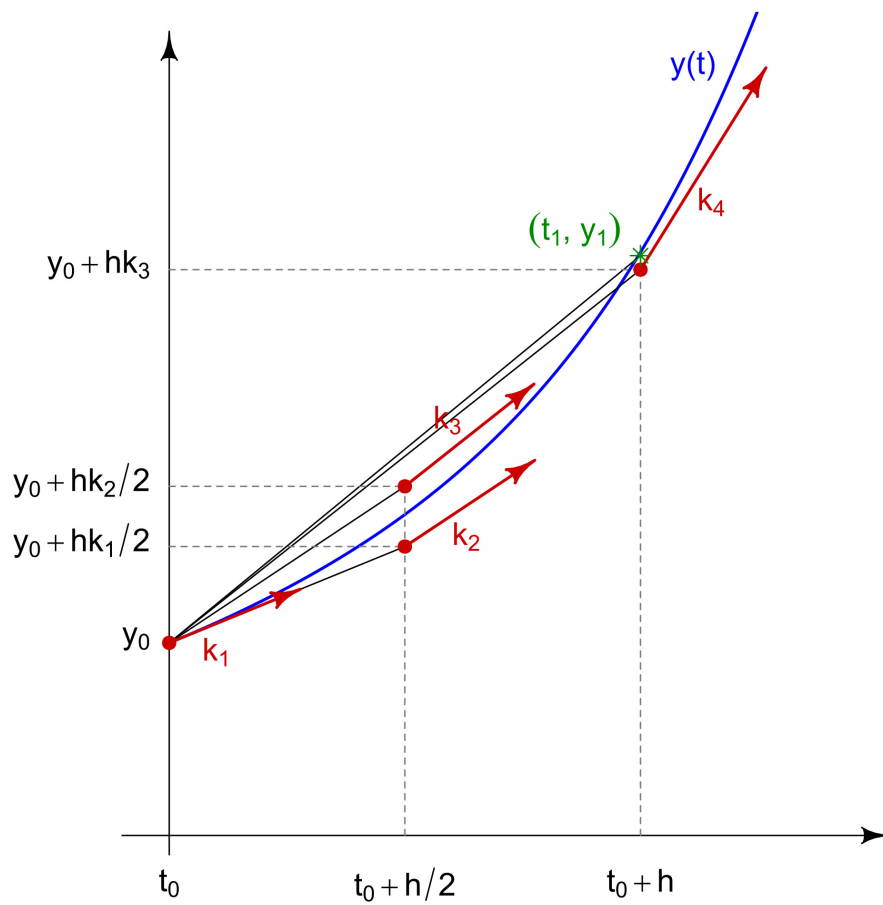


Final Project

Numerical Solutions to PDEs

Connor Emmons and Noah Wells

9 May 2025



1 Basic Heat Equation

$$\begin{aligned} u_t &= 2u_{xx} \\ u(0, t) &= u(3, t) = 0 \\ u(x, 0) &= \sin(\pi x) + \sin(2\pi x) \end{aligned}$$

First we must find the analytic solution to this equation. This is an extremely basic heat equation with neumann boundary conditions. The solution boils down to

$$\sum_{n=1}^{\infty} e^{-2\left(\frac{n\pi}{3}\right)^2 t} \sin\left(\frac{n\pi x}{3}\right)$$

which when given the initial conditions we see is equal to

$$e^{-2\pi^2 t} \sin(\pi x) + e^{-8\pi^2 t} \sin(2\pi x).$$

Using the central difference approximation for second derivatives gives:

$$\begin{aligned} \frac{du}{dt} &= 2 \frac{u_{i+1}^n - 2u_i^n + u_{i-1}^n}{\Delta x^2} \\ u(t_0) &= \sin(\pi x) + \sin(2\pi x) \end{aligned} \tag{1}$$

Note that the index i relates to the spatial steps and the index n relates to the time steps. The standard RK4 formulation is given as:

$$\begin{aligned} k_1 &= f(t_n, y_n) \\ k_2 &= f\left(t_n + \frac{h}{2}, y_n + \frac{h}{2}k_1\right) \\ k_3 &= f\left(t_n + \frac{h}{2}, y_n + \frac{h}{2}k_2\right) \\ k_4 &= f(t_n + h, y_n + hk_3) \\ y_{n+1} &= y_n + \frac{h}{6}(k_1 + 2k_2 + 2k_3 + k_4) \end{aligned} \tag{2}$$

Adapting this to the problem at hand gives:

$$\begin{aligned} k_1 &= \alpha \frac{u_{i+1}^{n-1} - 2u_i^{n-1} + u_{i-1}^{n-1}}{\Delta x^2} \\ k_2 &= \alpha \frac{\left(u_{i+1}^{n-1} + \frac{\Delta t}{2}k_1\right) - 2\left(u_i^{n-1} + \frac{\Delta t}{2}k_1\right) + \left(u_{i-1}^{n-1} + \frac{\Delta t}{2}k_1\right)}{\Delta x^2} \\ k_3 &= \alpha \frac{\left(u_{i+1}^{n-1} + \frac{\Delta t}{2}k_2\right) - 2\left(u_i^{n-1} + \frac{\Delta t}{2}k_2\right) + \left(u_{i-1}^{n-1} + \frac{\Delta t}{2}k_2\right)}{\Delta x^2} \\ k_4 &= \alpha \frac{\left(u_{i+1}^{n-1} + \Delta t k_3\right) - 2\left(u_i^{n-1} + \Delta t k_3\right) + \left(u_{i-1}^{n-1} + \Delta t k_3\right)}{\Delta x^2} \\ u_i^n &= u_i^{n-1} + \frac{\Delta t}{6}(k_1 + 2k_2 + 2k_3 + k_4) \end{aligned} \tag{3}$$

An example of the numerical solution plotted against the analytical solution is shown below. In order to view the full animation, please run the main.m file and select the corresponding option.

