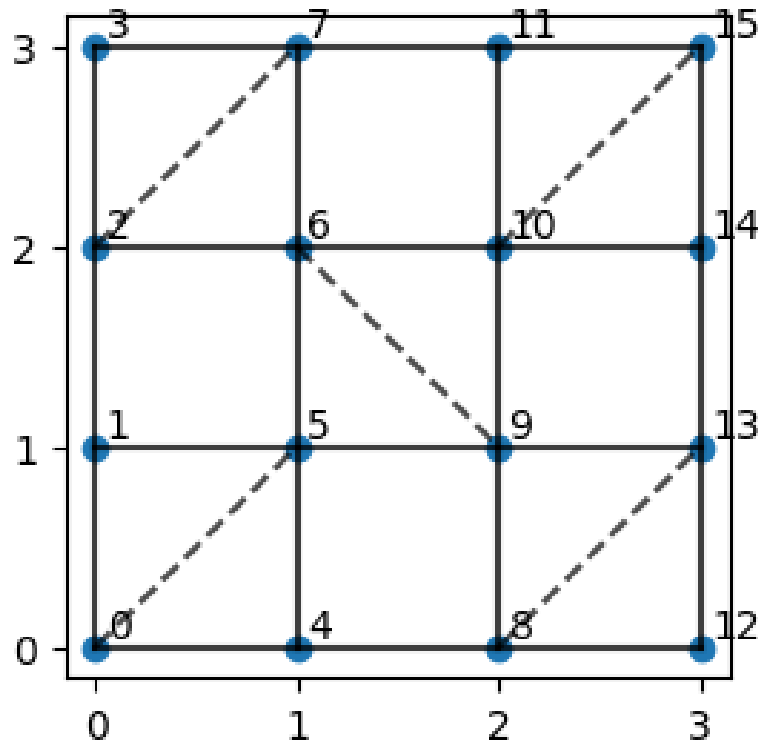


# Shastry-Sutherland model

Introduzione ai sistemi quantistici a molti corpi

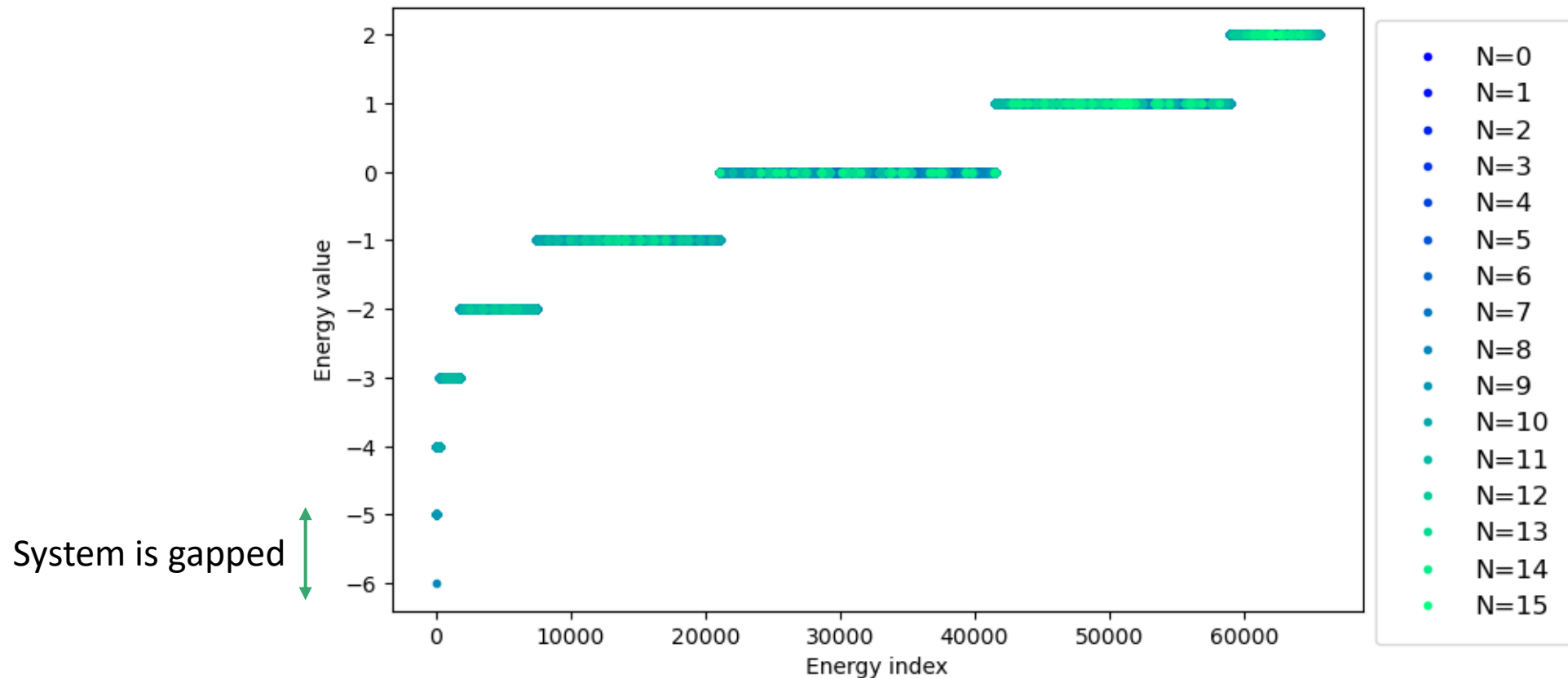
# The Hamiltonian

$$\mathcal{H} = J_1 \sum_{\langle i,j \rangle} S_i S_j + J_2 \sum_{\langle\langle i,j \rangle\rangle} S_i S_j$$



