

The Dynamical Origin of $m_{\phi,\text{crit}} \simeq 1.965$: Phase-Boundary Physics in Siamese Universes

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Abstract

In the Phase–Bifurcation Cosmogenesis scenario, the early universe is described by a CPT-symmetric pair of “Siamese” branches that share a quantum vacuum at a bifurcation surface and then acquire a relative phase difference $\Delta\phi$ as they evolve. The late-time fate partitions into three sectors: (A) synchrony ($\Delta\phi \rightarrow 0$), (B) antipodal (trapping near π), and (C) escape/oscillatory (recurrent barrier crossing). A sharp transition in the synchrony probability P_A at an effective phase mass $m_{\phi,\text{crit}} \simeq 1.965$ was reported previously. We reanalyse its origin and robustness with a controlled grid scan over (m_ϕ, k_{rot}) , using $N_\Delta = 40$ initial phases per grid point and uniform sampling in $\Delta\phi_{\text{ini}} \in [0, \pi]$. We confirm a clean monotonic transition in $P_A(m_\phi)$ with a global $m_{\phi,\text{crit}}$ and find that the probability band in (m_ϕ, k_{rot}) is nearly vertical: m_ϕ is the primary control parameter, while k_{rot} produces a modest modulation peaking around $k_{\text{rot}} \simeq 0.33$. Phase portraits at $m_{\phi,\text{crit}}$ show a single stable spiral at $\Delta\phi = 0$ and a smooth separatrix. These results establish $m_{\phi,\text{crit}}$ as a robust dynamical scale with implications for CPT-symmetric cosmology, baryogenesis and directional observables.

1 Introduction

CPT-symmetric cosmologies posit that our visible universe is mirrored by a partner evolving with the opposite arrow of time across the Big Bang [1–4]. In Siamese frameworks the relative phase between the branches becomes a physical degree of freedom. The phase mass m_ϕ controls whether the system locks in synchrony (A), is trapped antipodally (B), or escapes into oscillations (C), akin to driven-damped nonlinear oscillators [5–7]. Clarifying whether the previously reported value $m_{\phi,\text{crit}} \simeq 1.965$ is a numerical artefact or a structurally stable feature is essential because P_A feeds directly into phase-locked vs. misaligned cosmological histories and thereby into baryogenesis pathways and large-scale anisotropies.

2 Effective phase dynamics

We treat $\Delta\phi$ as an angular variable with periodic potential

$$V_{\text{eff}}(\Delta\phi) = m_\phi^2 [1 - \cos(\Delta\phi)], \quad (1)$$

leading, in $N = \ln a$, to

$$\Delta\phi'' + (3 + \xi(N)) \Delta\phi' + m_\phi^2 \sin \Delta\phi = S_{\text{rot}}(N; k_{\text{rot}}), \quad (2)$$

with a decaying rotational source $S_{\text{rot}}(N; k_{\text{rot}}) = k_{\text{rot}} e^{-\gamma N} f(N)$, $\gamma \simeq 2$. Initial conditions at N_{ini} adopt $\Delta\phi'(N_{\text{ini}}) = 0$ and $\Delta\phi_{\text{ini}} \in [0, \pi]$ uniformly.

3 Numerical setup and sector classification

We scan a 7×7 grid

$$m_\phi \in \{0.375, 0.75, 1.25, 1.625, 1.95, 2.5, 3.0\}, \quad (3)$$

$$k_{\text{rot}} \in \{0.0, 0.083, 0.167, 0.25, 0.333, 0.417, 0.5\}, \quad (4)$$

and draw $N_\Delta = 40$ values of $\Delta\phi_{\text{ini}}$ per point. Late-time, sector labels use time-averaged thresholds over $[N_f - \Delta N, N_f]$ with $\epsilon_\phi = 0.1$, $\epsilon_\pi = 0.1$, $\epsilon_{\phi'} = 0.05$.

4 Results

Global $P_A(m_\phi)$ and $m_{\phi,\text{crit}}$.

$P_A(m_\phi, k_{\text{rot}})$ and local $m_{\phi,\text{crit}}(k_{\text{rot}})$.