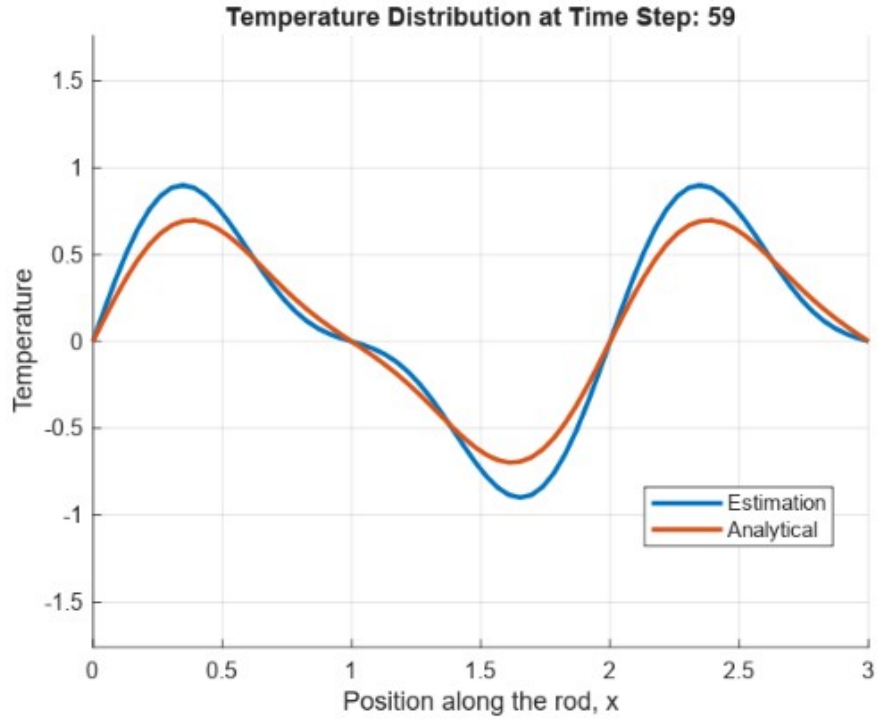


Figure 1: Iteration 59 of RK4 Solution to Heat Equation

With the proper function definition in MatLab, it is easy to use `ode45()` to obtain an equivalent solution. Note that for this comparison to work directly, the full vector of t values used by RK4 was passed to `ode45()` as an input. This means that while `ode45()` is still free to adjust the step size as necessary, the outputted solution will be on the specified values. Shown below is an example plot of all three solutions on top of each other. To view the full solution animated, please run the `main.m` file and select the corresponding option.

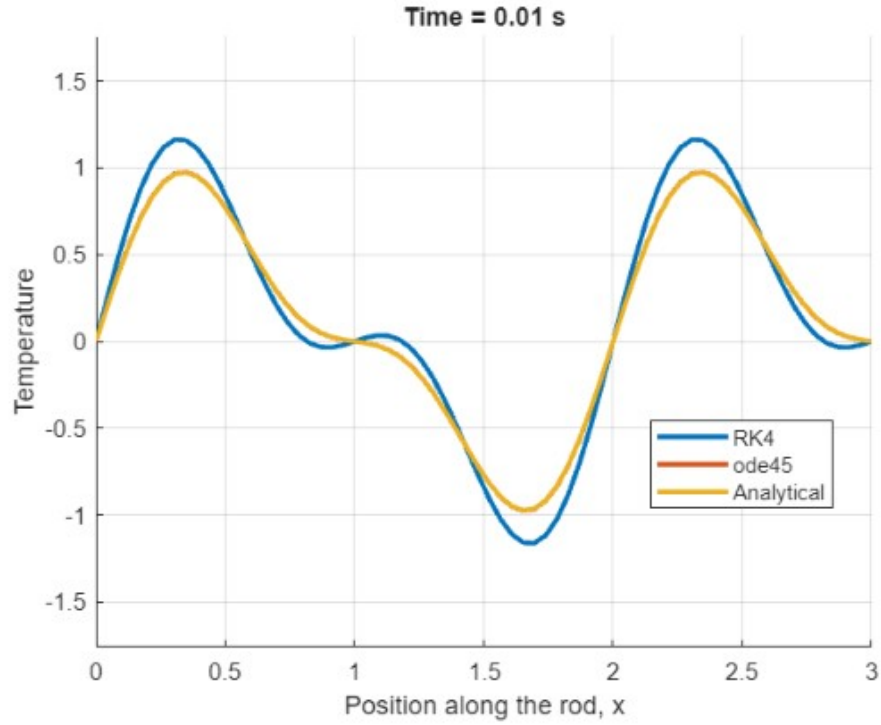


Figure 2: Iteration 34 of RK4 Solution and ode45() Solution vs. Analytical Solution to Heat Equation

We found the maximum absolute error to be equal to 0.2178 and the error in the L2 norm to be equal to 16.0653. You can see the error of the numerical solution in the graph below.

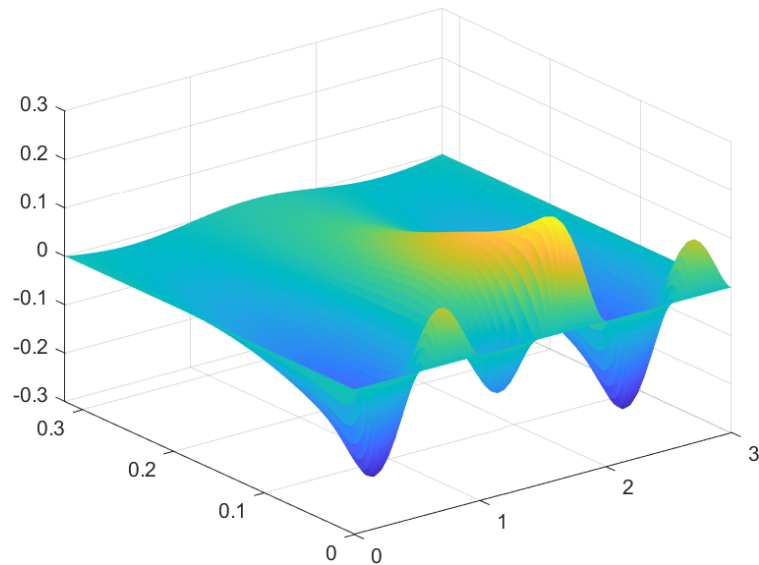


Figure 3: Error of Numerical solution

2 General Wave Equation

$$u_{tt} = 0.16u_{xx} + 0.02 \sin(x + t)$$

$$u(0, t) = 0.01 \sin(t)$$

$$u_x(2, t) = 0$$

$$u(x, 0) = \sin\left(\frac{\pi x}{2}\right)$$

$$u_t(x, 0) = \cos\left(\frac{\pi x}{2}\right)$$

Using the central difference approximation for second derivatives gives the following:

$$\begin{aligned} \frac{u_i^{n+1} - 2u_i^n + u_i^{n-1}}{\Delta t^2} &= 0.16 \left(\frac{u_{i+1}^n - 2u_i^n + u_{i-1}^n}{\Delta x^2} \right) + 0.02 \sin(x + t) \\ u_i^{n+1} &= \frac{0.16\Delta t^2}{\Delta x^2} (u_{i+1}^n - 2u_i^n + u_{i-1}^n) + 0.02\Delta t^2 \sin(x + t) + 2u_i^n - u_i^{n-1} \end{aligned} \quad (4)$$

Note that the i index represents the spatial steps and the n index represents the temporal steps. This gives a formulation for the next time step. Rewriting this gives a formulation for the current time step based on the previous time steps.

$$u_i^n = \frac{0.16\Delta t^2}{\Delta x^2} (u_{i+1}^{n-1} - 2u_i^{n-1} + u_{i-1}^{n-1}) + 0.02\Delta t^2 \sin(x + t) + 2u_i^{n-1} - u_i^{n-2} \quad (5)$$

This is the form which will be implemented to produce the numerical solution. Note that this equation requires knowledge of the two previous time steps, and thus a different method is required for initialization. In order to do this, begin with the forward difference equation in time and the given initial condition.

$$\frac{u_i^{n+1} - u_i^n}{\Delta t} = \cos\left(\frac{\pi x}{2}\right) \quad (6)$$

For the wave equation, the other initial condition gives the value, but the wave equation requires two previous time steps. Using this formulation, a "ghost step" can be obtained and used to generate the first iteration of the wave equation. This is given by:

$$u_i^0 = u_i^1 - \Delta t \cos\left(\frac{\pi x}{2}\right) \quad (7)$$

For the boundary condition at the left endpoint, it is readily apparent that the formulation for this must be:

$$u_{end}^n = u_{end-1}^n \quad (8)$$

Shown below is an image of the wave at one of the iterations. To see the full animation, please run the main.m file and select the corresponding option.

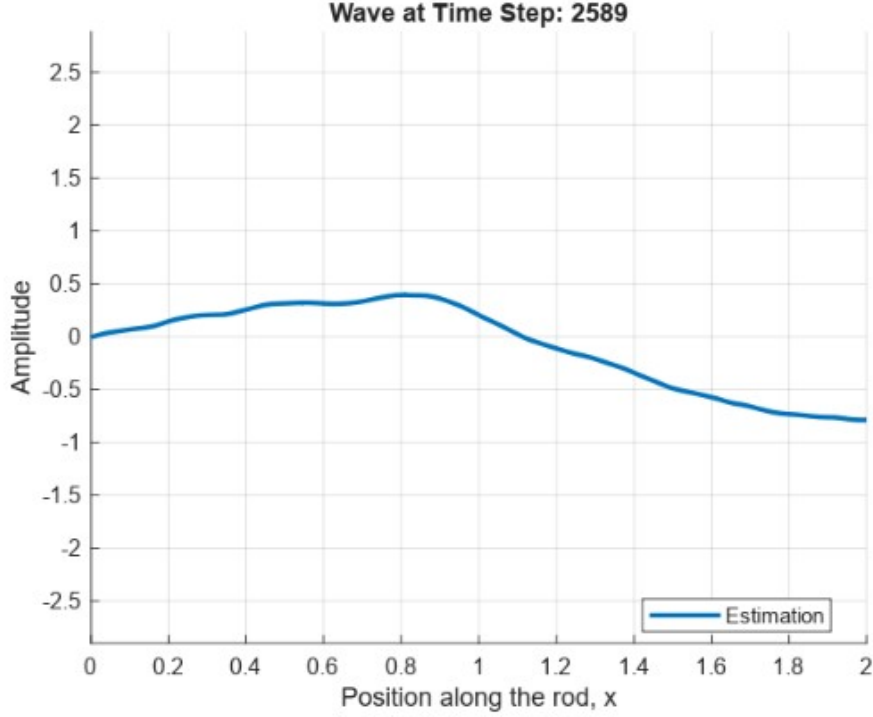


Figure 4: Iteration 2589 of Numerical Solution of Wave Equation

3 Two-Dimensional General Heat Equation

$$u_t = 0.5\nabla^2 u + e^{-t}u$$

$$u_x(0, y, t) = 0$$

$$u_y(x, 0, t) = 0$$

$$u(x, 4, t) = e^{-t}$$

$$u(4, y, t) = 0$$

$$u(x, y, 0) = \begin{cases} 1 & \text{for } y \geq x \\ 0 & \text{for } y < x \end{cases}$$

Using the forward difference approximation for first derivatives and the central difference approximation for second derivatives gives the following:

$$\begin{aligned} \frac{u_{i,j}^{n+1} - u_{i,j}^n}{\Delta t} &= 0.5 \left(\frac{u_{i+1,j}^n - 2u_{i,j}^n + u_{i-1,j}^n}{\Delta x^2} + \frac{u_{i,j+1}^n - 2u_{i,j}^n + u_{i,j-1}^n}{\Delta y^2} \right) + e^{-t} u_{i,j}^n \\ u_{i,j}^{n+1} &= 0.5\Delta t \left(\frac{u_{i+1,j}^n - 2u_{i,j}^n + u_{i-1,j}^n}{\Delta x^2} + \frac{u_{i,j+1}^n - 2u_{i,j}^n + u_{i,j-1}^n}{\Delta y^2} \right) + (e^{-t}\Delta t + 1) u_{i,j}^n \end{aligned} \quad (9)$$