# Limit $J_1=0$ : Dimer phase

$$\mathcal{H} = J_2 \sum_{\langle\langle i,j \rangle\rangle} S_i S_j$$

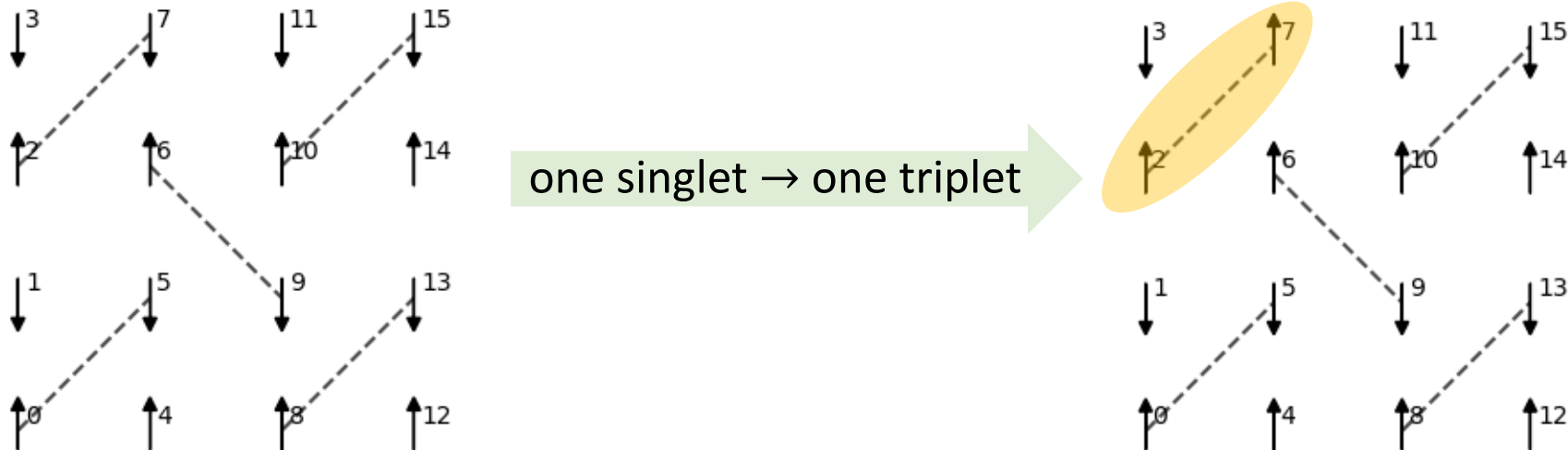


# Dimer phase

Ground state:  $|GS\rangle = \prod_{\langle\langle i,j \rangle\rangle} \frac{1}{\sqrt{2}} [| \uparrow_i \downarrow_j \rangle - | \downarrow_i \uparrow_j \rangle]$

Energy:

- $E_{GS} = -\frac{3}{4}J_2 N_{dimer}$
  - $E_1 = -\frac{3}{4}J_2(N_{dimer} - 1) + \frac{1}{4}J_2$
- }  $\Delta = E_1 - E_{GS} = J_2$



# Dimer phase: motion of triplets

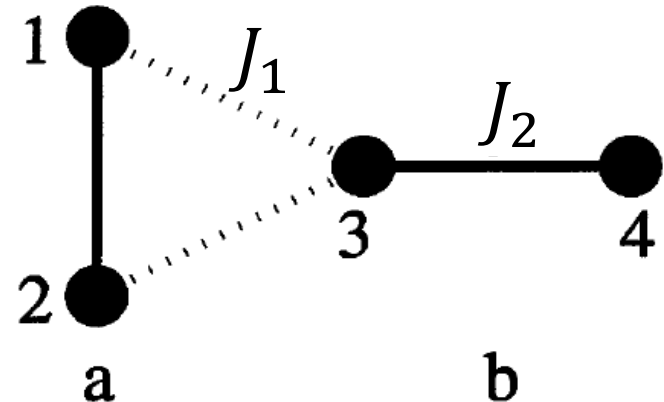
Let consider also a term  $J_1$  that connects dimers between themselves.

