

Searching for patterns in the Universe



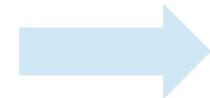
Gustavo Niz
Universidad de Guanajuato

Content

Our Universe

What we know...

What we don't know...



Perturbation theory in the “cosmic collider”

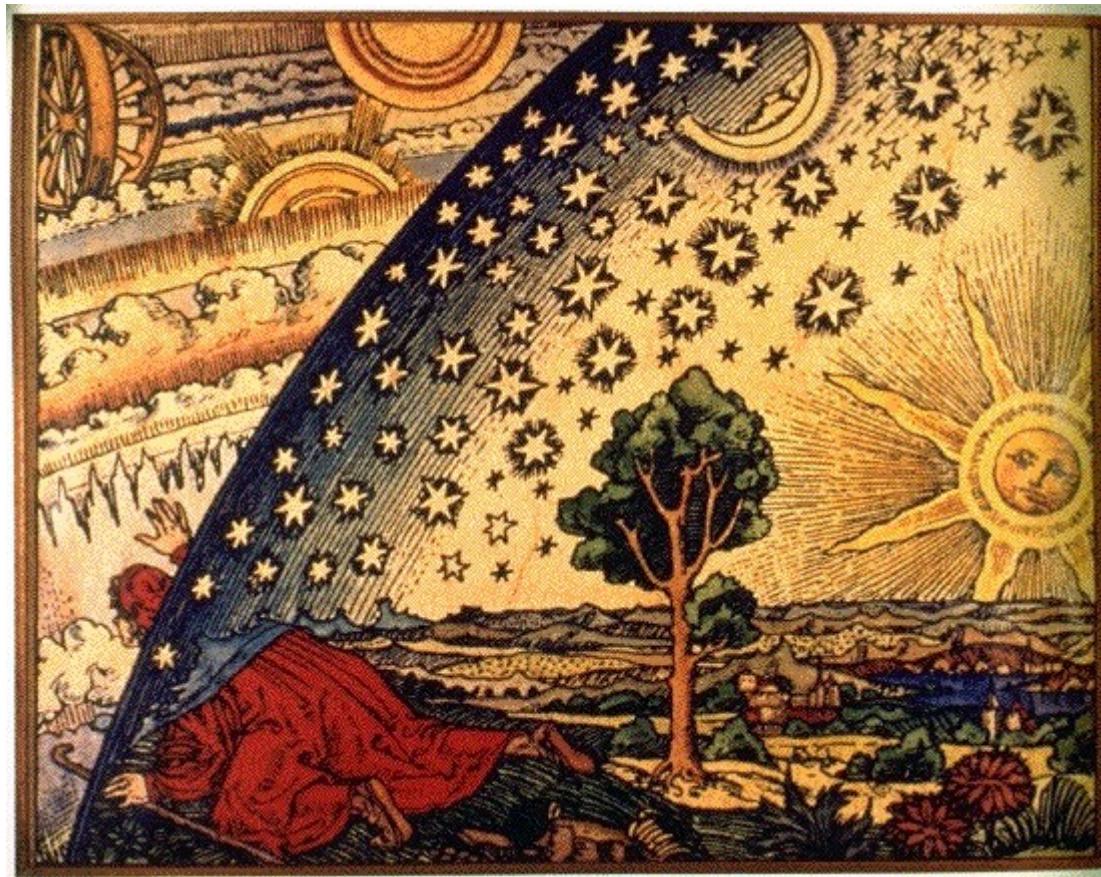
Early Universe

Cosmic Microwave Background (CMB)

Large Scale Structure (LSS)

Dark Energy Spectroscopic Instrument

What do we know about the Universe?



Flammarion, 1888

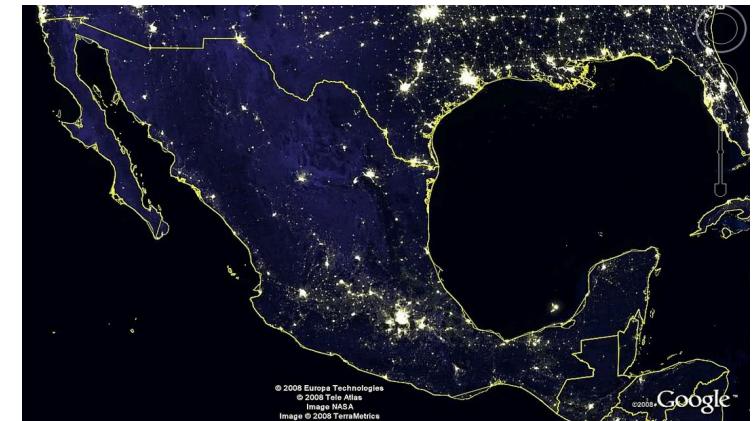
What do we know about the Universe?

We observe the Universe by

1) Mostly



(photons)



2) but also by: neutrinos, alpha particles, gravitational waves, ...

Note: We observe the past and only on the “light-cone”



What we know...



Sloan Digital Sky Survey

What we know...

Our visible Universe is BIG

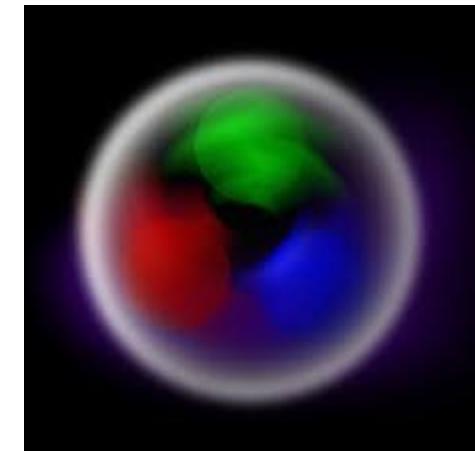
$$93 \text{ billion light years} = 10^{26} \text{ m}$$

Universe



Leon 10^3 m

Sun 10^8 m



Proton 10^{-15} m

What we know...

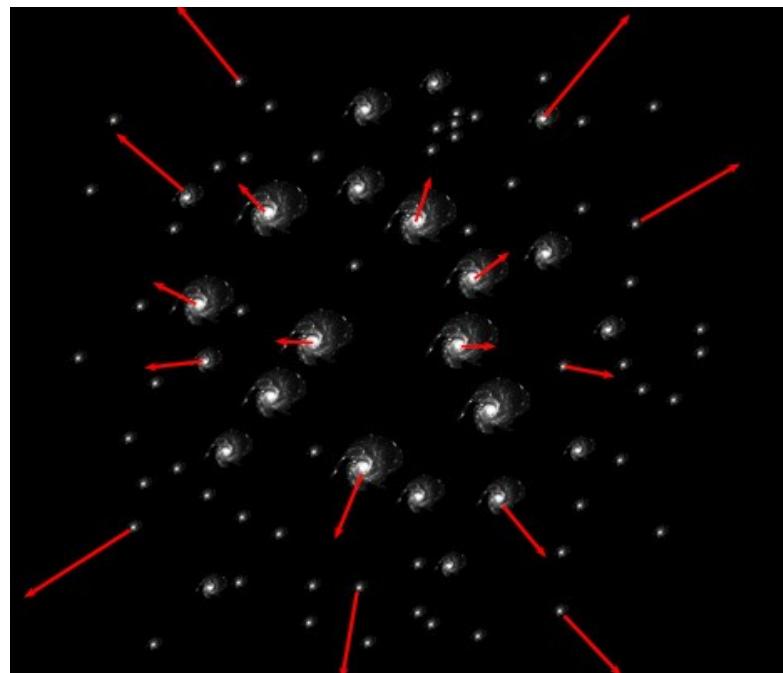
13.7 billion years



Oldest stars ~ 13.6 billion years

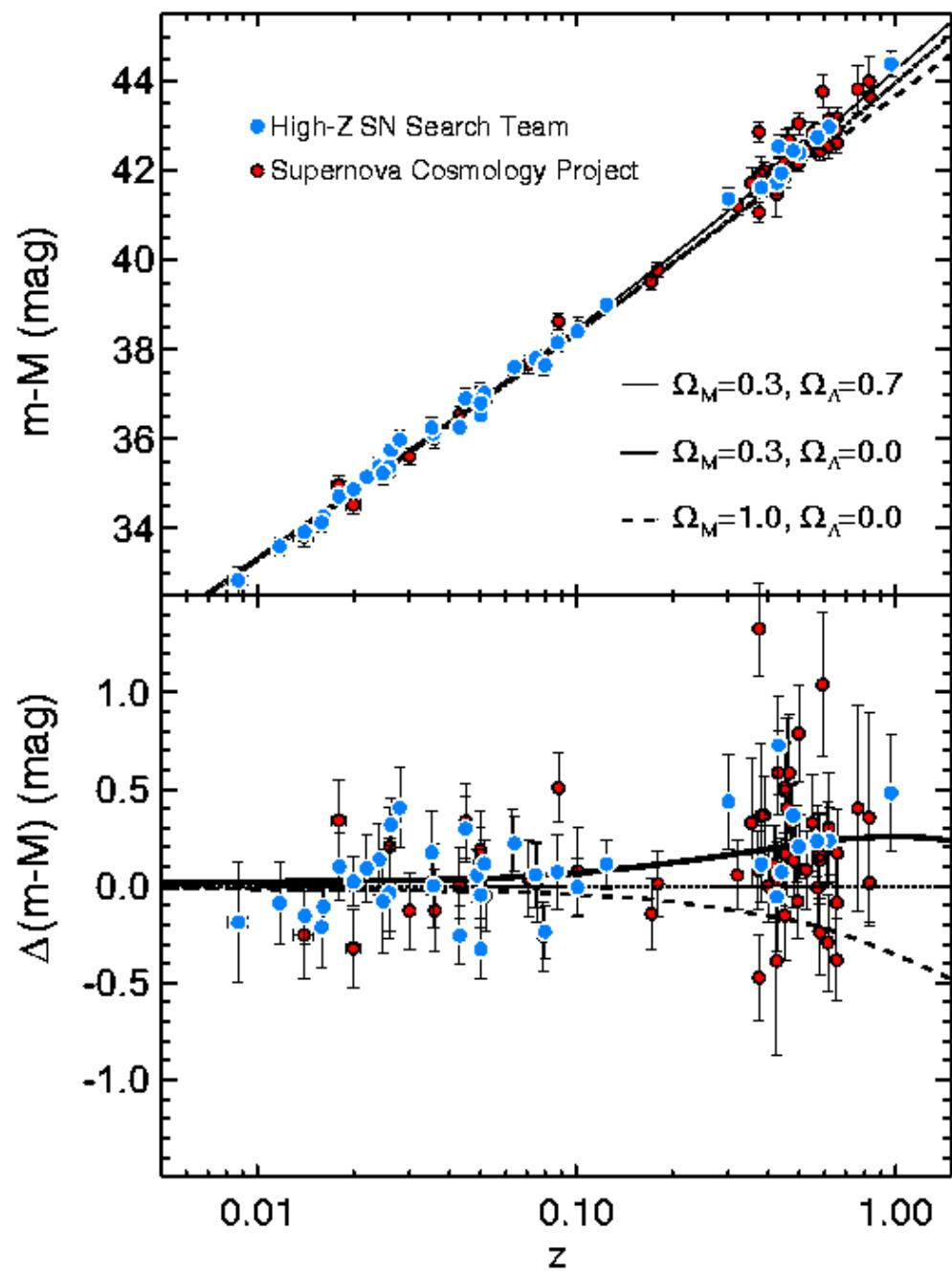
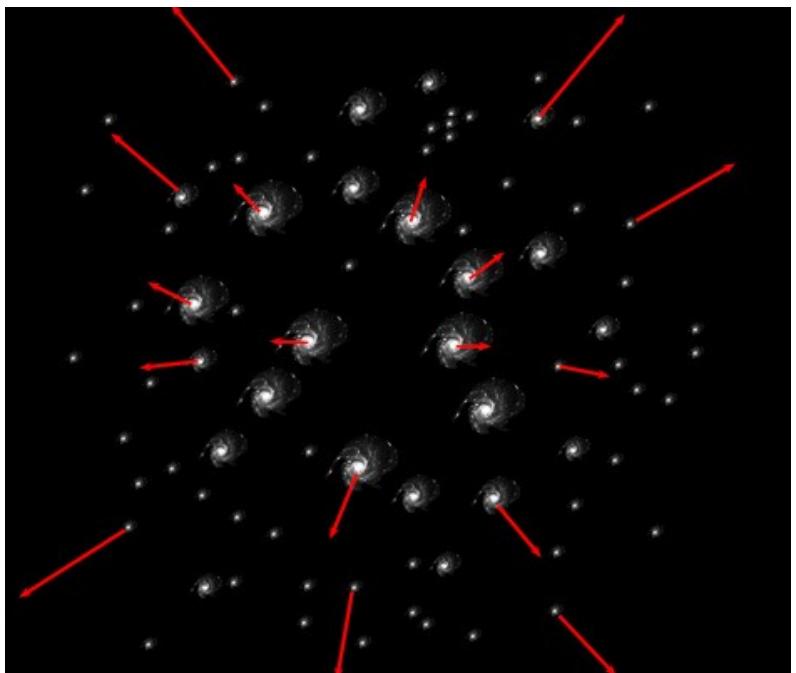
What we know...

The Universe expands



What we know...

The Universe expands



What we know...

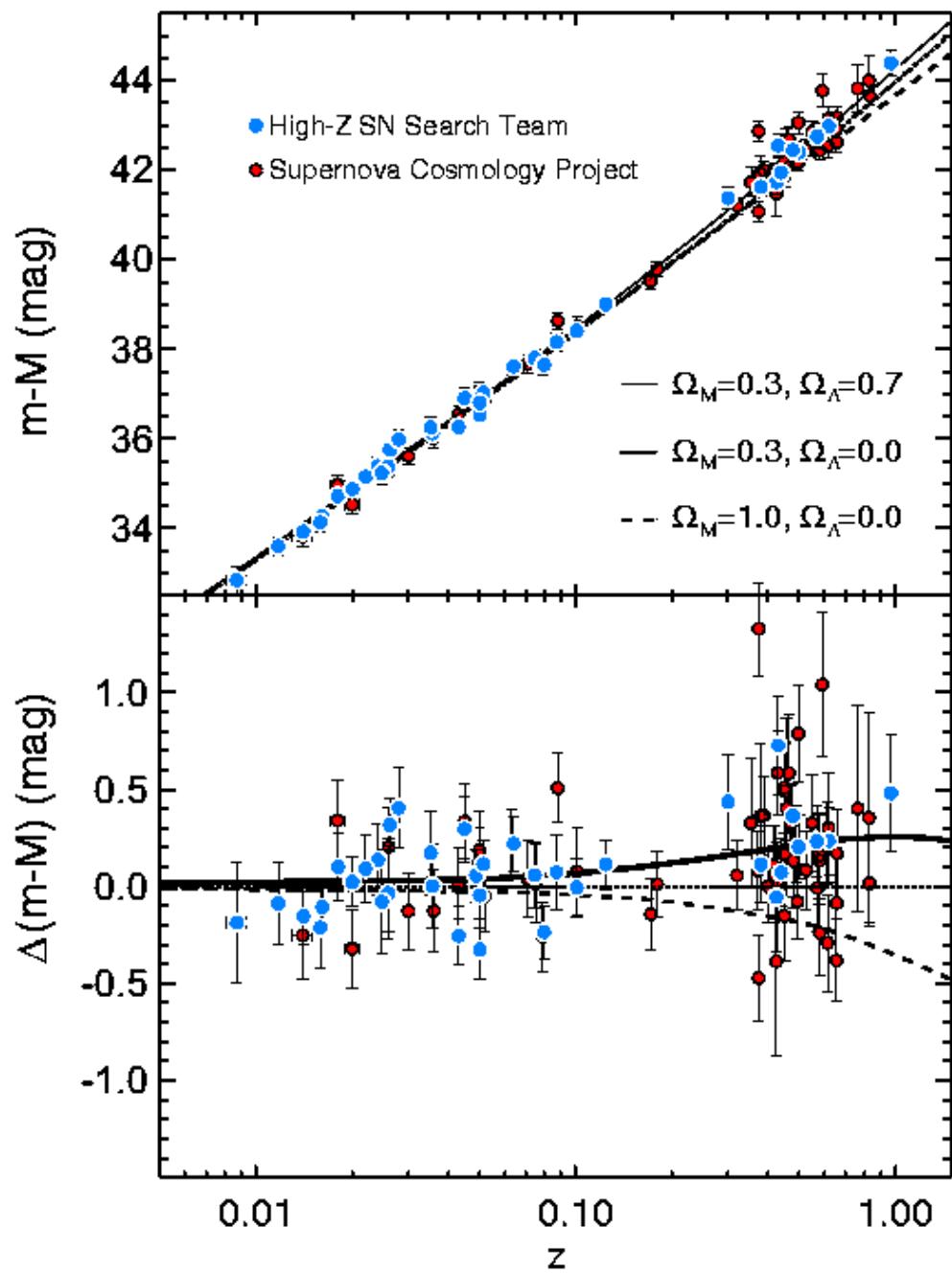
The Universe expands

Actually, it accelerates!

(Nobel 2011 to Perlmutter,
Riess & Schmidt).

Cause:

DARK ENERGY!



What we know...