Phase-sector geometry and trajectories.

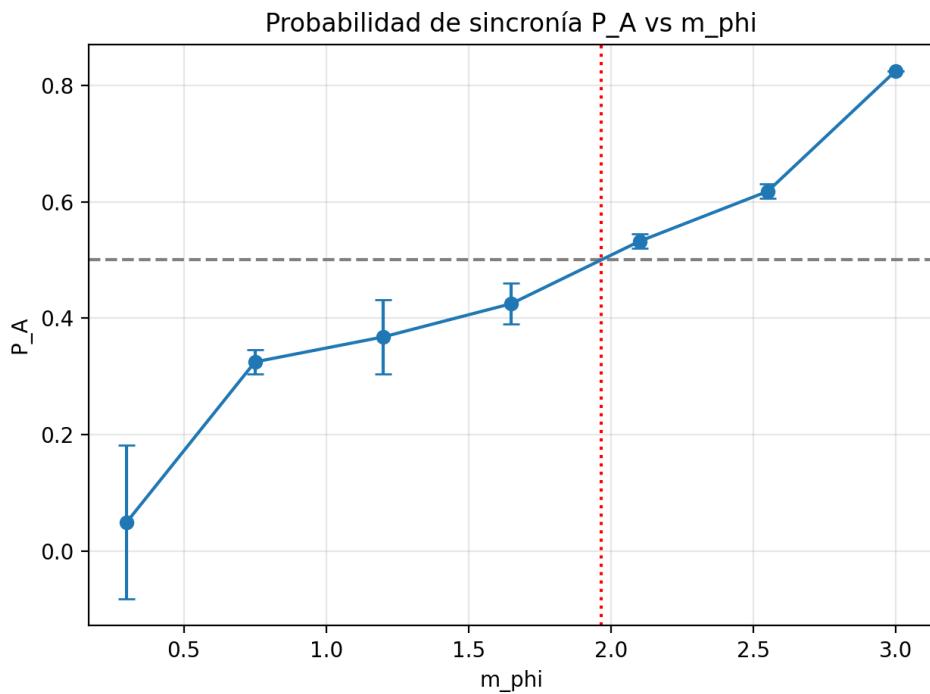


Figure 1: Synchrony probability $P_A(m_\phi)$ averaged over k_{rot} . Dashed line: $P_A = 0.5$. Dotted vertical: $m_{\phi,\text{crit}} \simeq 1.965$.