It can be shown that:

$$\mathcal{H}_{ab}^1 |t_m\rangle_a |s\rangle_b = \frac{J_1}{2} |t_m\rangle_a |t_0\rangle_b - \frac{J_1}{2} |t_0\rangle_a |t_m\rangle_b$$

$$\mathcal{H}_{ab}^1 |t_0\rangle_a |s\rangle_b = \frac{J_1}{2} |t_1\rangle_a |t_{-1}\rangle_b - \frac{J_1}{2} |t_{-1}\rangle_a |t_1\rangle_b$$

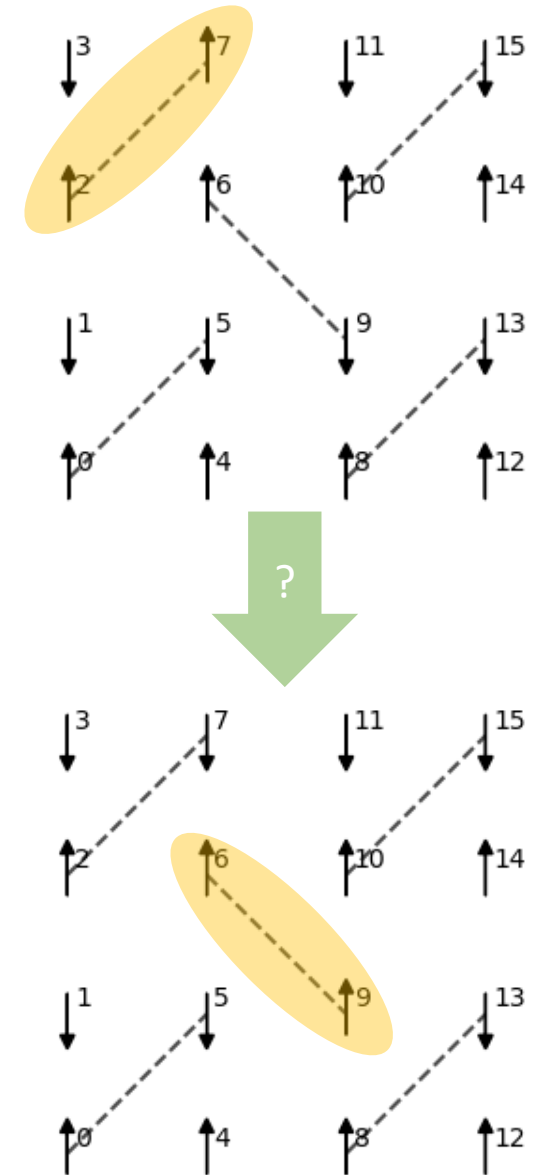
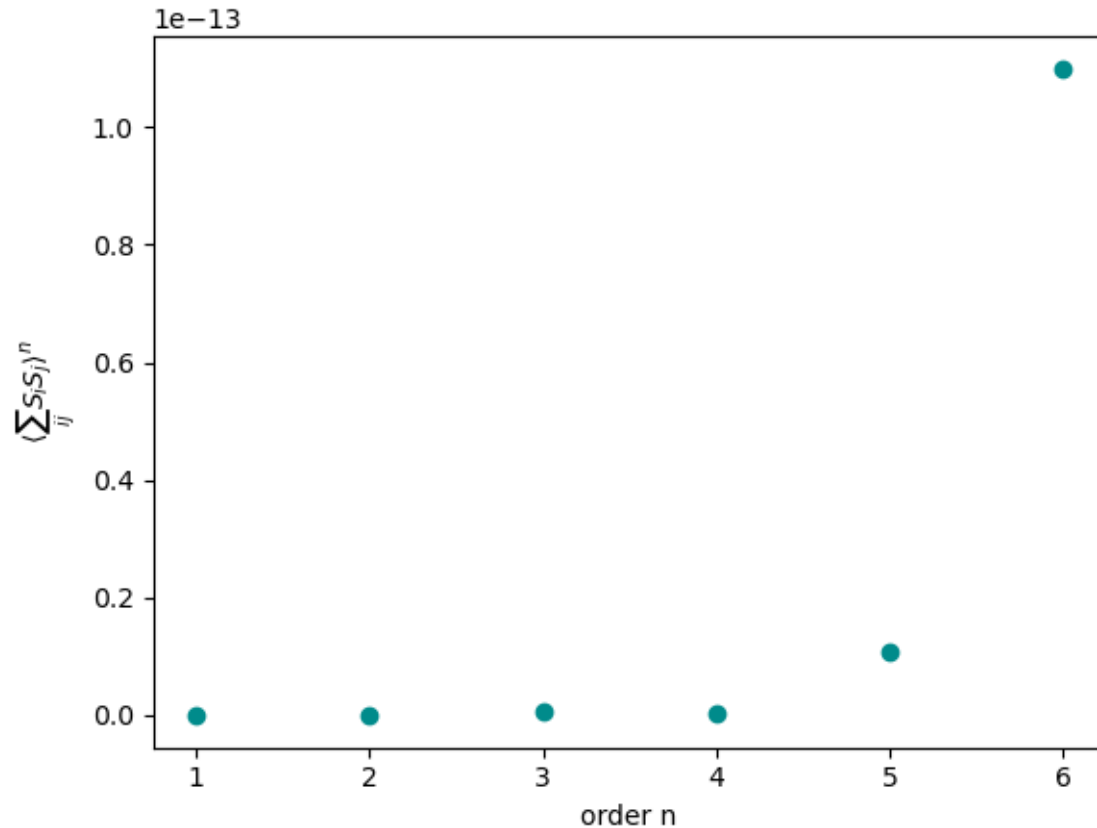
$$\mathcal{H}_{ab}^1 |s\rangle_a |t_m\rangle_b = 0$$