In this case, i and j represent the x and y spatial steps, respectively. This gives a formulation for the next time step. Assuming the same step size for both spatial dimensions allows the simplification $\Delta x = \Delta y = h$. Along with rewriting the formulation to give the current time step as a function of the previous time steps, this gives:

$$u_{i,j}^{n+1} = \frac{0.5\Delta t}{h^2} (u_{i+1,j}^{n-1} + u_{i-1,j}^{n-1} + u_{i,j+1}^{n-1} + u_{i,j-1}^{n-1} - 4u_{i,j}^{n-1}) + (e^{-t}\Delta t + 1) u_{i,j}^{n-1} \quad (10)$$

This is the form which will be implemented to produce the numerical solution. Similar to the wave equation above, for the derivative boundary conditions, it is readily apparent that the formulation for this must be:

$$\begin{aligned} u_{start,j}^n &= u_{start+1,j}^n \\ u_{i,start}^n &= u_{i,start+1}^n \end{aligned} \quad (11)$$

Shown below is an image of the heat on the surface at one of the iterations. To see the full animation, please run the main.m file and select the corresponding option.

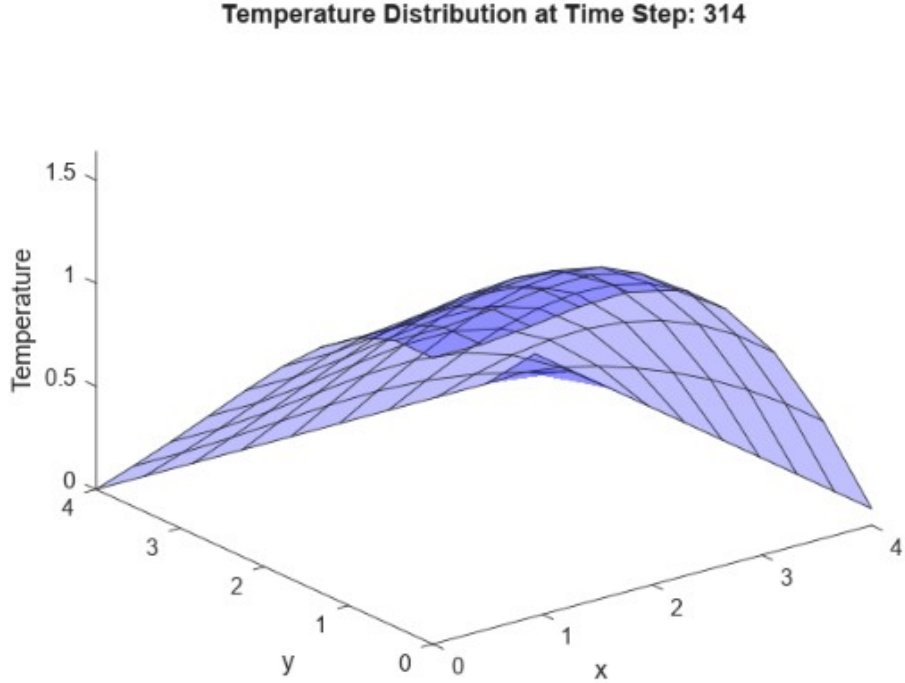


Figure 5: Iteration 314 of Numerical Solution of 2D Heat Equation

4 Poisson's Equation

$$\nabla^2 u = 1 + 0.2\delta_{(1,3)}(x, y) = f(x, y)$$

$$u(x, 4) = x$$

$$u(2, y) = 1$$

$$u(0, y) = 1 \text{ for } 2 \leq y \leq 4$$

$$u(1, y) = 0 \text{ for } 0 \leq y \leq 2$$

$$u(x, 2) = 0 \text{ for } 0 \leq x \leq 1$$

$$u(x, 0) = 1 \text{ for } 1 \leq x \leq 2$$

Using the central difference approximation for the second derivatives gives:

$$\left(\frac{u_{i+1,j}^n - 2u_{i,j}^n + u_{i-1,j}^n}{\Delta x^2} + \frac{u_{i,j+1}^n - 2u_{i,j}^n + u_{i,j-1}^n}{\Delta y^2} \right) = f(x, y) \quad (12)$$

Assuming that $h = \Delta x = \Delta y$ gives:

$$u_{i,j} = \frac{1}{4} (u_{i+1,j} + u_{i-1,j} + u_{i,j+1} + u_{i,j-1} - h^2 f(x, y)) \quad (13)$$

In order to solve this, use the method of iterative relaxation, where the matrix is initialized with zeros in all entries except the boundary conditions and iterations are performed until the change from step to step is within some specified tolerance. In some sense, this can be thought of as iterating through the transient part of the 2D heat equation until the steady-state solution is achieved. Shown below is an image of the heat on the surface as determined by the iteration. To have a figure which can be manipulated in space, please run the main.m file and select the corresponding option.

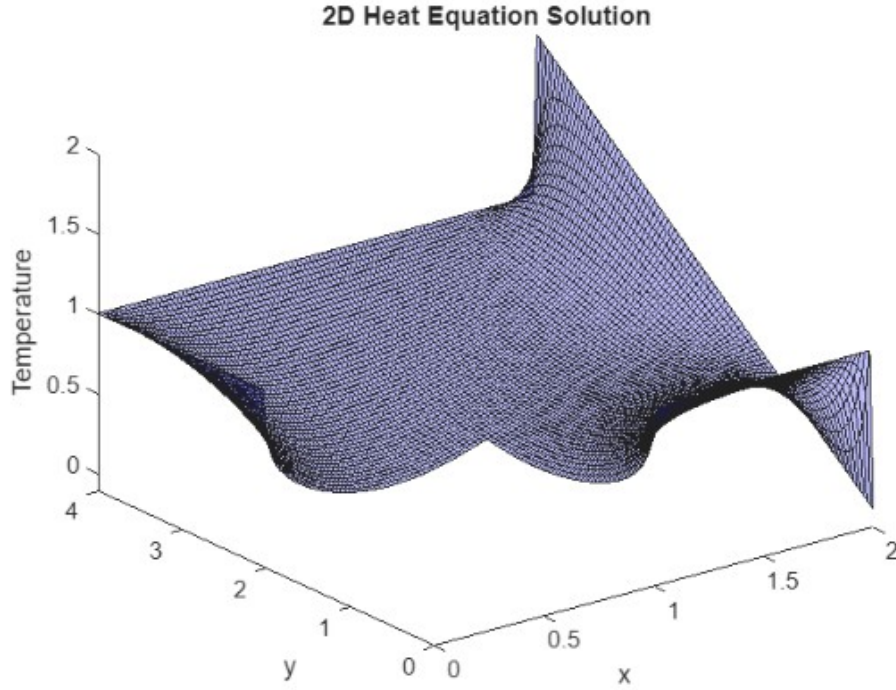


Figure 6: Final Numerical Solution of Poisson's Equation

5 Spacecraft Application

$$\begin{aligned}\nabla^2 u &= 0 \\ u(0, y) &= 0 \\ u(a, y) &= 0 \\ u(x, 0) &= 0 \\ u(x, b) &= \beta x(a - x)\end{aligned}$$