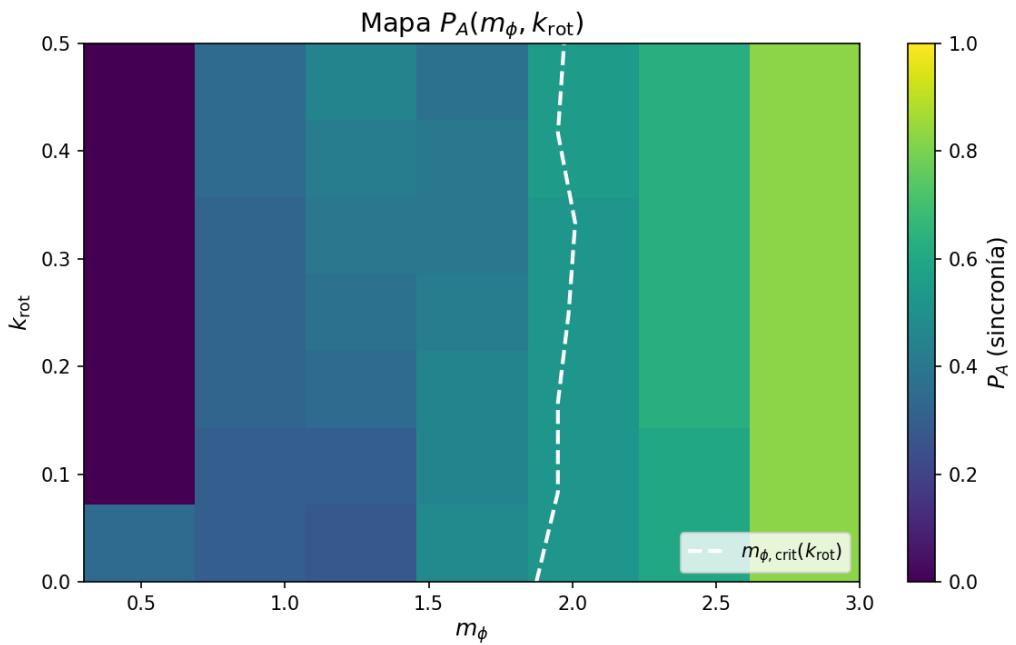


Figure 2: 2D map $P_A(m_\phi, k_{\text{rot}})$. White dashed: locus $P_A = 0.5$ i.e. local $m_{\phi,\text{crit}}(k_{\text{rot}})$. Nearly vertical transition band.

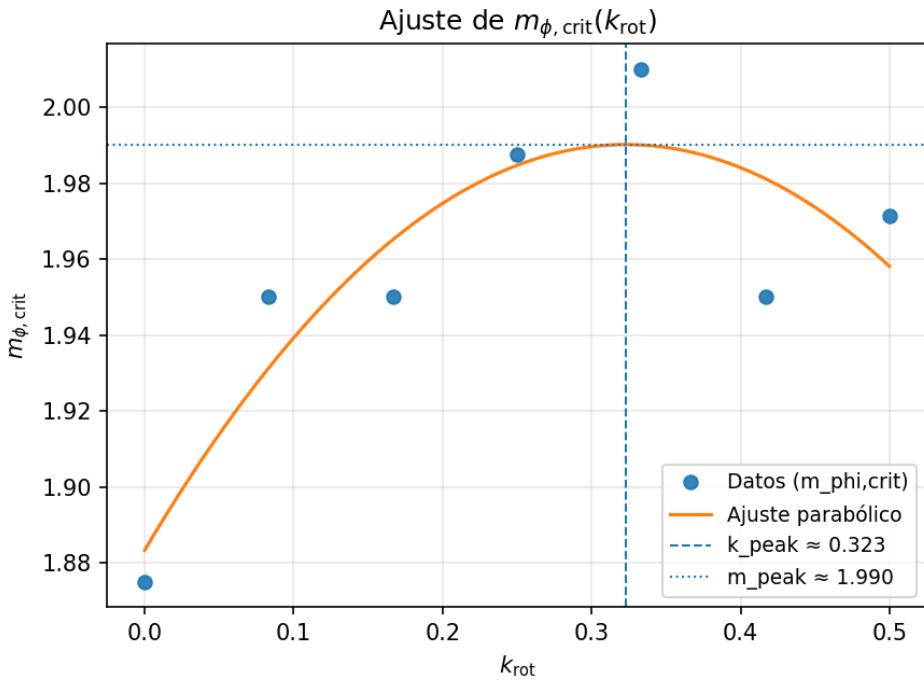


Figure 3: Quadratic fit to $m_{\phi, \text{crit}}(k_{\text{rot}})$ showing a weak maximum near $k_{\text{rot}} \simeq 0.33$.

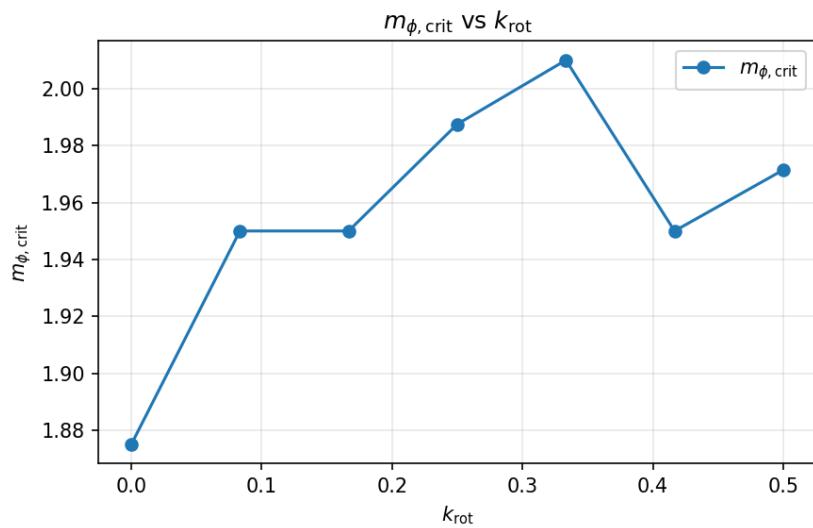


Figure 4: Discrete points used in the fit of Fig. 3.

