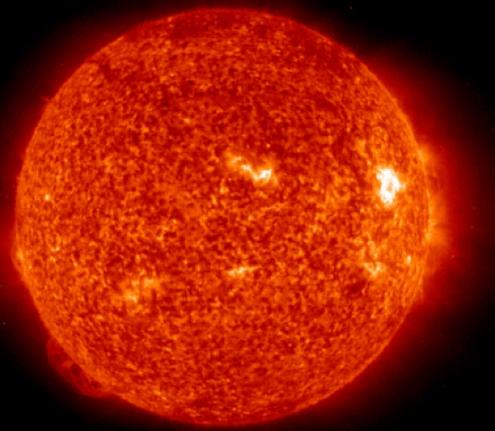
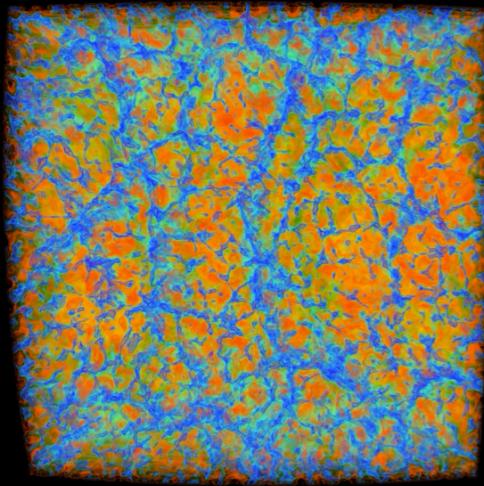


Simulations of Flux Emergence in Cool Stars: What's Convection, Rotation, and Stellar Structure got to do with it?



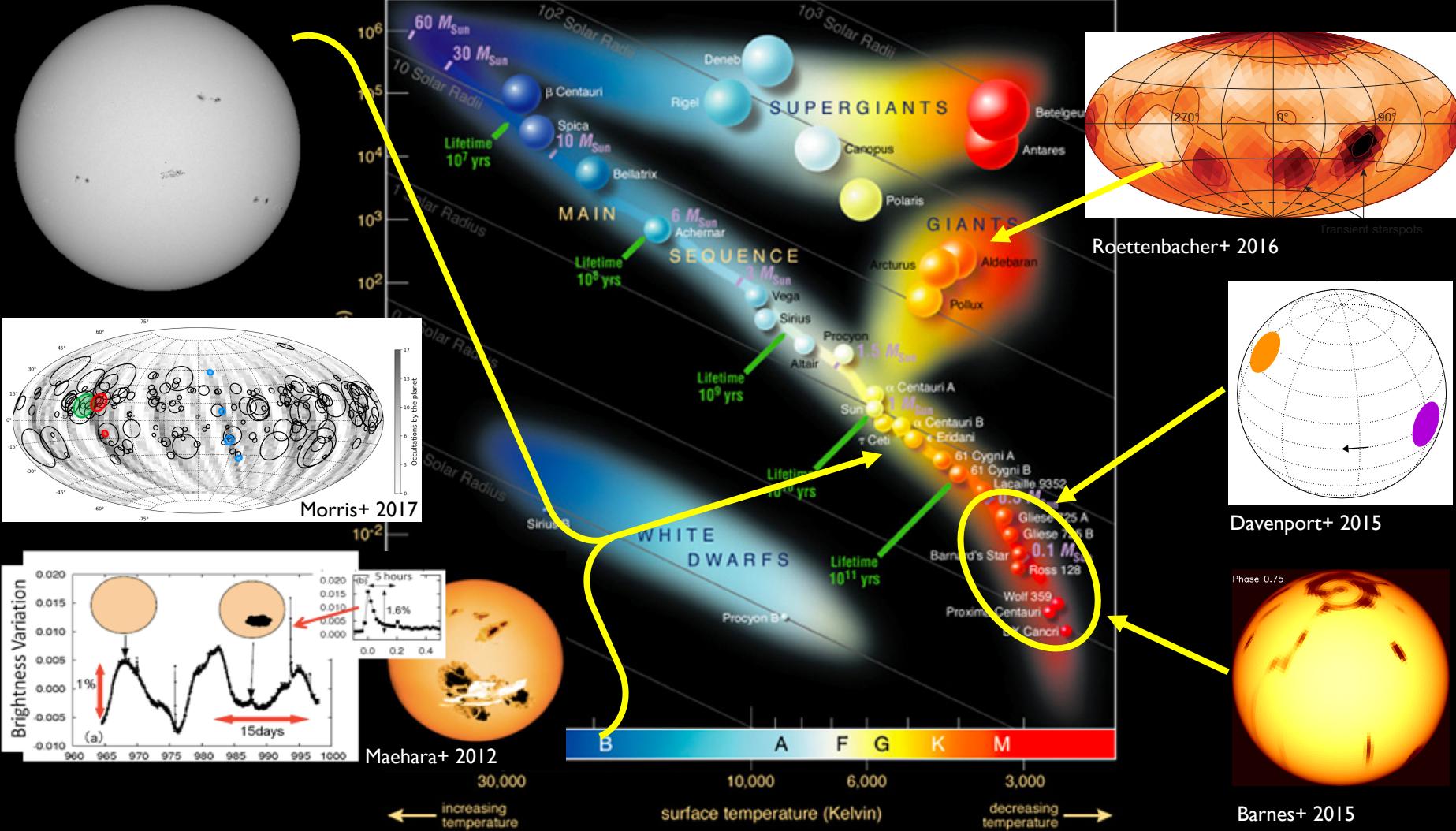
**Maria Weber^{1,2}, Matthew Browning³, Nicholas Nelson⁵, Yuhong Fan⁶, Mark Miesch⁷, Ben Brown^{8,9},
⁴Suzannah Boardman, ⁴Joshua Clarke, ⁴Samuel Pugsley, ⁴Edward Townsend**

