MATH 521 - Numerical Analysis of Differential Equations

Christoph Ortner, 03/2024

Assignment 4: Heat Equation

Name:

Student ID:

In this assignment we will solve the follow heat equation (prototype diffusion equation)

$$egin{aligned} u_t - \Delta u &= f, & x \in \Omega, t \in (0,T], \ u &= 0, & x \in \Omega, t \in [0,T], \ u &= u_0, & x \in \Omega, t = 0. \end{aligned}$$

where the solution is a function u(x,t) and T is the final time. Assume that $u_0 \in H^2(\Omega) \cap H^1_0(\Omega)$, and f = f(x,t) is arbitrarily smooth in space and time. Throughout, you may assume that the solution u(t,x) exists, is unique, and as smooth as needed to carry out your analysis.

We first discretize the problem using a P_1 -FEM: Let \mathcal{T}_h be a regular triangulation of the polygonal domain $\Omega \subset \mathbb{R}^2$ and let $V_h := \mathcal{P}_1(\mathcal{T}_h)$ be the corresponding FE space. Then the semi-discrete (continuous in time) formulation is to find $u_h \in C^1([0,T],V_h)$, such that $u_h(t=0) = I_h u_0$ (with I_h the nodal interpolation operator), and

$$(u_{h,t},v_h)+(
abla u_h,
abla v_h)=(f,\eta^{n+1}+\eta^n) \qquad orall v_h\in V_h, t\in (0,T].$$

with $(\cdot, \cdot) = (\cdot, \cdot)_{L^2(\Omega)}$. The Assignment is primarily concerned with the effect of using different discretization schemes in time.

Throughout this assignment we assume that the family \mathcal{T}_h , h>0 is uniformly shape-regular (no need to further mention this) and quasiuniform, i.e. $ch \leq h_T \leq h$ for all $T \in \mathcal{T}_h$ with c independent of h.

Q1: Crank-Nicholson Scheme [10]

For a system of ODEs, $\dot{u}=F(t,u)$ the CN scheme reads,

$$U^{n+1} = U^n + \Delta t \frac{F(t_n, U^n) + F(t_{n+1}, U^{n+1})}{2}$$

where $t_n = n\Delta t$, Δt is the time-step and U^n the approximation to $u(t_n)$.

- (a) Formulate the CN discretization of the semi-discrete FEM.
- (b) Under suitable regularity assumptions on the solution \emph{u}_{i} prove that

$$\max_{n=1,\dots,N} \|u(t_n) - u_h^n\|_{L^2} \leq C(h^2 + \Delta t^2)$$

with $N\Delta t \leq T$. How does C depend on N (or on T) in your result?

For part-b there is a part of the proof that is very repetitive (basically the same as in class) and I won't make you go through it again. I'll very briefly sketch this out, and you can then perform the last step.

Solution Q1a

$$\left(rac{u_h^{n+1}-u_h^n}{\Delta t},v_h
ight)+rac{1}{2}ig((
abla u_h^n,
abla v_h)+(
abla u_h^{n+1},
abla v_h)ig)=rac{1}{2}ig((f^n,v_h)+(f^{n+1},v_h)ig) \quad orall v_h\in V_h$$

Solution Q1b

Sketch of Part 1:

Let $u^n(x)=u(x,t_n)$ and $\tilde{u}^n_h\in V_h$ the Ritz projection of u^n i.e.

$$(\nabla (u^n - \tilde{u}_h^n), \nabla v_h) = 0 \quad \forall v_h \in V_h$$

We split the error

$$e_h^n = u_h^n - u^n = (u_h^n - \tilde{u}_h^n) + (\tilde{u}_h^n - u^n) = \eta^n + \epsilon^n.$$

Following the class notes almost verbatim we can then prove

$$(\eta^{n+1}-\eta^n,v_h)+rac{\Delta t}{2}(
abla\eta^{n+1}+
abla\eta^n,
abla v_h)=rac{\Delta t}{2}(g^n,v_h)$$

where

$$||g^n||_{L^2} \le C_1(h^2 + \Delta t^2)$$

with C_1 depending on $\|\nabla^2 u_t\|_{L^\infty(L^2)}$ and $\|u_{ttt}\|_{L^\infty(L^2)}$, but independent of $h, n, \Delta t$.