Expansion

gave rise to



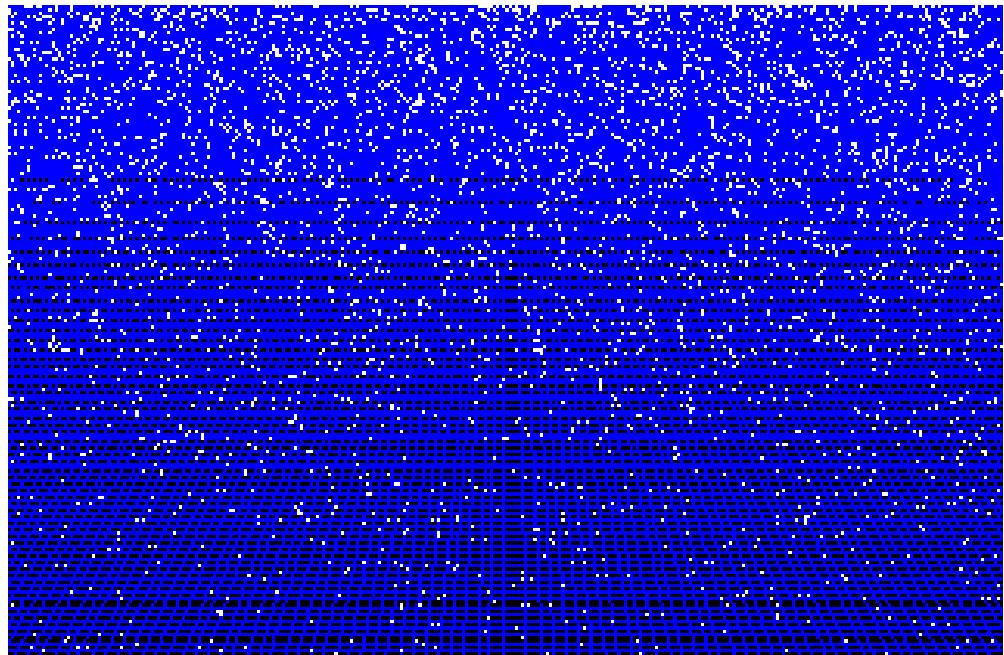
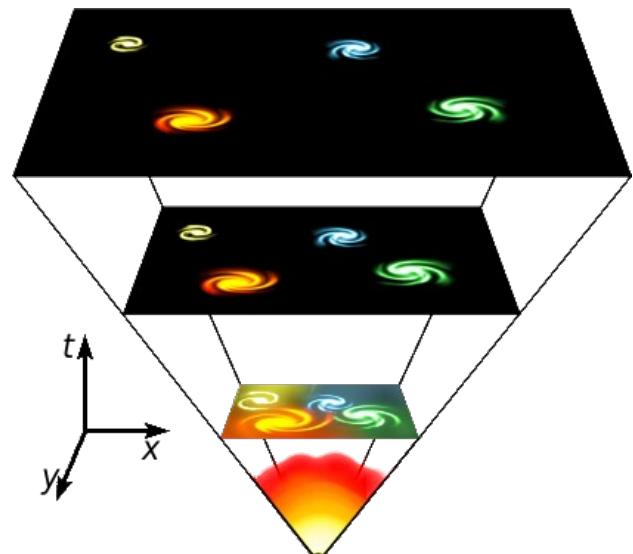
What we know...

Expansion

gave rise to



What we know...

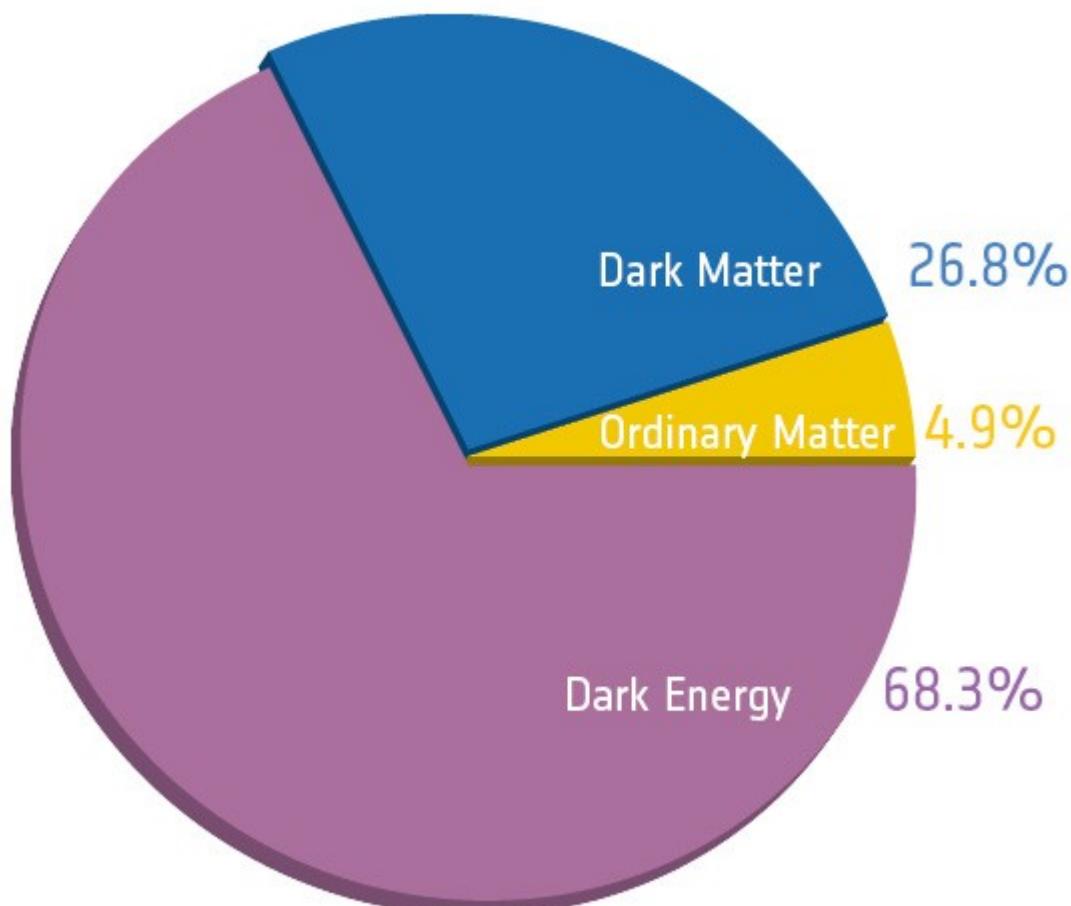


Observations

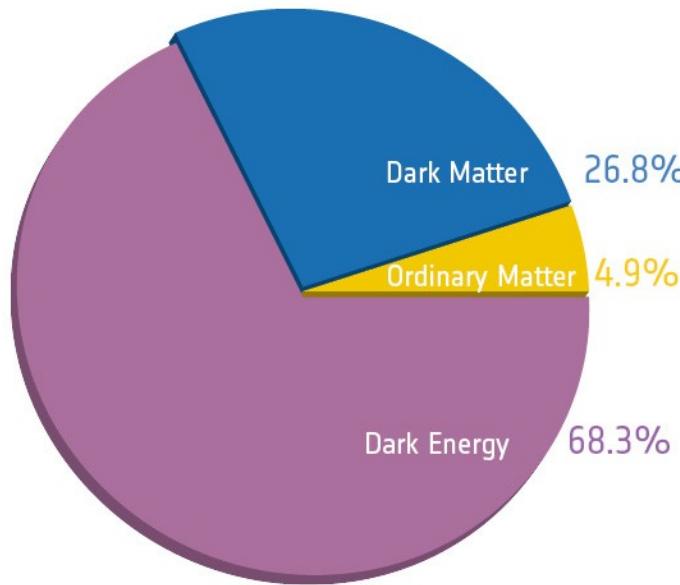
- There was no explosion
- No expansion into something else
- Expansion may be faster than speed of light

What we know...

Composed today by



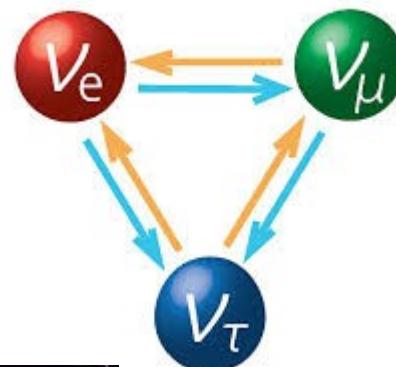
What we know...



Stars ~ 0.5%



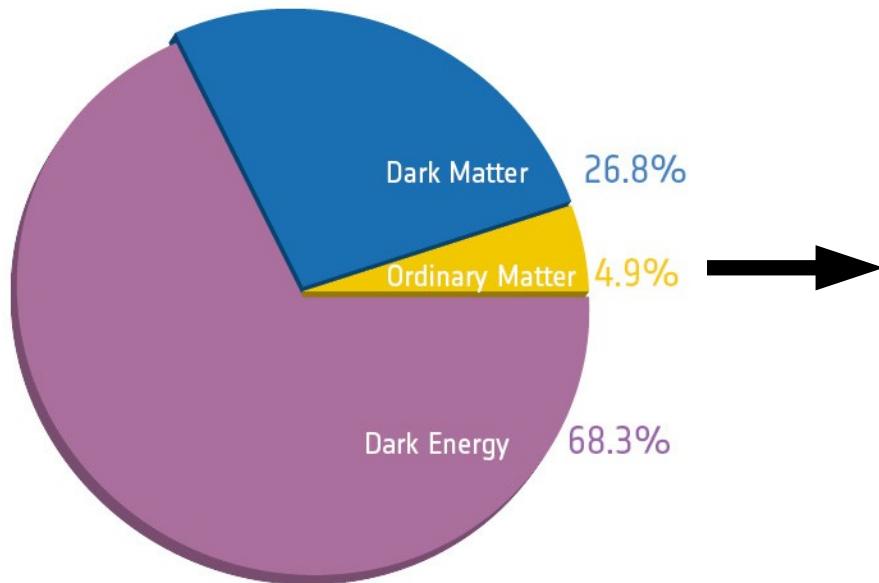
Heavy elements ~ 0.03%



Neutrinos ~ 0.3%

Free hydrogen & Helium ~ 4%

What we know...



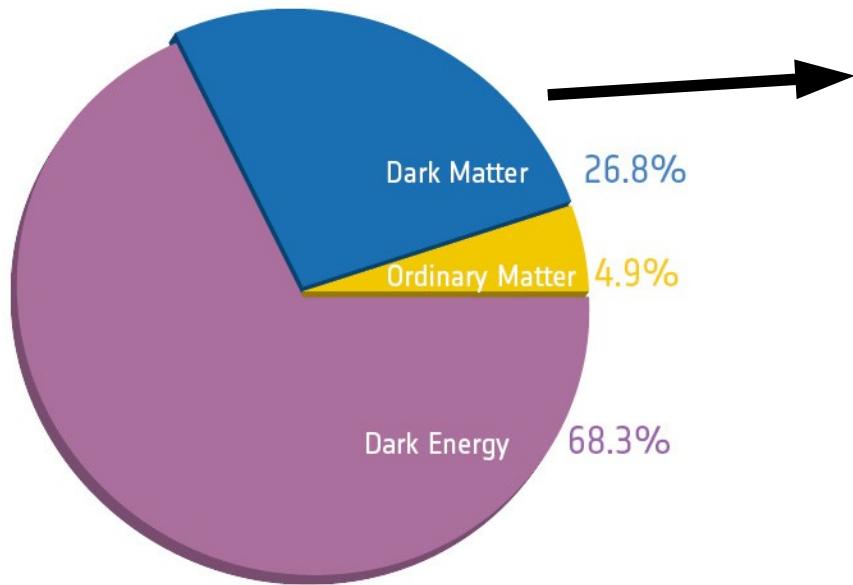
The infographic is titled "Otherworldly Elements" and features four sections about elements:

- Copernicium:** A metal that's a gas at room temperature. It is shown as a glowing, wispy cloud. Periodic table entry: 112 (285) Cn Copernicium.
- Curium:** So radioactive it glows in the dark. It is shown as a glowing sphere. Periodic table entry: 96 (247) Cm Curium.
- Caesium:** Explodes when it hits water. It is shown as a small glass tube containing a red liquid. Periodic table entry: 55 132.90547 Cs Cesium.
- Bismuth:** Can levitate a magnet between two pieces of itself. It is shown as a colorful, metallic crystal. Periodic table entry: 83 208.980 Bi Bismuth.

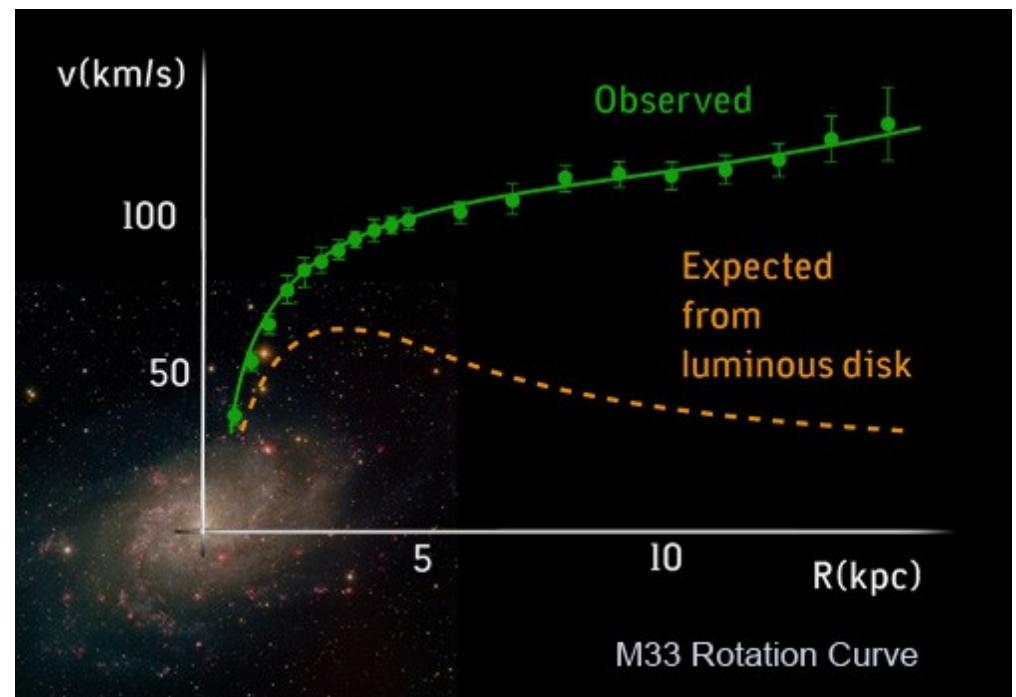
Learn more at curiosity.com

Source: PBS

What we do not know...



Dark Matter

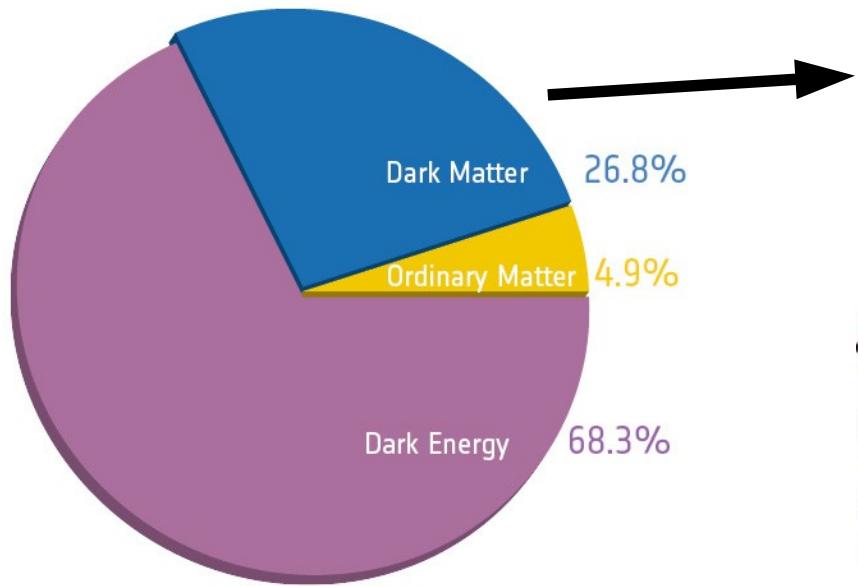


Best candidate:

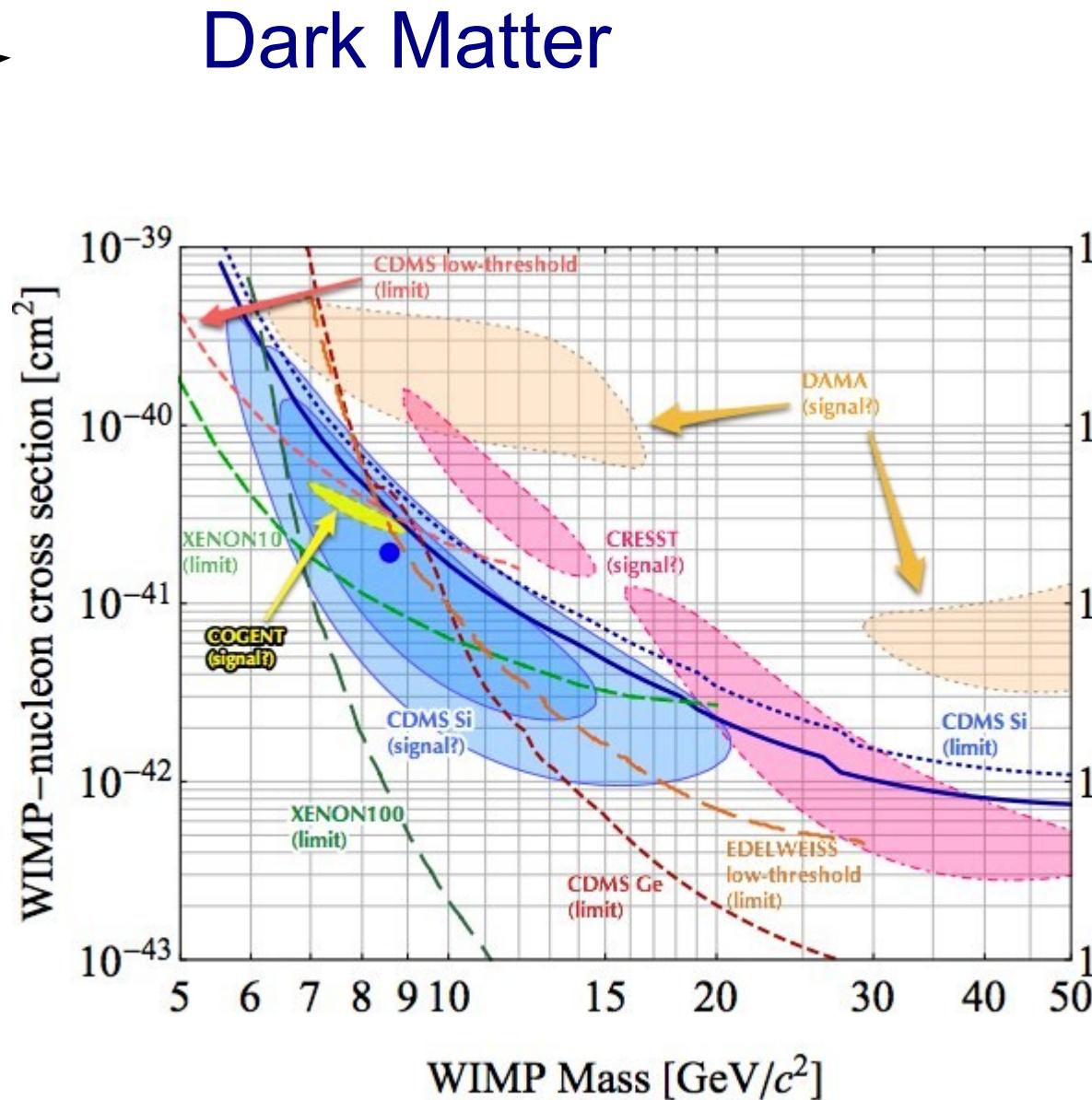
Massive particle with:

Null or weak interactions with the visible matter

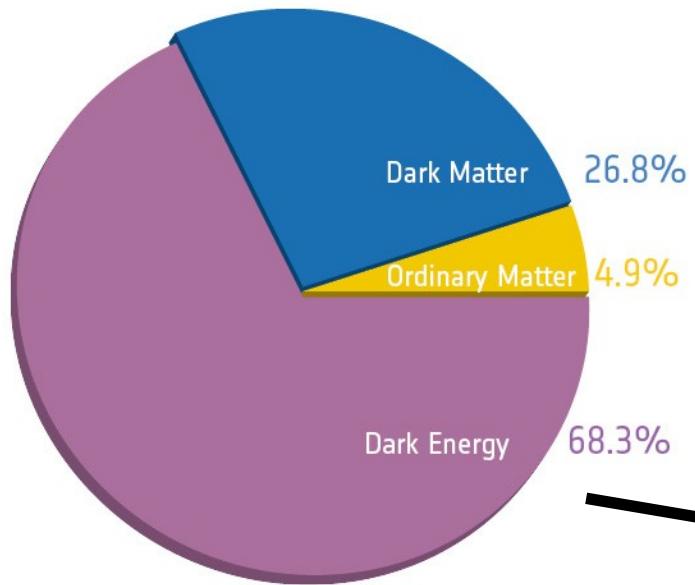
What we do not know...



