

MPS PRACTICE PROBLEMS 2

SOLUTIONS

$$1 \text{ (a)} \quad E = \frac{3}{2} N k_B T$$
$$\Rightarrow T = \frac{2}{3} \frac{E}{N k_B}$$

Then

$$S = N k_B \left(\frac{3}{2} \ln T + \ln V + \kappa \right)$$

$$= N k_B \left(\frac{3}{2} \ln \left(\frac{2}{3} \frac{E}{N k_B} \right) + \ln V + \kappa \right)$$

$$= N k_B \left(\frac{3}{2} \ln E - \frac{3}{2} \ln N \right. \\ \left. + \ln V + \ln \left(\frac{2}{3 k_B} \right) + \kappa \right)$$

$$= N k_B \left(\frac{3}{2} \ln E - \frac{3}{2} \ln N \right. \\ \left. + \ln V + \tilde{\kappa} \right)$$

$$\text{where } \tilde{\kappa} = \ln \left(\frac{2}{3 k_B} \right) + \kappa.$$

(b) In (a), we have $S(E, V, N)$, the entropy is expressed in terms of its "natural" variables. There are three equations of state that result from partial derivative of S with respect to its natural variables. They are

$$\begin{aligned} \bullet \quad \frac{1}{T} &= \frac{\partial S(E, V, N)}{\partial E} \\ &= N k_B \cdot \frac{3}{2} \cdot \frac{1}{E} \end{aligned}$$

$$\Rightarrow \boxed{E = \frac{3}{2} N k_B T}$$

$$\begin{aligned} \bullet \quad \frac{P}{T} &= \frac{\partial S(E, V, N)}{\partial V} \\ &= N k_B \cdot \frac{1}{V} \end{aligned}$$

$$\Rightarrow \boxed{P V = N k_B T}$$

$$\bullet \quad \frac{\mu}{T} = - \frac{\partial S(E, V, N)}{\partial N}$$

$$\begin{aligned} &= - k_B \left(\frac{3}{2} \ln E - \frac{3}{2} \ln N + \ln V + \tilde{\kappa} \right) \\ &\quad - N k_B \left(-\frac{3}{2} \right) \frac{1}{N} \end{aligned}$$

$$\begin{aligned} &= - k_B \left(\frac{3}{2} \ln E - \frac{3}{2} \ln N + \ln V \right. \\ &\quad \left. - \frac{3}{2} + \tilde{\kappa} \right) \end{aligned}$$

$$\Rightarrow \mu = k_B T \left(-\frac{3}{2} \ln E + \frac{3}{2} \ln N - \ln V + \frac{3}{2} - \tilde{\kappa} \right).$$

To completely determine μ , we need to know $\tilde{\kappa}$ — we will calculate it using quantum mechanics later in the course.

$$2 (a) \quad G = E - TS + pV$$

$$\Rightarrow dG = dE - d(TS) + d(pV)$$

$$= dE - SdT - TdS + Vdp + pdV.$$

Substituting the statement of conservation of energy

$$dE = \delta W + \delta Q + \sum_i \mu_i dN_i$$

$$= -pdV + TdS + \sum_i \mu_i dN_i,$$

we get

$$dG = -\cancel{pdV} + \cancel{TdS} + \sum_i \mu_i dN_i$$

$$- SdT - \cancel{TdS} + Vdp + \cancel{pdV}$$

$$= -SdT + Vdp + \sum_i \mu_i dN_i$$

$\Rightarrow G$ is naturally a function of T ,

p and \vec{N} as the change in G

is determined by the change in T ,

p and \vec{N} using only conservation of

energy (i.e. not using any equations of state, which are specific to a given system).

(b) From (a),

$$dG = -SdT + Vdp + \sum_i \mu_i dN_i, \quad (*)$$

On the other hand, via the definition of partial derivatives

$$\begin{aligned} dG(T, p, N) &= \frac{\partial G(T, p, \vec{N})}{\partial T} dT \\ &+ \frac{\partial G(T, p, \vec{N})}{\partial p} dp \\ &+ \sum_i \frac{\partial G(T, p, \vec{N})}{\partial N_i} dN_i \quad (**)$$

Comparing the coefficients of dT , dp and dN_i in (*) and (**) give the three equations of state

