

# Chapter 8 - Plane (2D) Problems

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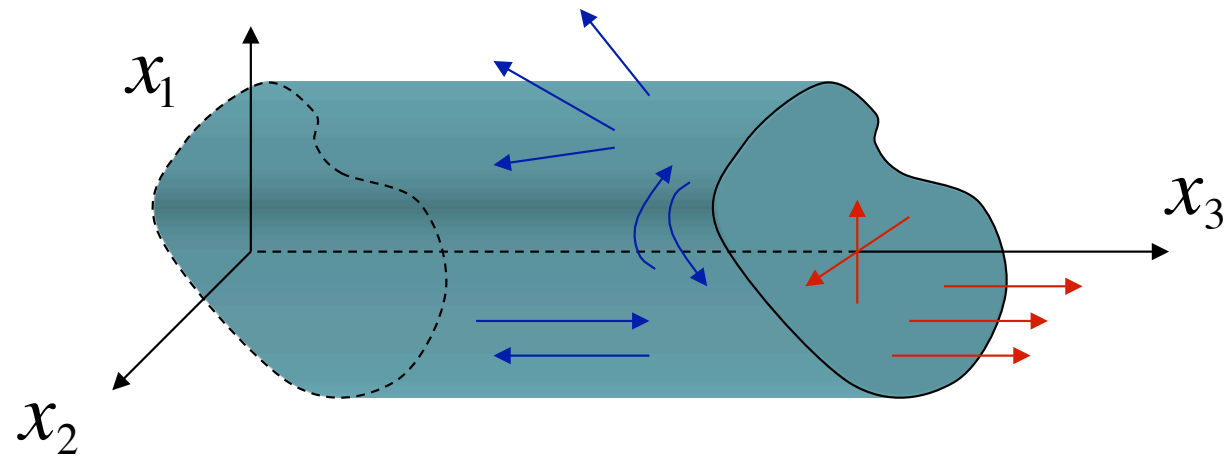
- ★ **Plane strain**
  - Plane strain formulation
  - Airy stress function

## Plane stress

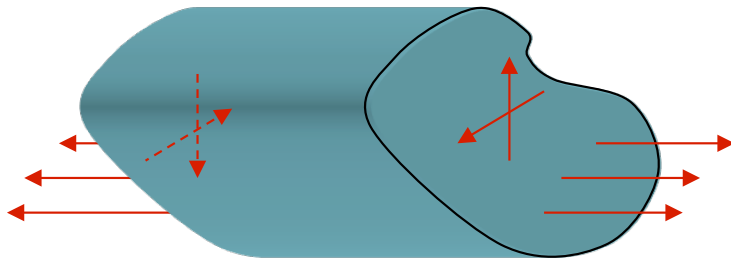
## Solutions

- Rectangular domains (beam problems)
- Axisymmetric domains
  - cylinder with internal/external pressure
  - hole under tension
- Half-space problems

- Recall:



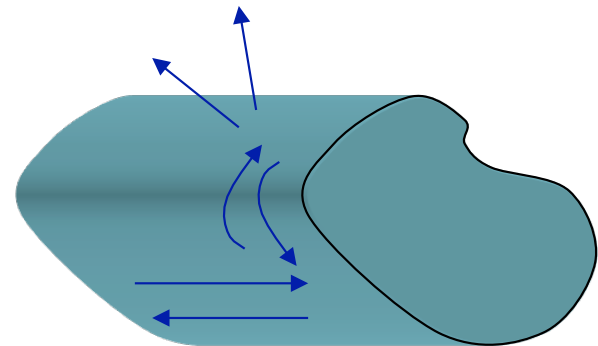
- Superposition



St. Venant type problems

A

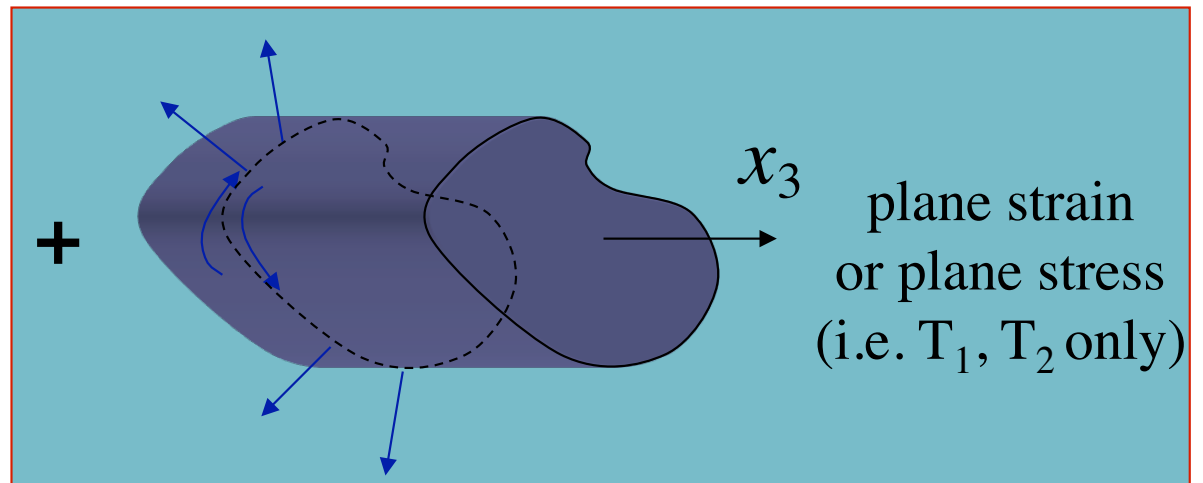
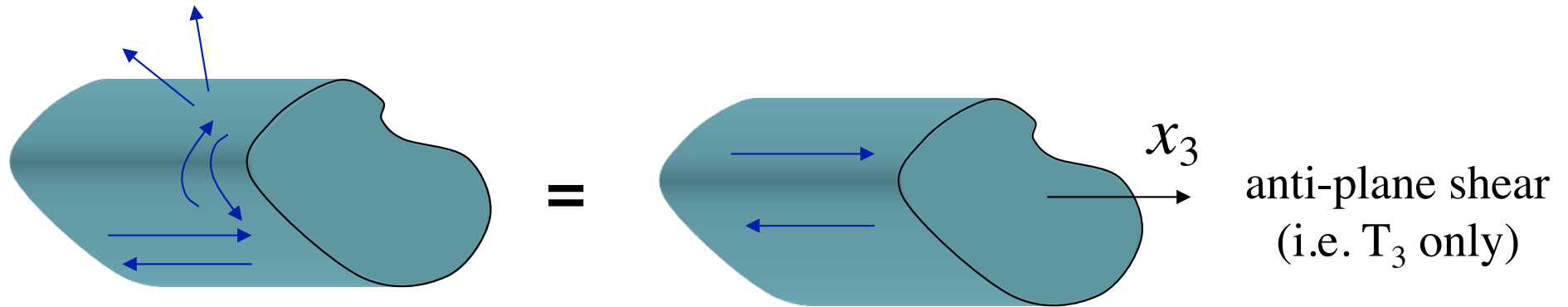
+



Load on lateral surface

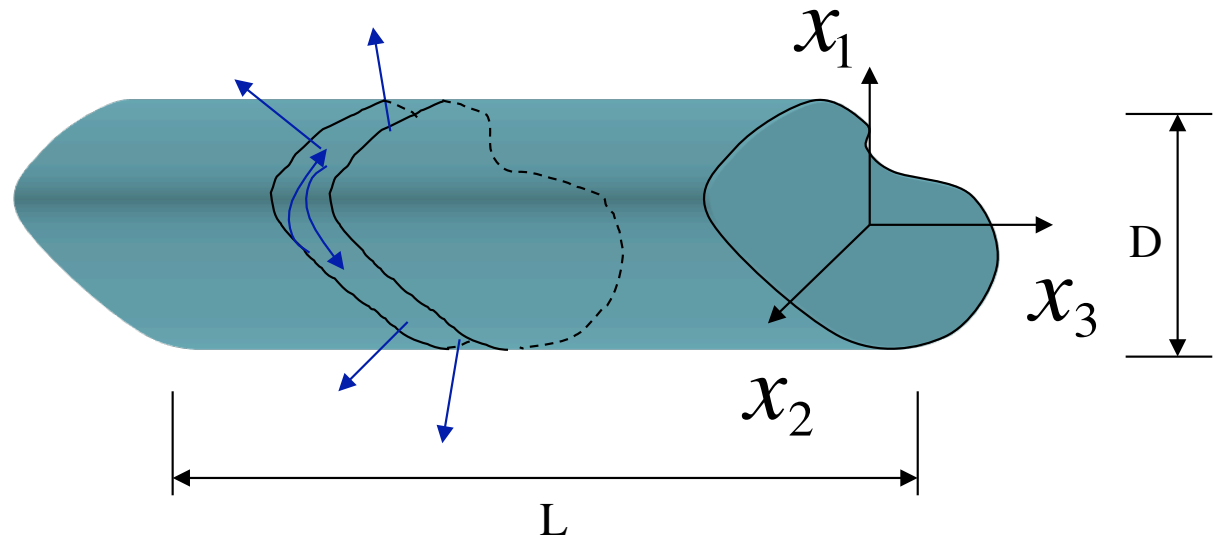
B

- Problem (B) :