# Dimer phase: motion of triplets

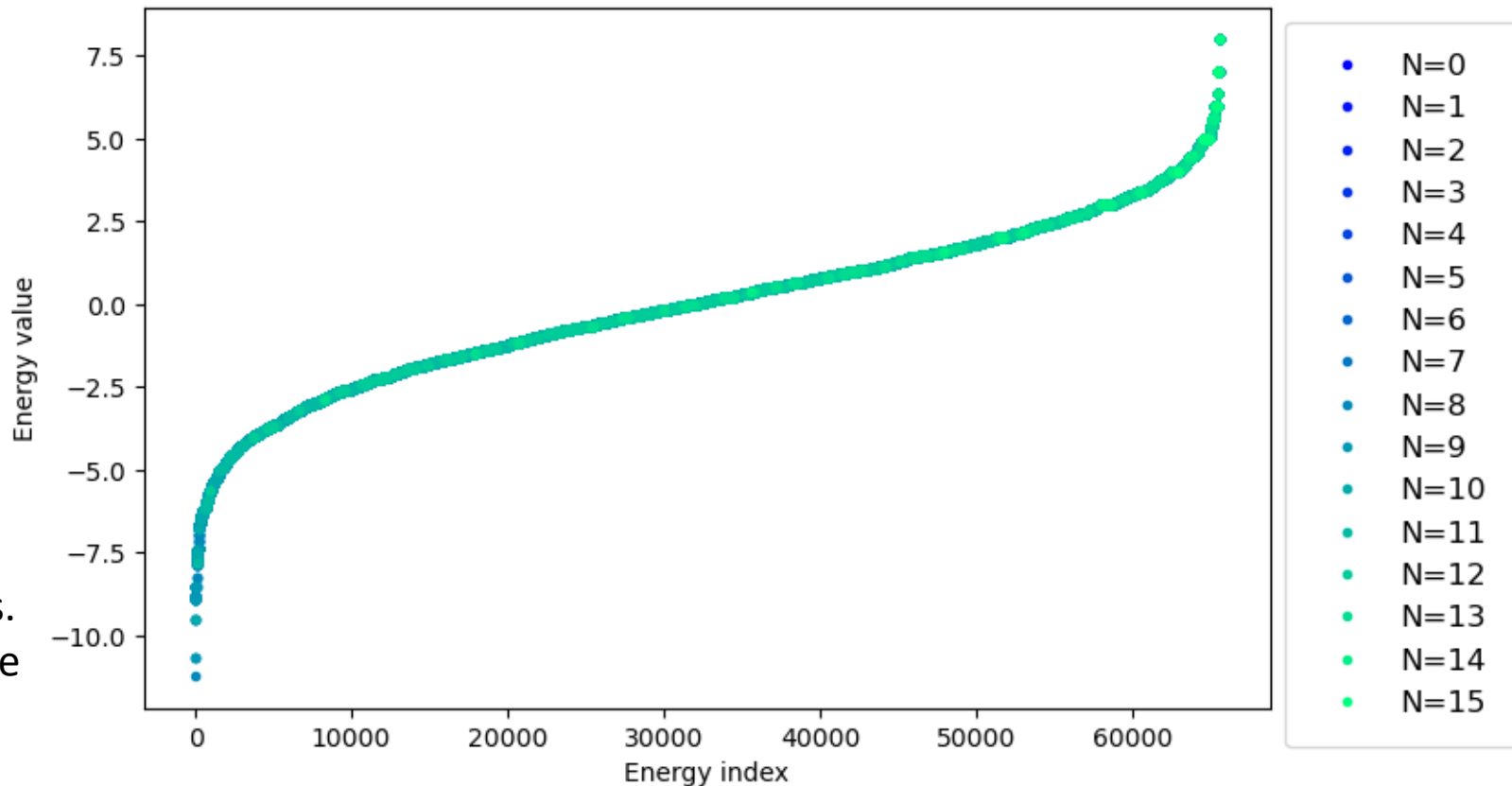
First order:  $\langle \psi_1' | \mathcal{H}_{J_1} | \psi_1 \rangle = 0$ , where  $|\psi_1\rangle$  is a state with one triplet.

At which order of  $\mathcal{H}_{J_1}$  is possible to move one triplet from one bond to a near bond?



# Limit $J_2=0$ : Heisenberg 2D

$$\mathcal{H} = J_1 \sum_{\langle i,j \rangle} S_i S_j$$



Gap due to finite-size effects.  
In the thermodynamic limit the  
system is gapless

# Heisenberg 2D

$$\mathcal{H} = J_1 \sum_{\langle i,j \rangle} S_i S_j = \dots = -2J_1 S^2 L + 4J_1 S \sum_k \omega_k \alpha_k^+ \alpha_k + 2SJ_1 \sum_k (\omega_k - 1)$$

$$\frac{E_{GS}}{N} = -2J_1 S^2 + 2SJ_1 \underbrace{\frac{1}{N} \sum_k (\omega_k - 1)}_{\simeq -0.158} \simeq -0.658J_1$$