¹University of Chicago, ²Adler Planetarium, ³University of Exeter, ⁴former Mphys students at the University of Exeter,

⁵California State University, Chico, ⁶HAO/NCAR, ⁷NOAA, ⁸LASP,

⁹Department of Astrophysical and Planetary Sciences, University of Colorado, Boulder

@SolarisMaria



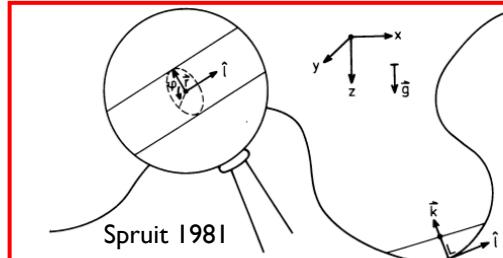
Model Schematic:

TFT
Approximation

Initial
Conditions

Stellar
Structure

Convective
Velocity Field

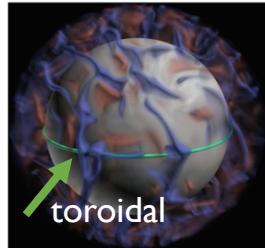


Equation of Motion:

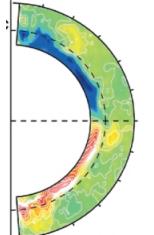
$$\rho \frac{d\mathbf{v}}{dt} = -2\rho(\boldsymbol{\Omega}_0 \times \mathbf{v}) - (\rho_e - \rho)\mathbf{g} + 1 \frac{\partial}{\partial s} \left(\frac{B^2}{8\pi} \right) + \frac{B^2}{4\pi} \mathbf{k}$$

$$- C_d \frac{\rho_e |(\mathbf{v} - \mathbf{v}_e)_\perp| |(\mathbf{v} - \mathbf{v}_e)_\perp|}{(\pi \Phi / B)^{1/2}}$$

