

TD1

1.1

1. $\rho = |\psi\rangle\langle\psi|$

~~Def~~ \hat{A} Positive semi-def $\Leftrightarrow \forall |\psi\rangle \in \mathcal{H}, \langle\psi|\hat{A}|\psi\rangle \geq 0$

Here: $\langle\psi|\rho|\psi\rangle = \langle\psi|\psi\rangle\langle\psi|\psi\rangle$
 $= |\langle\psi|\psi\rangle|^2 \geq 0$

Hermitian: $\rho^\dagger = |\psi\rangle\langle\psi|^\dagger = |\psi\rangle\langle\psi| = \rho$

$$\text{Tr}(\rho) = \sum_{|\psi\rangle \in \mathcal{B}(\mathcal{H})} \langle\psi|\psi\rangle\langle\psi|\psi\rangle$$

$\mathcal{B}(\mathcal{H})$ basis: choose $\{|\psi\rangle, |\psi_1\rangle, |\psi_2\rangle, \dots\}$ by

Gram Schmidt orthonormal:

$$\begin{aligned} \text{Tr}(\rho) &= \langle\psi|\psi\rangle\langle\psi|\psi\rangle + \sum_i \underbrace{\langle\psi_i|\psi\rangle\langle\psi|\psi_i\rangle}_{=0} \\ &= 1 \end{aligned}$$

2) $\partial_t |\psi\rangle = \frac{1}{i\hbar} \hat{H} |\psi\rangle$

$$\partial_t \langle\psi| = -\frac{1}{i\hbar} \langle\psi| \hat{H}$$

$$\begin{aligned} \partial_t \rho &= (\partial_t |\psi\rangle) \langle\psi| + |\psi\rangle (\partial_t \langle\psi|) \\ &= \frac{1}{i\hbar} [\hat{H} |\psi\rangle \langle\psi| - |\psi\rangle \langle\psi| \hat{H}] \\ &= \frac{1}{i\hbar} [\hat{H}, \rho] \end{aligned}$$

3)

$$\langle \hat{A} \rangle = \langle \psi | \hat{A} | \psi \rangle$$

$|\psi\rangle \in B(H)$

$$\langle \hat{A} \rangle = \langle \psi | \hat{A} | \psi \rangle$$

$$= \text{Tr}(\langle \psi | \hat{A} | \psi \rangle)$$

$$= \text{Tr}(|\psi\rangle\langle\psi| \hat{A})$$

$$= \text{Tr}(\hat{\rho} \hat{A})$$

1.2 4. LINEARITY

$$\begin{aligned} 5. \quad \text{Proof:} \quad \text{Tr}(\hat{\rho}^2) &= \text{Tr}(|\psi\rangle\langle\psi| |\psi\rangle\langle\psi|) \\ &= \text{Tr}(|\psi\rangle\langle\psi|) \\ &= 1 \end{aligned}$$

$$\text{Next:} \quad \text{Tr}(\hat{\rho}^2) = \text{Tr} \left(\left(\sum_{\mu} P_{\mu} |\mu\rangle\langle\mu| \right) \left(\sum_{\nu} P_{\nu} |\nu\rangle\langle\nu| \right) \right)$$

$$= \text{Tr} \sum_{\mu, \nu} P_{\mu} P_{\nu} |\mu\rangle\langle\mu| |\nu\rangle\langle\nu|$$

$$= \text{Tr} \sum_{\mu, \nu} P_{\mu} P_{\nu} |\mu\rangle\langle\mu| |\nu\rangle\langle\nu|$$

$$= \sum_{\mu, \nu} P_{\mu} P_{\nu} \begin{vmatrix} \text{Tr}(|\mu\rangle\langle\mu| |\nu\rangle\langle\nu|) \\ \text{Tr}(\langle\mu| |\nu\rangle\langle\nu| |\mu\rangle) \end{vmatrix}$$

$$= \text{Tr}(|\mu\rangle\langle\mu| |\nu\rangle\langle\nu|)$$

$$= |\langle\mu|\nu\rangle|^2$$

$$0 \leq \dots \leq 1$$

$$\begin{aligned} \text{Tr}(\rho^2) &\leq \sum_{r,s} P_r P_s \\ &\leq \sum_r P_r \sum_s P_s \\ &\leq 1 \end{aligned}$$