From the simulation:

$$\frac{E_{GS}}{N} \simeq -0.702J_1$$



## Heisenberg 2D: introducing a perturbation $J_2 \ll J_1$

$$\mathcal{H} = J_1 \sum_{\langle i,j \rangle} S_i S_j + J_2 \sum_{\langle\langle i,j \rangle\rangle} S_i S_j = \mathcal{H}_0 + V$$

$$E_{GS} = E_{GS}^{(0)} + E_{GS}^{(1)} \rightarrow E_{GS}^{(1)} = \langle \psi_0 | V | \psi_0 \rangle = J_2 \sum_{\langle\langle i,j \rangle\rangle} \langle \psi_0 | S_i S_j | \psi_0 \rangle$$

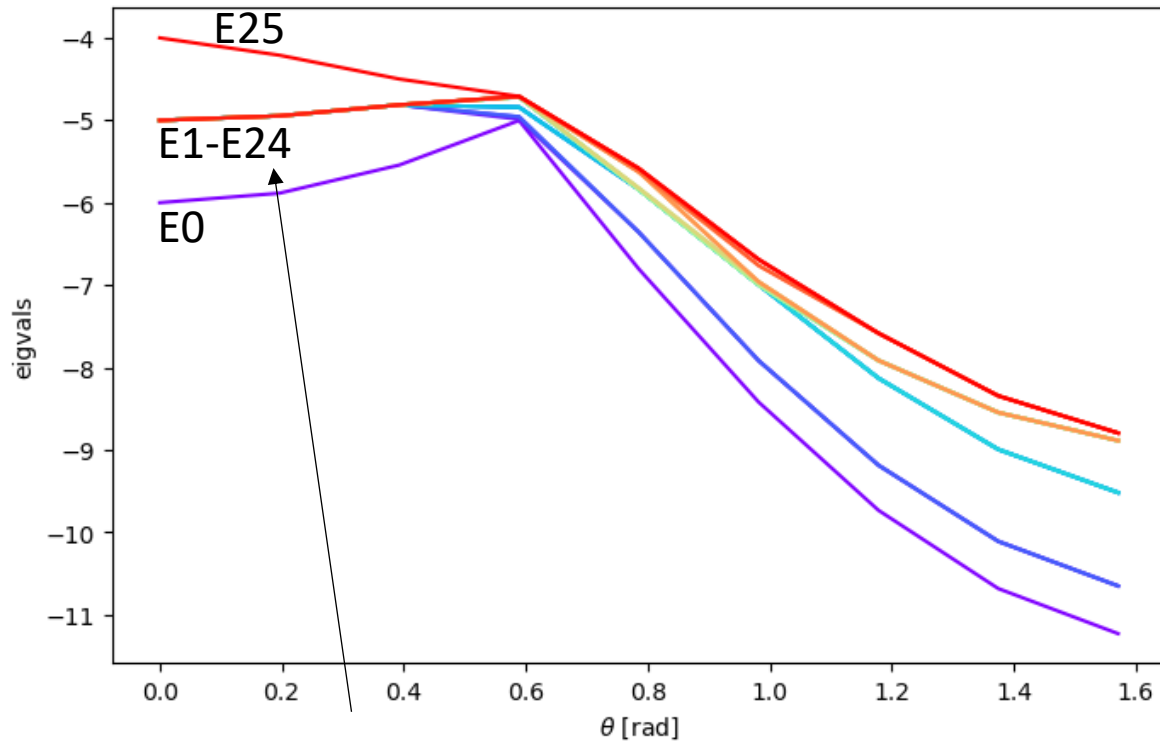
From Miyahara-Ueda:  $\frac{E_{AF}}{N} = -0.669J_1 + 0.102J_2 \rightarrow \text{Transition at : } \left(\frac{J_1}{J_2}\right)_c = 0.71$

From simulation:  $\frac{E_{AF}}{N} = -0.702J_1 + 0.107J_2 \rightarrow \text{Transition at: } \left(\frac{J_1}{J_2}\right)_c = 0.687$