Tens of experiments
...searching...



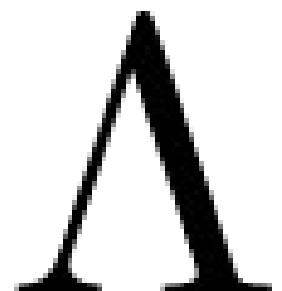
What we do not know...



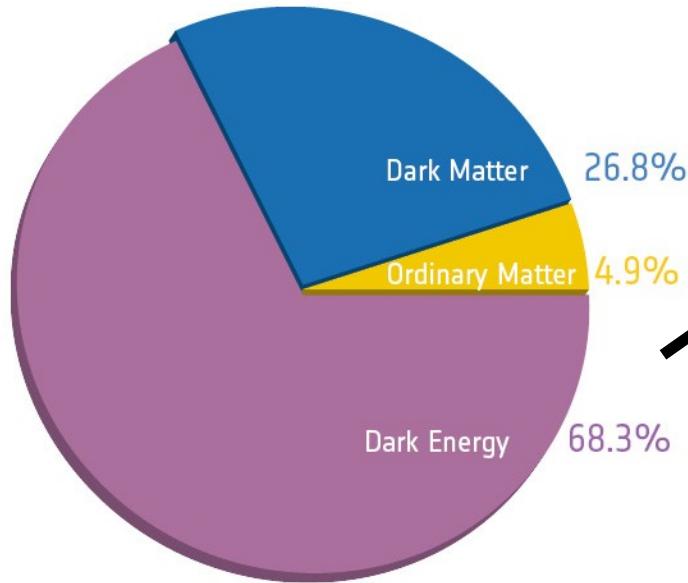
Dark Energy

Best candidate:

Cosmological Constant



What we do not know...

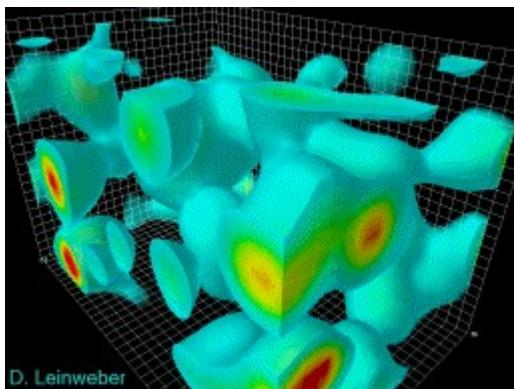


Dark Energy

Theoretical problem

Cosmological Constant value

$$\Lambda_{vac} = 10^{120} \Lambda_{obs}$$



Empty space perturbations

1000
00
00
00

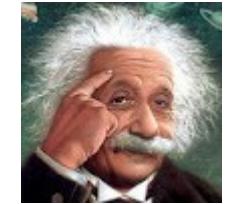
THE COSMIC COLLIDER

A Perturbed Universe

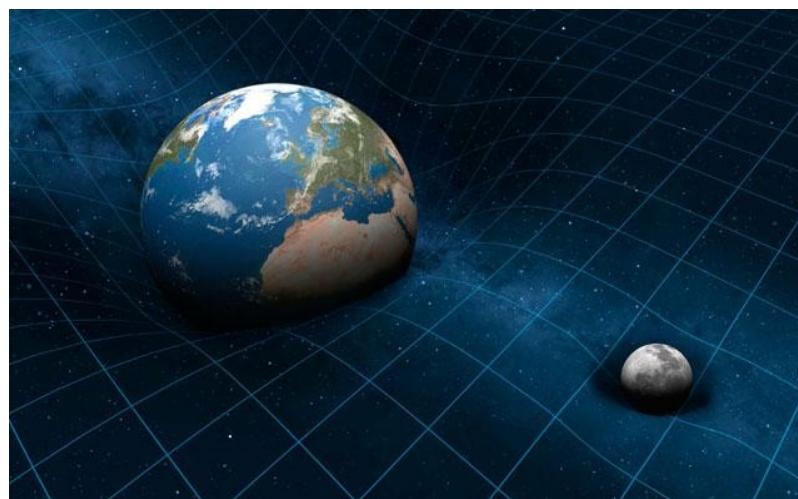
From Quantum Fluctuations to Human Beings...

Simple model

Einstein (1915)



$$G_{\mu\nu} = 8\pi G_N T_{\mu\nu}$$



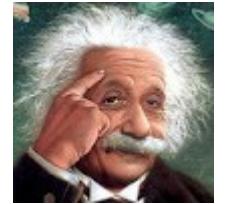
*Space-time
Geometry*

Matter/energy

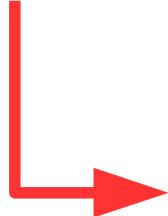
Non-linear theory - describes interactions!

$$G_{\mu\nu} = 8\pi G_N T_{\mu\nu}$$

Simple model



$\mu = \nu = 3$

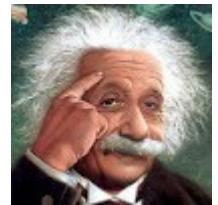


$$\begin{aligned}
 & \frac{1}{4} g_{00}^{-1} g_{11}^{-1} g_{22}^{-1} \frac{\partial g_{00}}{\partial y} \frac{\partial g_{11}}{\partial y} + \frac{1}{4} g_{00}^{-1} g_{11}^{-1} g_{33}^{-1} \frac{\partial g_{11}}{\partial t} \frac{\partial g_{33}}{\partial t} + \frac{1}{4} g_{00}^{-1} g_{11}^{-1} g_{33}^{-1} \frac{\partial g_{00}}{\partial x} \frac{\partial g_{33}}{\partial x} + \\
 & \frac{1}{4} g_{00}^{-1} g_{11}^{-1} g_{33}^{-1} \frac{\partial g_{00}}{\partial z} \frac{\partial g_{11}}{\partial z} + \frac{1}{2} g_{00}^{-1} g_{11}^{-1} \frac{\partial^2 g_{11}}{\partial t^2} - \frac{1}{4} (g_{00}^{-1})^2 g_{11}^{-1} \frac{\partial g_{00}}{\partial t} \frac{\partial g_{11}}{\partial t} - \\
 & \frac{1}{4} g_{00}^{-1} (g_{11}^{-1})^2 \left(\frac{\partial g_{11}}{\partial t} \right)^2 + \frac{1}{2} g_{00}^{-1} g_{11}^{-1} \frac{\partial^2 g_{00}}{\partial x^2} - \frac{1}{4} (g_{00}^{-1})^2 g_{11}^{-1} \left(\frac{\partial g_{00}}{\partial x} \right)^2 - \\
 & \frac{1}{4} g_{00}^{-1} (g_{11}^{-1})^2 \frac{\partial g_{00}}{\partial x} \frac{\partial g_{11}}{\partial x} + \frac{1}{4} g_{00}^{-1} g_{22}^{-1} g_{33}^{-1} \frac{\partial g_{00}}{\partial y} \frac{\partial g_{33}}{\partial y} + \frac{1}{2} g_{00}^{-1} g_{33}^{-1} \frac{\partial^2 g_{33}}{\partial t^2} - \\
 & \frac{1}{4} (g_{00}^{-1})^2 g_{33}^{-1} \frac{\partial g_{00}}{\partial t} \frac{\partial g_{33}}{\partial t} - \frac{1}{4} g_{00}^{-1} (g_{33}^{-1})^2 \left(\frac{\partial g_{33}}{\partial t} \right)^2 + \frac{1}{2} g_{00}^{-1} g_{33}^{-1} \frac{\partial^2 g_{00}}{\partial z^2} - \\
 & \frac{1}{4} (g_{00}^{-1})^2 g_{33}^{-1} \left(\frac{\partial g_{00}}{\partial z} \right)^2 - \frac{1}{4} g_{00}^{-1} (g_{33}^{-1})^2 \frac{\partial g_{00}}{\partial z} \frac{\partial g_{33}}{\partial z} + \frac{1}{4} g_{11}^{-1} g_{22}^{-1} g_{33}^{-1} \frac{\partial g_{11}}{\partial y} \frac{\partial g_{33}}{\partial y} + \\
 & \frac{1}{2} g_{11}^{-1} g_{33}^{-1} \frac{\partial^2 g_{33}}{\partial x^2} - \frac{1}{4} g_{11}^{-1} (g_{33}^{-1})^2 \left(\frac{\partial g_{33}}{\partial x} \right)^2 - \frac{1}{4} (g_{11}^{-1})^2 g_{33}^{-1} \frac{\partial g_{11}}{\partial x} \frac{\partial g_{33}}{\partial x} + \\
 & \frac{1}{2} g_{11}^{-1} g_{33}^{-1} \frac{\partial^2 g_{11}}{\partial z^2} - \frac{1}{4} g_{11}^{-1} (g_{33}^{-1})^2 \frac{\partial g_{11}}{\partial z} \frac{\partial g_{33}}{\partial z} - \frac{1}{4} (g_{11}^{-1})^2 g_{33}^{-1} \left(\frac{\partial g_{11}}{\partial z} \right)^2
 \end{aligned}$$

Non-linear theory

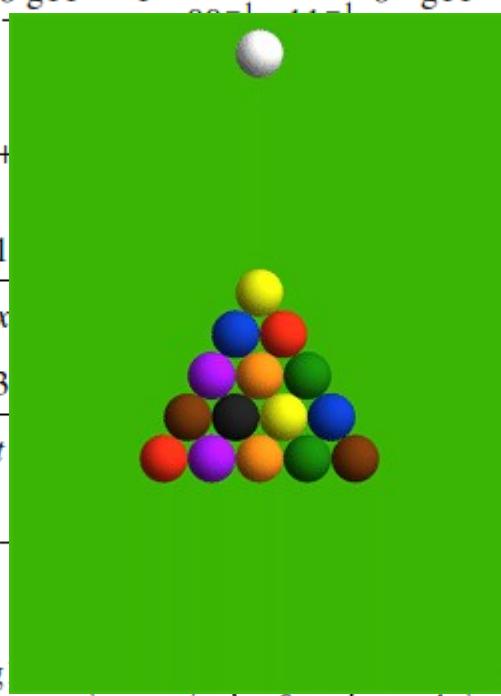
$$G_{\mu\nu} = 8\pi G_N T_{\mu\nu}$$

Simple model



$\mu = \nu = 3$

$$\begin{aligned}
 & \frac{1}{4} g_{00}^{-1} g_{11}^{-1} g_{22}^{-1} \frac{\partial g_{00}}{\partial y} \frac{\partial g_{11}}{\partial y} + \frac{1}{4} g_{00}^{-1} g_{11}^{-1} g_{33}^{-1} \frac{\partial g_{11}}{\partial t} \frac{\partial g_{33}}{\partial t} + \frac{1}{4} g_{00}^{-1} g_{11}^{-1} g_{33}^{-1} \frac{\partial g_{00}}{\partial x} \frac{\partial g_{33}}{\partial x} + \\
 & \frac{1}{4} g_{00}^{-1} g_{11}^{-1} g_{33}^{-1} \frac{\partial g_{00}}{\partial z} \frac{\partial g_{11}}{\partial z} + \frac{1}{4} g_{00}^{-1} \left(g_{11}^{-1} \right)^2 \left(\frac{\partial g_{11}}{\partial t} \right)^2 + \frac{1}{4} \left(g_{00}^{-1} \right)^2 g_{11}^{-1} \frac{\partial g_{00}}{\partial t} \frac{\partial g_{11}}{\partial t} - \\
 & \frac{1}{4} g_{00}^{-1} \left(g_{11}^{-1} \right)^2 \frac{\partial g_{00}}{\partial x} \frac{\partial g_{11}}{\partial x} + \frac{1}{4} \left(g_{00}^{-1} \right)^2 g_{33}^{-1} \frac{\partial g_{00}}{\partial t} \frac{\partial g_{33}}{\partial t} + \frac{1}{2} g_{00}^{-1} g_{33}^{-1} \frac{\partial^2 g_{33}}{\partial t^2} - \\
 & \frac{1}{4} \left(g_{00}^{-1} \right)^2 g_{33}^{-1} \frac{\partial g_{00}}{\partial z} \frac{\partial g_{33}}{\partial z} + \frac{1}{2} g_{00}^{-1} g_{33}^{-1} \frac{\partial^2 g_{00}}{\partial z^2} - \\
 & \frac{1}{4} \left(g_{00}^{-1} \right)^2 g_{33}^{-1} \left(\frac{\partial g_{00}}{\partial z} \right)^2 + \frac{1}{4} g_{11}^{-1} g_{22}^{-1} g_{33}^{-1} \frac{\partial g_{11}}{\partial y} \frac{\partial g_{33}}{\partial y} + \\
 & \frac{1}{2} g_{11}^{-1} g_{33}^{-1} \frac{\partial^2 g_{33}}{\partial x^2} - \frac{1}{4} g_{11}^{-1} g_{33}^{-1} \frac{\partial^2 g_{11}}{\partial x^2} + \frac{1}{4} g_{11}^{-1} g_{33}^{-1} \frac{\partial^2 g_{11}}{\partial z^2} - \frac{1}{4} g_{11}^{-1} \left(g_{33}^{-1} \right)^2 \frac{\partial g_{11}}{\partial z} \frac{\partial g_{33}}{\partial z} - \frac{1}{4} \left(g_{11}^{-1} \right)^2 g_{33}^{-1} \left(\frac{\partial g_{11}}{\partial z} \right)^2
 \end{aligned}$$



Describes interactions!

Simple model

$$G_{\mu\nu} = 8\pi G_N T_{\mu\nu}$$

Cosmological principle: Isotropy and homogeneity

$$ds^2 = -dt^2 + a^2(t)(dr^2 + f(r)^2 d\Omega_2^2)$$

FRWL

(Perfect) fluids $\rho = \bar{\rho}$ $p = \bar{p} = \omega \bar{\rho}$

Describes well the averaged Universe

Fractional components

$$H = H_0 \left(\frac{\Omega_b^0}{a^3} + \frac{\Omega_{dm}^0}{a^3} + \frac{\Omega_r^0}{a^4} + \frac{\Omega_k^0}{a^2} + \Omega_\Lambda^0 \right)$$

