MECH 503 Final Exam

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$$BC_{1} = \sigma_{xy}(x, y = \pm \frac{h}{2}) = 0$$

$$BC_{2} = \sigma_{y}(x, y = \frac{h}{2}) = 0$$

$$BC_{3} = \sigma_{y}(x, y = \frac{-h}{2}) = -q$$

$$BC_{4} = \int_{-\frac{h}{2}}^{\frac{h}{2}} \sigma_{x}(x = \pm L, y) = 0$$

$$BC_{5} = \int_{-\frac{h}{2}}^{\frac{h}{2}} \sigma_{x}(x = \pm L, y)y = 0$$

$$BC_{6} = \int_{-\frac{h}{2}}^{\frac{h}{2}} \sigma_{xy}(x = \pm L, y) = \pm Lq$$

One could observe the boundary conditions are satisfied in classical and weak sense. For the normal stress in x-direction:

$$\sigma_x(x = \pm L, y) = \pm \frac{qy(3h^2 - 20y^2)}{5h^3}$$

which is symmetric about the y-axis, thus, the moment is balanced. The distribution of shear

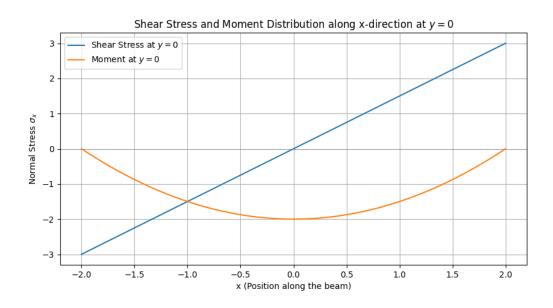


Figure 1: Shear Stress and Moment Distribution along x-direction at y = 0

stress and moment is shown in figure 1, since no specific property values are given, one could conclude that the general shape of the distribution agrees with the theory.

The Python code is provided in the Appendix file.

Part 1

$$\sigma_{\theta\theta}(r=a) = 0$$
 $\sigma_{r\theta}(r=a) = 0$

Part 2

$$\sigma_{rr}(\frac{r}{a} = \infty) = \sigma_0 \sin^2(\theta)$$

$$\sigma_{r\theta}(\frac{r}{a} = \infty) = -\frac{\sigma_0 \sin(2\theta)}{2}$$

$$\sigma_{\theta\theta}(\frac{r}{a} = \infty) = \sigma_0 \sin^2(\theta)$$

Part 3

$$\sigma_{xx}(\frac{r}{a} = \infty) = -\sigma_0 \sin^2(\theta) \cos(2\theta)$$

$$\sigma_{yy}(\frac{r}{a} = \infty) = -\frac{\sigma_0 \sin(4\theta)}{4}$$

$$\sigma_{xy}(\frac{r}{a} = \infty) = -\sigma_0 \sin^2(\theta) \cos(2\theta)$$

Part 4

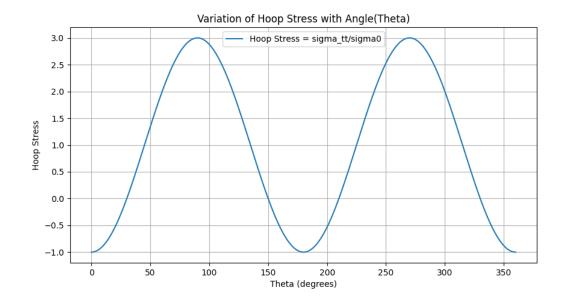


Figure 2: Hoop Stress vs. $Angle(\theta)$

The maximum value of hoop stress is 3, which occurs at $\theta = 90.0$ degrees.

2.1 Part 5

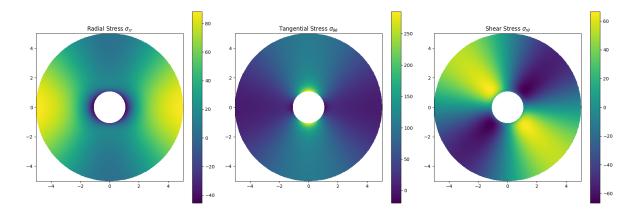


Figure 3: Stress Distribution around the Hole

The radius is assumed to be 1m which is different from the given value, however, the general idea is the same.

3.1 Part A

By equilibrium equation and divergence theorem:

$$\rho \frac{\partial^2 u_i}{\partial t^2} = \sigma_{ij,j}$$

The stress tensor is given by:

$$\sigma_{ij} = C_{ijkl} \epsilon_{kl}$$

where C_{ijkl} are the elastic stiffness coefficients, and the infinitesimal strain tensor is given by:

$$\epsilon_{kl} = \frac{1}{2}(u_{k,l} + u_{l,k})$$

Substituting:

$$\sigma_{ij,j} = (C_{ijkl}\epsilon_{kl})_{,j} = C_{ijkl}\epsilon_{kl,j}$$

Since C_{ijkl} are constants with respect to spatial coordinates:

$$\sigma_{ij,j} = C_{ijkl} \frac{1}{2} (u_{k,lj} + u_{l,kj})$$

Since $C_{ijkl} = C_{ijlk}$ and the first equation,

$$C_{ijkl}u_{k,lj} = \rho \frac{\partial^2 u_i}{\partial t^2}$$

$$\rho \frac{\partial^2 u_i}{\partial t^2} = C_{ijkl} \frac{\partial^2 u_k}{\partial x_i \partial x_l}$$