# Plane Strain Formulation

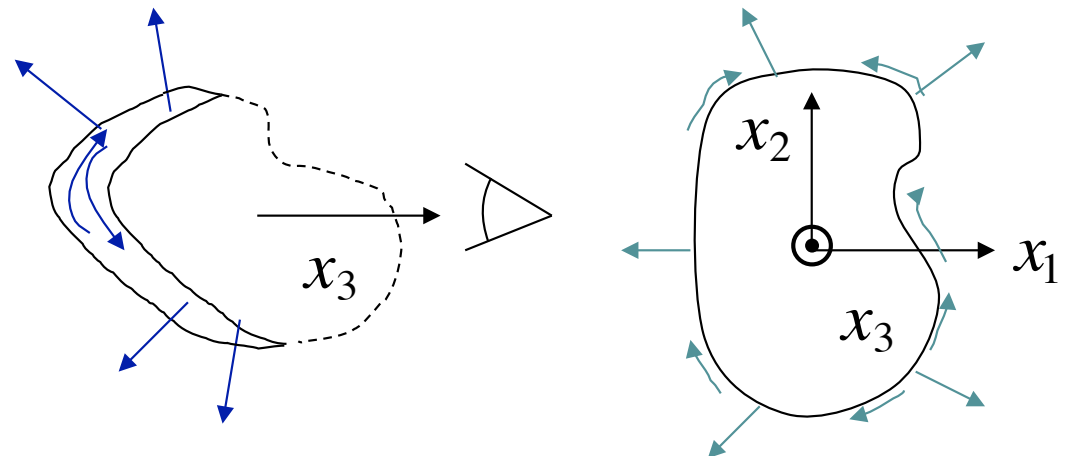
- Assume  $L \gg D$



- Assume no  $x_3$  dependence of tractions

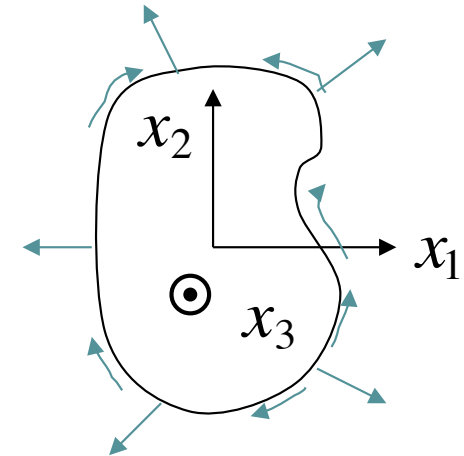
$$T_\alpha = T_\alpha(x_1, x_2) \quad \alpha = 1, 2$$

- Consider one “slice” of cylinder



# Plane Strain Formulation (cont.)

- Guess:  $u_1 = u_1(x_1, x_2)$   
 $u_2 = u_2(x_1, x_2)$   
 $u_3 = 0$  (i.e. no axial strain since  $L \gg D$ )



- Strains:  $\epsilon_{3i} = 0$   
 $\left. \begin{aligned} \epsilon_{11} &= \epsilon_{11}(x_1, x_2) \\ \epsilon_{22} &= \epsilon_{22}(x_1, x_2) \\ \epsilon_{12} &= \epsilon_{12}(x_1, x_2) \end{aligned} \right\} \text{ (i.e. 2D only)}$

- Stresses:  $E\epsilon_{11} = \sigma_{11} - \nu(\sigma_{22} + \sigma_{33})$   
 $E\epsilon_{22} = \sigma_{22} - \nu(\sigma_{11} + \sigma_{33})$   
 $E\epsilon_{33} = 0 = \sigma_{33} - \nu(\sigma_{11} + \sigma_{22}) \Rightarrow \sigma_{33} = \nu(\sigma_{11} + \sigma_{22})$
- (i.e. axial stress present in pl-ε)

## Plane Strain Formulation (cont.)

$$E\varepsilon_{11} = \sigma_{11} - \nu(\sigma_{22} + \nu(\sigma_{11} + \sigma_{22}))$$

$$\Rightarrow E\varepsilon_{11} = (1 - \nu^2)\sigma_{11} - \nu(1 + \nu)\sigma_{22}$$

$$\therefore \left. \begin{aligned} \frac{E}{1 - \nu^2} \varepsilon_{11} &= \sigma_{11} - \frac{\nu}{1 - \nu} \sigma_{22} \\ \frac{E}{1 - \nu^2} \varepsilon_{22} &= \sigma_{22} - \frac{\nu}{1 - \nu} \sigma_{11} \end{aligned} \right\}$$

$$\text{Setting: } \left. \begin{aligned} E' &= \frac{E}{1 - \nu^2} \\ \nu' &= \frac{\nu}{1 - \nu} \end{aligned} \right\} \Rightarrow$$

$$\begin{aligned} E' \varepsilon_{11} &= \sigma_{11} - \nu' \sigma_{22} \\ E' \varepsilon_{22} &= \sigma_{22} - \nu' \sigma_{11} \end{aligned}$$

$$\text{Also, } \sigma_{13} = \sigma_{23} = 0 \quad \text{and}$$

$$\sigma_{12} = 2\mu\varepsilon_{12}$$

# Plane Strain Formulation (cont.)

- Equilibrium  
(no body forces)

$$\sigma_{11,1} + \sigma_{12,2} + \cancel{\sigma_{13,3}}^0 = 0$$

$$\sigma_{12,1} + \sigma_{22,2} + \cancel{\sigma_{23,3}}^0 = 0$$

$$\cancel{\sigma_{13,1}}^0 + \cancel{\sigma_{23,2}}^0 + \sigma_{33,3} = 0$$

$$0 \text{ since } \sigma_{33} = \nu(\sigma_{11} + \sigma_{22}) = \sigma_{33}(x_1, x_2)$$

- Thus, in plane strain:

## Displacements

$$\begin{aligned} u_1 &= u_1(x_1, x_2) \\ u_2 &= u_2(x_1, x_2) \\ u_3 &= 0 \end{aligned}$$

## Stress-strain

$$\begin{aligned} E' \epsilon_{11} &= \sigma_{11} - \nu' \sigma_{22} \\ E' \epsilon_{22} &= \sigma_{22} - \nu' \sigma_{11} \\ \sigma_{12} &= 2\mu \epsilon_{12} \\ \sigma_{33} &= \nu(\sigma_{11} + \sigma_{22}) \end{aligned}$$

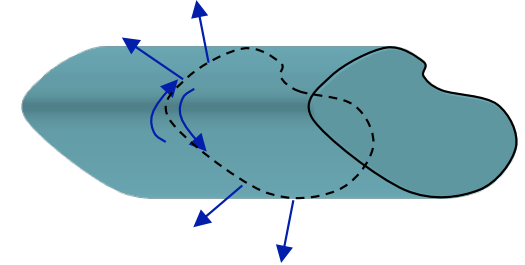
## Equilibrium

$$\begin{aligned} \sigma_{11,1} + \sigma_{12,2} &= 0 \\ \sigma_{12,1} + \sigma_{22,2} &= 0 \end{aligned}$$

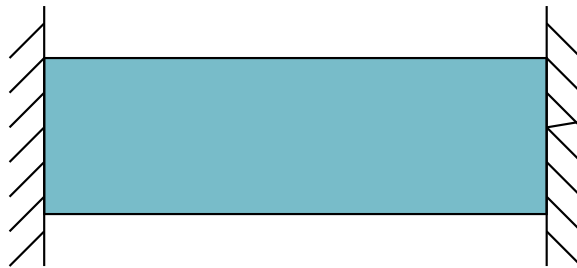
→ i.e. we have now a 2D class of problems