B_0
 r_0
 θ_0
 Φ

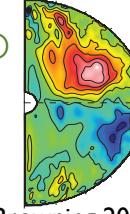


Solar $1\Omega_\odot, 5\Omega_\odot$
Tachocline interface dynamo

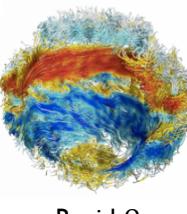


Browning+ 2006

M dwarf, Solar $3\Omega_\odot$
Distributed dynamo

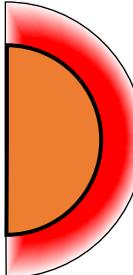


Browning 2008
M Dwarf

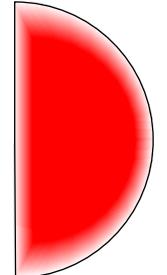


Rapid Ω_0
Brown+ 2011

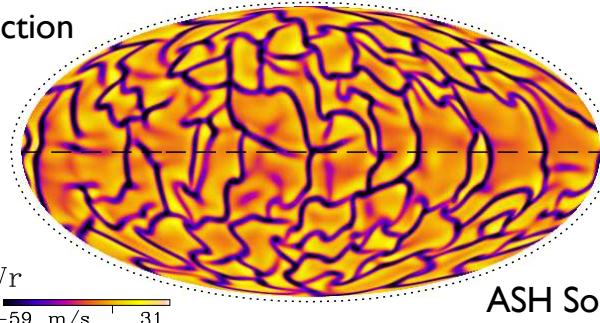
Sun



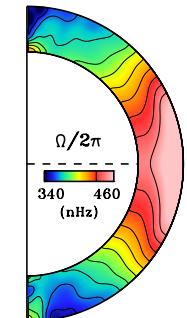
$0.3M_\odot$ M Dwarf



Convection