By separation of variables we obtain an analytic temperature profile of

$$u(x, y) = \sum_{n=1}^{\infty} A_n \sin(8nx) \sinh(8ny)$$

where

$$A_n = \frac{625 \left(4 \sin \left(\frac{\pi n}{2} \right)^2 - \pi n \sin(\pi n) \right)}{2n^3 \pi \sinh \left(\frac{\pi n}{2} \right)}.$$

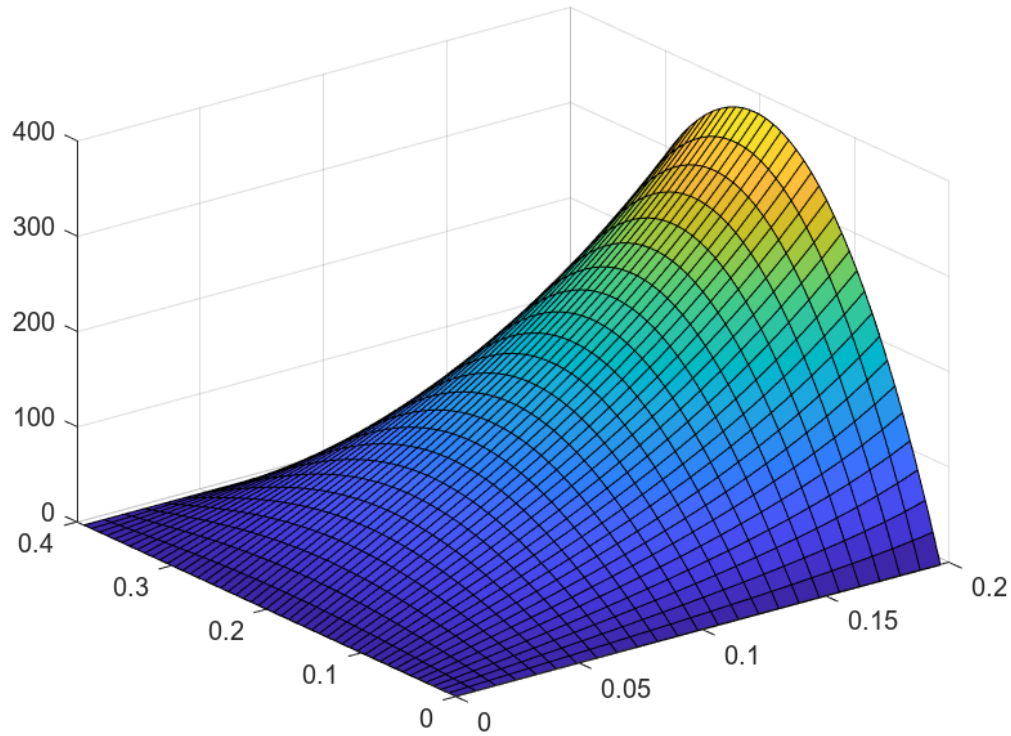


Figure 7: Analytical Solution of Radiator Fin

We can see quite clearly how temperature barely drops off the further away we are from the heat source so as such it would make sense to reduce the rectangle just to a parabola around the main heat source. This would provide the most amount of heat dissipation without greatly increasing weight. Utilizing our formula from question 4, we can find that our answer seems to drop off quicker than the analytical solution. This is likewise also seen even more exaggerated in our MOL solution (which utilizes a very similar formula found in question 1) as seen in the figures below.