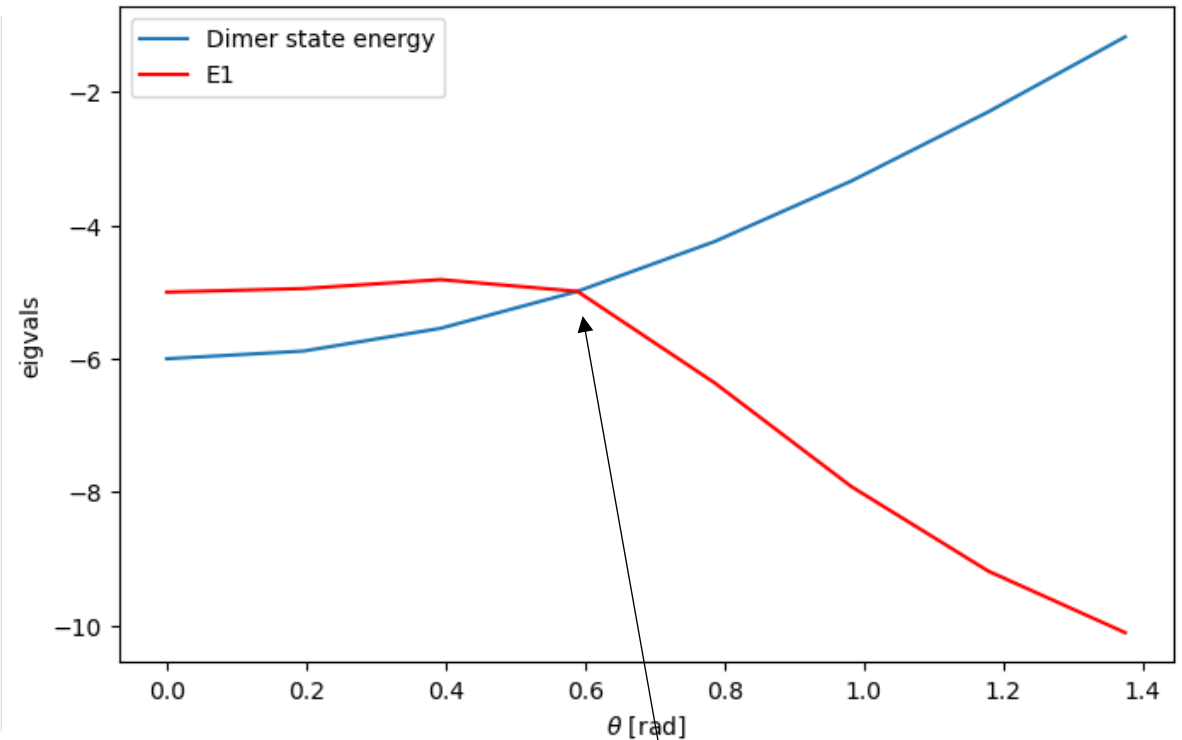
# Changing $J_1$ and $J_2$

We can define the coupling constants through trigonometric functions:

$$J_1 = \sin \theta, J_2 = \cos \theta$$



deg=8 for  $S_z=+1$  |  $\uparrow\uparrow$   
deg=8 for  $S_z=-1$  |  $\uparrow\downarrow$  +  $\downarrow\uparrow$   
deg=8 for  $S_z=0$  |  $\downarrow\downarrow$



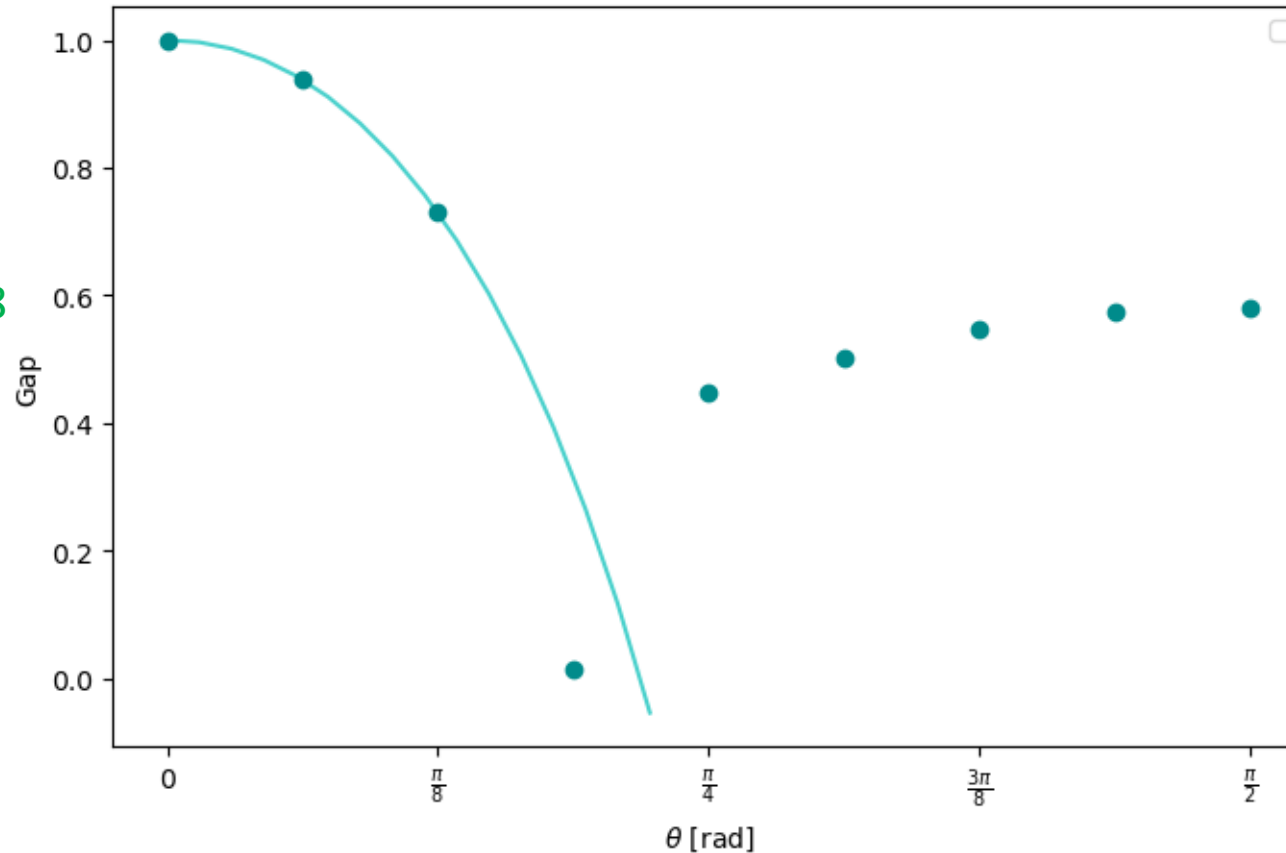
$J_1/J_2=0.66818$

# Changing $J_1$ and $J_2$ : gap

We can define the coupling constants through trigonometric functions:  $J_1/J_2 = \tan \theta$

$$\Delta = \cos(\theta) [1 - \tan^2(\theta) - 1/2 \tan^3(\theta) - 1/8 \tan^4(\theta)]$$