3.2 Part b

$$\begin{aligned} u_i &= a_i \exp[i\omega(t - \frac{x_j p_j}{c})] \\ u_{i,j} &= -\frac{i\omega p_j}{c} a_i \exp[i\omega(t - \frac{x_j p_j}{c})] \\ u_{i,jk} &= -\frac{i\omega p_j}{c} (-\frac{i\omega p_k}{c}) a_i \exp[i\omega(t - \frac{x_j p_j}{c})] \\ &= \frac{\omega^2}{c^2} p_j p_k a_i \exp[i\omega(t - \frac{x_j p_j}{c})] \\ C_{ijkl} \frac{\omega^2}{c^2} p_l p_k a_k \exp[i\omega(t - \frac{x_j p_j}{c})] &= \rho \omega^2 a_i \exp[i\omega(t - \frac{x_j p_j}{c})] \\ C_{ijkl} p_l p_k a_k &= \rho c^2 a_i \end{aligned}$$

This last equation states that the product of C_{ijkl} with p_lp_k must equal ρc^2 times the identity matrix when contracted with a_k . This implies that a_i and p_i must be eigenvectors of the tensor equation, with c^2 being related to the corresponding eigenvalue. This shows that a plane wave solution is feasible under certain conditions on the wave vector p_i , frequency ω , and phase velocity c.

4 Question 4

This question in coded in Python entirely, thus, for detail information, please refer to the code provided in Appendix file. And also note that the computation process is modified according to the MATLAB code posted on Canvas. Mode 1 here is the speed in longitudinal and mode 2&3 are the speed in shear.

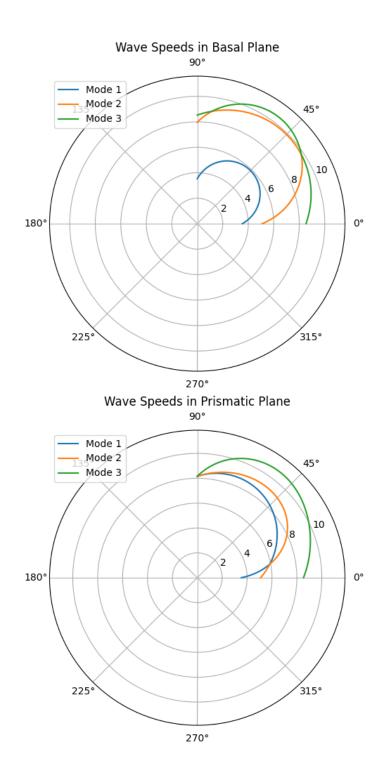


Figure 4: Wave Speed Distribution in Polar Plot

$$u_z = \frac{b}{2\pi}\theta$$

$$\varepsilon_{z\theta} = \frac{1}{2} \left(\frac{\partial u_z}{\partial \theta} \right) = \frac{b}{4\pi}$$

By Hooke's Law:

$$\sigma_{ij} = 2\mu\varepsilon_{ij} + \lambda\delta_{ij}\varepsilon_{kk}$$

$$\sigma_{z\theta} = 2\mu\varepsilon_{z\theta} = \frac{\mu b}{2\pi}$$

The strain energy density W per unit volume can be expressed as:

$$W = \frac{1}{2}\sigma_{ij}\varepsilon_{ij}$$

$$W = \frac{1}{2}\sigma_{z\theta}\varepsilon_{z\theta} = \frac{1}{2}\left(\frac{\mu b}{2\pi}\right)\left(\frac{b}{4\pi}\right) = \frac{\mu b^2}{16\pi^2}$$

The strain energy per unit length of the dislocation line, U, is:

$$U = \int_{A} W \, dA$$

$$U = \int_{0}^{2\pi} \int_{0}^{R} \frac{\mu b^{2}}{16\pi^{2}} r \, dr \, d\theta = \frac{\mu b^{2} R^{2}}{16\pi^{2}}$$

Strain Energy distribution is plotted below, notice some properties are assumed to some number for plotting successfully.

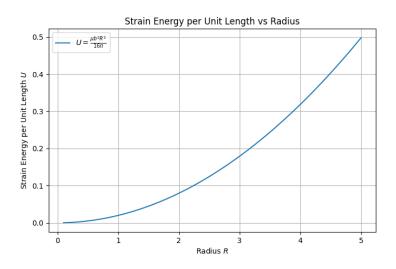


Figure 5: Strain Energy per Unit vs. Radius

$$\mathbf{u} = \begin{pmatrix} -Cx_2 + Bx_3 \\ Cx_1 - Ax_3 \\ -Bx_1 + Ax_2 \end{pmatrix}$$

$$\nabla \mathbf{u} = \begin{pmatrix} 0 & -C & B \\ C & 0 & -A \\ -B & A & 0 \end{pmatrix}$$

The vorticity tensor is given by:

$$\mathbf{w} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} - \frac{\partial u_j}{\partial x_i} \right)$$
$$= \begin{pmatrix} 0 & -C & B \\ C & 0 & -A \\ -B & A & 0 \end{pmatrix}$$

since w is skew-symmetric,

$$\mathbf{w} = -\mathbf{w}^{T}$$

$$\mathbf{E} = \frac{1}{2} \left(\frac{\partial u_{i}}{\partial x_{j}} + \frac{\partial u_{j}}{\partial x_{i}} \right)$$

$$= \begin{pmatrix} 0 & -C & B \\ C & 0 & -A \\ -B & A & 0 \end{pmatrix} + \begin{pmatrix} 0 & C & -B \\ -C & 0 & A \\ B & -A & 0 \end{pmatrix}$$

$$= \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

Thus, since all the components in **E** is zero, it is a rigid body rotation.