# Plane Strain Formulation (cont.)

- Note: end faces should be traction free  
 $\rightarrow$  problem, since we have  $\sigma_{33} = \nu(\sigma_{11} + \sigma_{22})!!!$



- Plane strain would be an exact solution for:

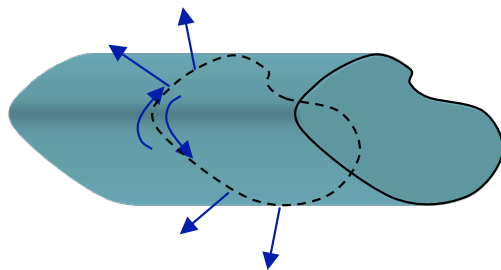


Lubricated wall  $\sigma_{13} = \sigma_{23} = 0$

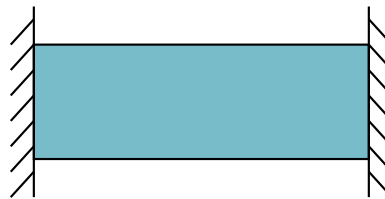
and  $u_3 = 0 \Rightarrow \epsilon_{33} = 0$

$$\Rightarrow \sigma_{33} = \nu(\sigma_{11} + \sigma_{22})$$

- In general we would need one more superposition:



=



plane strain

+



St. Venant type

$$\sigma_{33} = -\nu(\sigma_{11} + \sigma_{22})$$



# Airy Stress Function

- Assume a function  $\Phi(x_1, x_2)$  : **Airy Stress Function**

$$\begin{aligned}\sigma_{11} &= \frac{\partial^2 \Phi}{\partial x_2^2} \\ \sigma_{22} &= \frac{\partial^2 \Phi}{\partial x_1^2} \\ \sigma_{12} &= -\frac{\partial^2 \Phi}{\partial x_1 \partial x_2}\end{aligned}$$

- Equilibrium:

$$\left. \begin{aligned}\sigma_{11,1} + \sigma_{12,2} &= 0 \\ \sigma_{12,1} + \sigma_{22,2} &= 0\end{aligned}\right\} \Rightarrow$$

$$\begin{aligned}\frac{\partial^3 \Phi}{\partial x_2^2 \partial x_1} + \left( -\frac{\partial^3 \Phi}{\partial x_1 \partial x_2^2} \right) &= 0 \\ -\frac{\partial^3 \Phi}{\partial x_1^2 \partial x_2} + \frac{\partial^3 \Phi}{\partial x_1^2 \partial x_2} &= 0\end{aligned}$$

i.e. equilibrium identically satisfied!

## Airy Stress Function (cont.)

- Strains:  $E' \varepsilon_{11} = \sigma_{11} - \nu' \sigma_{22}$

$$E' \varepsilon_{22} = \sigma_{22} - \nu' \sigma_{11}$$

$$\sigma_{12} = 2\mu \varepsilon_{12}$$

$$\sigma_{33} = \nu(\sigma_{11} + \sigma_{22})$$

- Compatibility needed:

$$2\varepsilon_{12,12} = \varepsilon_{11,22} + \varepsilon_{22,11}$$

$$\Rightarrow \frac{\partial^4 \Phi}{\partial x_1^4} + 2 \frac{\partial^4 \Phi}{\partial x_1^2 \partial x_2^2} + \frac{\partial^4 \Phi}{\partial x_2^4} = 0$$

$$\Rightarrow \nabla^4 \Phi(x_1, x_2) = 0 \quad \text{Bi-harmonic equation}$$

- Any  $\Phi$  that satisfies this and gives  $\sigma$  as above is a possible pl- $\varepsilon$  solution. We have solved the entire class of pl- $\varepsilon$  problems. What distinguishes between problems are the b.c.s.

## Airy Stress Function (cont.)

- If body forces are present,  $f_i(x_1, x_2)$

$$\sigma_{11} = \frac{\partial^2 \Phi}{\partial x_2^2} + V$$

$$\sigma_{22} = \frac{\partial^2 \Phi}{\partial x_1^2} + V$$

$$\sigma_{12} = -\frac{\partial^2 \Phi}{\partial x_1 \partial x_2}$$

where

$$\frac{\partial V}{\partial x_1} = -f_1$$

$$\frac{\partial V}{\partial x_2} = -f_2$$

$V$ : “potential” energy

- Equilibrium is again identically satisfied

- Compatibility becomes:
- $$\nabla^4 \Phi(x_1, x_2) = -\left(\frac{1-2\nu}{1-\nu}\right) \nabla^2 V$$

# Chapter 8 - Plane (2D) Problems

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Plane strain formulation

Airy stress function

### ★ **Plane stress**

### **Solutions**

Rectangular domains (beam problems)

Axisymmetric domains

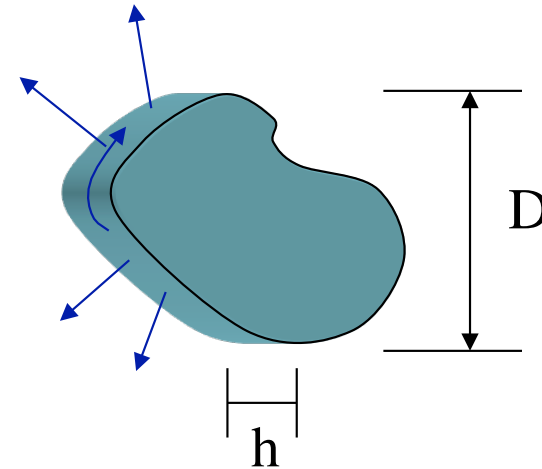
- cylinder with internal/external pressure
- hole under tension

Half-space problems

# Plane Stress

- “Short” cylinder, i.e. thin plate!

$$h \ll D$$



- Assume  $\sigma_{13} = \sigma_{23} = \sigma_{33} = 0$  and
 
$$\sigma_{11} = \sigma_{11}(x_1, x_2)$$

$$\sigma_{22} = \sigma_{22}(x_1, x_2)$$

$$\sigma_{12} = \sigma_{12}(x_1, x_2)$$

- Strains:
 

$$E\epsilon_{11} = \sigma_{11} - \nu(\sigma_{22} + \cancel{\sigma_{33}}^0)$$

$$E\epsilon_{22} = \sigma_{22} - \nu(\sigma_{11} + \cancel{\sigma_{33}}^0)$$

$$E\epsilon_{33} = \cancel{\sigma_{33}}^0 - \nu(\sigma_{11} + \sigma_{22})$$

$$\sigma_{13} = \sigma_{23} = 0$$

$$\sigma_{12} = 2\mu\epsilon_{12}$$

← Note: there is an  $\epsilon_{33}$

## Plane Stress (cont.)

- Equilibrium (no body forces):
 
$$\begin{aligned} \sigma_{11,1} + \sigma_{12,2} + \cancel{\sigma_{13,3}}^0 &= 0 \\ \sigma_{12,1} + \sigma_{22,2} + \cancel{\sigma_{23,3}}^0 &= 0 \\ \cancel{\sigma_{13,1}}^0 + \cancel{\sigma_{23,2}}^0 + \cancel{\sigma_{33,3}}^0 &= 0 \end{aligned}$$

- Plane stress:      Stress-strain                      Equilibrium

$$E\varepsilon_{11} = \sigma_{11} - \nu\sigma_{22}$$

$$E\varepsilon_{22} = \sigma_{22} - \nu\sigma_{11}$$

$$\sigma_{12} = 2\mu\varepsilon_{12}$$

$$\sigma_{11,1} + \sigma_{12,2} = 0$$

$$\sigma_{12,1} + \sigma_{22,2} = 0$$

- Thus plane stress is also a two dimensional problem with identical equations, but different “coefficients” in the stress-strain equations. Solve both problems using Airy stress function.

- Note: In pl- $\sigma$  bi-harmonic (i.e. compatibility) we neglect the effect of  $\varepsilon_{33}(\neq 0)$ . This is small and as  $h \rightarrow 0$  it becomes smaller.