2D Heat Equation Solution

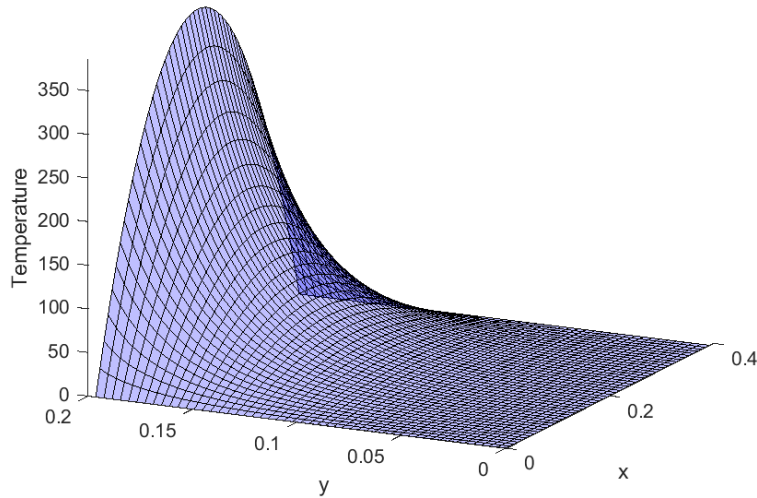


Figure 8: Poisson solution method of Radiator Fin

2D Heat MOL Solution

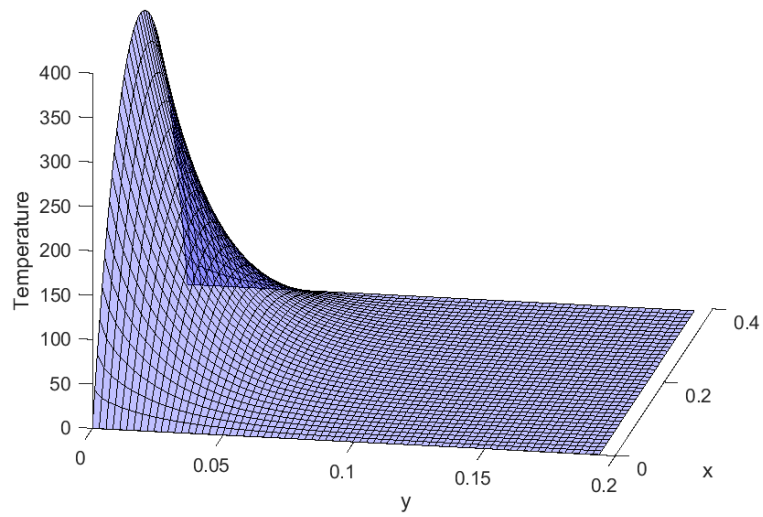


Figure 9: MOL solution method of Radiator Fin

While these methods generally followed the structure of the analytical solution, its especially clear in the MOL solution that it trended to 0 too quickly and as such didn't model the solution as good as it could have.

```
% These control the number of iterations in both space and time. Note that
% this is not the number of steps, which would be either of these values
% minus 1. The output matrices will be exactly these dimensions.
t_iter_heat = 5000;
x_iter_heat = 70;

% The t-values used for analysis are contained in t_vals, and the x-values
% used for analysis are contained in x_vals. X_analytical contains the
% computed values for the analytical solution. The first row contains the
% initial condition, and each subsequent row contains the rod analyzed at
% the x-values at the next timestep. X and U_sol are set up in the same way
% and are the calculated values for RK4 and ode45 respectively. The step
% sizes for both x and t are given by delta_x and delta_t respectively.
disp("Basic Heat Equation:")
[t_vals_heat, x_vals_heat, X_analytical_heat, X_heat, U_sol_heat, t_sol_heat,
delta_x_heat, delta_t_heat] = basic_heat(t_iter_heat, x_iter_heat);

% Wave Code
t_iter_wave = 5000;
x_iter_wave = 100;
c_wave = 0.4;
disp("General Wave Equation:")
[t_vals_wave, x_vals_wave, X_wave, delta_x_wave, delta_t_wave] = general_wave
(t_iter_wave, x_iter_wave, c_wave);

% 2D Heat Code
t_iter_2D = 1000;
x_iter_2D = 10;
y_iter_2D = 10;
alpha_2D = 0.5;
disp("General 2D Heat Equation:")
[t_vals_2D, x_vals_2D, y_vals_2D, X_2D, delta_xy_2D, delta_t_2D] = general_heat_2D
(t_iter_2D, x_iter_2D, y_iter_2D, alpha_2D);

% Poisson's Equation Code
x_iter_pois = 101;
y_iter_pois = x_iter_pois*2 - 1;
TOL_pois = 10^-8;
max_iter_pois = 100000;
disp("Poisson's Equation:")
[x_vals_pois, y_vals_pois, X_pois, delta_x_pois, delta_y_pois] = general_poisson
(x_iter_pois, y_iter_pois, TOL_pois, max_iter_pois);
```

```

function [t_vals, x_vals, X_analytical, X, U_sol, t_sol, delta_x, delta_t] =
basic_heat(t_iter, x_iter)
    t_0 = 0;
    t_f = 2;
    t_vals = linspace(t_0, t_f, t_iter);

    L = 3;
    X = zeros(t_iter, x_iter);
    x_0 = @(x) sin(pi*x) + sin(2*pi*x);
    x_vals = linspace(0,L,x_iter);
    X(1, :) = x_0(x_vals);
    X(:,1) = 0;
    X(:,end) = 0;

    K = zeros(4, x_iter);
    delta_t = t_vals(2) - t_vals(1);
    delta_x = x_vals(2) - x_vals(1);
    alpha = 2;

    if delta_t > 0.5*delta_x^2/alpha
        fprintf('Upper Bound: %g\n', 0.5*delta_x^2/alpha)
        fprintf('Delta t: %g\n', delta_t)
        disp("Unstable, increase t_iter or decrease x_iter")
        return
    end

    for i = 2 : t_iter
        K(1,2:x_iter-1) = alpha*(X(i-1,3:x_iter) - 2*X(i-1,2:x_iter-1) + X(i-1,1:
x_iter-2))/delta_x^2;
        K(2,2:x_iter-1) = alpha*((X(i-1,3:x_iter) + (delta_t/2)*K(1,3:x_iter)) - 2*(X
(i-1,2:x_iter-1) + (delta_t/2)*K(1,2:x_iter-1)) + (X(i-1,1:x_iter-2) + (delta_t/2)*K
(1,1:x_iter-2)))/delta_x^2;
        K(3,2:x_iter-1) = alpha*((X(i-1,3:x_iter) + (delta_t/2)*K(2,3:x_iter)) - 2*(X
(i-1,2:x_iter-1) + (delta_t/2)*K(2,2:x_iter-1)) + (X(i-1,1:x_iter-2) + (delta_t/2)*K
(2,1:x_iter-2)))/delta_x^2;
        K(4,2:x_iter-1) = alpha*((X(i-1,3:x_iter) + (delta_t)*K(3,3:x_iter)) - 2*(X
(i-1,2:x_iter-1) + (delta_t)*K(3,2:x_iter-1)) + (X(i-1,1:x_iter-2) + (delta_t)*K(3,1:
x_iter-2)))/delta_x^2;

        X(i,:) = X(i-1,:) + (delta_t/6)*(K(1,:) + K(2,:) + K(3,:) + K(4,:));
    end

    X_analytical = zeros(t_iter, x_iter);

    f = @(x,t) exp(-2*pi^2*t)*sin(pi*x) + exp(-8*pi^2*t)*sin(2*pi*x);

    for i = 1:length(t_vals)
        for j = 2:length(x_vals) - 1