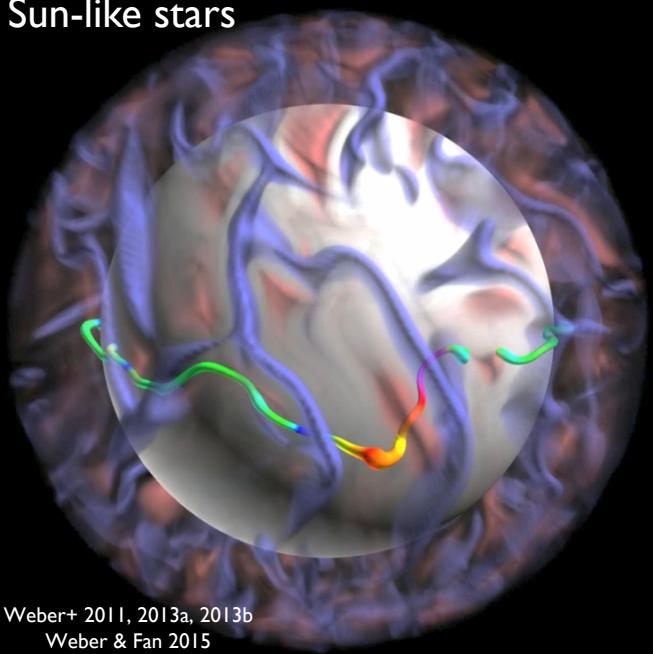


ASH Solar



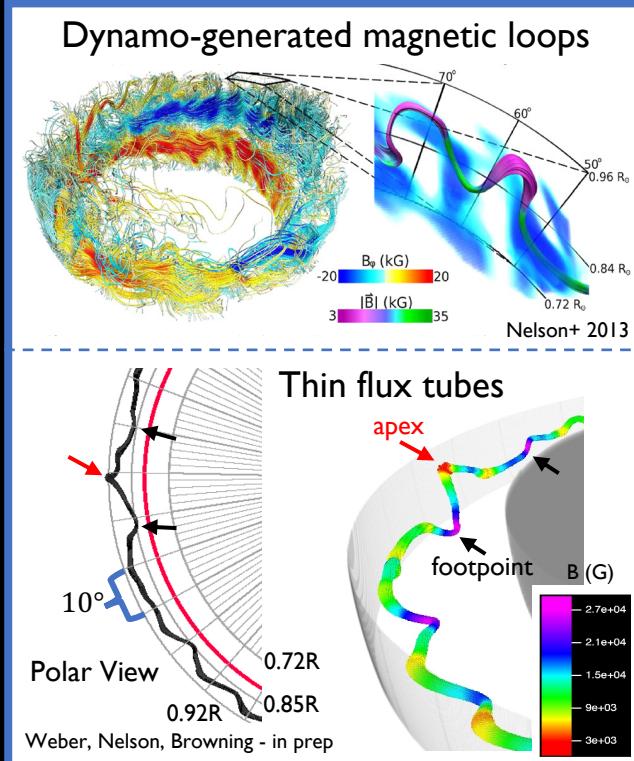
Convection modulates flux emergence

Sun-like stars

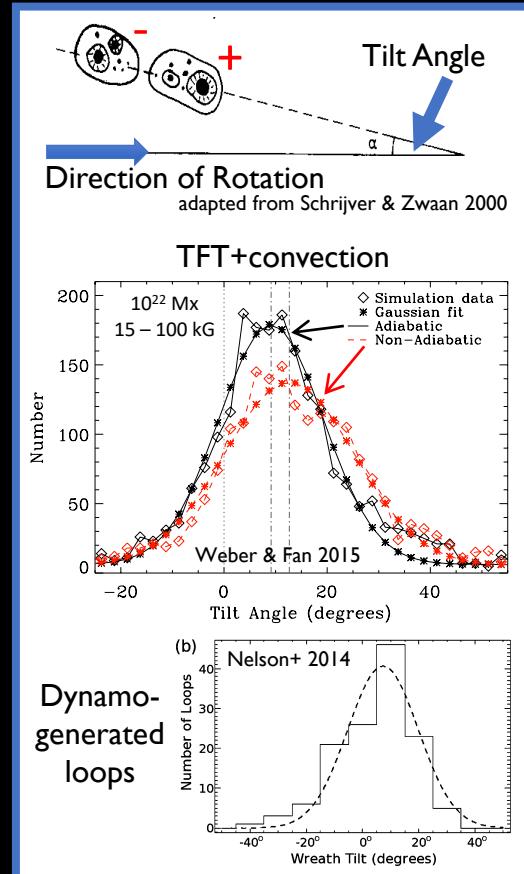


Weber+ 2011, 2013a, 2013b
Weber & Fan 2015

- Convection & magnetic buoyancy work in concert to promote flux emergence



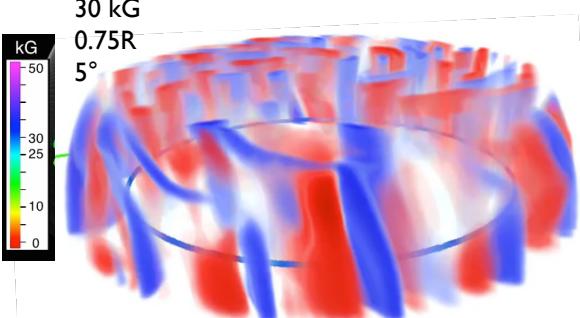
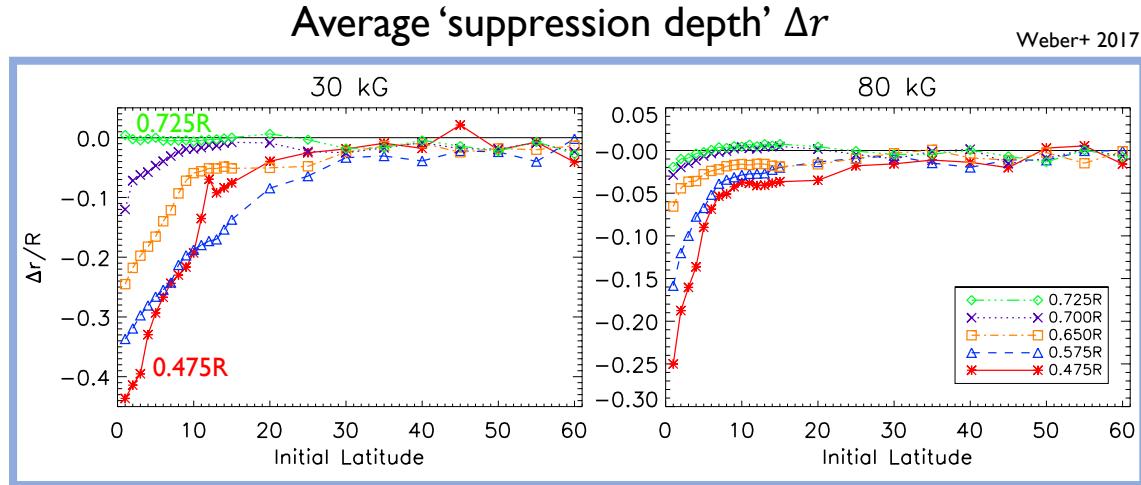
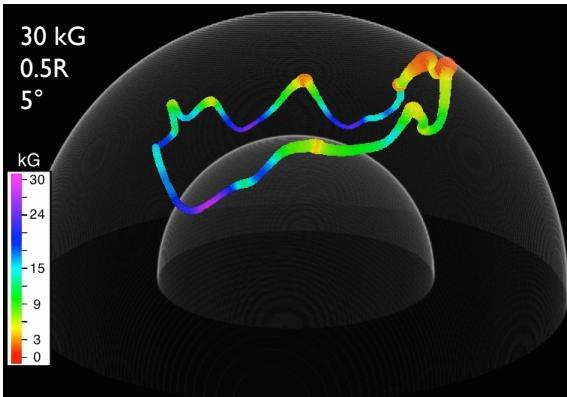
- Downflows naturally induce loops $\sim 15^\circ - 20^\circ$ apart