## Plane Stress (cont.)

- BVP for two-dimensional plane problems (both pl- $\sigma$  and pl- $\varepsilon$ ):

Compatibility

$$\nabla^4 \Phi(x_1, x_2) = 0$$

Stress-strain

$$E' \varepsilon_{11} = \sigma_{11} - \nu' \sigma_{22}$$

$$E' \varepsilon_{22} = \sigma_{22} - \nu' \sigma_{11}$$

$$\sigma_{12} = 2\mu \varepsilon_{12}$$

$$E' = \begin{cases} E & \text{for } pl-\sigma \\ E/(1-\nu^2) & \text{for } pl-\varepsilon \end{cases}$$

$$\nu' = \begin{cases} \nu & \text{for } pl-\sigma \\ \frac{\nu}{1-\nu} & \text{for } pl-\varepsilon \end{cases}$$

Equilibrium (identically satisfied)

$$\sigma_{11} = \frac{\partial^2 \Phi}{\partial x_2^2}$$

$$\sigma_{22} = \frac{\partial^2 \Phi}{\partial x_1^2}$$

$$\sigma_{12} = -\frac{\partial^2 \Phi}{\partial x_1 \partial x_2}$$

$$\sigma_{33} = \begin{cases} 0 & \text{for } pl-\sigma \\ \nu(\sigma_{11} + \sigma_{22}) & \text{for } pl-\varepsilon \end{cases}$$

- Solution methods:
  - Trial and error
  - Separation of variables
  - Transforms (e.g. Fourier)
  - Green's functions

# Chapter 8 - Plane (2D) Problems

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### **Plane strain**

Plane strain formulation

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### **Plane stress**

### ★ **Solutions**

Rectangular domains (beam problems)

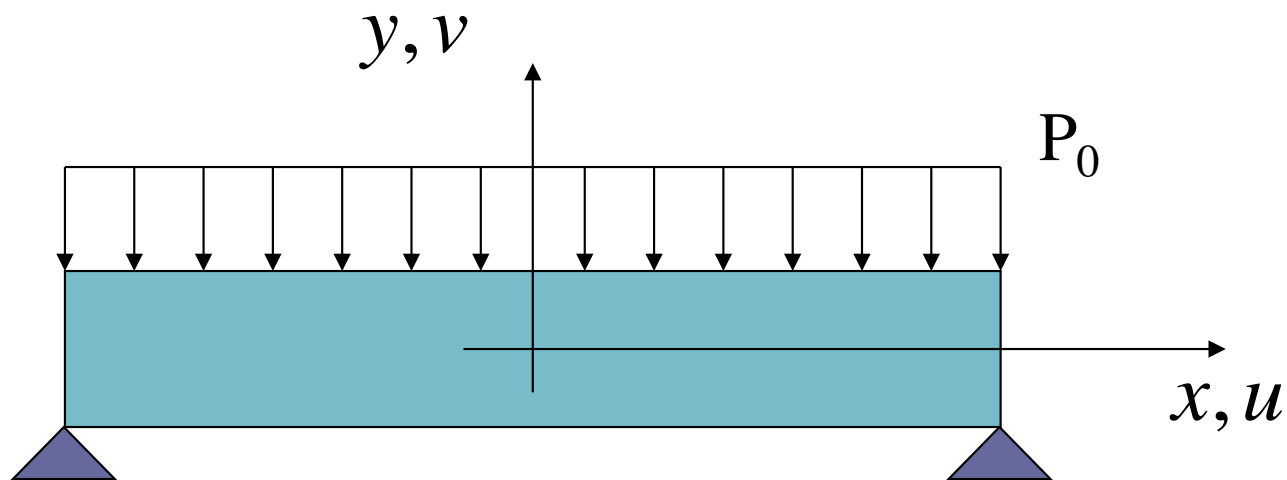
Axisymmetric domains

- cylinder with internal/external pressure
- hole under tension

Half-space problems



# Rectangular Domains



- Simply supported beam
- No body forces
- Plane strain

- Solution (Done in class... **Example 8-3** from textbook)
- See also examples 8-1, 8-2

## Rectangular Domains (cont.)

- Solution

$$\sigma_{xx} = -\frac{3P_0}{2} \left\{ \left( \frac{l}{h} \right)^2 - 4 \left( \frac{x}{l} \right)^2 \right\} \frac{y}{l} + P_0 \left\{ \frac{3}{5} - 4 \left( \frac{y}{h} \right)^2 \right\} \frac{y}{l}$$

$$\sigma_{xy} = \frac{3P_0}{2} \left\{ 1 - 4 \left( \frac{y}{h} \right)^2 \right\} \frac{x}{h}$$

$$\sigma_{yy} = -\frac{P_0}{2} - \frac{3P_0}{2} \frac{y}{h} + 2 \left( \frac{y}{h} \right)^3$$

Strength of Materials solution

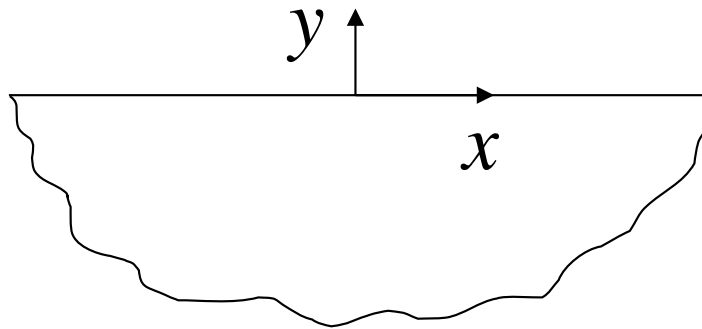
- Note: Reaction forces come in through a net shear force  $S^*$  at supports

$$S^* = \int_{-h/2}^{h/2} \sigma_{xy} dy = \frac{3P_0}{2} \int_{-h/2}^{h/2} \left\{ 1 - 4 \left( \frac{y}{h} \right)^2 \right\} \frac{l}{2h} dy$$

$$\Rightarrow S^* = \frac{1}{2} P_0 l$$

## Rectangular Domains (cont.)

- Other rectangular domains:

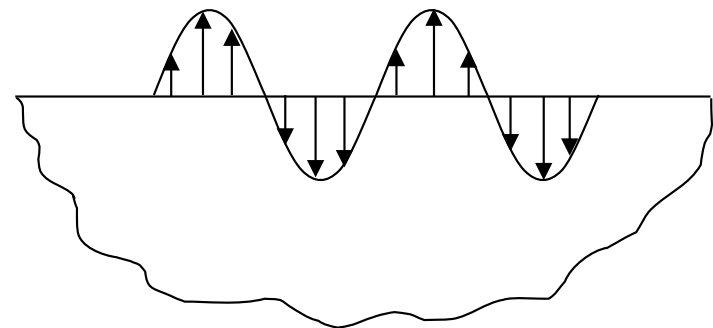
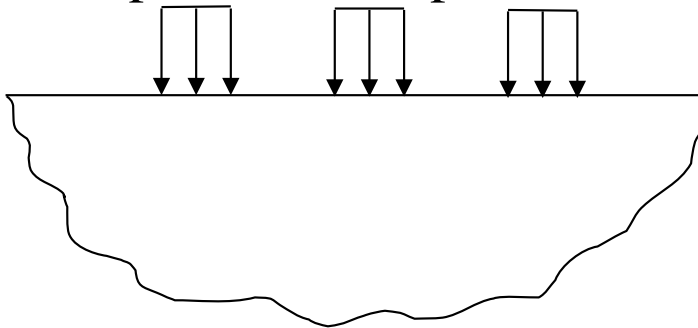


### Half Space problems

- Fourier transform or series
- Separation of variables

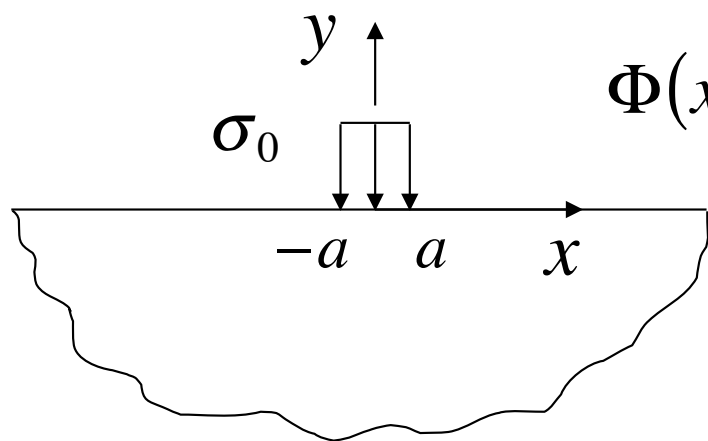
$$\Phi(x, y) = \sum_{n=1}^{\infty} \underbrace{(a_n \cos \alpha_n x + b_n \sin \alpha_n x)}_{f(x)} \underbrace{(A_n e^{\alpha_n y} + B_n e^{-\alpha_n y})}_{f(y)}$$

- Can solve periodic b.c. problems:

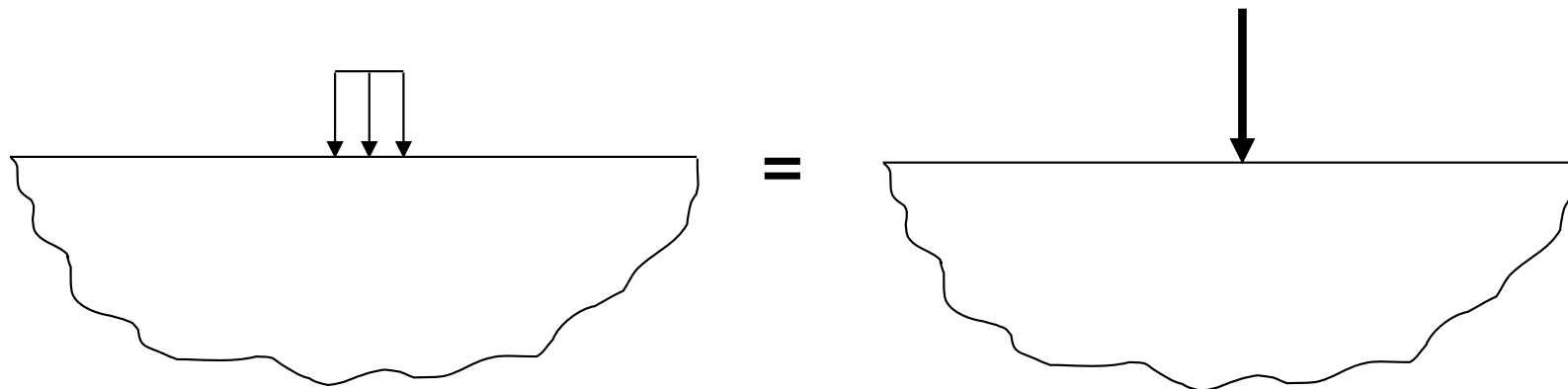


## Rectangular Domains (cont.)

- For non-periodic b.c.s use **Fourier Transform**


$$\Phi(x, y) = -\frac{\sigma_0}{\pi} \int_{-\infty}^{\infty} \frac{\sin \alpha a}{\alpha^3} \cos \alpha x (1 - \alpha y) e^{\alpha y} d\alpha$$

- Note: According to St Venant's principle



# Axisymmetric Domains

- Equilibrium in cylindrical coordinates:

$$\begin{aligned} \frac{\partial \sigma_r}{\partial r} + \frac{1}{r} \frac{\partial \tau_{r\theta}}{\partial \theta} + \frac{\partial \tau_{rz}}{\partial z} + \frac{\sigma_r - \sigma_\theta}{r} + R &= 0 \\ \frac{\partial \tau_{r\theta}}{\partial r} + \frac{1}{r} \frac{\partial \sigma_\theta}{\partial \theta} + \frac{\partial \tau_{\theta z}}{\partial z} + \frac{2\tau_{r\theta}}{r} + \Theta &= 0 \\ \frac{\partial \tau_{zr}}{\partial z} + \frac{\partial \tau_{\theta z}}{r \partial \theta} + \frac{\partial \sigma_z}{\partial z} + \frac{\tau_{zr}}{r} + Z &= 0 \end{aligned}$$

- Reduce to 2D pl- $\sigma$  and pl- $\epsilon$  as before and use:

$$\begin{aligned} \sigma_{rr} &= \frac{1}{r} \frac{\partial \Phi}{\partial r} + \frac{1}{r^2} \frac{\partial^2 \Phi}{\partial \theta^2} + V \\ \sigma_{\theta\theta} &= \frac{\partial^2 \Phi}{\partial r^2} + V \\ \sigma_{r\theta} &= -\frac{\partial}{\partial r} \left( \frac{1}{r} \frac{\partial \Phi}{\partial \theta} \right) \end{aligned}$$

With:

$$\{R, \Theta\} = \left\{ -\frac{\partial V}{\partial r}, -\frac{1}{r} \frac{\partial V}{\partial \theta} \right\}$$

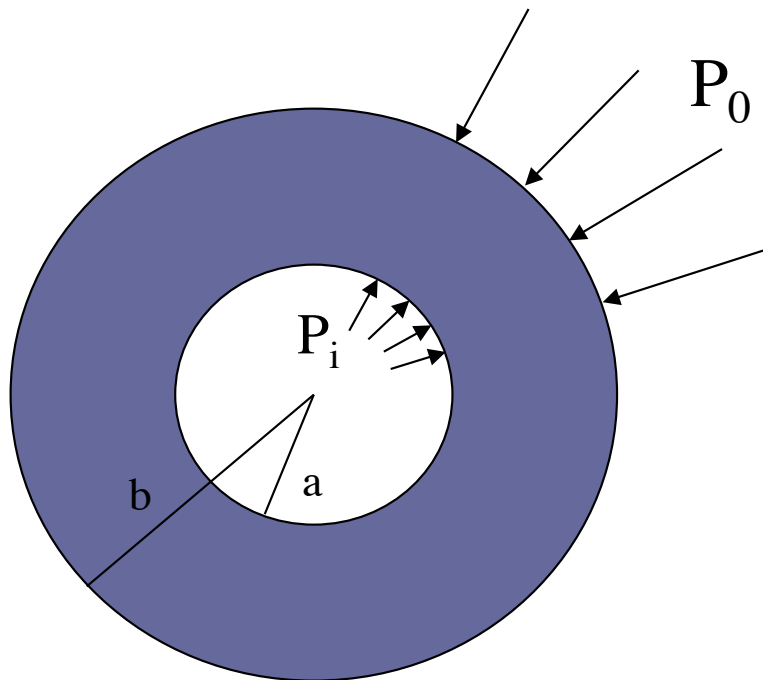
## Axisymmetric Domains (cont.)

- Compatibility:  $\nabla^4 \Phi(r, \theta) = 0$   $\nabla^2 \equiv \frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2}$
- General solution:

$$\begin{aligned} \Phi(r, \theta) = & a_0 \log r + b_0 r^2 + c_0 r^2 \log r + d_0 r^2 \theta + a'_0 \theta \\ & + \frac{a_1}{2} r \theta \sin \theta + \left( b_1 r^3 + a'_1 r^{-1} + b'_1 r \log r \right) \cos \theta \\ & - \frac{c_1}{2} r \theta \cos \theta + \left( d_1 r^3 + c'_1 r^{-1} + d'_1 r \log r \right) \sin \theta \\ & + \sum_{n=2}^{\infty} \left( a_n r^n + b_n r^{n+2} + a'_n r^{-n} + b'_n r^{-n+2} \right) \cos n\theta \\ & + \sum_{n=2}^{\infty} \left( c_n r^n + d_n r^{n+2} + c'_n r^{-n} + d'_n r^{-n+2} \right) \sin n\theta \end{aligned}$$

## Axisymmetric Domains (cont.)

- Thick walled cylinder with internal and external pressure:



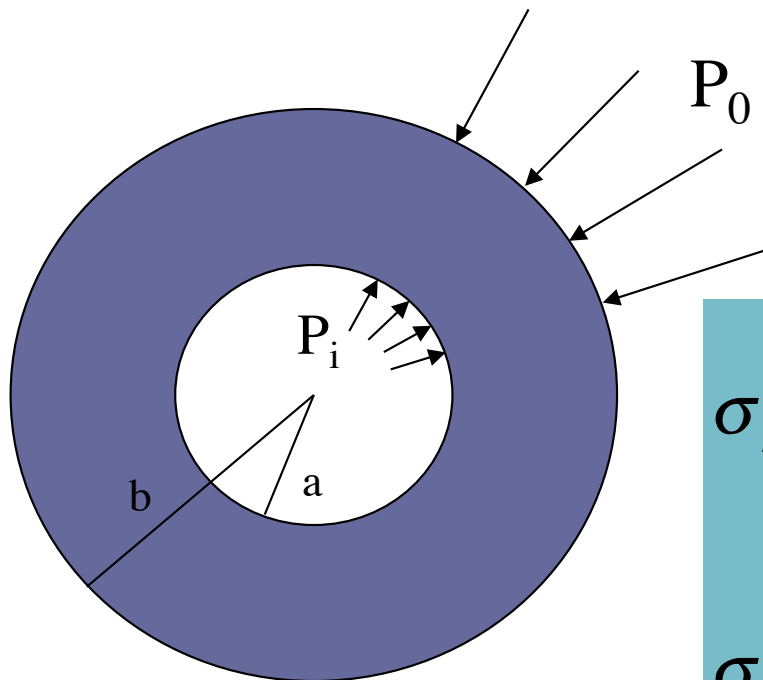
- **Axisymmetry**, i.e. no  $\theta$  dependence

$$\frac{\partial}{\partial \theta} \equiv 0$$

- Solution (Done in class...**Example 8-6**)