```

```

        X_analytical(i,j) = f(x_vals(j),t_vals(i));
    end
end

U0 = x_0(x_vals);
tspan = t_vals;

% Solve using ode45
[t_sol, U_sol] = ode45(@(t,U) heatODE(t, U, alpha, delta_x, x_iter), tspan, U0);

str = "";
if t_iter >= 1000
    str = "Caution: At current number of iterations, animation may take a long
time. ";
end
decide = input(str + "Press enter to exit. Type 1 for RK4 vs. Analytical, Type 2
for RK4 vs ode45, Type 3 for all three.", "s");
switch decide
    case "1"
        animation_speed = 0.05;
        figure;
        hold on
        h = plot(x_vals, X(1,:), 'LineWidth', 2);
        g = plot(x_vals, X_analytical(1,:), 'LineWidth', 2);
        xlabel('Position along the rod, x');
        ylabel('Temperature');
        title('Heat Equation Animation');
        legend('Estimation', 'Analytical', 'Location', 'best')
        grid on;

        % Fix y-axis limits to avoid recalculating each frame
        ylim([min(X(:)), max(X(:))]);

        % Improve rendering performance by reducing overhead
        set(gcf, 'Renderer', 'painters');

        % Efficiently animate without redrawing the full figure
        for timestep = 1:size(X,1)
            h.YData = X(timestep,:); % only updating data, very efficient
            g.YData = X_analytical(timestep,:);
            title(sprintf('Temperature Distribution at Time Step: %d',
timestep));
            drawnow limitrate; % significantly improves performance
            pause(animation_speed); % adjust animation speed
        end
    case "2"
        animation_speed = 0;

        figure;

```

```

    hold on
    g = plot(x_vals, X(1,:), 'LineWidth', 2);
    h = plot(x_vals, U_sol(1,:), 'LineWidth', 2);
    ylim([min(U_sol(:)) max(U_sol(:))]);
    grid on;
    xlabel('Position along the rod, x');
    ylabel('Temperature');
    legend('RK4', 'ode45', 'Location', 'best')
    for k = 1:length(t_sol)
        set(h, 'YData', U_sol(k,:));
        set(g, 'YData', X(k,:));
        title(sprintf('Time = %.2f s', t_sol(k)));
        drawnow limitrate;
        pause(animation_speed);
    end
case "3"
    animation_speed = 0;

    figure;
    hold on
    g = plot(x_vals, X(1,:), 'LineWidth', 2);
    h = plot(x_vals, U_sol(1,:), 'LineWidth', 2);
    i = plot(x_vals, X_analytical(1,:), 'LineWidth', 2);
    ylim([min(U_sol(:)) max(U_sol(:))]);
    grid on;
    xlabel('Position along the rod, x');
    ylabel('Temperature');
    legend('RK4', 'ode45', 'Analytical', 'Location', 'best')
    for k = 1:length(t_sol)
        set(h, 'YData', U_sol(k,:));
        set(g, 'YData', X(k,:));
        set(i, 'YData', X_analytical(k,:));
        title(sprintf('Time = %.2f s', t_sol(k)));
        drawnow limitrate;
        pause(animation_speed);
    end
case isempty(decide)
    return
end
end

function dUdt = heatODE(t, U, alpha, delta_x, x_iter)
    dUdt = zeros(x_iter,1);

    % Boundary conditions (example: Dirichlet - fixed endpoints)
    U(1) = 0;           % u(0,t) boundary (fixed)
    U(x_iter) = 0;      % u(L,t) boundary (fixed)

    % Interior points (second-order finite difference)

```

```
for i = 2:x_iter-1
    dUdt(i) = alpha*(U(i+1)-2*U(i)+U(i-1))/delta_x^2;
end
end
```

```

function [t_vals, x_vals, X, delta_x, delta_t] = general_wave(t_iter, x_iter, c)

    t_0 = 0;
    t_f = 120;
    t_vals = linspace(t_0, t_f, t_iter);

    L = 2;
    X = zeros(t_iter, x_iter);
    x_0 = @(x) sin(pi*x/2);
    x_vals = linspace(0,L,x_iter);
    X(1, :) = x_0(x_vals);
    BC1 = @(t) 0.01*sin(t);
    X(:,1) = BC1(t_vals);
    X(1,end) = X(1,end-1);

    delta_t = t_vals(2) - t_vals(1);
    delta_x = x_vals(2) - x_vals(1);

    if delta_t > delta_x/c
        fprintf('Upper Bound: %g\n', delta_x/c)
        fprintf('Delta t: %g\n', delta_t)
        disp("Unstable, increase t_iter or decrease x_iter")
        return
    end

    ghost_step = X(1,:) - delta_t*cos(pi*x_vals/2);
    i = 2;
    X(2,2:x_iter-1) = (c*delta_t/delta_x)^2*(X(i-1,3:x_iter) - 2*X(i-1,2:x_iter-1) + X(i-1,1:x_iter-2)) + 0.02*delta_t^2*sin(x_vals(2:x_iter-1)+t_vals(i-1)) + 2*X(i-1,2:
x_iter-1) - ghost_step(2:x_iter-1);

    for i = 3:t_iter
        X(i,2:x_iter-1) = (c*delta_t/delta_x)^2*(X(i-1,3:x_iter) - 2*X(i-1,2:x_iter-
1) + X(i-1,1:x_iter-2)) + 0.02*delta_t^2*sin(x_vals(2:x_iter-1)+t_vals(i-1)) + 2*X(i-
1,2:x_iter-1) - X(i-2, 2:x_iter-1);
        X(i,end) = X(i,end-1);
    end

    str = "";
    if t_iter >= 1000
        str = "Caution: At current number of iterations, animation may take a long
time. ";
    end
    decide = input(str + "Press enter to exit. Type 1 for Finite Difference
Animation", "s");
    switch decide
        case "1"
            animation_speed = 0.0;
            figure;