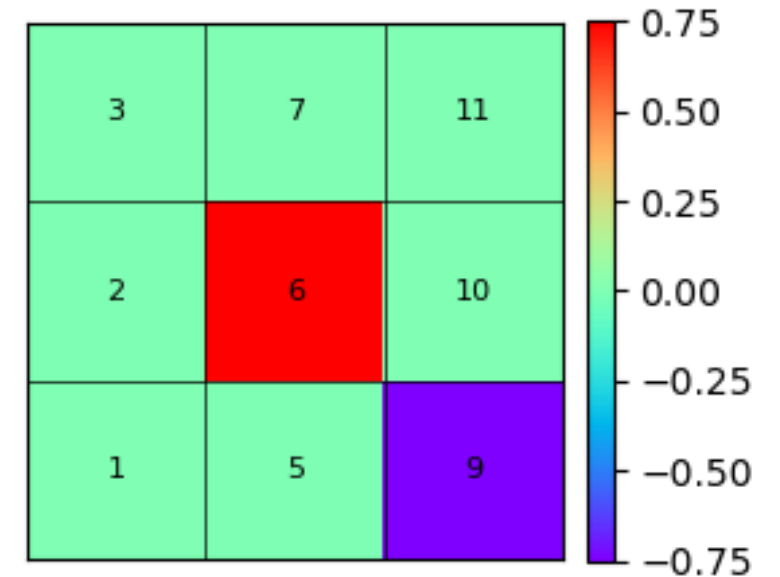
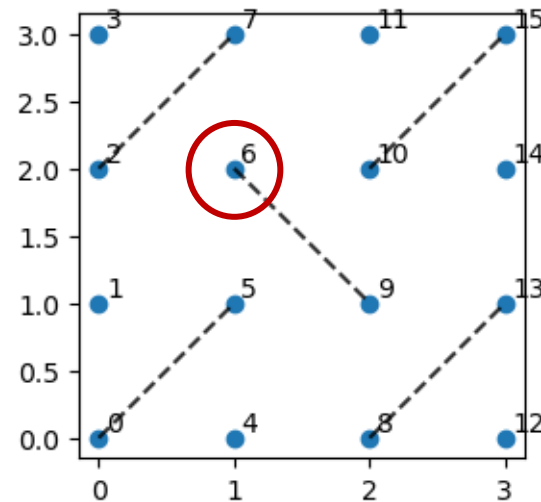
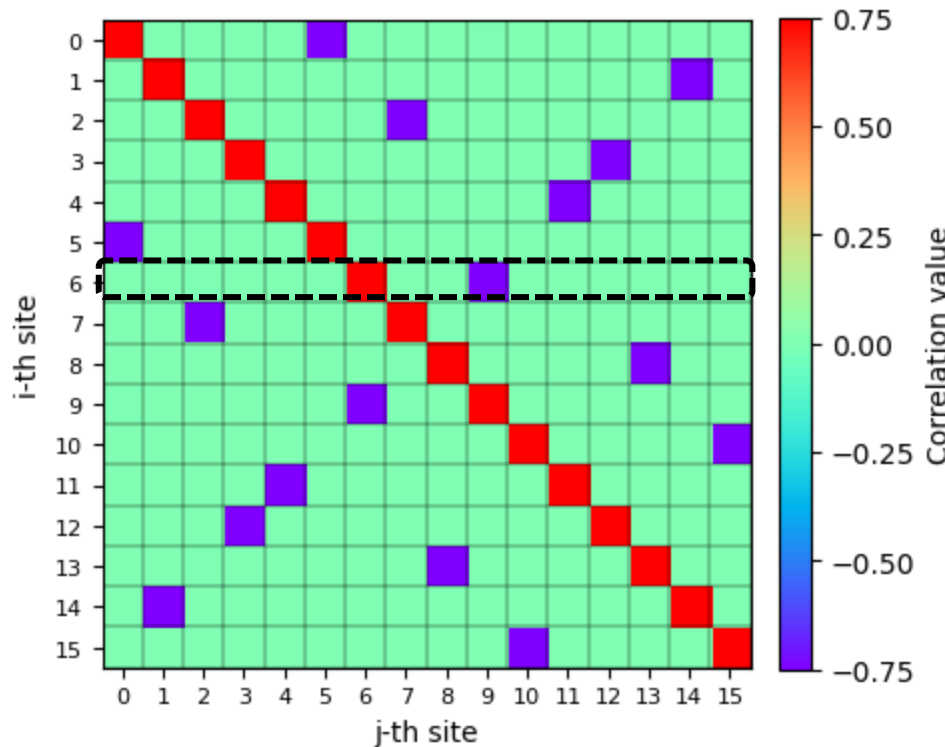
Transition between the  
gapless and the gaped  
phase:  $J_1/J_2 = 0.66818$