Egalité : CS. ineq : only
if one $P_r = 1$ and others
 $= 0$

We saw pure $\Rightarrow \text{Tr}(\rho^2) = 1$

Now $\text{Tr}(\rho^2) = 1 \Rightarrow ?$

~~$$\text{Tr} \sum_{r,s} P_r P_s |r\rangle\langle r| \otimes |s\rangle\langle s| = 1$$~~

~~$$\Rightarrow \sum_{r,s} P_r P_s \langle r| \otimes \langle s| \otimes |r\rangle \otimes |s\rangle = 1$$~~

~~$$\Rightarrow \left[\sum_{r,s} P_r P_s \langle r| \otimes |s\rangle \right]^\dagger = 1$$~~

populations diag

coherences hors-diag

II

$$1) \mathcal{H} = \mathcal{H}_1 \otimes \mathcal{H}_2$$

$$L_1 \text{ basis } : B = \{ | \uparrow \uparrow \rangle, | \uparrow \downarrow \rangle, | \downarrow \uparrow \rangle, | \downarrow \downarrow \rangle \}$$

$$| \psi \rangle = \alpha | \uparrow \uparrow \rangle + \beta | \uparrow \downarrow \rangle + \gamma | \downarrow \uparrow \rangle + \delta | \downarrow \downarrow \rangle.$$

$$\text{We note: } | \psi \rangle = \sum_{a,b} \psi_{ab} | a b \rangle, \quad a, b \in \uparrow, \downarrow$$

$$| \psi \rangle = \begin{pmatrix} \psi_{\uparrow\uparrow} \\ \psi_{\uparrow\downarrow} \\ \psi_{\downarrow\uparrow} \\ \psi_{\downarrow\downarrow} \end{pmatrix} = \begin{pmatrix} \alpha \\ \beta \\ \gamma \\ \delta \end{pmatrix} \text{ in basis } B \quad (4)$$

$$\rho = | \psi \rangle \langle \psi | = \sum_{a_1, b_1, a_2, b_2} \psi_{a_1 b_1} \psi_{a_2 b_2}^* | a_1 b_1 \rangle \langle a_2 b_2 |$$

In basis B :

$$\rho = \begin{pmatrix} \psi_{\uparrow\uparrow} \psi_{\uparrow\uparrow}^* & \psi_{\uparrow\uparrow} \psi_{\uparrow\downarrow}^* & \dots \\ \psi_{\uparrow\downarrow} \psi_{\uparrow\uparrow}^* & \psi_{\uparrow\downarrow} \psi_{\uparrow\downarrow}^* & \\ \vdots & & \ddots \end{pmatrix}$$

$$2) \hat{A} = \begin{pmatrix} a_{11} & a_{12} \\ a_{21} & a_{22} \end{pmatrix}, \quad \hat{B} = \begin{pmatrix} b_{11} & b_{12} \\ b_{21} & b_{22} \end{pmatrix}$$

$$\hat{A} \otimes \hat{B} = \begin{pmatrix} a_{11} \hat{B} & a_{12} \hat{B} \\ a_{21} \hat{B} & a_{22} \hat{B} \end{pmatrix} = \begin{pmatrix} a_{11} b_{11} & a_{11} b_{12} & a_{11} b_{21} & a_{11} b_{22} \\ a_{12} b_{11} & a_{12} b_{12} & a_{12} b_{21} & a_{12} b_{22} \\ a_{21} b_{11} & a_{21} b_{12} & a_{21} b_{21} & a_{21} b_{22} \\ a_{22} b_{11} & a_{22} b_{12} & a_{22} b_{21} & a_{22} b_{22} \end{pmatrix}$$

3) $\hat{O} = \hat{A} \otimes \hat{B}$ Only acting on subsystem A : $\hat{B} = 1$

$$\hat{O} = \hat{A} \otimes 1$$

$$\hookrightarrow \hat{O} = \begin{pmatrix} a_{11} & 0 & | & a_{11} & 0 \\ 0 & a_{11} & | & 0 & a_{11} \\ \vdots & \vdots & | & \vdots & \vdots \\ a_{21} & 0 & | & a_{21} & 0 \\ 0 & a_{21} & | & 0 & a_{21} \end{pmatrix} \quad (5)$$

$$\langle \hat{O} \rangle = \text{Tr}(\hat{\rho} \hat{O})$$

4) We start with:

$$\text{Tr}(\hat{\rho} \hat{O}) = \sum_{a_1 b_1} \langle a_1 b_1 | \hat{\rho} \hat{O} | a_1 b_1 \rangle$$