Part 2: please complete this.

Hint:
$$(u-v, u+v)_H = ||u||_H^2 - ||v||_H^2$$

Choose $v_h = \eta^{n+1} + \eta^n$, yielding:

$$(\eta^{n+1}-\eta^n,\eta^{n+1}+\eta^n)+\frac{\Delta t}{2}(\nabla\eta^{n+1}+\nabla\eta^n,\nabla\eta^{n+1}+\nabla\eta^n)=\frac{\Delta t}{2}(g^n,\eta^{n+1}+\eta^n)$$

Simplifying the left-hand side using the identity provided in the hint, and the C-S inequality gives:

$$||\eta^{n+1}||^2 - ||\eta^n||^2 + \frac{\Delta t}{2}||\nabla \eta^{n+1}||^2 - ||\nabla \eta^n||^2 \le \frac{\Delta t}{2}||g^n|| ||\eta^{n+1} + \eta^n||.$$

Above is what I got for Q2, I realize that next steps should be related to the Discrete Gronwall's Lemma, but I do not know how to process from that. Here is the Lemma I searched from Google.

Gronwall's Lemma (Discrete Version)

The recurrence relation above, if summed from n=0 to n=N-1 and applying the discrete version of Gronwall's lemma, suggests that:

$$||\eta^N||^2 \leq \expigg(C\Delta t \sum_{n=0}^{N-1} 1igg) \left(||\eta^0||^2 + C_1^2 (h^2 + \Delta t^2)^2 T
ight),$$

where C is a constant depending on the domain and coefficients of the equation, not on N or T. Given that $\eta^0=0$ (since $u_h^0=I_hu_0$ and $u^0=u_0$), this simplifies to:

$$\|\eta^N\| \leq C(h^2 + \Delta t^2),$$

where now the constant may depend on T but not explicitly on N, other than through the product $N\Delta t \leq T$.

Though the detail of the Lemma is not clear to me, but observe from the result, one can conclude: Combining the bounds for η^n and ϵ^n :

$$\max_{n=1,\dots,N}\|u(t_n)-u_h^n\|\leq \max_{n=1,\dots,N}(\|\eta^n\|+\|\epsilon^n\|)\leq C(h^2+\Delta t^2).$$

Thus, the overall constant C in the error estimate reflects dependencies from the norms of higher derivatives of u and the final time T, but crucially it does not depend explicitly on N. This result ensures that the Crank-Nicolson scheme provides a robust and effective method for the numerical approximation of the heat equation in both space and time.

Q2: Inverse Estimate [5]

Under the assumptions on \mathcal{T}_h stated at the beginning of the assignment, prove that there exists $c^i>0$ such that

$$\|\nabla v_h\|_{L^2} \le c^i h^{-1} \|v_h\|_{L^2} \qquad \forall v_h \in \mathcal{P}_1(\mathcal{T}_h).$$

Hint: transform to the reference element.

Solution Q2

Each element T in the triangulation \mathcal{T}_h can be mapped to a reference element \hat{T} by:

$$x = B_T \hat{x} + b$$

And the area of each element ${\cal T}$

$$|T| = |\det(B_T)||\hat{T}|,$$

Thus, the gradients transform as follows:

$$\nabla v_h(x) = B_T^{-1} \nabla \hat{v}_h(\hat{x})$$
.

Therefore, the norm of the gradient scales as:

$$\|
abla v_h \|_{L^2(T)}^2 = \int_T |
abla v_h |^2 dx = \int_{\hat{T}} | B_T^{-1}
abla \hat{v}_h |^2 |\det(B_T)| \, d\hat{x}.$$

Using the norm of a matrix and its inverse, particularly $||B_T^{-1}||$ where the norm is the operator norm (maximum stretching factor), we find:

$$||B_T^{-1}|| \le Ch^{-1},$$

assuming that the transformation linearly scales with the diameter h of the elements (since $h_T \le h$ and is quasiuniform). Thus, we estimate:

$$\|\nabla v_h\|_{L^2(T)}^2 \le \|B_T^{-1}\|^2 |\det(B_T)| \|\nabla \hat{v}_h\|_{L^2(\hat{T})}^2.$$

Thus

$$\|\nabla v_h\|_{L^2(\Omega)}^2 \leq C^2 h^{-2} \|v_h\|_{L^2(\Omega)}^2$$
.