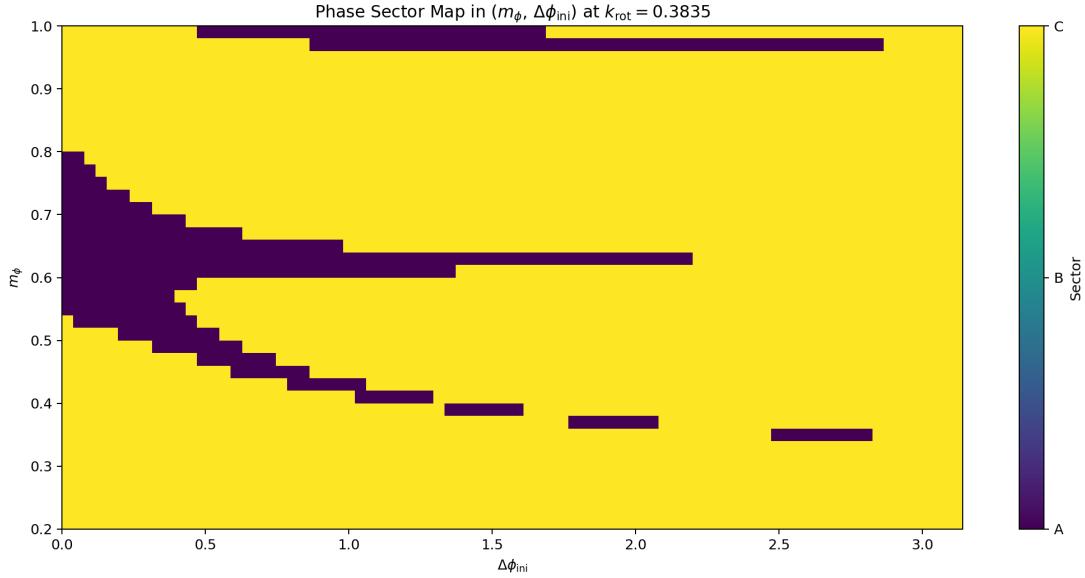


Figure 5: Phase-sector map in $(m_\phi, \Delta\phi_{\text{ini}})$ at representative k_{rot} . Dark: synchrony (A). Yellow: escape (C).

