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Finding time evolution operator for interaction picture Hamiltonian in resonant case
\hat{H}_{int} = \lambda \hbar (\hat{a}\hat{\sigma}_{t} + \hat{a}^{\dagger}\hat{\sigma}_{-})
Use basis | e, n-1) (excited atom, n-1 photons) and | g, n) (ground atom, n photons)
\hat{H}_{int}|e,n-1\rangle = 2 \pm \sqrt{n}|g,n\rangle

\hat{H}_{int}|g,n\rangle = 2 \pm \sqrt{n}|e,n-1\rangle
Time evolution operator: \hat{\mathcal{U}}(t) = e^{-\frac{i\hat{H}_{int}t}{\hbar}}
                                                             = \sum_{k=0}^{\infty} \frac{1}{k!} \left( -\frac{1}{k!} \hat{H}_{ink} t \right)^{k}
\widehat{\mathcal{U}}(t)|e,n-1\rangle = \frac{\infty}{t} \frac{(-1)^{\frac{k}{2}} (\lambda \sqrt{n} t)^{\kappa} |e,n-1\rangle - i \sum_{k=1}^{\infty} \frac{(-1)^{\frac{k-1}{2}} (\lambda \sqrt{n} t)^{\kappa} |g,n\rangle}{k!}
                             = cos(\lambda \sqrt{n} t)|e,n-1\rangle - isin(\lambda \sqrt{n} t)|g,n\rangle
\hat{\mathcal{U}}(t)|g,n\rangle = -i\sum_{k=1,3\cdots}^{\infty} \frac{(-1)^{\frac{k-1}{2}}(\lambda |n|t)^{k}|e,n-1\rangle}{k!} + \sum_{k=n,2\cdots}^{\infty} \frac{(-1)^{\frac{k}{2}}(\lambda |n|t)^{k}|g,n\rangle}{k!} + \sum_{k=n,2\cdots}^{\infty} \frac{(-1)^{\frac{k}{2}}(\lambda |n|t)^{k}|g,n\rangle}{k!}
                         = -isin(\lambda \sqrt{n} t)|c,n-1> + \cos(\lambda \sqrt{n} t)|g,n>
= \hat{\mathcal{U}}(t) = \cos(\lambda \sqrt{n} t) |e, n-1\rangle \langle e, n-1| - i\sin(\lambda \sqrt{n} t) |e, n-1\rangle \langle g, n| - \cdots
              \cdots - isin(\lambda \sqrt{n} t)|g,n>\(e,n-1| + cos(\lambda \sqrt{n} t)|g,n>\(g,n|
If the system starts in the state |\Psi(0)\rangle = |g, n\rangle
Then |\Psi(t)\rangle = \hat{\mathcal{U}}(t)|g,n\rangle = -i\sin(\lambda \sqrt{n} t)|e,n-1\rangle + \cos(\lambda \sqrt{n} t)|g,n\rangle
=> Probability of measuring atom in excited state: P_e(t) = |\langle e, n-1 | \psi(0) \rangle|^2 = \sin^2(\lambda | n | t)
       Similar for ground state: P_g(t) = \cos^2(\lambda \ln t)
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