And finally,

$$\|\nabla v_h\|_{L^2(\Omega)} \le c^i h^{-1} \|v_h\|_{L^2(\Omega)}.$$

Q3: Explicit Euler [15]

For a system of ODEs, $\dot{u}=F(t,u)$ the Explicit Euler (EE) scheme reads,

$$U^{n+1} = U^n + \Delta t F(t_n, U^n),$$

where $t_n = n\Delta t$, Δt is the time-step and U^n the approximation to $u(t_n)$.

- (a) Formulate the EE discretization of the semi-discrete FEM.
- (b) For f=0 and under a suitable restriction on Δt and h that you should derive, prove the **CONDITIONAL STABILITY** result

$$||u_h^n||_{L^2} \le q_{\Delta t} ||u_h^{n+1}||_{L^2}$$

where $q_{\Delta t} < 1$.

(NOTE: a full error estimate is a bit more involved, so we only prove this stability estimate instead. This is still surprisingly hard, so I will give you the outline of the proof so you can follow it.)

Solution Q3a

Solution Q3b

Step 1: By testing with $v_h = u_h^{n+1} + u_h^n$ (or otherwise) show that

$$||u_h^{n+1}||^2 + \frac{\Delta t}{2}||\nabla u_h^{n+1}||^2 = ||u_h^n||^2 - \frac{\Delta t}{2}||\nabla u_h^n||^2 + \frac{\Delta t}{2}||\nabla u_h^{n+1} - \nabla u_h^n||^2$$

HINT: $(u,v)_H = \frac{1}{2} ||u||_H^2 + \frac{1}{2} ||v||_H^2 - \frac{1}{2} ||u-v||_H^2$

Step 2: By testing with $v_h=u_h^{n+1}-u_h^n$ and applying the inverse inequality prove that

$$||u_b^{n+1} - u_b^n||^2 \le \Delta t \mu ||\nabla u_b^n||^2,$$

where you should determine μ in terms of $c^i, h, \Delta t$.

Step 3: Apply the inverse inequality again to the result Step-1, then apply Step-2, then collect your terms. The result should now follow under a restriction on μ that you should specify.

Q4: Implementation of Heat Equation [20]

For this assignment you may choose any code-base you like, our first implementation of P1-FEM, the Ferrite implementation, or any other code in Julia or Python or Matlab. This is a little harder than the previous coding assignments in that I'm giving you a lot less help.

(a) Select one of the three discretizations of the heat equation that we covered: P1-FEM in space and IE, EE or CN in time. Implement this scheme to solve the heat equation on the time interval $t \in [0,1]$ with $\Omega = (-1,1)^2$, f(x,t) = 1 and $u_0(x) = 0$.

Visualize the solution at the final time T=1 and print out the value of $\max_{x\in\Omega}u(x,t=1)$. (e.g. put this in the title of the figure)

(b) Design and implement a numerical test that demonstrates the convergence rate we proved numerically.

Solution Q4a

```
In [ ]:
In [ ]:
```

Solution Q4b

```
In []: # to use the method of manufactured solutions the following # may be useful. (or some variant...)  \begin{array}{l} u_-ex = (x,\,t) \rightarrow t * \cos(x[1] * t + x[2] * \sin(t)) * (x[1]^2 - 1) * (x[2]^2 - 1) \\ \nabla^2 u_-ex = (x,\,t) \rightarrow \text{ForwardDiff.hessian}(x \rightarrow u_-ex(x,\,t),\,x) \\ \Delta u_-ex = (x,\,t) \rightarrow \text{tr}(\nabla^2 u_-ex(x,\,t)) \\ \partial_t u_-ex = (x,\,t) \rightarrow \text{ForwardDiff.derivative}(t \rightarrow u_-ex(x,\,t),\,t) \\ f_-ex = (x,\,t) \rightarrow \partial_t u_-ex(x,\,t) - \Delta u_-ex(x,\,t) \end{array}  In []:
```