$$= \sum_{\substack{a_1 b_1 \\ a_2 b_2}} \langle a_1 b_1 | \hat{\rho} | a_2 b_2 \rangle \langle a_2 b_2 | \hat{O} | a_1 b_1 \rangle$$

$$= \sum_{\substack{a_1 b_1 \\ a_2 b_2}} \psi_{a_1 b_1} \psi_{a_2 b_2}^* \langle a_2 | A | a_1 \rangle \delta_{b_1 b_2}$$

$$= \sum_{\substack{a_1 b_1 \\ a_2}} \psi_{a_1 b_1} \psi_{a_2 b_1}^* \langle a_2 | A | a_1 \rangle.$$

And compare with :

$$\text{Tr}(\rho_A A) = \sum_{a_3} \langle a_3 | \rho_A A | a_3 \rangle$$

$$= \sum_{a_3 a_4} \langle a_3 | \rho_A | a_4 \rangle \langle a_4 | A | a_3 \rangle.$$

$$= \sum_{\substack{a_1, b_1 \\ a_2, b_2 \\ a_3, a_4}} \psi_{a_1 b_1} \psi_{a_2 b_2}^* \underbrace{\langle a_3 | a_1 \rangle \langle a_2 | a_4 \rangle}_{\delta_{a_3 a_1} \delta_{a_2 a_4}} \underbrace{\langle a_4 | A | a_3 \rangle \text{Tr}(b_1 \times b_2)}_{\delta_{b_1 b_2}}.$$

$$= \sum_{\substack{a_1, b_1 \\ a_2}} \psi_{a_1 b_1} \psi_{a_2 b_1}^* \langle a_2 | A | a_1 \rangle.$$

(6)

$$= \text{Tr}(\rho \hat{O}) = \langle \hat{O} \rangle.$$

TD Wigner

$$1) \quad \mathbb{A} = \int dp |p\rangle\langle p| \quad \tilde{A}(x,p) = \int e^{-ipy/\hbar} \langle x+y/2 | \hat{A} | x-y/2 \rangle dy$$

$$\tilde{A}(x,p) = \iiint e^{-ipy/\hbar} \langle x+y/2 | p_1 \rangle \langle p_1 | \hat{A} | p_2 \rangle \langle p_2 | x-y/2 \rangle dy dp_1 dp_2$$

$$\langle x | p \rangle = \frac{1}{\sqrt{2\pi\hbar}} e^{ipx/\hbar}$$

(1)

$$= \frac{1}{\sqrt{2\pi\hbar}} \iiint dy dp_1 dp_2 e^{-ipy/\hbar} e^{ip_1(x+y/2)/\hbar} e^{-ip_2(x-y/2)/\hbar} \langle p_1 | \hat{A} | p_2 \rangle$$

$$= \frac{1}{\sqrt{2\pi\hbar}} \iint dp_1 dp_2 \langle p_1 | \hat{A} | p_2 \rangle \int dy e^{\frac{i}{\hbar}(-py + p_1(x+y/2) - p_2(x-y/2))}$$

$$= \frac{1}{\sqrt{2\pi\hbar}} \int dy e^{\frac{i}{\hbar}(-py + p_1(x+y/2) - p_2(x-y/2))}$$

$$= \frac{1}{\sqrt{2\pi\hbar}} \iint dp_1 dp_2 \langle p_1 | \hat{A} | p_2 \rangle e^{i(p_1-p_2)x/\hbar} \int dy e^{i(\frac{p_1+p_2}{2}-p)y}$$

$$= \frac{1}{\sqrt{2\pi\hbar}} \iint dp_1 dp_2 \langle p_1 | \hat{A} | p_2 \rangle e^{i(p_1-p_2)x/\hbar} 2\pi\hbar \delta(\frac{p_1+p_2}{2}-p)$$

$$= \iint dp_1 dp_2 \delta\left(\frac{p_1+p_2}{2} - p\right) \langle p_1 | A | p_2 \rangle e^{i(p_1-p_2)x/\hbar}$$

$$u = p_1 - p_2, \quad v = \frac{p_1+p_2}{2}$$

$$p_1 = v + \frac{u}{2}, \quad p_2 = v - \frac{u}{2}$$

$$du dv = \cancel{dp_1 dp_2} = \cancel{dp_1} \cancel{dp_2} \left(\cancel{dp_1} \cancel{dp_2} \right) \left(\cancel{dp_1} \cancel{dp_2} \right)$$

$$= \iint du dv \delta(v - p) \langle v + \frac{u}{2} | A | v - \frac{u}{2} \rangle e^{iux/\hbar}$$

$$= \int du e^{iux/\hbar} \langle p + \frac{u}{2} | A | p - \frac{u}{2} \rangle$$