# Axisymmetric Domains (cont.)

- Thick walled cylinder with internal and external pressure:



$$\sigma_r = \frac{(P_0 - P_i)a^2b^2}{(b^2 - a^2)} \frac{1}{r^2} - \frac{(P_0b^2 - P_ia^2)}{(b^2 - a^2)}$$
$$\sigma_\theta = -\frac{(P_0 - P_i)a^2b^2}{(b^2 - a^2)} \frac{1}{r^2} - \frac{(P_0b^2 - P_ia^2)}{(b^2 - a^2)}$$
$$\sigma_{r\theta} = 0$$

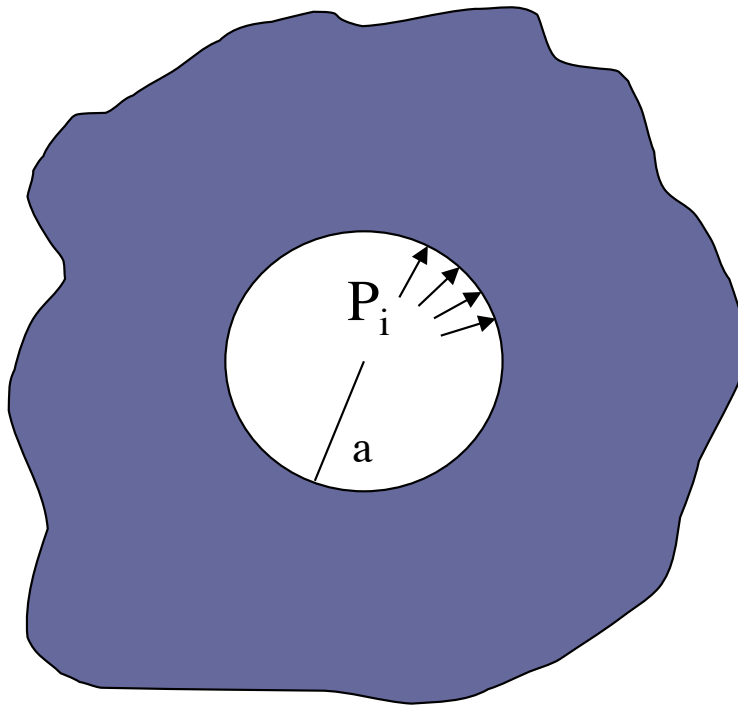
- Note: in the thin walled cylinder  $\sigma_\theta$  does NOT depend on  $r$ .



## Axisymmetric Domains (cont.)

### Special Cases:

- “Infinite cylinder” with internal pressure:  $P_0 = 0$  and  $b \rightarrow \infty$



$$\left. \begin{aligned} \sigma_r &= \frac{-P_i a^2 b^2}{(b^2 - a^2)} \frac{1}{r^2} - \frac{-P_i a^2}{(b^2 - a^2)} \\ \sigma_\theta &= -\frac{-P_i a^2 b^2}{(b^2 - a^2)} \frac{1}{r^2} - \frac{-P_i a^2}{(b^2 - a^2)} \end{aligned} \right\}$$

 $\Rightarrow$ 

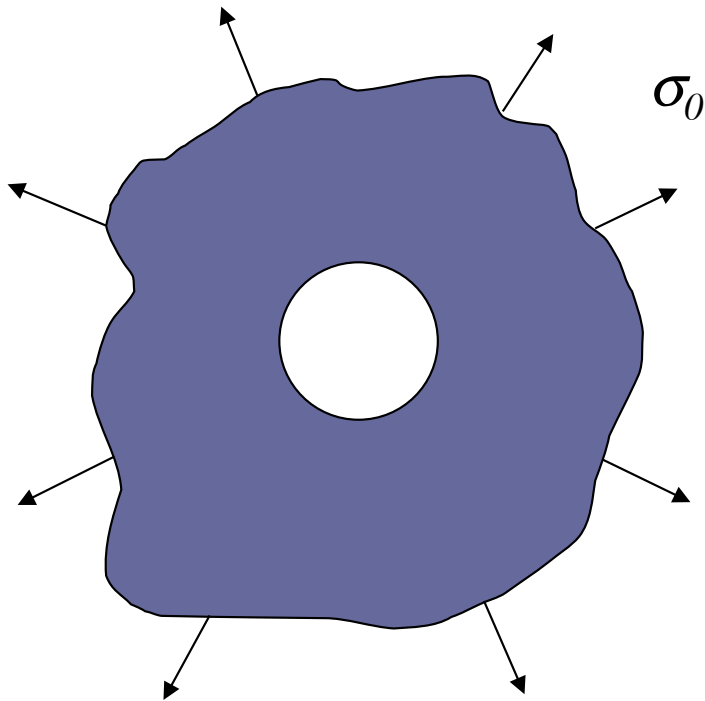
$$\begin{aligned} \sigma_r &= -P_i \left( \frac{a}{r} \right)^2 \\ \sigma_\theta &= P_i \left( \frac{a}{r} \right)^2 \end{aligned}$$

- Stresses decay to 0 far away from the hole since load is self equilibrating

 $\Rightarrow$  St. Venant!

## Axisymmetric Domains (cont.)

- “Infinite cylinder” with external tension:  $P_i = 0$ ,  $P_0 = -\sigma_0$ , and  $b \rightarrow \infty$



$$\left. \begin{aligned} \sigma_r &= -\frac{\sigma_0 a^2 b^2}{(b^2 - a^2)} \frac{1}{r^2} + \frac{\sigma_0 b^2}{(b^2 - a^2)} \\ \sigma_\theta &= \frac{\sigma_0 a^2 b^2}{(b^2 - a^2)} \frac{1}{r^2} + \frac{\sigma_0 b^2}{(b^2 - a^2)} \end{aligned} \right\}$$

$$\Rightarrow \begin{cases} \sigma_r = \sigma_0 \left\{ 1 - \frac{a^2}{r^2} \right\} \\ \sigma_\theta = \sigma_0 \left\{ 1 + \frac{a^2}{r^2} \right\} \end{cases}$$

- As  $r \rightarrow \infty$   $\sigma_r = \sigma_0$   
 $\sigma_\theta = \sigma_0$

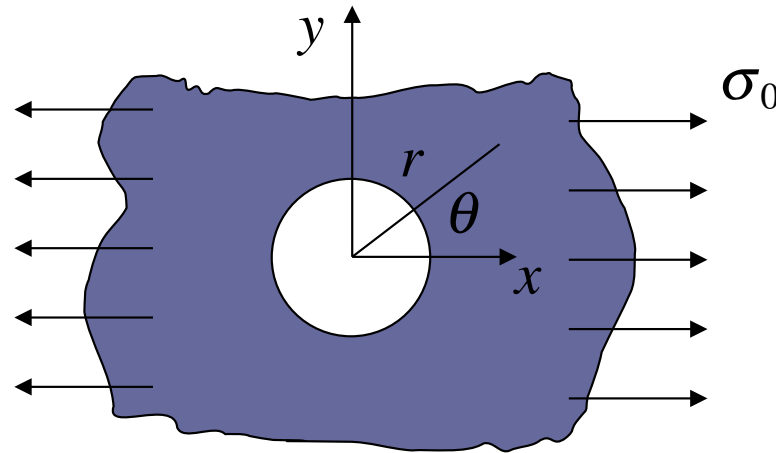
- At  $r=a$   $\sigma_r = 0$   $\Rightarrow$  **Stress concentration**  
 $\sigma_\theta = 2\sigma_0$