Beyond average



Universe is not fully homogeneous and isotropic

Trick: Use perturbation theory

$$ds^2 = ds_{FRWL}^2 + ds_{pert}^2$$

$$\rho = \bar{\rho}(1 + \delta) \quad \dots$$

Beyond average

$$\rho = \bar{\rho}(1 + \delta)$$

$$\bar{\rho} \sim 10^{-26} \text{Kg/m}^3$$

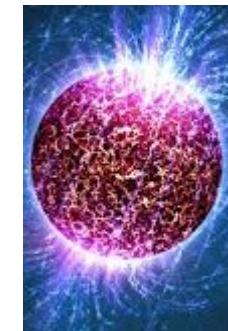
Note that



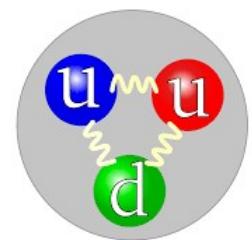
$$\delta \sim 100$$



$$\delta \sim 10^{29}$$

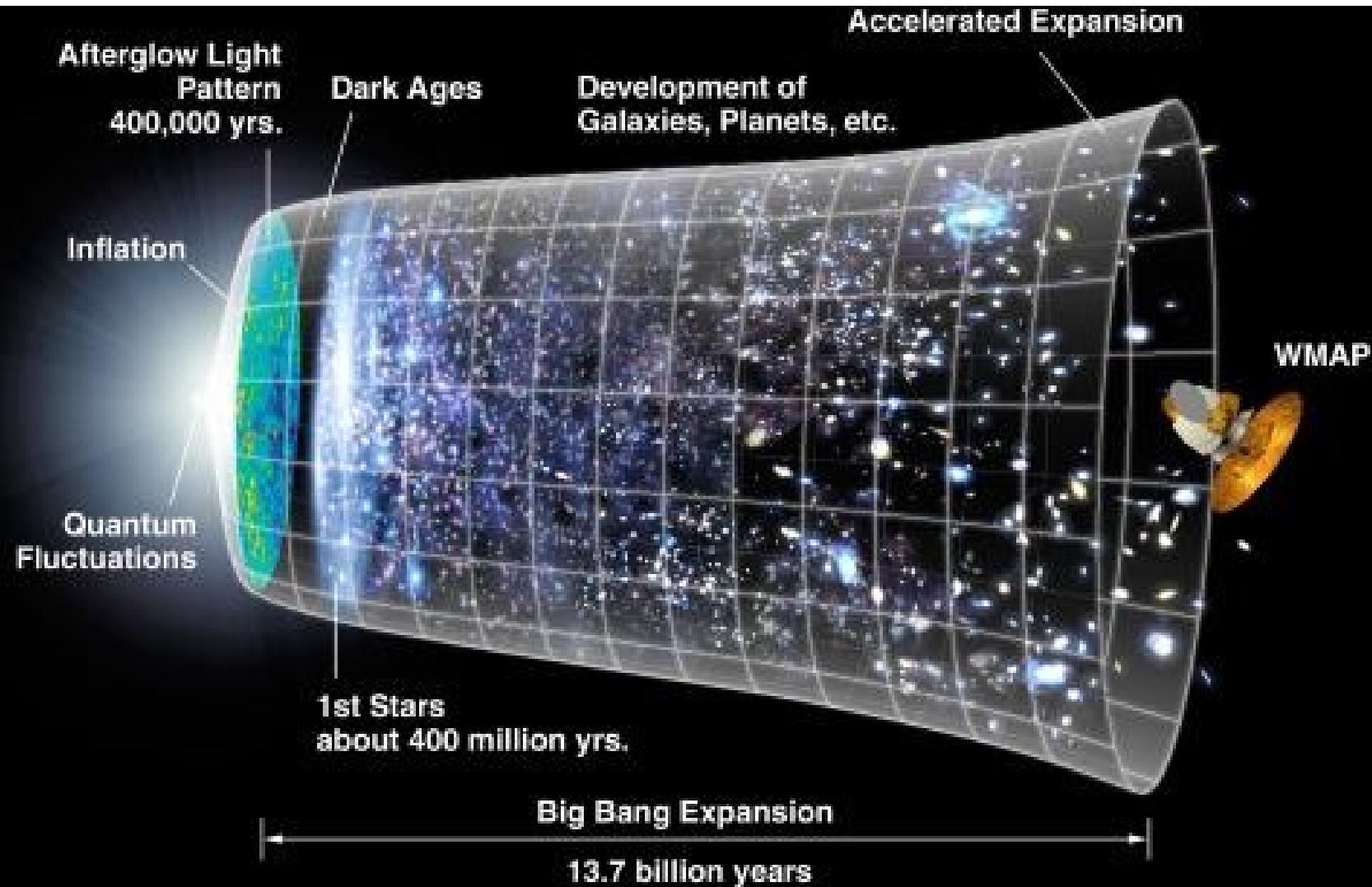


$$\delta \sim 10^{43}$$

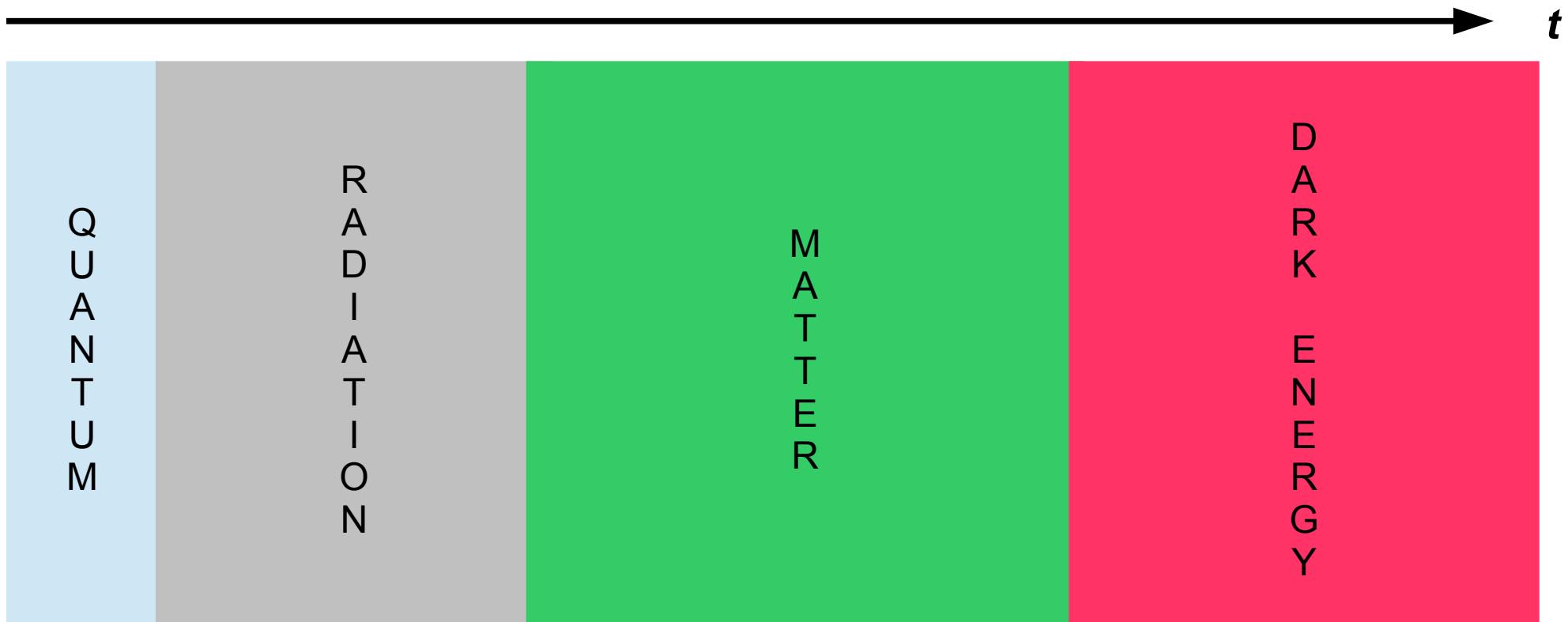


Could it really describe our Universe?

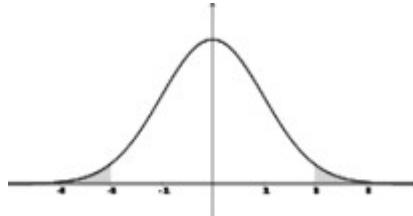
Basic idea



The story of “pertuby”



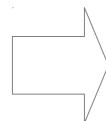
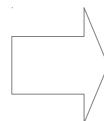
The story of “pertuby”



“Classical” pert.



Interactions!



t

Q
U
A
N
T
U
M

R
A
D
I
A
T
I
O
N

M
A
T
T
E
R

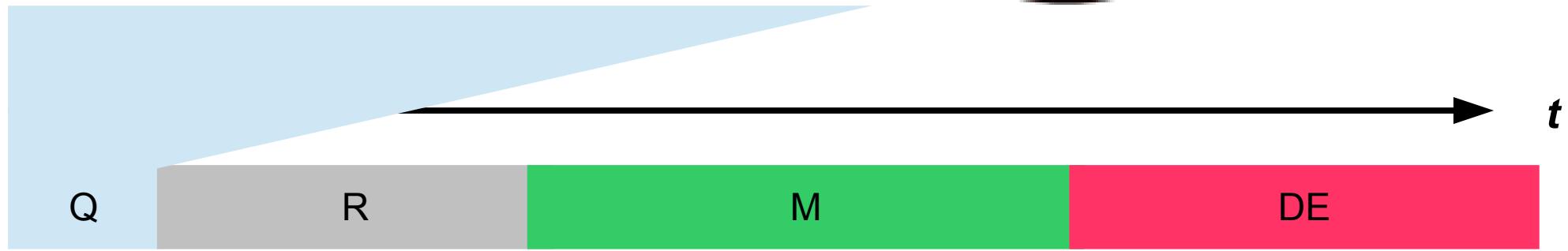
D
A
R
K
E
N
E
R
G
Y

The story of “pertuby”

Formally (the inflationary mechanism)

$$t \sim 10^{-32}$$

$$10^{50} \rightarrow$$



The story of “pertuby”

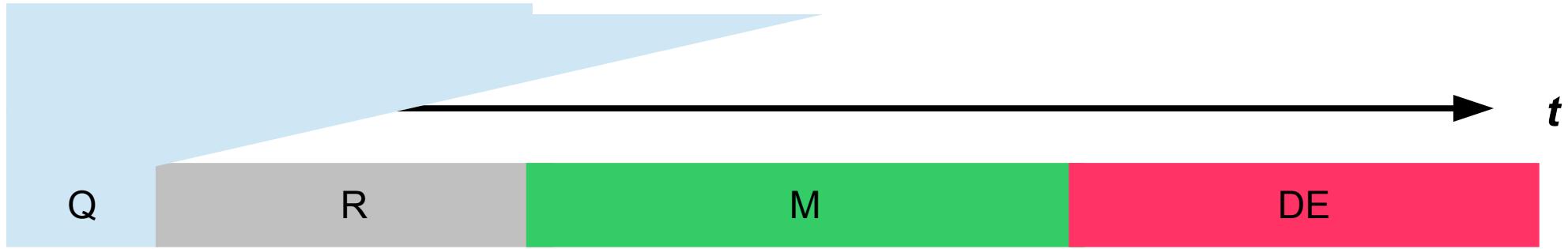
Formally (EFT perspective)

A single “clock” stopping de Sitter (accelerated) phase
that spontaneously breaks time symmetry

$$\pi \sim \delta t \sim \frac{\delta\phi}{\dot{\phi}}$$

Goldstone mode related to metric perturbations

$$\zeta \sim \frac{\delta a}{a} \sim H\pi$$

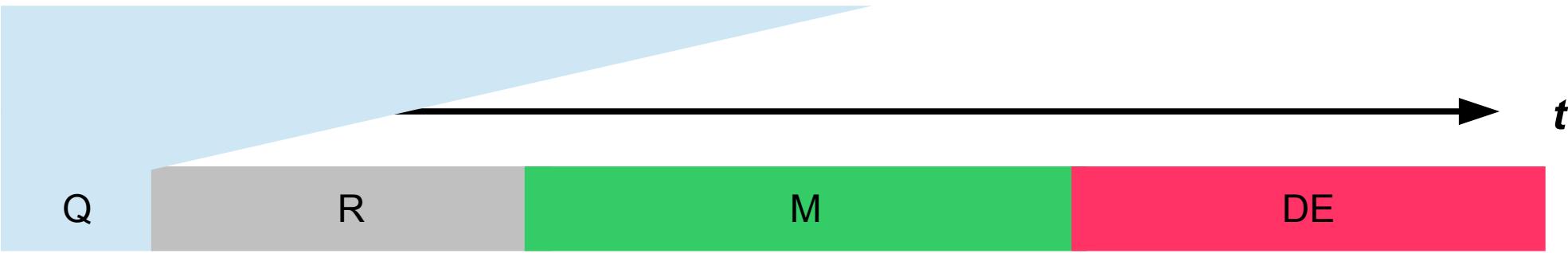


The story of “pertuby”

Most general Lagrangian which breaks time symmetry

$$S = \int d^4x \sqrt{-g} \left[\frac{1}{2}M_{\text{Pl}}^2 R - M_{\text{Pl}}^2 \left(3H^2(t + \pi) + \dot{H}(t + \pi) \right) + \right.$$

Cheung et al, ...

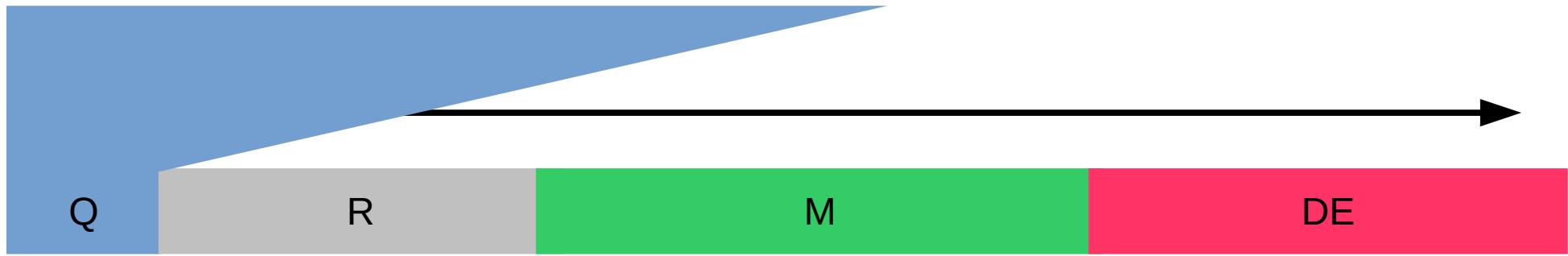
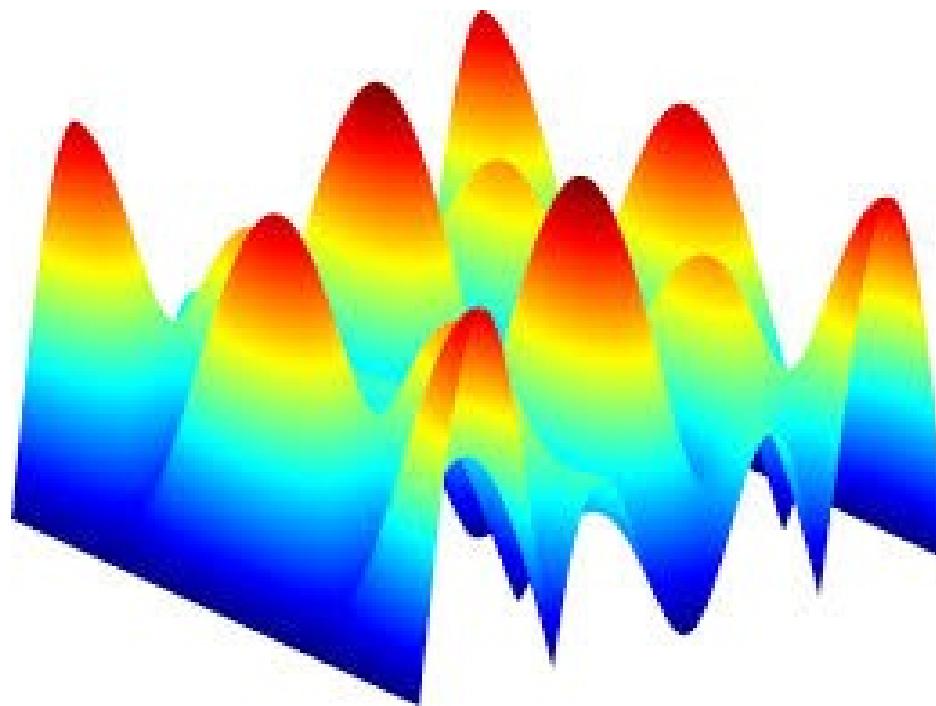
$$\begin{aligned} &+ M_{\text{Pl}}^2 \dot{H}(t + \pi) \left((1 + \dot{\pi})^2 g^{00} + 2(1 + \dot{\pi}) \partial_i \pi g^{0i} + g^{ij} \partial_i \pi \partial_j \pi \right) + \\ &\frac{M_2(t + \pi)^4}{2!} \left((1 + \dot{\pi})^2 g^{00} + 2(1 + \dot{\pi}) \partial_i \pi g^{0i} + g^{ij} \partial_i \pi \partial_j \pi + 1 \right)^2 + \\ &\left. \frac{M_3(t + \pi)^4}{3!} \left((1 + \dot{\pi})^2 g^{00} + 2(1 + \dot{\pi}) \partial_i \pi g^{0i} + g^{ij} \partial_i \pi \partial_j \pi + 1 \right)^3 + \dots \right] \end{aligned}$$


- *Unifies most single field models
- *Avoids questions about inflation

Tensors &
Higher der.

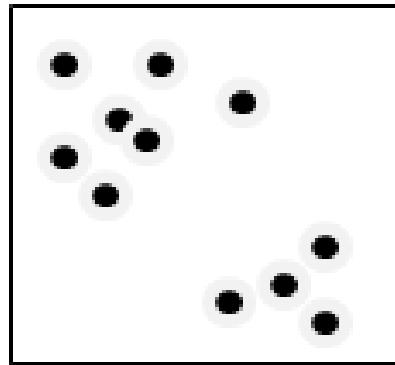
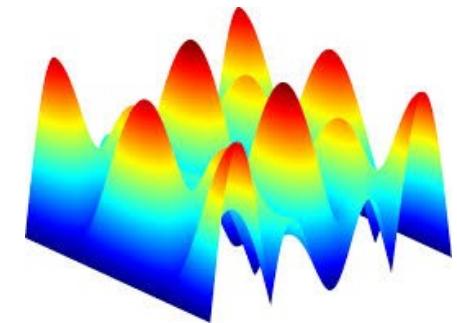
The story of “pertuby”