```



```
hold on
h = plot(x_vals, X(1,:), 'LineWidth', 2);
xlabel('Position along the rod, x');
ylabel('Amplitude');
title('Wave Equation Animation');
legend('Estimation', 'Location', 'best')
grid on;

% Fix y-axis limits to avoid recalculating each frame
ylim([min(X(:)), max(X(:))]);
xlim([0 2]);

% Improve rendering performance by reducing overhead
set(gcf, 'Renderer', 'painters');

% Efficiently animate without redrawing the full figure
for timestep = 1:size(X,1)
    h.YData = X(timestep,:); % only updating data, very efficient
    title(sprintf('Wave at Time Step: %d', timestep));
    drawnow limitrate; % significantly improves performance
    pause(animation_speed); % adjust animation speed
end
case isempty(decide)
    return
end

end
```

```

function [t_vals, x_vals, y_vals, X, delta_xy, delta_t] = general_heat_2D(t_iter, x_iter, y_iter, alpha)

    t_0 = 0;
    t_f = 8;
    t_vals = linspace(t_0, t_f, t_iter);

    L = 4;
    H = 4;
    X = zeros(x_iter, y_iter, t_iter);
    x_0 = @(x, y) double(y >= x);
    x_vals = linspace(0, L, x_iter);
    y_vals = linspace(0, H, y_iter);
    for i = 1:length(x_vals)
        for j = 1:length(y_vals)
            X(i, j, 1) = x_0(x_vals(i), y_vals(j));
        end
    end

    X(1, :, 1) = X(2, :, 1);
    X(end, :, :) = 0;
    X(:, 1, 1) = X(:, 2, 1);
    for x_num = 1:x_iter
        X(x_num, end, :) = exp(-t_vals);
    end

    delta_t = t_vals(2) - t_vals(1);
    delta_xy = x_vals(2) - x_vals(1);

    if delta_t > 0.25*delta_xy^2/alpha
        fprintf('Upper Bound: %g\n', 0.25*delta_xy^2/alpha)
        fprintf('Delta t: %g\n', delta_t)
        disp("Unstable, increase t_iter or decrease x_iter")
        return
    end

    for i = 2:t_iter
        X(2:x_iter-1, 2:x_iter-1, i) = alpha*delta_t/delta_xy^2*...
            (X(3:x_iter, 2:x_iter-1, i-1) + X(1:x_iter-2, 2:x_iter-1, i-1) + ...
            X(2:x_iter-1, 3:x_iter, i-1) + X(2:x_iter-1, 1:x_iter-2, i-1) - 4*X(2:
x_iter-1, 2:x_iter-1, i-1)) + ...
            (exp(-t_vals(i))*delta_t + 1)*X(2:x_iter-1, 2:x_iter-1, i-1);
        X(1, :, i) = X(2, :, i);
        X(:, 1, i) = X(:, 2, i);
    end

    str = "";
    if t_iter >= 1000
        str = "Caution: At current number of iterations, animation may take a long

```

```

time. ";
end
decide = input(str + "Press enter to exit. Type 1 for Finite Difference
Animation", "s");
switch decide
    case "1"
        animation_speed = 0.0;

        % Create meshgrid for surface plot
        [X_grid, Y_grid] = meshgrid(x_vals, y_vals); % Assuming x and y are 1D
vectors

        figure;
        hold on
        h = surf(X_grid, Y_grid, X(:,:,1), "FaceAlpha", 0.25, 'FaceColor', 'b');
        xlabel('x');
        ylabel('y');
        zlabel('Temperature');
        title('2D Heat Equation Solution');
        % Fix z-axis limits to avoid recalculating each frame
        zlim([min(X(:)), max(X(:))]);
        view(3);

        % Improve rendering performance by reducing overhead
        set(gcf, 'Renderer', 'painters');

        for timestep = 1:size(X,3)
            set(h, 'ZData', X(:,:,timestep));
            title(sprintf('Temperature Distribution at Time Step: %d',
timestep));
            drawnow limitrate;
            pause(animation_speed);
        end
        case isempty(decide)
            return
    end
end
end

```

```

function [x_vals, y_vals, X, delta_x, delta_y] = general_poisson(x_iter, y_iter, TOL, max_iter)

    L = 2;
    H = 4;
    X = zeros(x_iter, y_iter, 1);
    x_vals = linspace(0,L,x_iter);
    y_vals = linspace(0,H,y_iter);

    X(1:(x_iter+1)/2,1:(y_iter+1)/2,1) = 0;
    X((x_iter+1)/2:end,1,1) = 1;
    X(1,(y_iter+1)/2:end,1) = 1;
    X(end,:,1) = 1;
    X(:,end,1) = x_vals;

    delta_x = x_vals(2) - x_vals(1);
    delta_y = y_vals(2) - y_vals(1);
    f = @(x, y) 1 + 0.2 * double(x == 1 & y == 3);
    f_mat = zeros(x_iter,y_iter);
    for i = 1:x_iter
        for j = 1:y_iter
            f_mat(i,j) = f(x_vals(i), y_vals(j));
        end
    end

    for i = 2:max_iter
        X(1:(x_iter+1)/2,1:(y_iter+1)/2,1) = 0;
        X((x_iter+1)/2:end,1,i) = 1;
        X(1,(y_iter+1)/2:end,i) = 1;
        X(end,:,i) = 1;
        X(:,end,i) = x_vals;
        X(2:x_iter-1,2:y_iter-1,i) = 0.25*(X(3:x_iter,2:y_iter-1,i-1) + X(1:x_iter-2,2:y_iter-1,i-1) + X(2:x_iter-1,3:y_iter,i-1) + X(2:x_iter-1,1:y_iter-2,i-1) - delta_x^2*f_mat(2:x_iter-1,2:y_iter-1));
        if max(abs(X(:, :, i) - X(:, :, i-1)), [], "all") <= TOL
            break
        end
    end
    X(1:(x_iter+1)/2,1:(y_iter+1)/2,end) = NaN;

    str = "";
    decide = input(str + "Press enter to exit. Type 1 for Finite Difference Solution", "s");
    switch decide
        case "1"

            % Create meshgrid for surface plot
            [X_grid, Y_grid] = meshgrid(x_vals, y_vals(1:2:end)); % Assuming x and y are 1D vectors

```

```
        figure;
        hold on
        h = surf(X_grid, Y_grid, X(:,1:2:end,end), "FaceAlpha", 0.25, \
'FaceColor', 'b');
        xlabel('x');
        ylabel('y');
        zlabel('Temperature');
        title('2D Heat Equation Solution');
        % Fix z-axis limits to avoid recalculating each frame
        zlim([min(X(:)), max(X(:))]);
        view(3);
    case isempty(decide)
        return
    end

end

end
```