# Axisymmetric Domains (cont.)

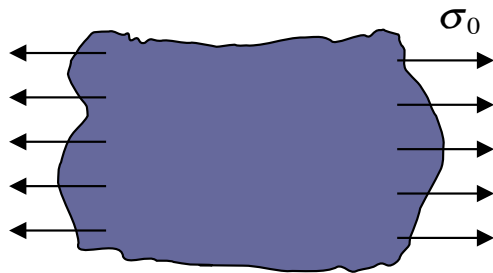
- Hole in uniaxial tension

Example 8-7 in textbook

**NOT an axisymmetric problem!**

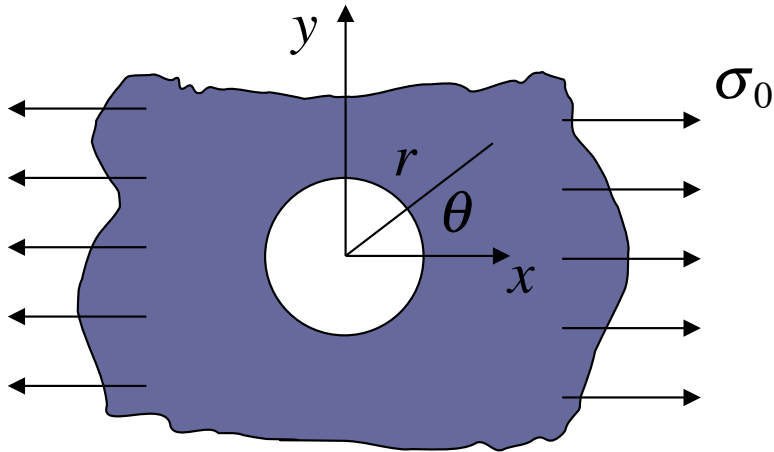


- At far field we get uniaxial tension:  $\sigma_x = \sigma_0$  or  $\sigma_r = \frac{\sigma_0}{2}(1 + \cos 2\theta)$   
 $\sigma_y = \sigma_{xy} = 0$   $\sigma_\theta = \frac{\sigma_0}{2}(1 - \cos 2\theta)$   
 $\sigma_{r\theta} = -\frac{\sigma_0}{2}\sin 2\theta$



$$\Rightarrow \Phi_I(x, y) = \frac{\sigma_0}{2} y^2 = \frac{\sigma_0}{2} r^2 \sin^2 \theta = \frac{\sigma_0}{4} r^2 (1 - \cos 2\theta)$$

## Axisymmetric Domains (cont.)



- Try:

$$\Phi = f(r)\cos 2\theta + \text{axisymmetry}$$

$$\Rightarrow \Phi(r, \theta) = f(r)\cos 2\theta + \cancel{Ar^2} \ln r + B \ln r + Cr^2 + D$$

0

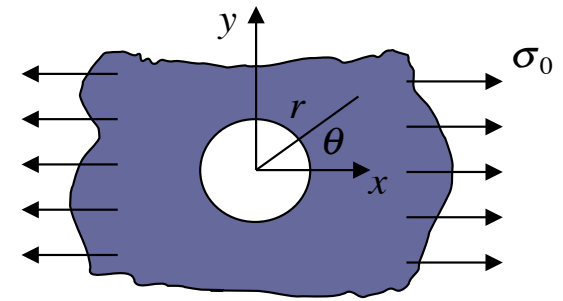
- Solution ...

## Axisymmetric Domains (cont.)

$$\sigma_r = \frac{\sigma_0}{2} \left[ 1 - \left( \frac{a}{r} \right)^2 \right] + \frac{\sigma_0}{2} \left[ 1 - 4 \left( \frac{a}{r} \right)^2 + 3 \left( \frac{a}{r} \right)^4 \right] \cos 2\theta$$

$$\sigma_\theta = \frac{\sigma_0}{2} \left[ 1 + \left( \frac{a}{r} \right)^2 \right] - \frac{\sigma_0}{2} \left[ 1 + 3 \left( \frac{a}{r} \right)^4 \right] \cos 2\theta$$

$$\sigma_{r\theta} = -\frac{\sigma_0}{2} \left[ 1 + 2 \left( \frac{a}{r} \right)^2 - 3 \left( \frac{a}{r} \right)^4 \right] \sin 2\theta$$



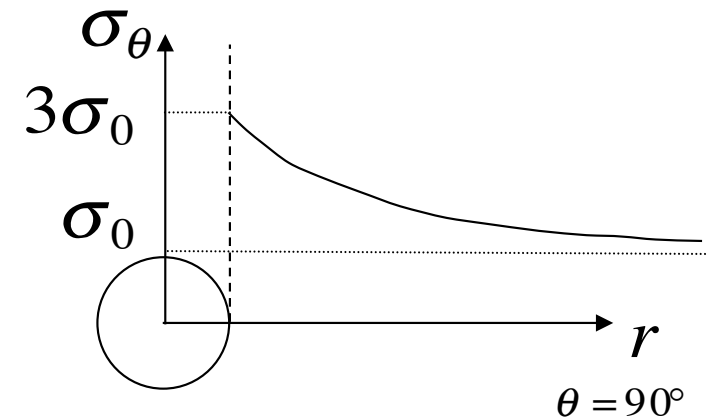
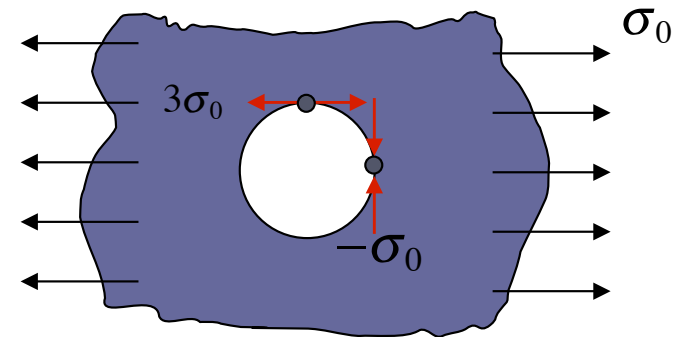
- As  $r \rightarrow \infty$  we recover the uniaxial stress state

$$\left. \begin{array}{l} \text{• But at the hole, } r=a : \\ \sigma_r = 0 \\ \sigma_\theta = \sigma_0 - 2\sigma_0 \cos 2\theta \\ \sigma_{r\theta} = 0 \end{array} \right\} \Rightarrow \left. \begin{array}{l} \sigma_\theta^{\max} \left( a, \frac{\pi}{2} \right) = 3\sigma_0 \\ \sigma_\theta^{\min} (a, 0) = -\sigma_0 \end{array} \right\} \Rightarrow$$

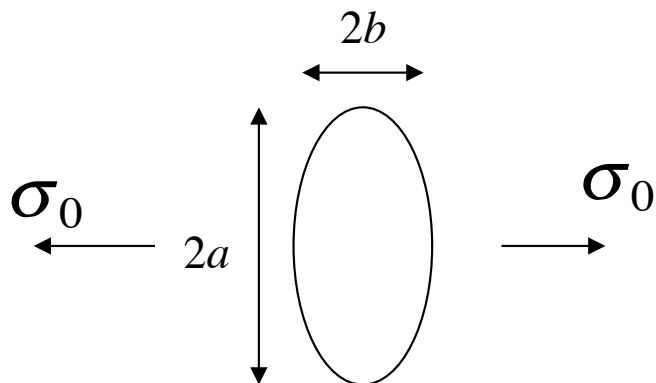
# Axisymmetric Domains (cont.)