$$\frac{\partial G(T, p, \vec{N})}{\partial T} = -S$$

$$\frac{\partial G(T, p, \vec{N})}{\partial p} = V$$

$$\frac{\partial G(T, p, \vec{N})}{\partial N_i} = \mu_i$$

$$3. \quad S(E, V) = \frac{4}{3} a V^{1/4} E^{3/4}$$

$$(a) \quad \frac{1}{T} = \frac{\partial S(E, V)}{\partial E}$$

$$= \frac{4}{3} a V^{1/4} \frac{3}{4} E^{-1/4}$$

$$= a V^{1/4} E^{-1/4}$$

$$\Rightarrow E^{1/4} = a V^{1/4} T$$

$$\Rightarrow E = a^4 V T^4$$

$$\Rightarrow \text{Energy density } \underline{\frac{E}{V} = a^4 T^4}$$

$$(b) \quad \frac{P}{T} = \frac{\partial S(E, V)}{\partial V}$$

$$= \frac{4}{3} a \cdot \frac{1}{4} V^{-3/4} E^{3/4}$$

$$= \frac{1}{3} a \left(\frac{E}{V} \right)^{3/4}$$

$$= \frac{1}{3} a (a^4 T^4)^{3/4} \text{ from (a)}$$

$$= \frac{1}{3} a \cdot a^3 T^3$$

$$= \frac{1}{3} a^4 T^3$$

$$\Rightarrow \underline{P = \frac{1}{3} a^4 T^4}$$

(c) Free energy

$$F = E - TS$$

$$= a^4 V T^4 - T \cdot \frac{4}{3} a V^{\frac{1}{4}} E^{3/4}$$

$$= a^4 V T^4 - T \cdot \frac{4}{3} a V^{\frac{1}{4}} (a^4 V T^4)^{3/4}$$

using (a)

$$= a^4 V T^4 - \frac{4}{3} a V^{\frac{1}{4}} \cdot T \cdot a^3 V^{3/4} \cdot T^3$$

$$= a^4 V T^4 - \frac{4}{3} a^4 V T^4$$

$$= -\frac{1}{3} a^4 V T^4$$

$$(d) \quad -p = \frac{\partial F(T, V)}{\partial V}$$

$$= -\frac{1}{3} a^4 T^4$$

$$\Rightarrow p = \frac{1}{3} a^4 T^4, \quad \text{consistent with (b).}$$

$$(e) \quad pV = \frac{1}{3} a^4 T^4 \cdot V$$

$$\text{From (a), } T^4 = \frac{E}{a^4 V}$$

$$\Rightarrow pV = \frac{1}{3} a^4 V \frac{E}{a^4 V}$$

$$= \frac{1}{3} E$$

$$(f) \quad dE = dQ + dW$$

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For an adiabatic process, $dQ = 0$.

If the compression is quasi-static,

$$dW = -p dV$$

$$\Rightarrow dE = -p dV \quad \text{--- (*)}$$

From (e), $E = 3 pV$.

Substituting into (*)

$$3p dV + 3V dp = -p dV$$

$$\Rightarrow 0 = 4p dV + 3V dp$$

Divide by $3pV$

$$0 = \frac{4}{3} \frac{dV}{V} + \frac{dp}{p}$$

$$= \frac{4}{3} d \ln V + d \ln p$$

$$= d (\ln V^{4/3} + \ln p)$$

$$= d \ln (p V^{4/3})$$

$$\Rightarrow \ln (p V^{4/3}) = \text{constant}$$

$$\Rightarrow p V^{4/3} = \text{constant.}$$

$$(g) \quad G = E - TS + pV$$

$$= aVT^4 - T \frac{4}{3} a V^{\frac{1}{4}} (a^4 VT^4)^{3/4}$$

$$+ \frac{1}{3} a^4 VT^4 \quad \text{using (a) and (e)}$$

$$= a^4 VT^4 - \frac{4}{3} a^4 VT^4 + \frac{1}{3} a^4 VT^4$$

$$= 0.$$

4. Equation of state

$$\frac{1}{T} = \frac{\partial S(E, A)}{\partial E}$$

$$\begin{aligned} S(E, A) &= N k_B \left(\ln \left(\frac{m A E}{2 \pi \hbar^2 N} \right) + 2 \right) \\ &= 2 N k_B + N k_B \ln E \\ &\quad + N k_B \ln \left(\frac{m A}{2 \pi \hbar^2 N} \right) \end{aligned}$$

$$\Rightarrow \frac{\partial S(E, A)}{\partial E} = N k_B \frac{1}{E}$$

$$\text{i.e. } \frac{1}{T} = N k_B \frac{1}{E}$$

$$\Rightarrow \underline{E = N k_B T}$$

The equipartition theorem says there is $\frac{1}{2} k_B T$ of energy on average associated with each degree of freedom. Each atom in the gas has two degrees of freedom (motion in the x and y directions) \Rightarrow average kinetic energy of $\frac{1}{2} k_B T$ in each of these directions \Rightarrow average energy of $k_B T$ per particle \Rightarrow energy $N k_B T$ for N particles

5 (a) Conservation of energy

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$$dE = TdS - pdV + \sum_i \mu_i dN_i \quad (*)$$

\Rightarrow the change in the energy is determined by the changes in S , V and \vec{N} for all systems

$\Rightarrow E = E(S, V, \vec{N})$, the energy is "naturally" a function of S , V and \vec{N}

(b) Using the ~~pro~~ definition of partial derivatives

$$\begin{aligned} dE(S, V, \vec{N}) &= \frac{\partial E(S, V, \vec{N})}{\partial S} dS \\ &+ \frac{\partial E(S, V, \vec{N})}{\partial V} dV \\ &+ \sum_i \frac{\partial E(S, V, \vec{N})}{\partial N_i} dN_i \quad (**)$$

Comparing the coefficients of dS , dV and dN_i in (*) and (**) give the equations of state

$$\begin{aligned} T &= \frac{\partial E(S, V, \vec{N})}{\partial S} \\ -p &= \frac{\partial E(S, V, \vec{N})}{\partial V} \\ \mu_i &= \frac{\partial E(S, V, \vec{N})}{\partial N_i} \end{aligned}$$