(2)

$$\overline{T_n}[\hat{A}, \hat{B}] =$$

$$2) \frac{1}{h} \iint dx dp \tilde{A}(x, p) \tilde{B}(x, p) = \frac{1}{h} \iint dx dp \int dy_1 e^{-i p y_1 / \hbar} \langle x + y_1/2 | A | x - y_1/2 \rangle \int dy_2 e^{-i p y_2 / \hbar} \langle x + y_2/2 | B | x - y_2/2 \rangle$$

$$= \frac{1}{h} \iiint dx dp dy_1 dy_2 e^{-i p (y_1 + y_2) / \hbar} \langle x + y_1/2 | A | x - y_1/2 \rangle \langle x + y_2/2 | B | x - y_2/2 \rangle$$

$$= \iiint dx dy_1 dy_2 \delta(y_1 + y_2) \langle x + y_1/2 | A | x - y_1/2 \rangle \langle x + y_2/2 | B | x - y_2/2 \rangle$$

$$= \iint dx dy_1 \langle x + y_1/2 | A | x - y_1/2 \rangle \langle x - y_1/2 | B | x + y_1/2 \rangle$$

$$\stackrel{\text{if } dx}{=} \int dx \quad u = x - \frac{y_1}{2}, \quad v = x + \frac{y_1}{2} \quad (3)$$

$$\Rightarrow dx dy_1 = du dv$$

$$= \iint du dv \langle v | A | u \rangle \langle u | B | v \rangle$$

$$= \int dv \langle v | A B | v \rangle$$

$$= T_n(\hat{A} \hat{B})$$

$$\tilde{A}(x, p) = \int e^{-i p y / \hbar} \langle x + y/2 | \hat{A} | x - y/2 \rangle dy$$

$$3) \langle \hat{A} \rangle = \text{Tr}(\hat{\rho} \hat{A}) = \frac{1}{h} \iint dx dp \tilde{\rho}(x,p) \tilde{A}(x,p) = \iint dx dp \mathcal{W}(x,p) \tilde{A}(x,p)$$

$$4) \int dp \mathcal{W}(x,p) = \frac{1}{h} \iint dp \int du e^{ixu/h} \langle p+u/2 | \hat{\rho} | p-u/2 \rangle$$

$$= \frac{1}{h} \iint dp du e^{ixu/h} \varphi(p+u/2) \varphi^*(p-u/2)$$

(4)

$$\cancel{\frac{1}{h} \iint dp du} \quad \alpha = p + \frac{u}{2}, \quad \beta = p - \frac{u}{2}, \quad dx dp = dx du$$

$$= \frac{1}{h} \iint dx d\beta e^{ix(\alpha-\beta)/h} \varphi(\alpha) \varphi^*(\beta)$$

$$= \left(\frac{1}{\sqrt{h}} \int dx e^{ix\alpha/h} \varphi(\alpha) \right) \left(\frac{1}{\sqrt{h}} \int d\beta e^{-ix\beta/h} \varphi^*(\beta) \right)$$

$$= \psi(x) \psi^*(x)$$

$$= |\psi(x)|^2$$

Same for $\int dx \mathcal{W}(x,p)$ w/ initial def of x .

$$\begin{aligned}
 5) \quad \mathcal{W}^*(x, p) &= \frac{1}{h} \int e^{+ip y / \hbar} \left(\langle x + y/2 | \hat{\rho} | x - y/2 \rangle \right)^* dy \\
 &= \frac{1}{h} \int dy \quad e^{+ip y / \hbar} \left(\psi(x + y/2) \psi^*(x - y/2) \right)^* \\
 &= \frac{1}{h} \int dy \quad e^{-ip y / \hbar} \quad \psi(x - y/2) \psi^*(x + y/2)
 \end{aligned}$$

$$y \rightarrow -y$$

$$= \frac{1}{h} \int dy \quad e^{-ip y / \hbar} \quad \psi(x + y/2) \psi^*(x - y/2)$$