By the end of inflation we're left with a distribution of perturbations

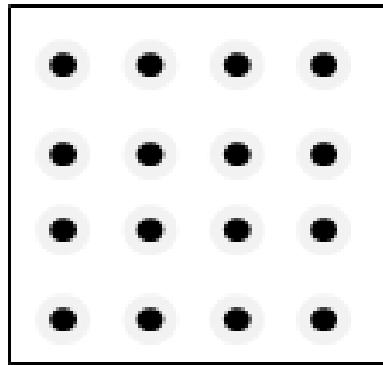


Detour

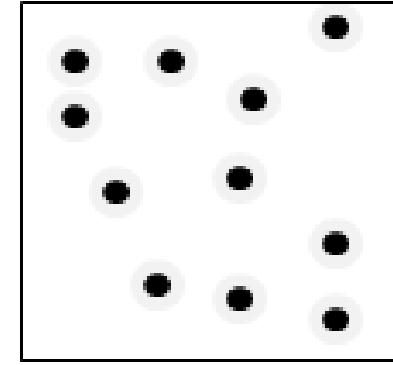
Characterising distributions



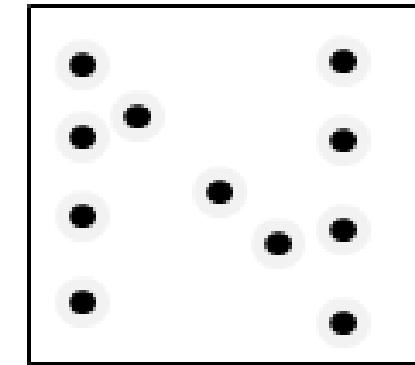
clustered



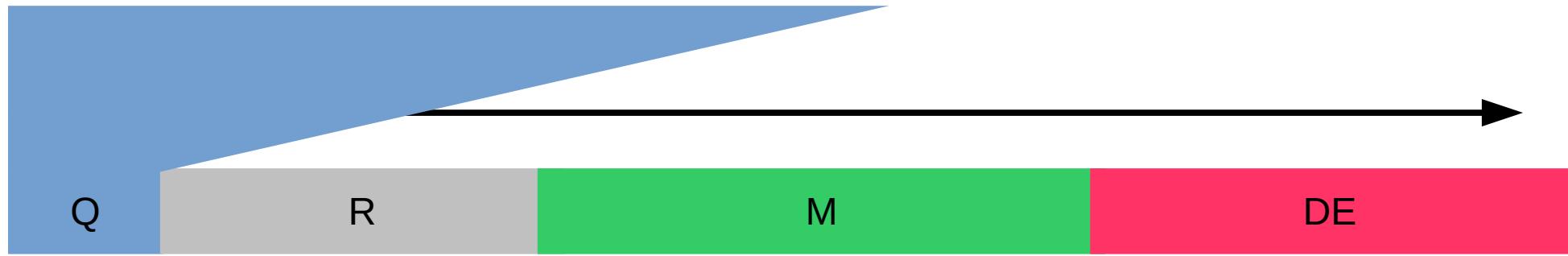
ordered



random



pattered

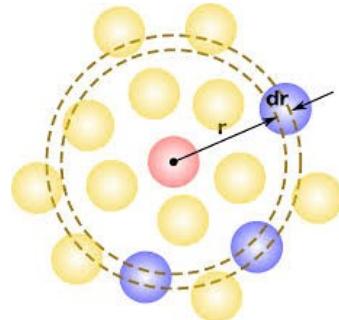


Detour

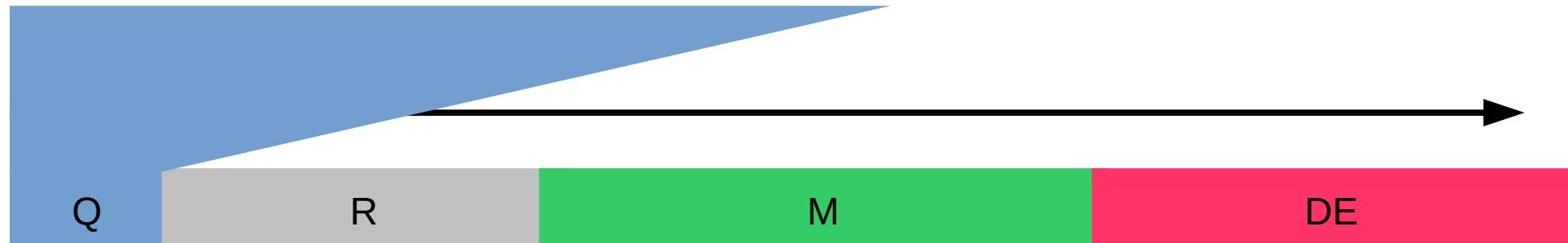
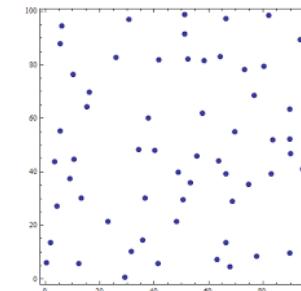
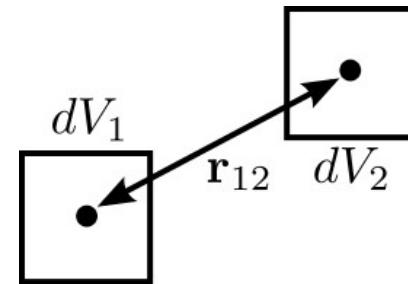
Usual to consider N-point correlation functions

- **N=2 (2pcf)**

$$dP = \bar{n}^2(1 + \xi^{(2)}(\mathbf{r}_{12}))dV_1 dV_2$$

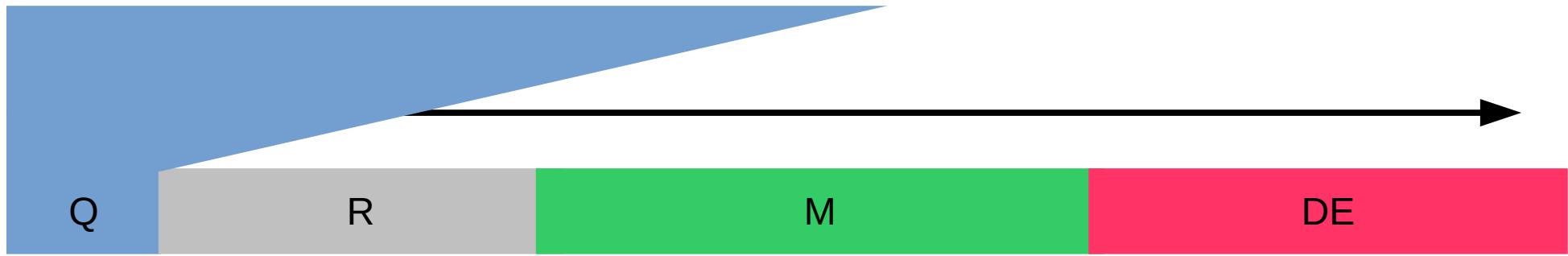
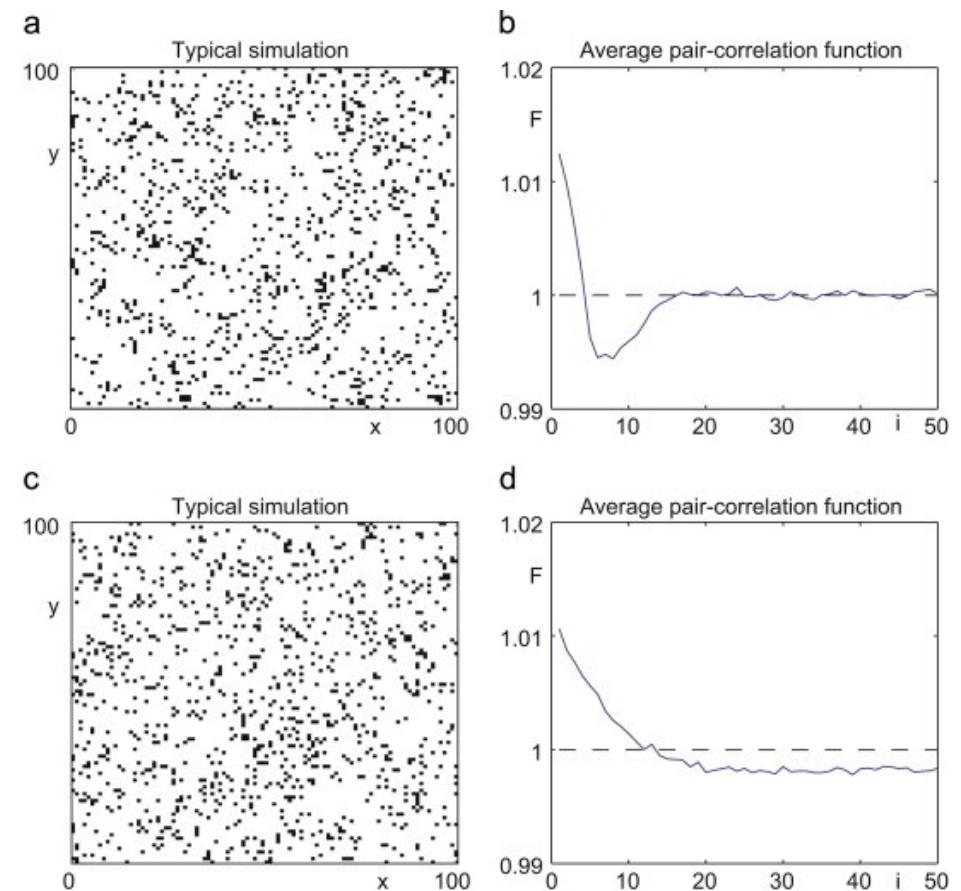
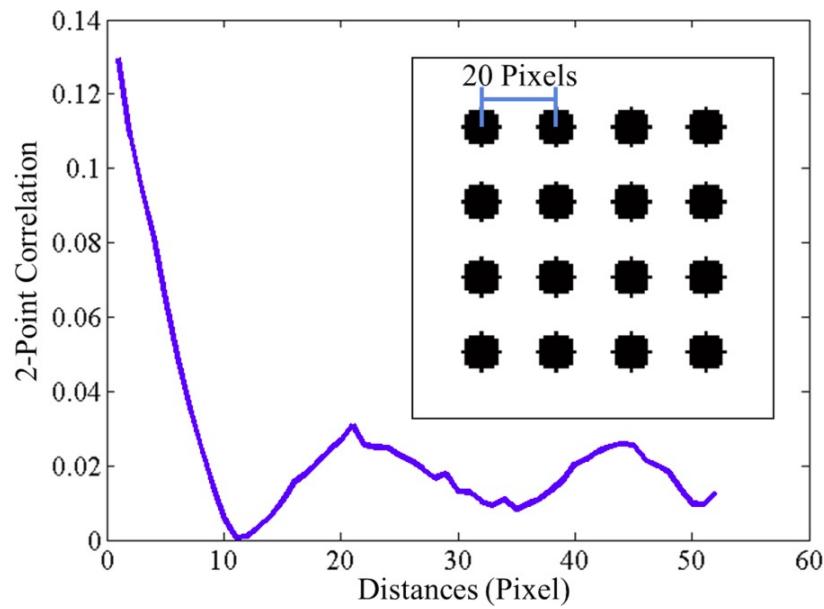


Excess correlation
over the random pairs



Detour

Examples



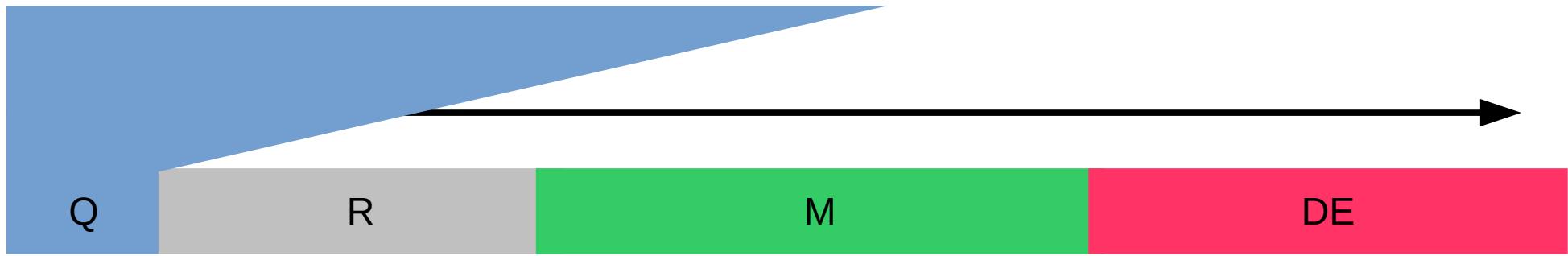
Detour

Real space $\xleftarrow{e^{ik \cdot x}}$ Fourier space

$$\xi^{(2)}(r) = \langle \delta(r)\delta(r') \rangle$$

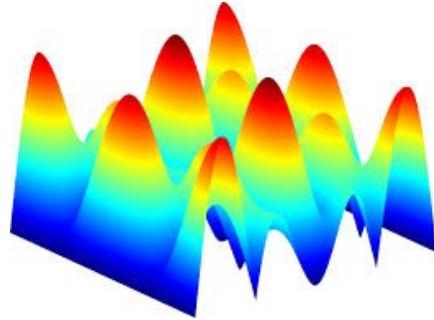
$$\langle \delta(k)\delta(k') \rangle = \delta(k - k')P(k)$$

Power Spectrum



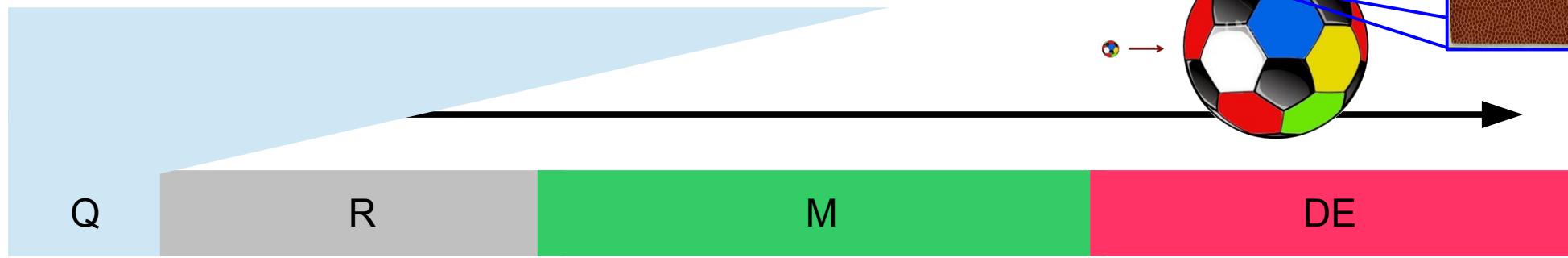
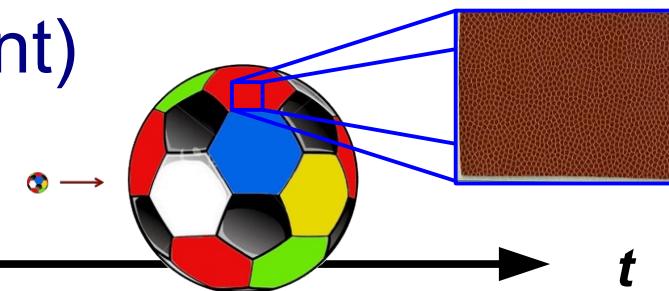
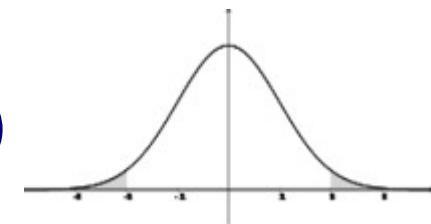
The story of “pertub”

Returning to inflation, we're left with a distribution

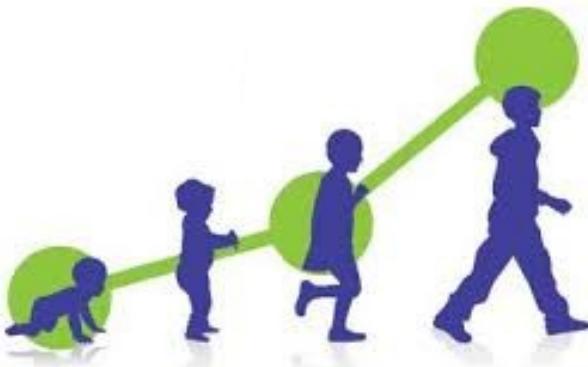


$$\xi \sim \leftarrow \begin{array}{c} \text{two images of a green, wavy, textured object} \\ P(k)\delta(k - k') \sim <\hat{\delta}(k)\hat{\delta}(k')> \end{array} \rightarrow$$

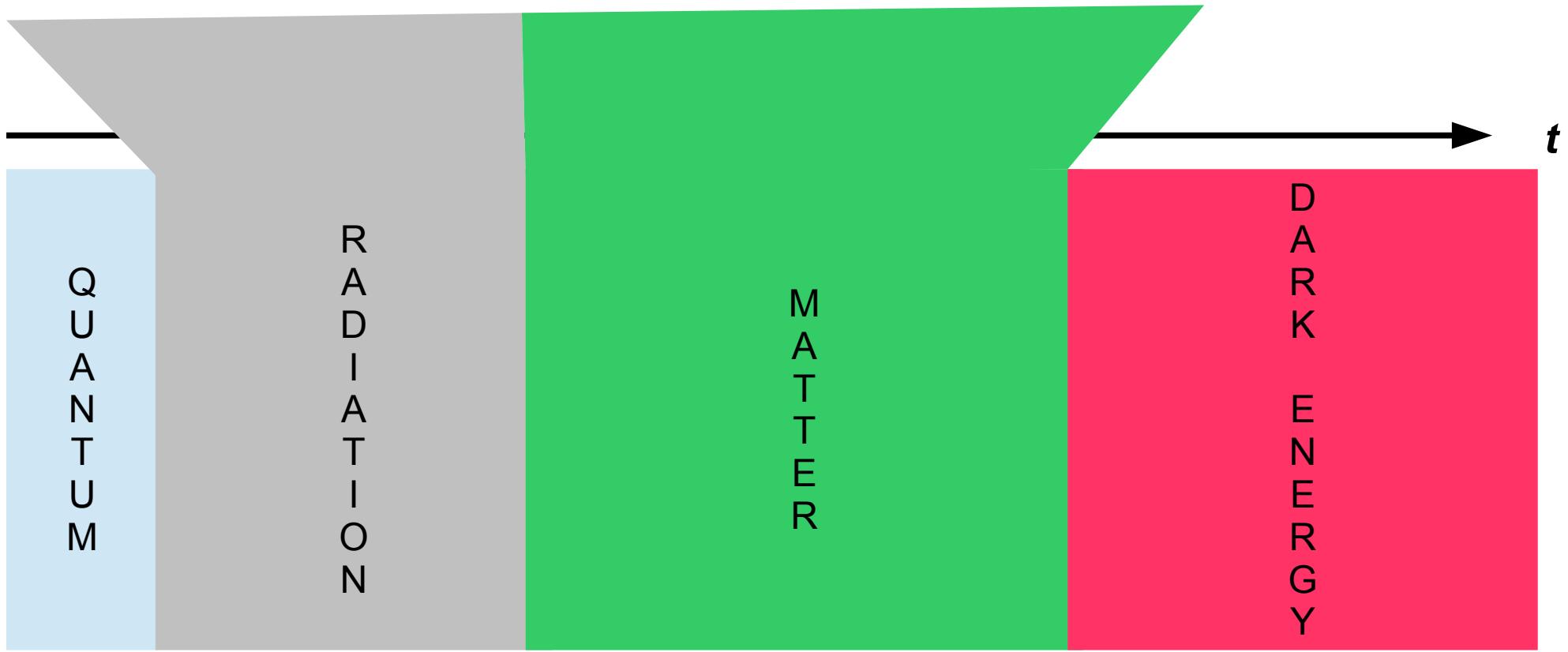
- *Almost scale invariant
- *Very Gaussian (2pt correlations are enough)
- *Input amplitude of (0.00001).
- *Gravitational waves (model dependent)