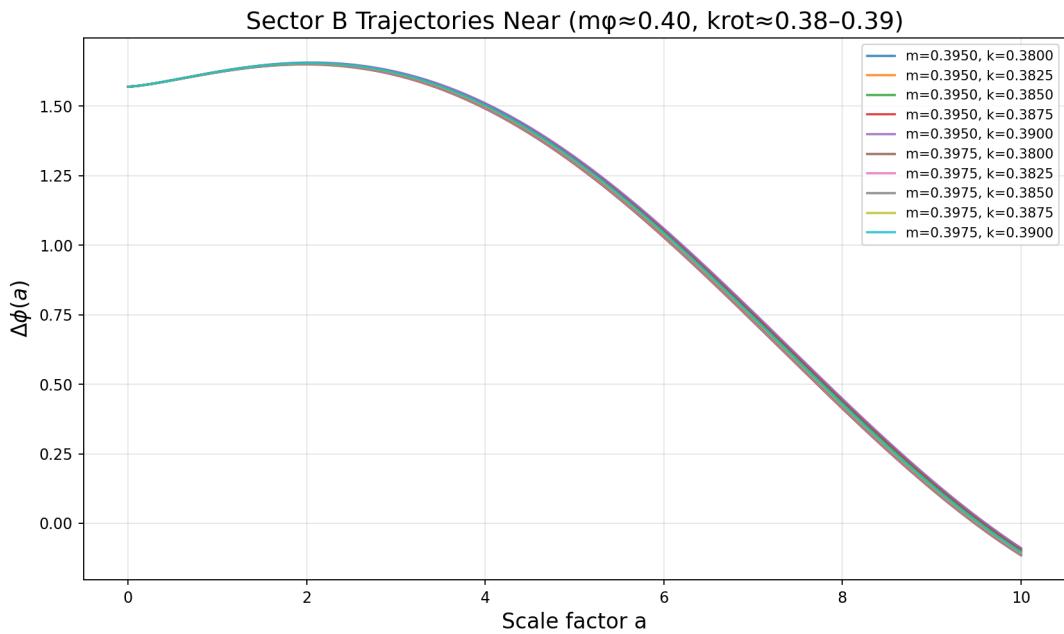


Figure 6: Representative trajectories in sector B (antipodal trapping) near $(m_\phi \approx 0.40, k_{\text{rot}} \approx 0.38-0.39)$.

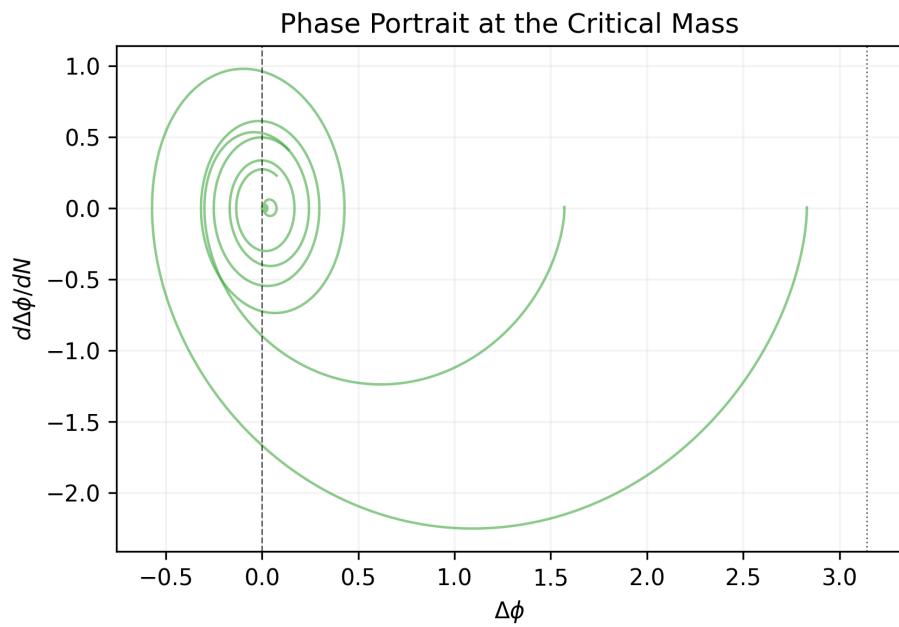


Figure 7: Phase portrait at $m_{\phi,\text{crit}}$: estable espiral en $\Delta\phi = 0$ y separatrix suave.