$$= \mathcal{W}(x, p) \quad \rightarrow \text{real}.$$

$$6) \quad \iint dx dp \quad \mathcal{W}(x, p) = \int dx \quad |\psi(x)|^2 = 1.$$

$$\mathcal{W}^*(x, p)$$

(5)

$$6) \iint dx dp \mathcal{W}(x, p) = \int dx |\psi(x)|^2 = 1.$$

⑥

$$7) \text{ ~~x~~ } x \rightarrow x-s :$$

$$\mathcal{W}'(x, p) = \frac{1}{h} \int dy e^{-i p y / \hbar} \psi(x - \frac{y}{2} - s) \psi^*(x + \frac{y}{2} - s) = \mathcal{W}(x-s, p)$$

$$p \rightarrow p - \hbar p :$$

$$\mathcal{W}''(x, p) = \frac{1}{h} \int du e^{i u x / \hbar} \psi(p - \frac{u}{2} - s_p) \psi^*(p + \frac{u}{2} - s_p) = \mathcal{W}(x, p - \hbar p).$$

$$8) |\langle \psi_a | \psi_b \rangle|^2 = \langle \psi_b | \psi_a \rangle \langle \psi_a | \psi_b \rangle$$

$$= \langle p_a \rangle_b$$

$$= \iint dx dp \mathcal{W}_b(x, p) \tilde{p}_a$$

$$= h \iint dx dp \mathcal{W}_b(x, p) \mathcal{W}_a(x, p).$$

Take 2 orthogonal state $|\psi_a\rangle, |\psi_b\rangle$: $|\langle \psi_a | \psi_b \rangle|^2 = 0$

so $\iint dx dp \mathcal{W}_a(x, p) \mathcal{W}_b(x, p) = 0$: one of them has to be < 0 somewhere

$$g) \text{ pose } \psi_1(y) = \frac{1}{\sqrt{h}} e^{-ipy/h} \psi(x+y/2)$$

$$\psi_2(y) = \frac{1}{\sqrt{h}} \psi(x-y/2)$$

Normalized ! why?

$$\int dy \psi_2^*(y) \psi_2(y) = \frac{1}{2} \int dy \psi^*(x-y/2) \psi(x-y/2)$$

~~$$\int dy \psi^*(x-y/2) \psi(x-y/2)$$~~

$$u = x - \frac{y}{2}, \quad dy = 2du$$

$$= \int dy \psi^*(u) \psi(u)$$

$$= 1.$$

$$\text{Then: } \psi(x,p) = \frac{2}{h} \int dy \psi_1(y) \psi_2^*(y)$$

$$= \frac{2}{h} \frac{1}{2} \langle \psi_1 | \psi_2 \rangle$$

$$|\psi(x,p)| \leq \frac{2}{h} |\langle \psi_1 | \psi_1 \rangle| \leq \frac{2}{h}$$

(7)

$$b) \tilde{A} = \int dy e^{-i p y / \hbar} \langle x + y/2 | \tilde{A}(\hat{z}) | x - y/2 \rangle$$

~~$$= \int dy e^{-i p y / \hbar} A(x - y/2) \langle x + y/2 | x - y/2 \rangle$$~~

$$= \int dy e^{-i p y / \hbar} A(x - y/2) \langle x + y/2 | x - y/2 \rangle$$

$$= \int dy e^{-i p y / \hbar} A(x - y/2) \delta(y)$$

$$= A(x).$$

8

$$\tilde{B} = B(\hat{p}) : \text{iden w/ other formula (4)}$$

$$\langle \hat{H} \rangle = \iint W(x, p) \tilde{H} dx dp \quad \tilde{H} = \frac{\tilde{p}^2}{2m} + \tilde{V}(x)$$

$$= \frac{p^2}{2m} + V(x) = H(x, p).$$

$$= \iint W(x, p) H(x, p) dx dp.$$

Example of HO:

2.4: Tool for HO.

$$H = \frac{p^2}{2m} + \frac{1}{2} m \omega^2 x^2$$

Lowest energy es ~~(=)~~ (solve Schröd eq):

Find eig ψ of 2 lowest eig states: ~~(show they respect eq (10))~~ ^(10.1)

$$W(x, p) = \frac{1}{h} \int e^{-ipx/\hbar} \psi(x+y/2) \psi^*(x-y/2) dy$$

$$\frac{\delta W}{\delta r} = \frac{1}{h} \int e^{-ipx/\hbar} \left[\frac{\delta \psi^*(x-y/2)}{\delta r} \psi(x+y/2) + \frac{\delta \psi(x+y/2)}{\delta r} \psi^*(x-y/2) \right] dy$$