In the thermodynamic limit,  
gap  $\rightarrow 0$

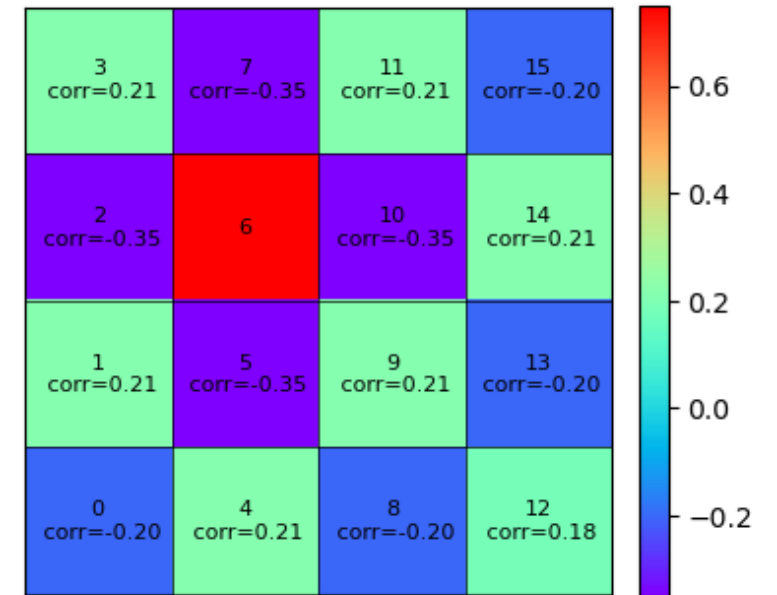
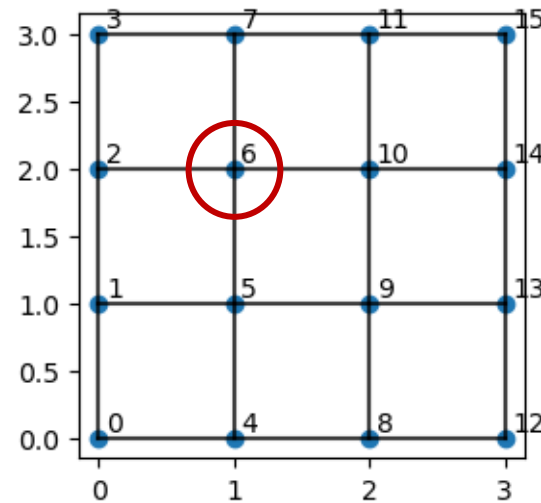
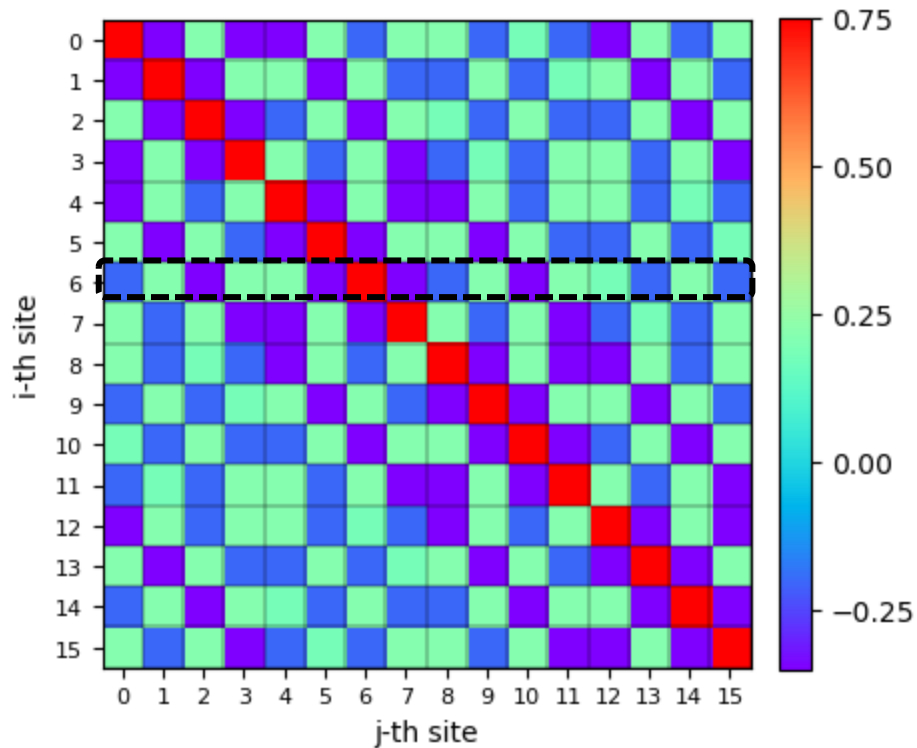
# Correlations: dimer phase

$$\langle GS | S_i S_j | GS \rangle$$



Correlations with nearest neighbors are 0, the only correlation survived is the one between the dimer sites

# Correlations: Heisenberg 2D



NN correlation = -0.35  $\rightarrow$  antiferromagnetic order is still present