- Convection introduces a statistical spread in tilt angles

Convection can also suppress flux emergence

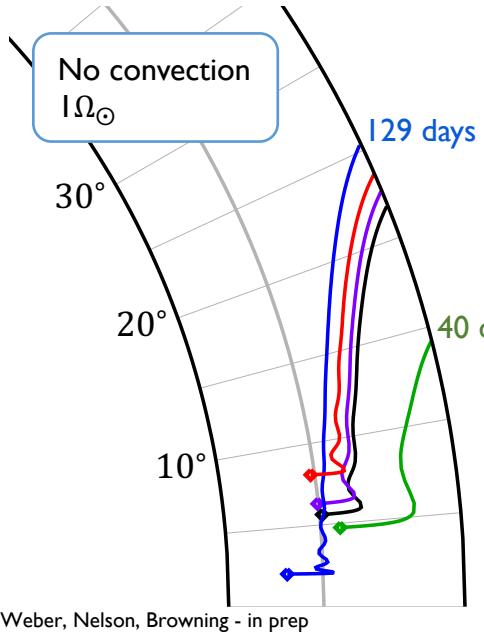
Fully convective M dwarf



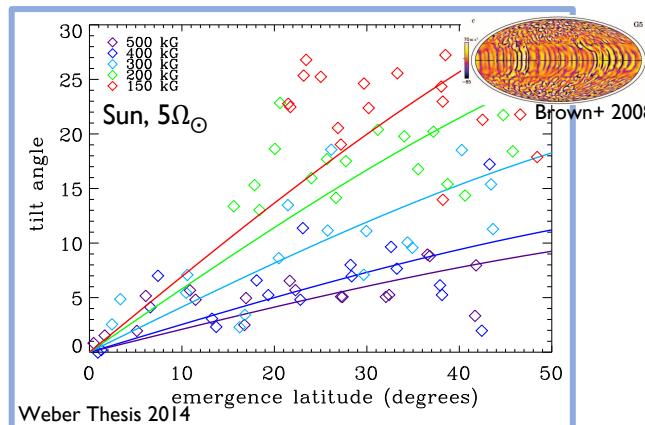
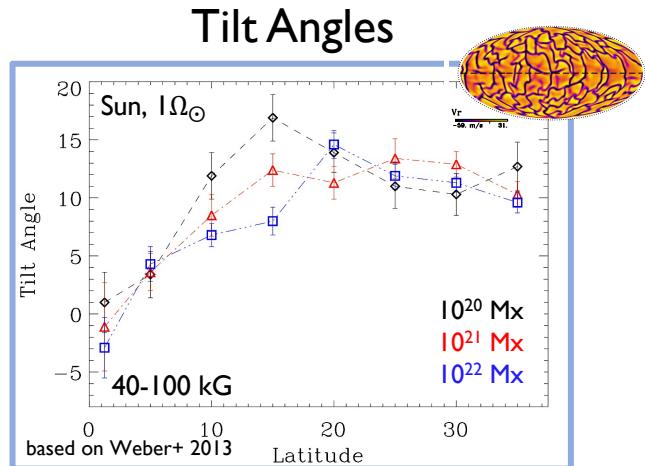
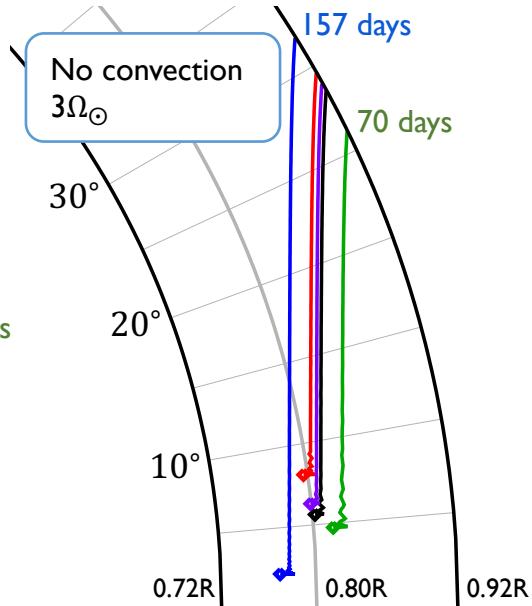
Weber & Browning 2016, Weber+ 2017