SE: $i\hbar \frac{\delta \psi}{\delta r} = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \psi + V(x) \psi$ (1)

So:
$$\begin{cases} \frac{\delta \psi}{\delta r} = -\frac{\hbar^2}{2im} \frac{\partial^2}{\partial x^2} \psi + \frac{1}{i\hbar} V(x) \psi \\ \frac{\delta \psi^*}{\delta r} = \frac{\hbar^2}{2im} \frac{\partial^2}{\partial x^2} \psi^* - \frac{1}{i\hbar} V(x) \psi^* \end{cases}$$

$$\frac{\delta W}{\delta r} = \frac{\delta W_1}{\delta r} + \frac{\delta W_2}{\delta r} \quad w/$$

$$(*) \quad \frac{\delta W_1}{\delta r} = \frac{1}{4\pi i m} \int e^{-ipx/\hbar} \left[\frac{\partial^2 \psi^*(x-y/2)}{\partial x^2} \psi(x+y/2) + \frac{\partial^2 \psi(x+y/2)}{\partial x^2} \psi^*(x-y/2) \right] dy$$

~~$$(**) \quad \frac{\delta W_2}{\delta r} = \frac{2\pi}{i\hbar^2} \int e^{-ipx/\hbar} [V(x+y/2) - V(x-y/2)] \psi^*(x-y/2) \psi(x+y/2) dy$$~~

$$\frac{\delta \mathcal{W}_0}{\delta t} = \frac{2\pi}{i\hbar^2} \int e^{-ipg/\hbar} \left[V(x+y/2) - V(x-y/2) \right] \psi^*(x-y/2) \times \psi(x+y/2) dy \quad (**)$$

Astuce :
(*) :

$$\frac{\partial^2 \psi^*(x-y/2)}{\partial x^2} = -2 \frac{\partial^2 \psi^*(x-y/2)}{\partial x \partial y}$$

$$\text{car : } \frac{\partial \psi^*(x-y/2)}{\partial x} = -2 \frac{\partial \psi^*(x-y/2)}{\partial y}$$

(2)

~~IPP?~~

~~$\frac{\delta \mathcal{W}_0}{\delta t} = -2 \int e^{-ipg/\hbar} \frac{\partial^2 \psi^*(x-y/2)}{\partial y \partial x} \psi(x+y/2) dy$~~

~~$\frac{\delta \mathcal{W}_0}{\delta t} = -\frac{2ip}{\hbar} \int e^{-ipg/\hbar} \frac{\partial \psi^*(x-y/2)}{\partial x} \frac{\partial \psi(x+y/2)}{\partial x} dy$~~

~~$+ \int e^{-ipg/\hbar} \frac{\partial \psi^*(x-y/2)}{\partial x} \frac{\partial \psi(x+y/2)}{\partial x} dy$~~

$$\int e^{-ipy/\hbar} \frac{\partial^2 \psi^*(x-y/\hbar)}{\partial x^2} \psi(x+y/\hbar) dy$$

$$= -2 \int e^{-ipy/\hbar} \frac{\partial^2 \psi^*(x-y/\hbar)}{\partial x \partial y} \psi(x+y/\hbar) dy \quad (3)$$

$$\stackrel{\text{I.P.P.}}{=} -\frac{2ip}{\hbar} \int e^{-ipy/\hbar} \frac{\partial \psi^*(x-y/\hbar)}{\partial x} \psi(x+y/\hbar) dy$$

$$+ 2 \int e^{-ipy/\hbar} \frac{\partial \psi^*(x-y/\hbar)}{\partial x} \underbrace{\frac{\partial \psi(x+y/\hbar)}{\partial y}}_{\text{transfer again.}} dy$$

$$= -\frac{2ip}{\hbar} \int e^{-ipy/\hbar} \frac{\partial \psi^*(x-y/\hbar)}{\partial x} \psi(x+y/\hbar) dy$$

$$+ \int e^{-ipy/\hbar} \frac{\partial \psi^*(x-y/\hbar)}{\partial x} \frac{\partial \psi(x+y/\hbar)}{\partial x} dy$$