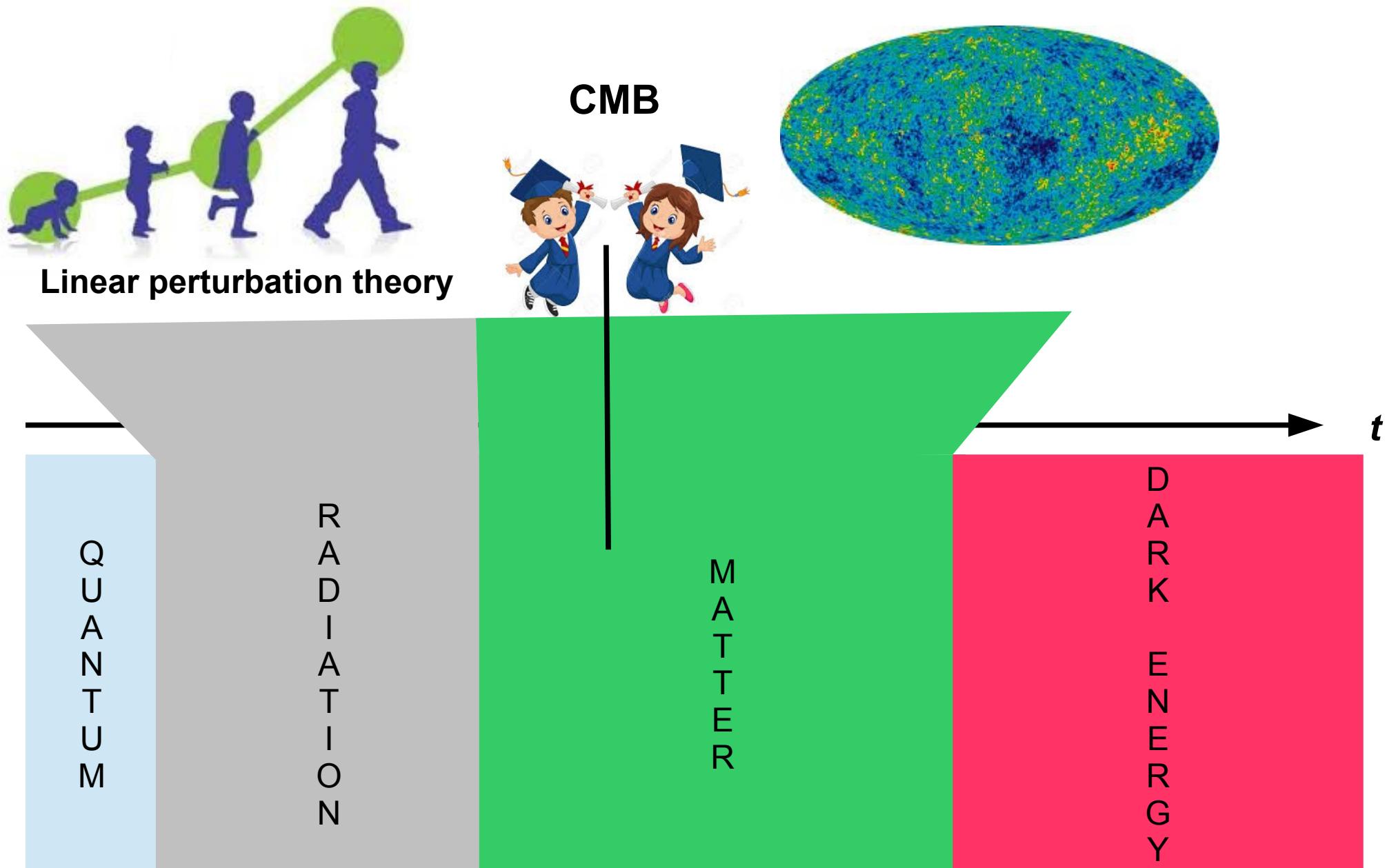
The story of “pertub”



Linear perturbation theory

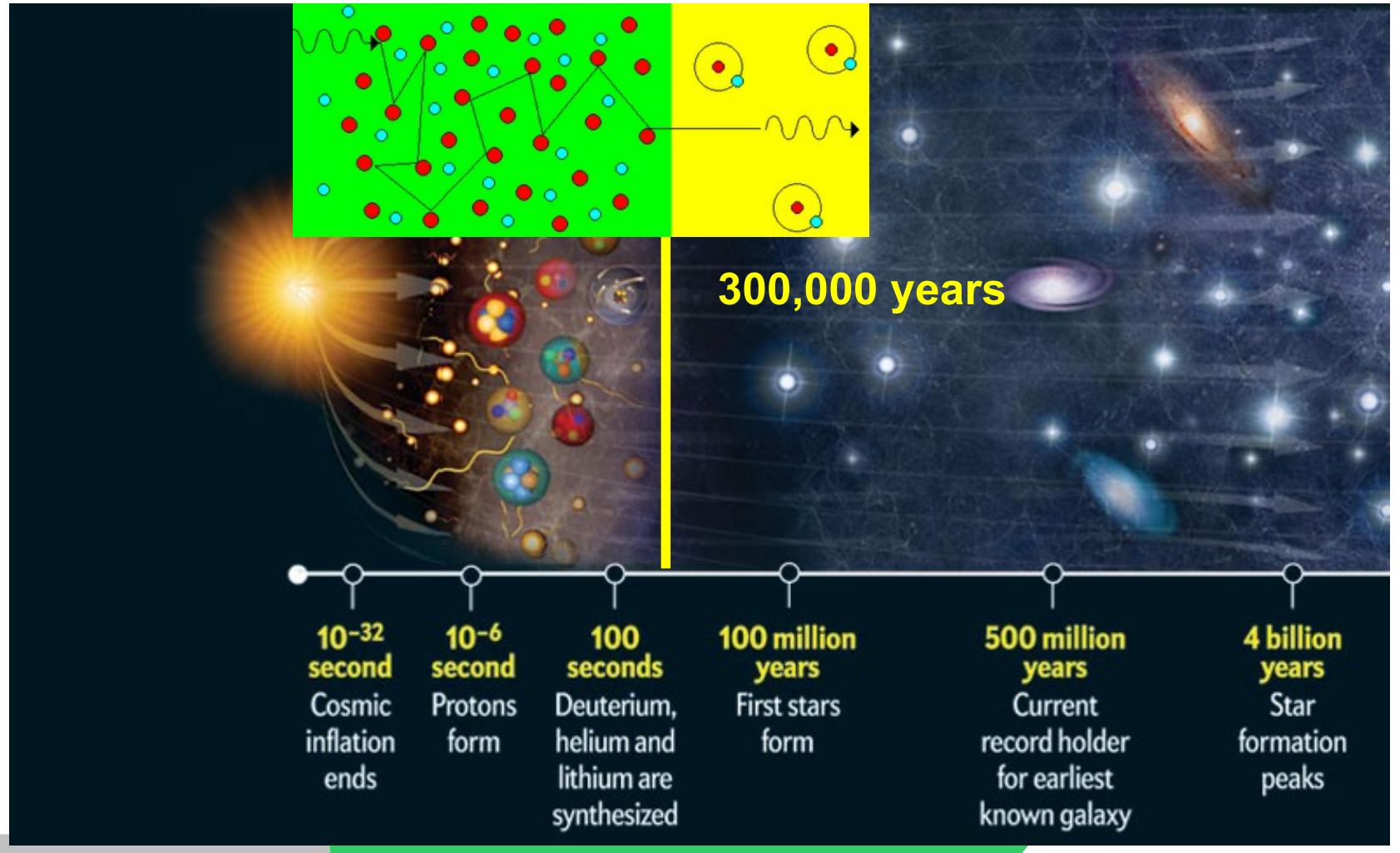


The story of “pertub”



The story of “pertub”

CMB



Q

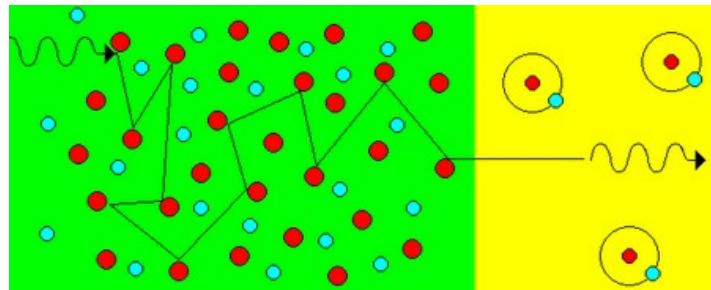
R

M

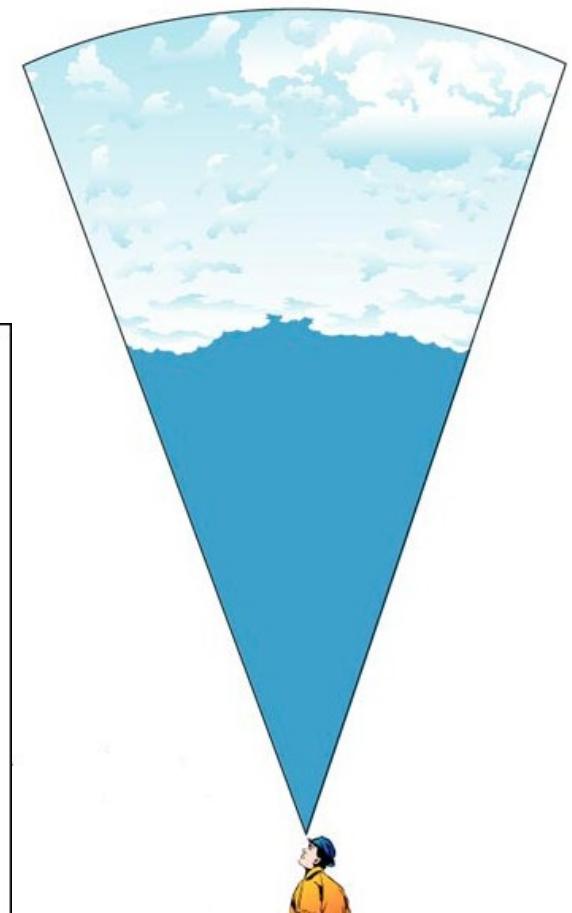
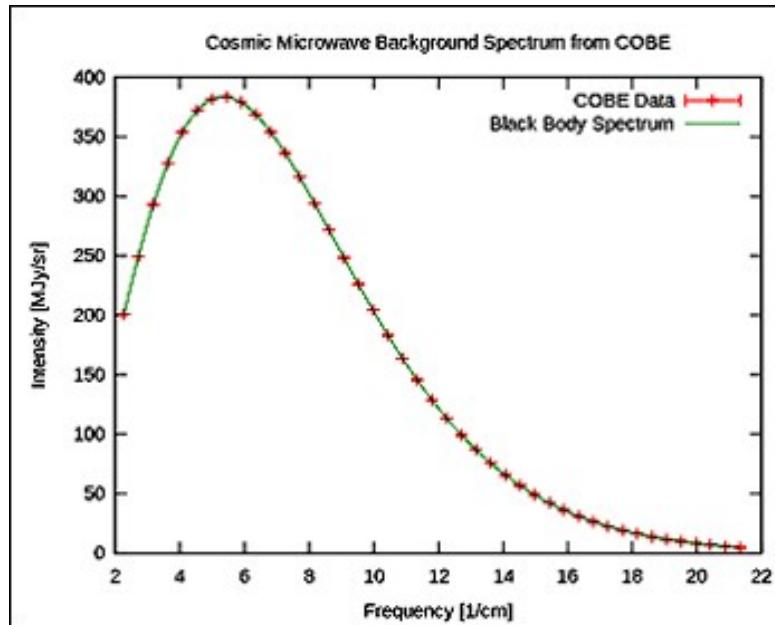
DE

t

The story of “pertub”



**CMB
(First observable)**



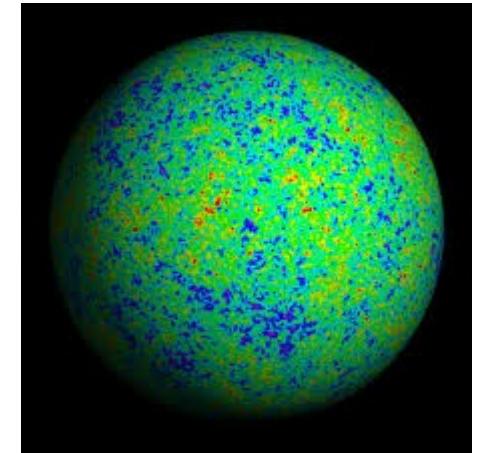
The story of “pertuby”

CMB Anisotropies

Natural to use spherical harmonics

$$\frac{\Delta T}{\bar{T}} = \sum_{lm} T_{lm} Y_{lm}$$

Power spectrum

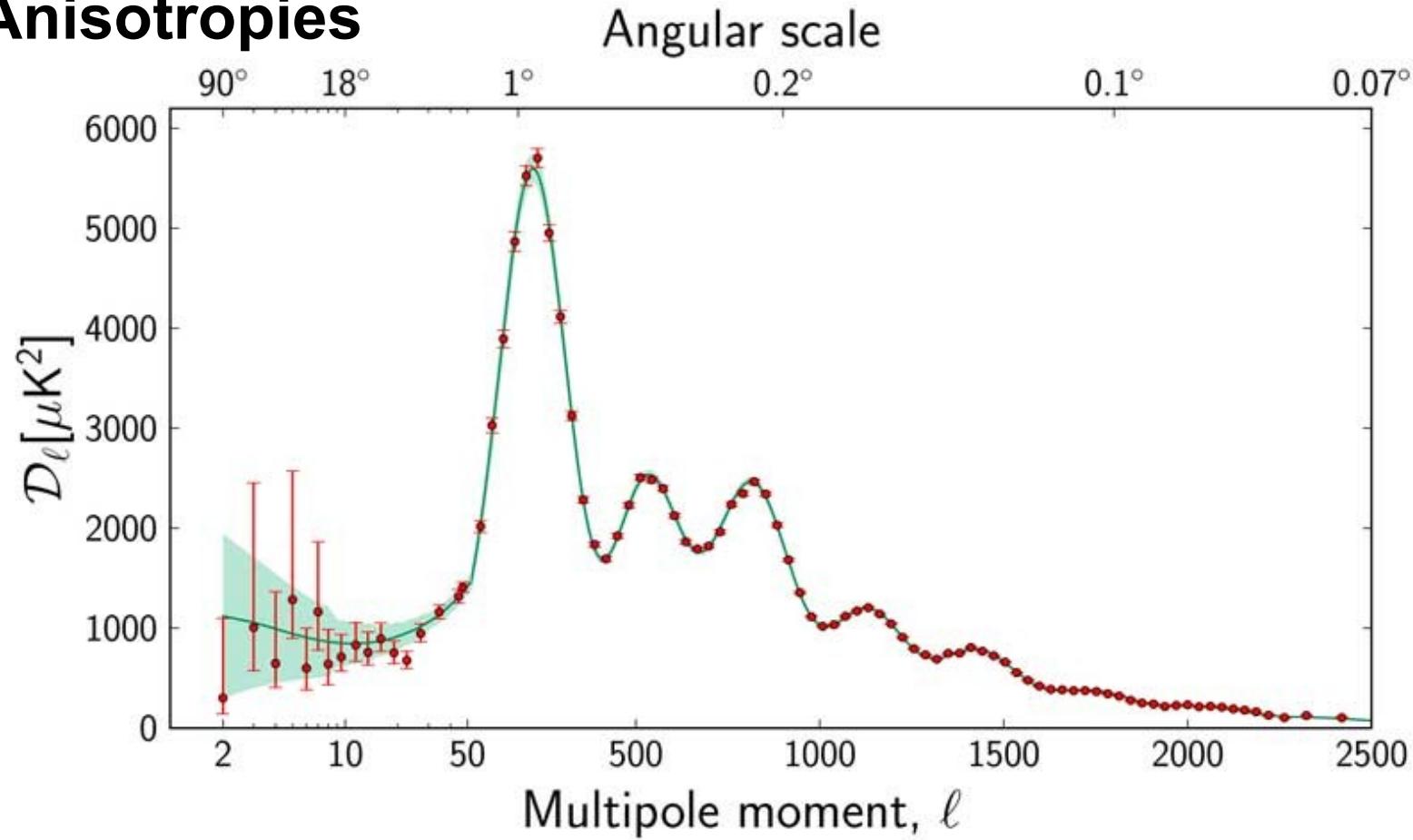


$$\langle T_{lm} T_{l'm'} \rangle = 2\pi \mathcal{D}_l \delta_{ll'} \delta_{mm'}$$



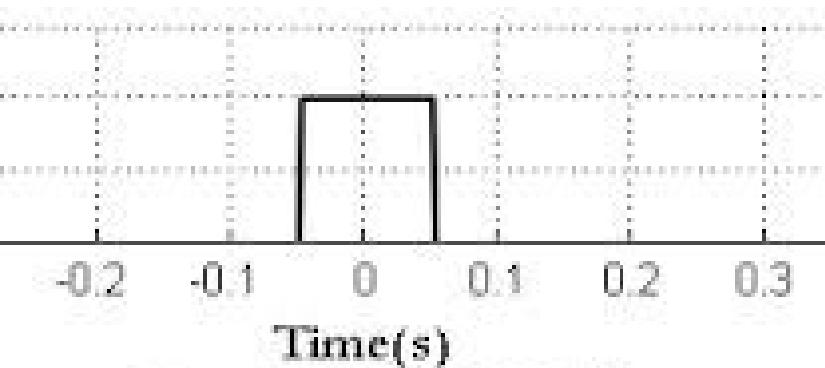
The story of “pertuby”

CMB Anisotropies

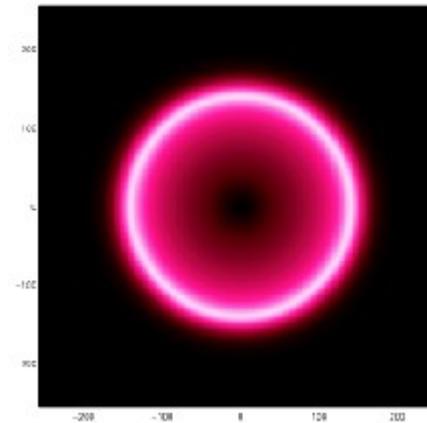
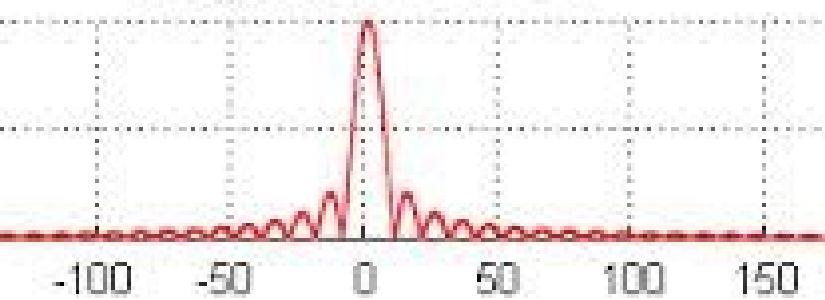


The story of “pertuby”