- The global rise of TFTs is more strongly suppressed by convective flows when the flux tube is initiated:
 - in the deeper interior
 - at lower latitudes
 - with a weaker magnetic field strength

Rotation alters emergence properties



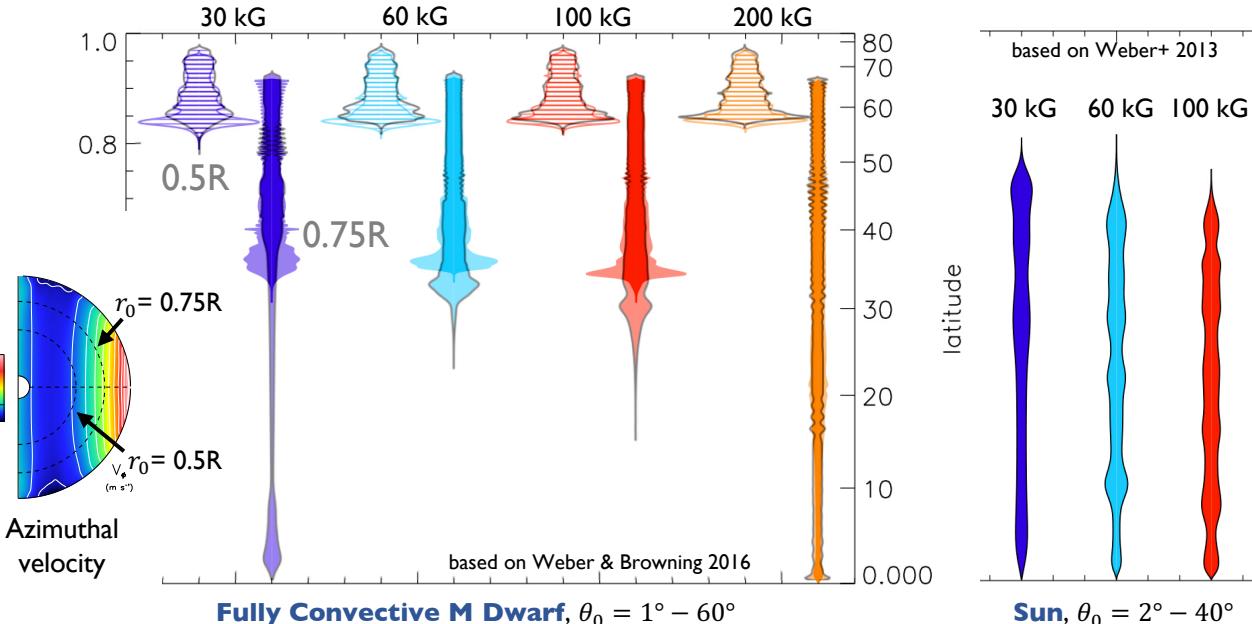
Weber, Nelson, Browning - in prep



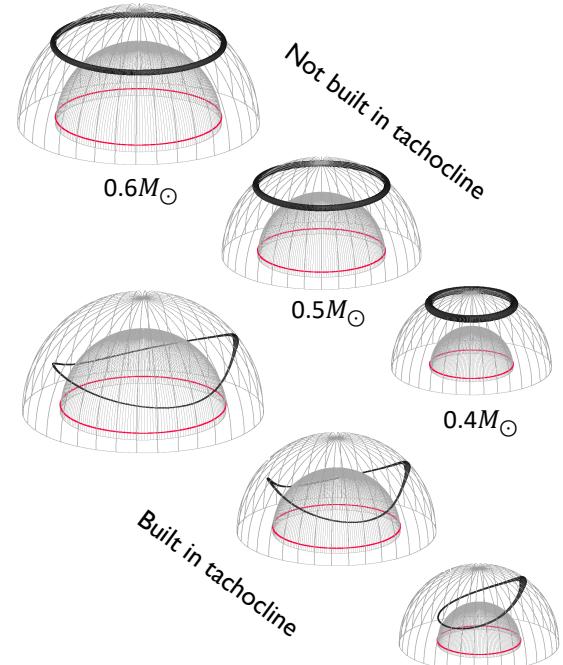
- Due to the Coriolis force, more rapid rotation:
 - Lengthens the rise time
 - Leads to poleward deflection
 - Increases tilt angles

Stellar structure impacts emergence latitudes and more

Emergence Latitude Probability Functions



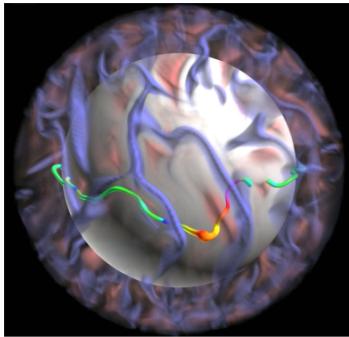
Tachocline or not?



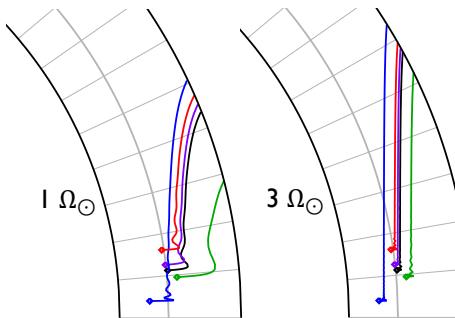
- Unlike solar case, in M dwarf there is a tendency for high latitude emergence ($> 30^\circ$)
- Exceptions when flux tubes initiated closer to the surface and of sufficiently weak (≤ 30 kG) or strong field strengths (≥ 200 kG)
- Increased density in M dwarfs leads to longer flux tube rise times by $\leq 10\times$
- Assumption of flux tube generating region, and thereby initial thermodynamic properties, matter

Summary

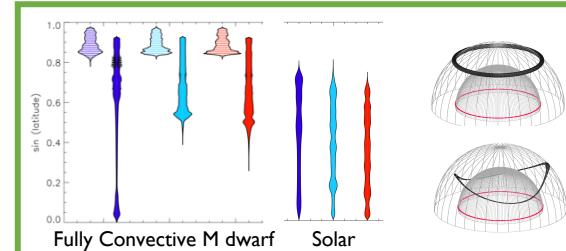
Convection, rotation, and stellar structure are all important contributing factors to the overall trend of flux emergence.



- Convection modulates flux emergence
- Fluid motions both suppress and promote the rise of magnetism
- Convection introduces a statistical spread in emergence properties



- Due to the Coriolis force, rapid rotation:
 - Lengthens the rise times
 - Leads to poleward emergence
 - Increases tilt angle



- Tendency for polar flux emergence in M dwarfs, unlike in solar-like stars
- Increased density in M dwarfs leads to longer rise times
- Assumptions about flux tube generating region (i.e. tachocline or not) has consequences for flux emergence

This work is a step toward linking magnetic flux emergence, convection, and dynamo action along the lower end of the main sequence.