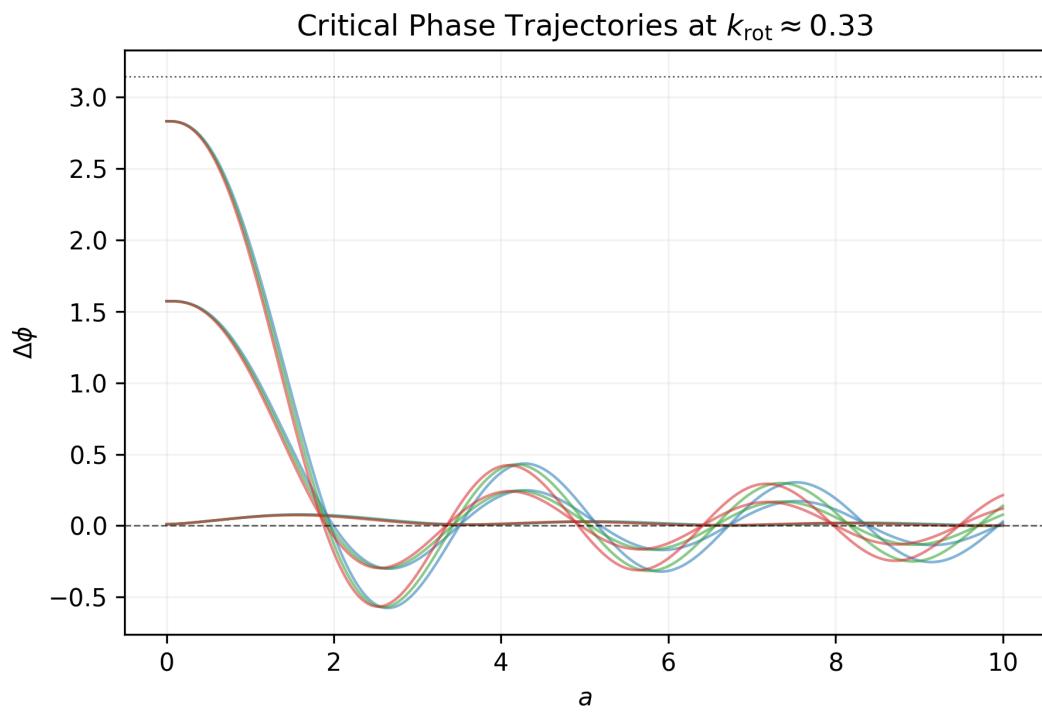


Figure 8: Near-critical families of trajectories around $k_{\text{rot}} \simeq 0.33$ ilustrando el origen del salto suave en $P_A(m_{\phi})$.

5 Data & Reproducibility

All figures are generated from artefacts in `results_phase_sectors/`. **Figures (PNG):** `PA_vs_mphi.png`, `PA_map_mphi_krot_with_mcrit.png`, `mphi_crit_vs_krot_fit.png`, `mphi_crit_vs_krot.png`, `phase_sector_map_delta.png`, `sectorB_trajectories.png`, `crit_phase_portrait.png`, `crit_trajectories.png`. **Grid:** m_ϕ and k_{rot} values as above; $N_\Delta = 40$; $\Delta\phi_{\text{ini}} \sim \mathcal{U}[0, \pi]$. **Integrator:** RK45 with $\text{rtol} = 10^{-7}$, $\text{atol} = 10^{-9}$; $\Delta N_{\text{max}} = 10^{-2}$. **Window:** $N_{\text{ini}} = -4$, $N_f = +10$; late averages on $\Delta N = 1$. **Thresholds:** $\epsilon_\phi = 0.1$, $\epsilon_\pi = 0.1$, $\epsilon_{\phi'} = 0.05$. **Seed:** 42. **Post-processing:** (i) $P_A = N_A/N_\Delta$; (ii) interpolate $P_A(m_\phi)$ to $P_A = 0.5$; (iii) compute $m_{\phi, \text{crit}}(k_{\text{rot}})$ and fit $ak_{\text{rot}}^2 + bk_{\text{rot}} + c$.

6 Discussion and Outlook

Our grid scan confirms a smooth transition yielding a well-defined $m_{\phi, \text{crit}}$. The near-vertical band in (m_ϕ, k_{rot}) and the modest peak of $m_{\phi, \text{crit}}(k_{\text{rot}})$ around 0.33 indicate that rotation slightly shifts the separatrix but does not alter the basic fate map. The portrait at $m_{\phi, \text{crit}}$ exhibits a single stable spiral at $\Delta\phi = 0$, consistent with a damped nonlinear oscillator. Next steps include embedding the dynamics in a calibrated Λ CDM background and replacing the uniform prior on $\Delta\phi_{\text{ini}}$ by a quantum-motivated distribution; on the observational side, FRB and CMB anisotropies provide falsifiable tests.

Acknowledgements

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References

- ### References
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