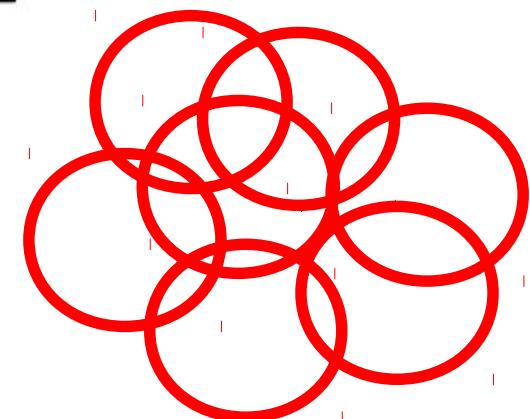
CMB Oscilations



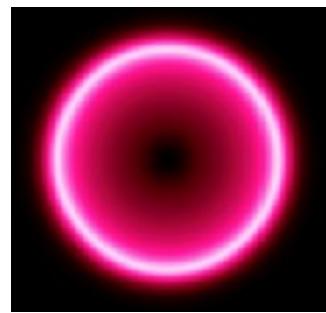
Magnitude of FFT



Remnant of photon's pressure

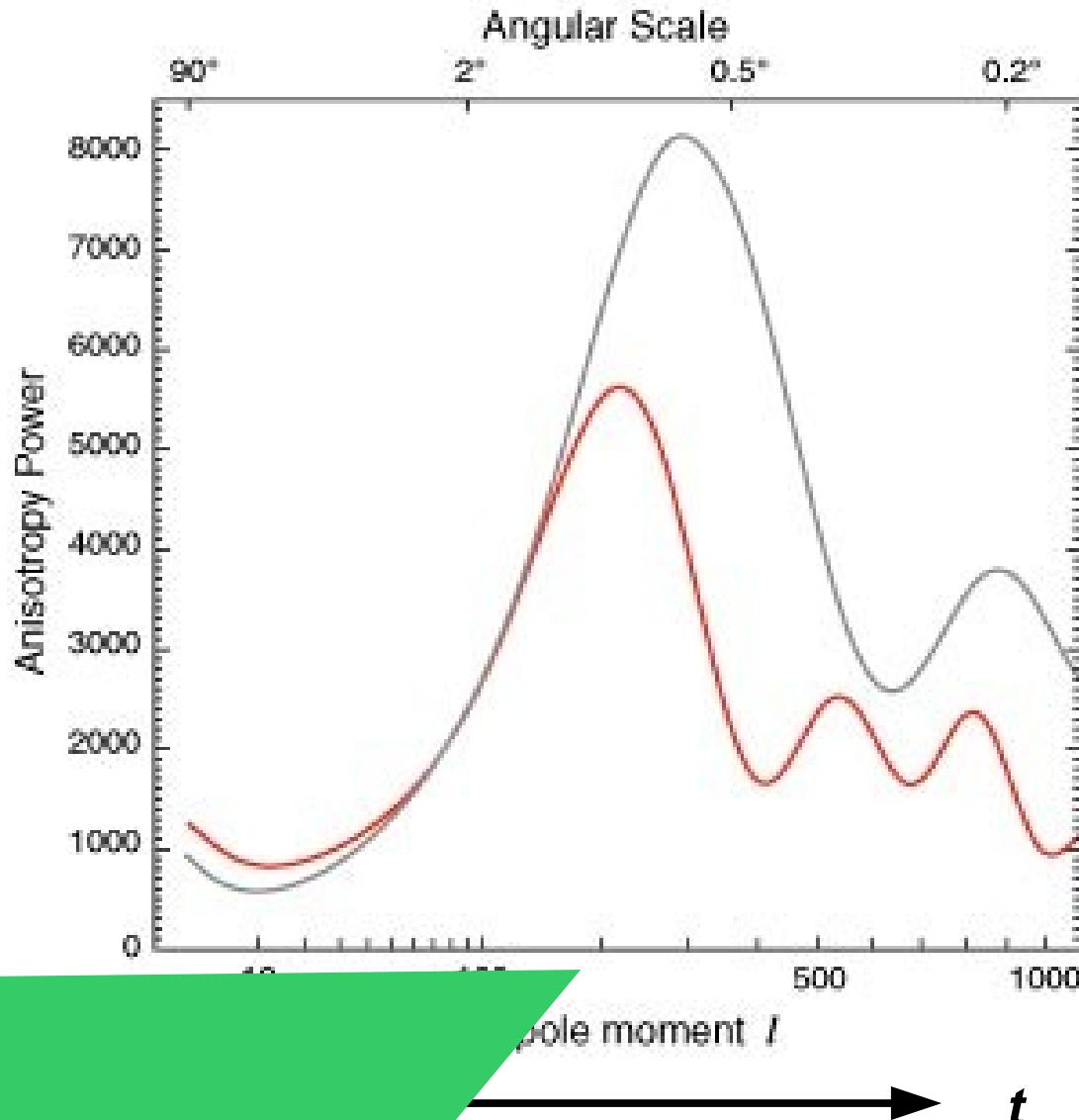
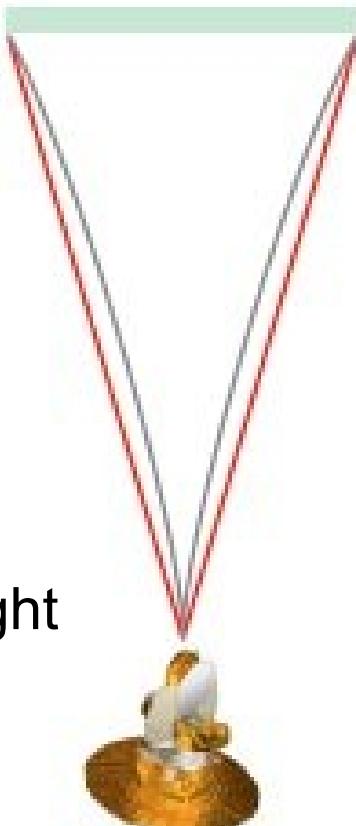


The story of “pertuby”



Standard Ruler:
1° arc measurement of
dominant energy spike

Planck
To 0.1% the flat
Universe model is right



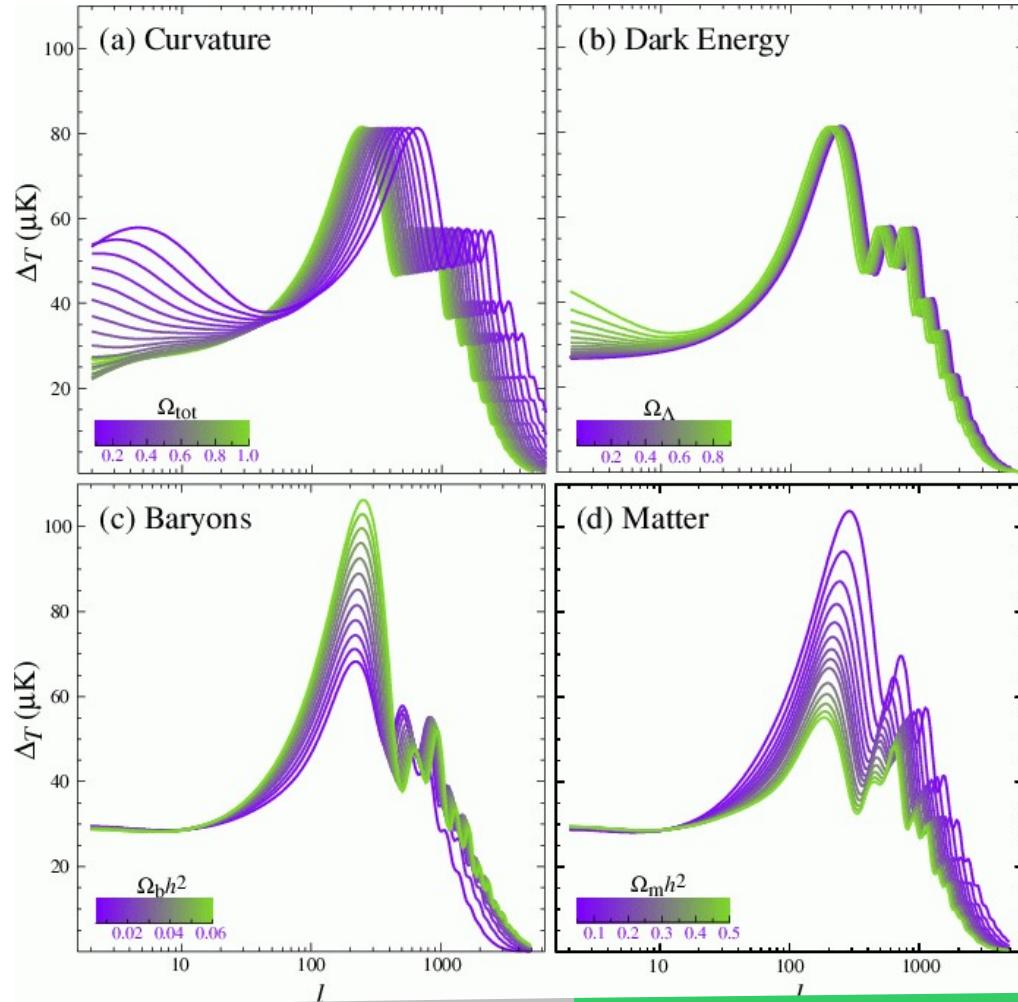
Q

R

M

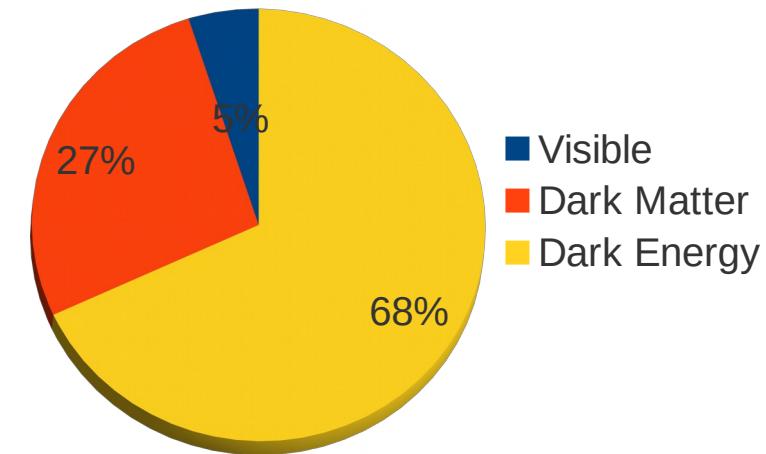
DE

The story of “pertub”



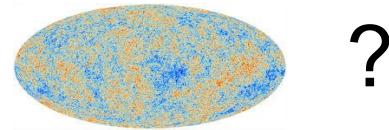
Planck satellite

Flat Universe to 0.1%



The story of “pertuby”

Is there more information in

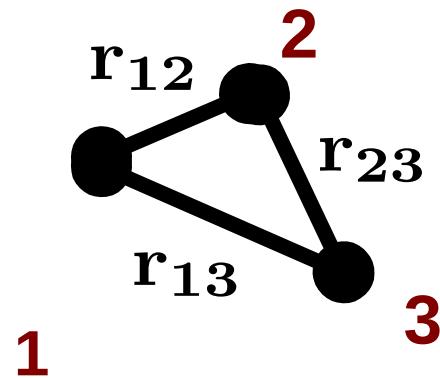


?

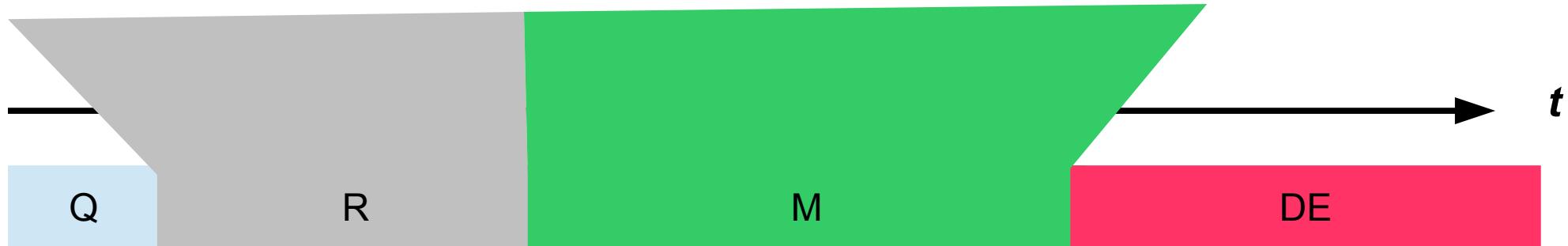
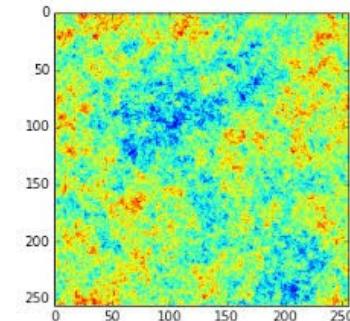
Higher N-point correlation functions

- $N=3$ (3pcf)

$$dP = \bar{n}^2(1 + \xi^{(3)}(\mathbf{r}_{12}, \mathbf{r}_{13}, \mathbf{r}_{23}))dV_1 dV_2 dV_3$$



Excess correlation
over the random triples



The story of “pertuby”

- Isotropy

$$\xi^{(3)}(\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3) = \xi^{(3)}(r_1, r_2, r_3)$$

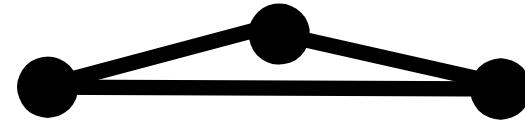
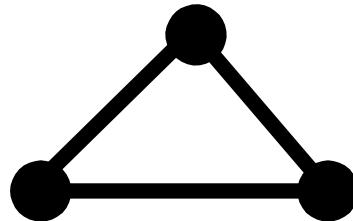
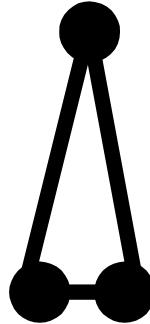
– 3 parameter function



4d plot!

Options

A) Consider particular shapes

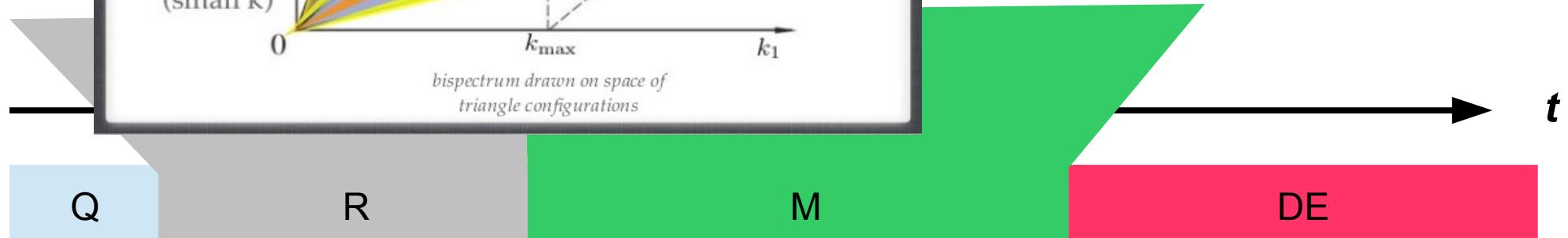
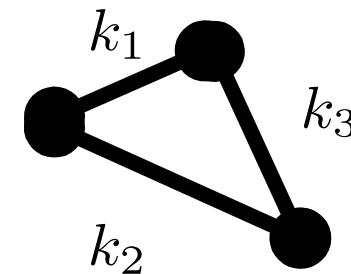
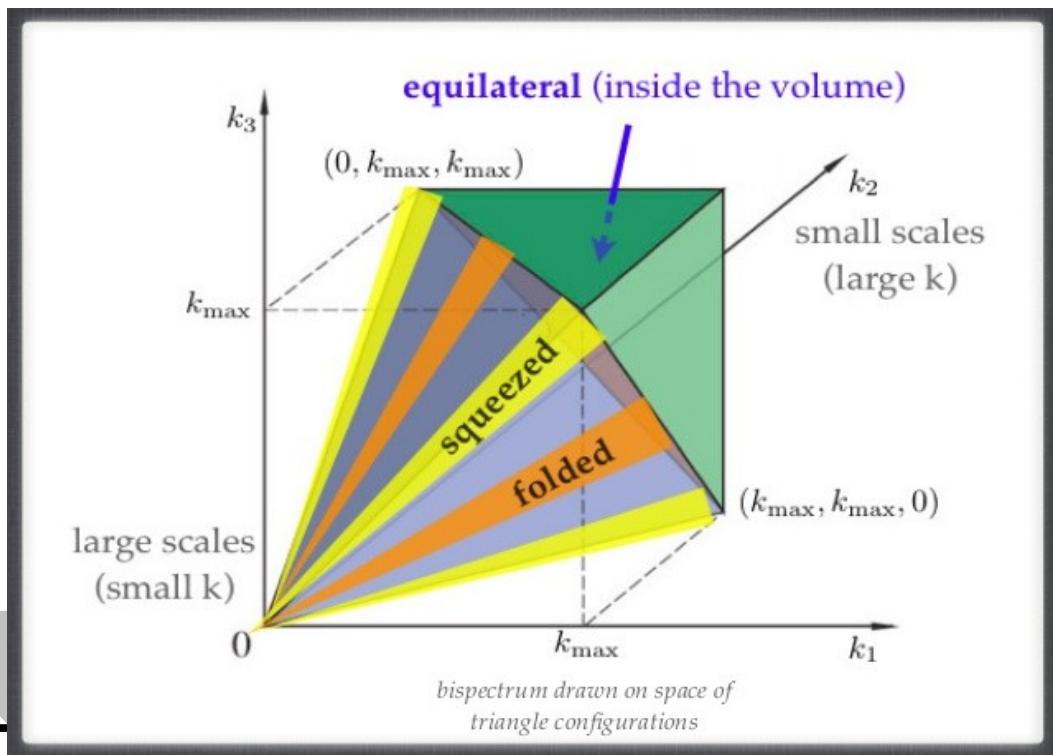


The story of “pertuby”

Real space $\xi^{(3)} = \langle \delta\delta\delta \rangle$

Fourier $\langle \delta(k_1)\delta(k_2)\delta(k_3) \rangle = \delta(\sum_i k_i)B(k_1, k_2, k_3)$