- Define **Stress Concentration Factor**,  $k$

$$k = \frac{\sigma_{\theta}^{\max}}{\sigma_0} = 3$$



- For an elliptical hole



$$k = \frac{\sigma_{\max}}{\sigma_0} = 1 + 2\frac{a}{b}$$

Note: As  $b \rightarrow 0$   
 $k \rightarrow \infty$

# Axisymmetric Domains (cont.)

De Havilland DH-106 Comet - First commercial jet airliner



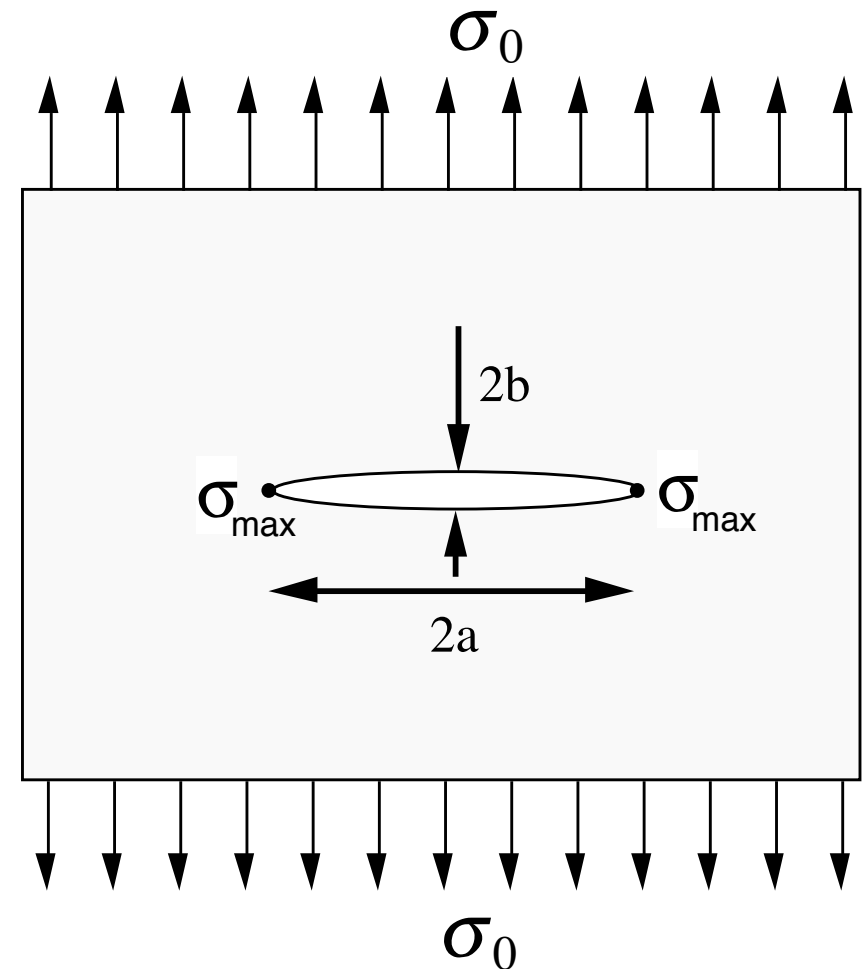
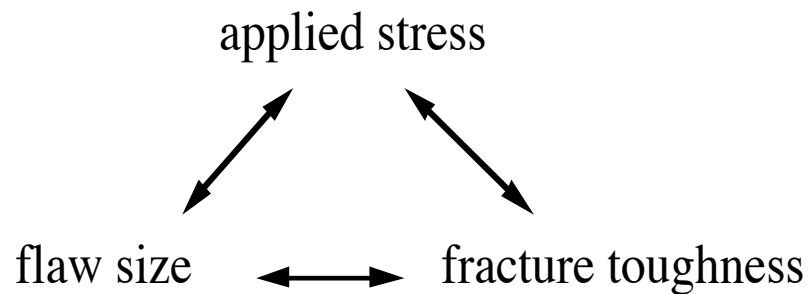
- 3/3/53 Canadian-Pacific Airlines,
- 5/2/53 British Overseas Airways,
- 1/10/54 British Overseas Airways,
- 4/8/54 British Overseas Airways,

Karachi, Pakistan  
Jalalogori, India  
Elba, Italy  
Stromboli, Italy

## Axisymmetric Domains (cont.)

Fracture Mechanics based  
approach in design

$$\frac{\sigma_{\max}}{\sigma_0} = 1 + 2 \frac{a}{b}$$

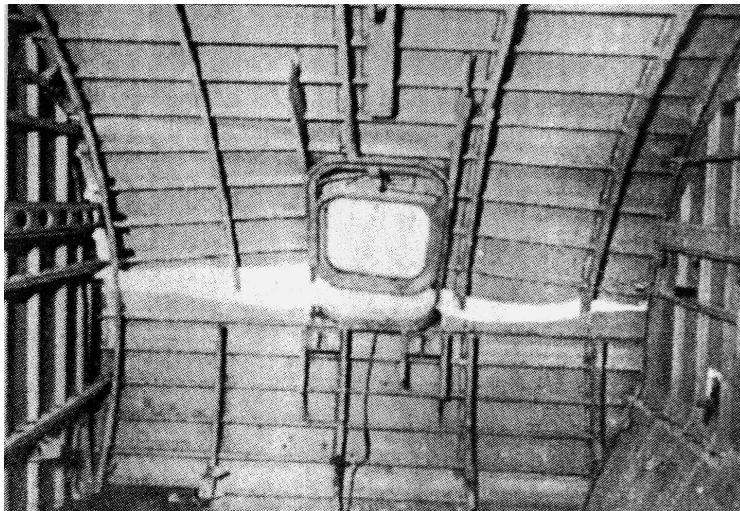




## Axisymmetric Domains (cont.)



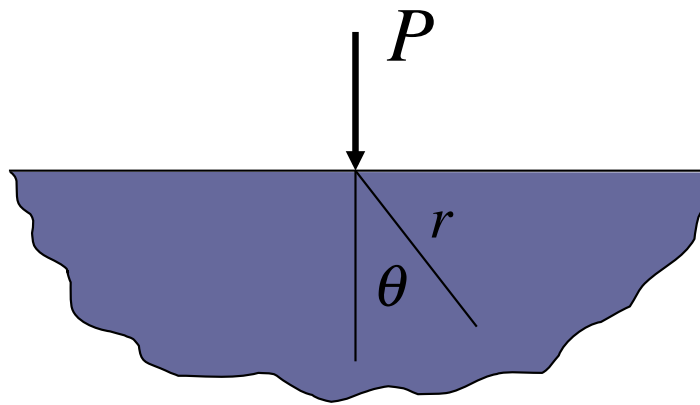
Old design



New design

# Half-space Problems

- Point load on half-space:



$$\Phi(r, \theta) = -\frac{P}{\pi} r \theta \sin \theta$$

 $\Rightarrow$ 

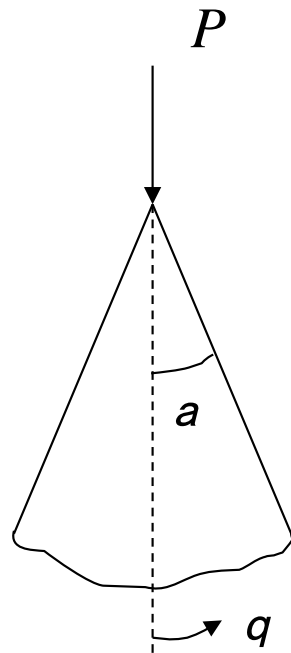
$$\left. \begin{aligned} \sigma_r &= -\frac{2P}{\pi} \frac{\cos \theta}{r} \\ \sigma_\theta &= \sigma_{r\theta} = 0 \end{aligned} \right\}$$

$$u_r = -\frac{2P}{\pi E'} \cos \theta \ln r - \left( \frac{1-\nu'}{\nu'} \right) \frac{P}{\pi} \theta \sin \theta + R.B.$$

$$u_\theta = \frac{2P}{\pi E'} (\nu' + \ln r) \sin \theta - \left( \frac{1-\nu'}{E'} \right) \frac{P}{\pi} \theta \cos \theta + \left( \frac{1-\nu'}{E'} \right) \frac{P}{\pi} \sin \theta + R.B.$$

# Half-space Problems (cont.)

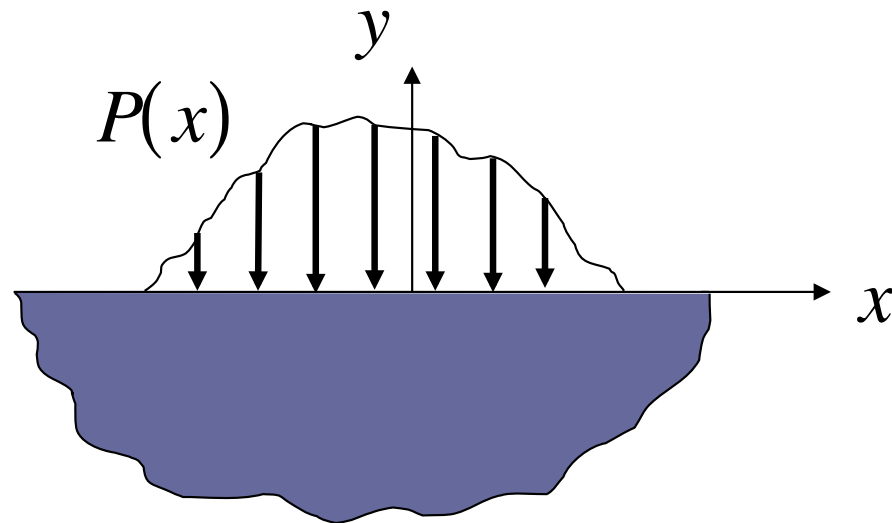
- Wedge with load at tip (Example 8-8):



- Solution...

## Half-space Problems (cont.)

- Arbitrary pressure distribution:



- Solution...

END OF CHAPTER 8