Same w/ 2nd term of eq (4):

$$-\int e^{-ipy/\hbar} \psi^*(x-y/\hbar) \frac{\partial^2 \psi(x+y/\hbar)}{\partial x^2} dy =$$

$$= -2 \int e^{-ipy/\hbar} \psi^*(x-y/\hbar) \frac{\partial^2 \psi(x+y/\hbar)}{\partial x \partial y} dy$$

$$= -\frac{2ip}{\hbar} \int e^{-ipy/\hbar} \psi^*(x-y/\hbar) \frac{\partial \psi(x+y/\hbar)}{\partial x} dy$$

$$- \frac{2ip}{\hbar} \int e^{-ipy/\hbar} \frac{\partial \psi^*(x-y/\hbar)}{\partial x} \frac{\partial \psi(x+y/\hbar)}{\partial x} dy$$

(*) \Rightarrow

(4)

$$\frac{\partial \mathcal{W}_T}{\partial T} = \frac{1}{4\pi m} x - \frac{2ip}{\hbar} \int e^{-ipy/\hbar} \frac{\partial}{\partial x} \left(\psi^* (x-y/2) \psi (x+y/2) \right)$$

$$= -\frac{p}{\hbar m} \frac{\partial \mathcal{W}_T}{\partial x}$$

(*) : ~~$U(x+y/2) = \sum_n \frac{1}{n!} \frac{\partial^n U(x)}{\partial x^n} \left(\frac{y}{2} \right)^n$~~

$$U(x+y/2) = \sum_n \frac{1}{n!} \frac{\partial^n U(x)}{\partial x^n} \bigg|_x \left(\frac{y}{2} \right)^n$$

$$U(x+y/2) - U(x-y/2) = \sum_n \frac{1}{n!} \frac{\partial^n U(x)}{\partial x^n} \bigg|_x \left[\left(\frac{y}{2} \right)^n - \left(-\frac{y}{2} \right)^n \right]$$

$$= \sum_s \frac{1}{(2s+1)!} \frac{\partial^{2s+1} U(x)}{\partial x^{2s+1}} 2 \left(\frac{y}{2} \right)^{2s+1}$$

$$= \sum_s \frac{1}{(2s+1)!} \frac{\partial^{2s+1} U(x)}{\partial x^{2s+1}} \left(\frac{1}{2} \right)^{2s} y^{2s+1}$$

$$\frac{\partial \mathcal{W}_0}{\partial T} = \frac{2\pi}{i\hbar^2} \sum_s \frac{1}{(2s+1)!} \frac{\partial^{2s+1} U(x)}{\partial x^{2s+1}} \left(\frac{1}{2} \right)^{2s} \frac{\partial^{2s+1} U(x)}{\partial x^{2s+1}} \times$$

$$\times \int e^{-ipy/\hbar} y^{2s+1} \psi^* (x-y/2) \psi (x+y/2) dy$$

$$= \frac{2\pi}{i\hbar^2} \sum_s \frac{1}{(2s+1)!} \left(\frac{1}{2} \right)^{2s} \frac{\partial^{2s+1} U(x)}{\partial x^{2s+1}} \frac{\partial^{2s+1} \mathcal{W}(x,p)}{\partial p^{2s+1}}$$

$$= \sum_s (-\hbar^2)^s \frac{1}{(2s+1)!} \left(\frac{1}{2} \right)^{2s} \frac{\partial^{2s+1} U(x)}{\partial x^{2s+1}} \frac{\partial^{2s+1} \mathcal{W}(x,p)}{\partial p^{2s+1}}$$

HO or for ref:

(5)

$$\frac{\partial U}{\partial x_m} = 0 \quad \text{for } m > 1.$$

$$L \quad \frac{\partial U(x)}{\partial t} = \cancel{L} \quad \frac{\partial U(x)}{\partial x} \quad \frac{\partial U(x, p)}{\partial p}$$

$$S_0 : \quad \frac{\partial W(x, p)}{\partial t} = - \frac{p}{m} \frac{\partial W(x, p)}{\partial x} + \frac{\partial V(x)}{\partial x} \quad \frac{\partial W(x, p)}{\partial p}$$

Classical point-like part:

$$x_0 = x \cos(\omega t) - \frac{p}{m\omega} \sin(\omega t)$$

$$p_0 = p \cos(\omega t) + m\omega x \sin(\omega t)$$

So:

$$W(x, p, t) = W\left(x \cos(\omega t) - \frac{p}{m\omega} \sin(\omega t), \right.$$

$$\left. p \cos(\omega t) + m\omega x \sin(\omega t), 0\right)$$