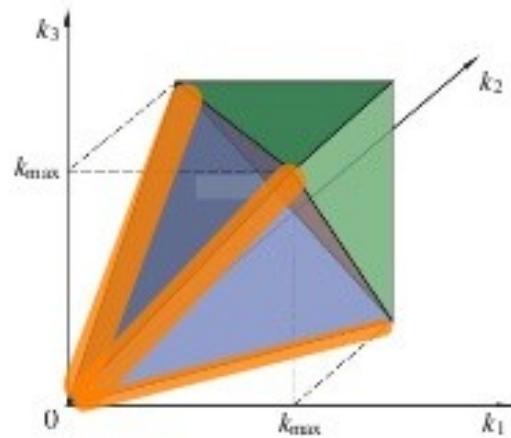
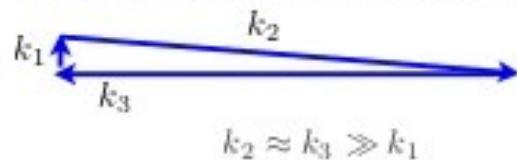
Bi-spectrum



BISPECTRUM SHAPES

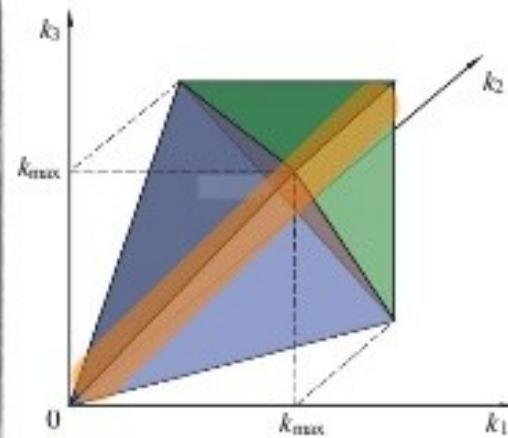
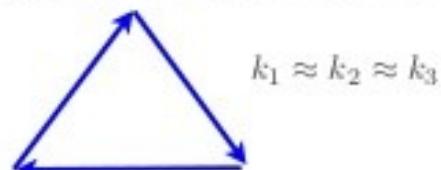
Different inflation models induce different momentum dependencies (shapes) of $B_\Phi(k_1, k_2, k_3)$

Squeezed triangles (local shape)



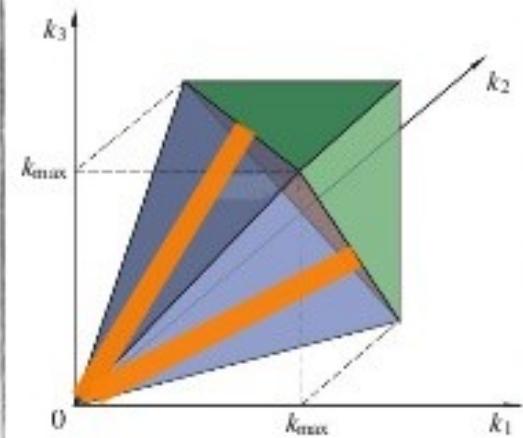
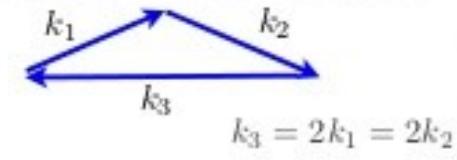
Arises in multifield inflation;
detection would rule out all
single field models!

Equilateral triangles



Typically higher derivative
kinetic terms, e.g. DBI inflation

Folded triangles

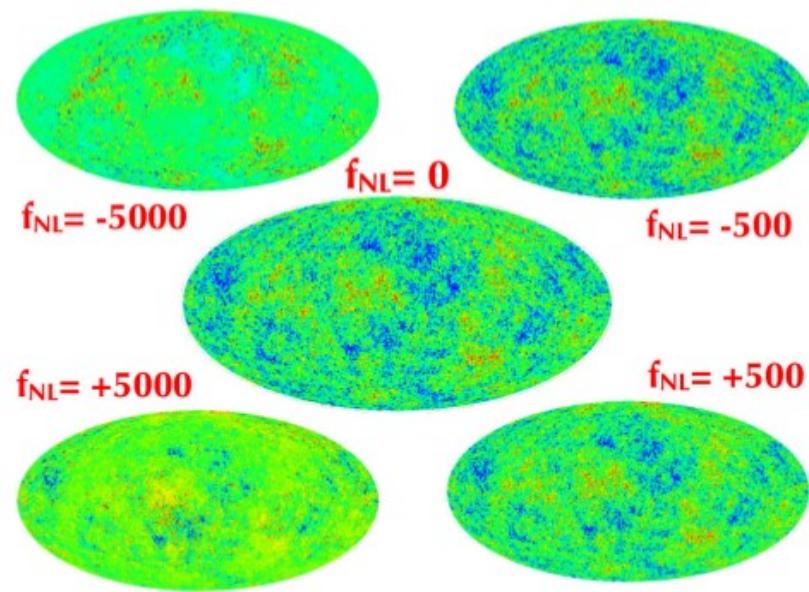


E.g. non-Bunch-Davies vacuum

The story of “pertub”

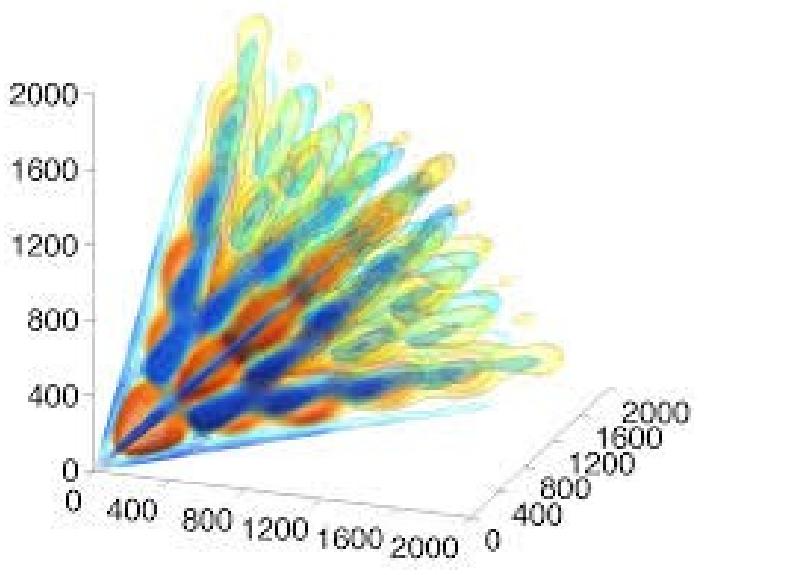
$$\phi = \phi_G + \phi_{NG}$$

Inflaton Gaussian Part Non-Gaussian Part



Simplest case: non-local NG

$$\phi = \phi_G + f_{NL}(\phi_G^2 - \langle \phi_G \rangle^2)$$



The story of “pertuby”

CMB

- Consistent with EFT of inflation (Gaussian & almost scale invariant distribution)
- Not much about interactions

$$\langle \delta^n \rangle$$

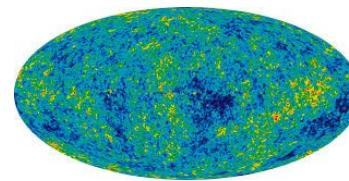
consistent with zero for $n>2$

- Cannot gain more mode statistics to achieve better accuracy than with Planck

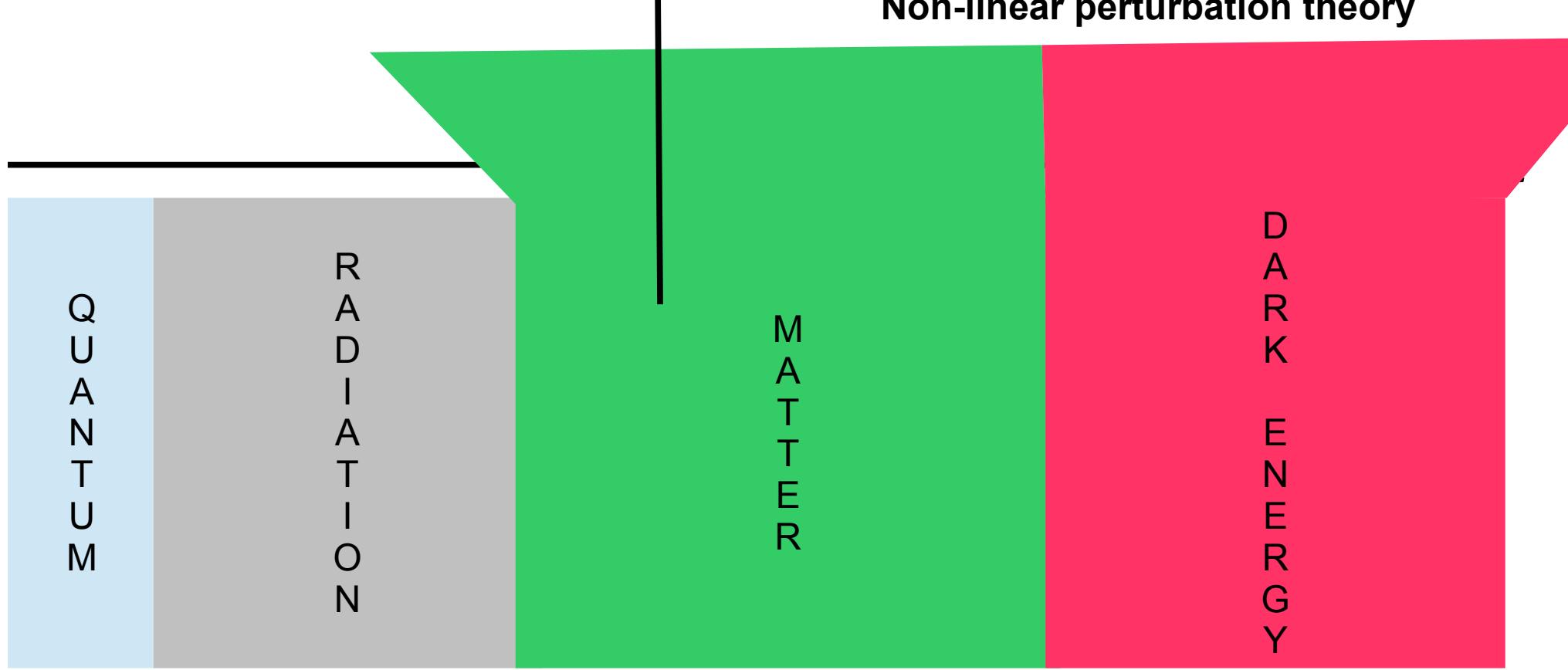
Future: explore the LSS



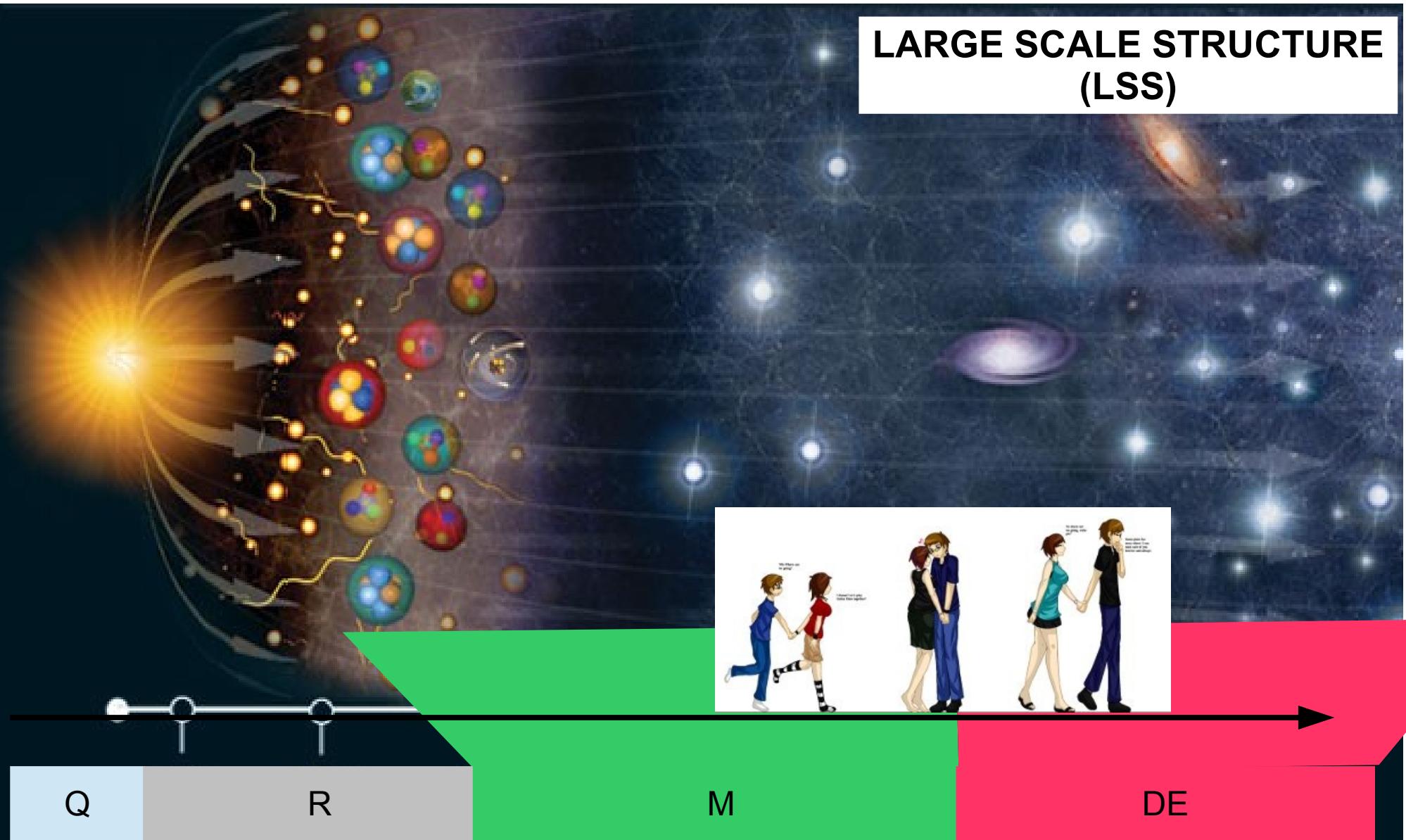
The story of “pertuby”



Non-linear perturbation theory

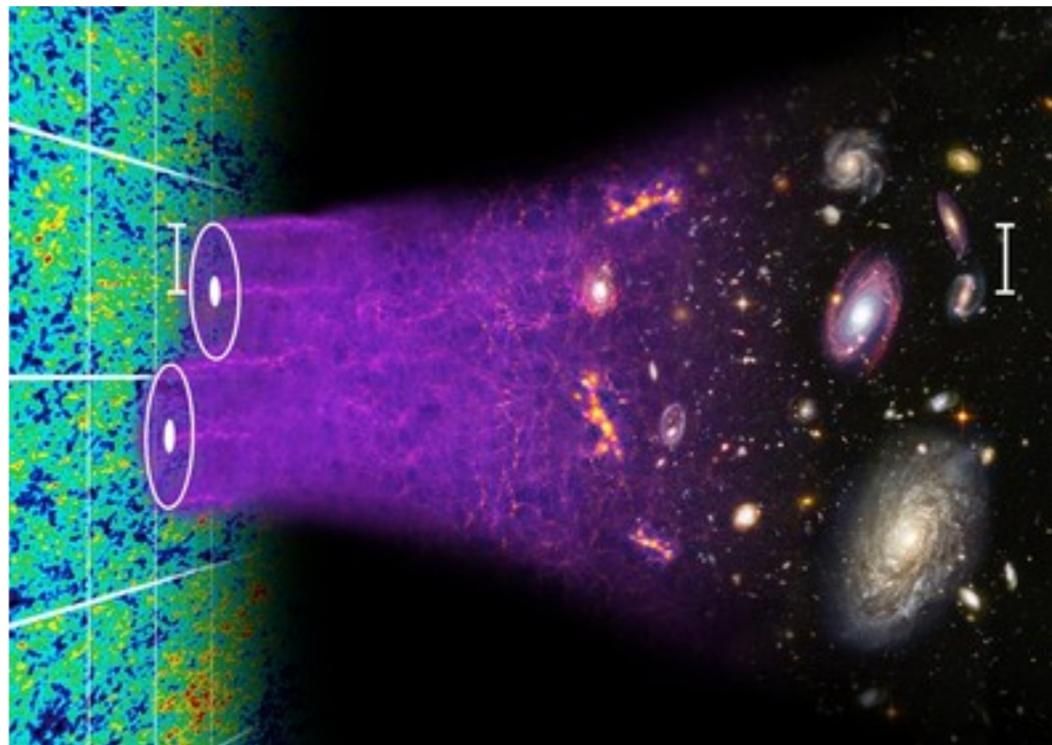


The story of “pertuby”

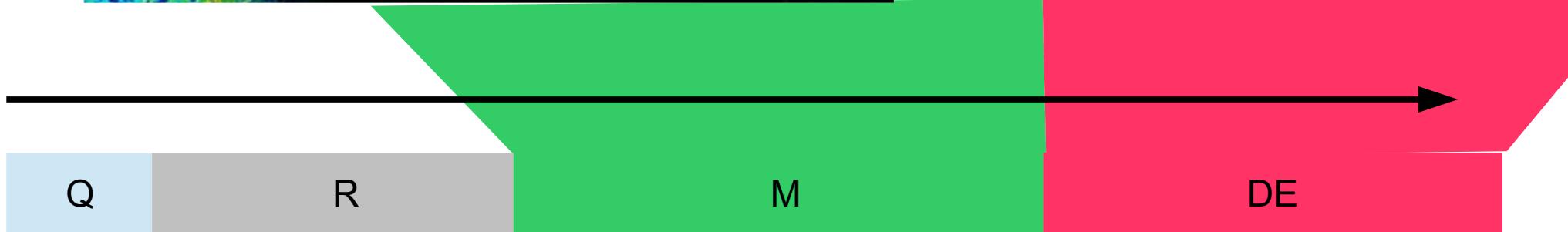
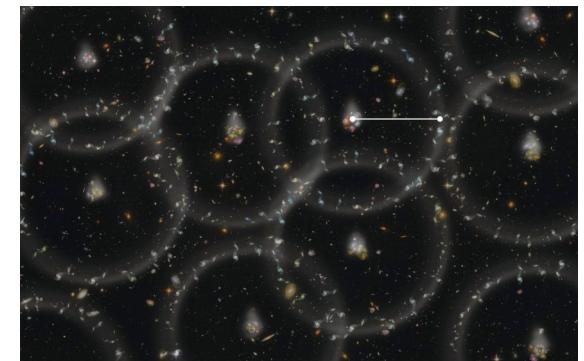


The story of “pertub”

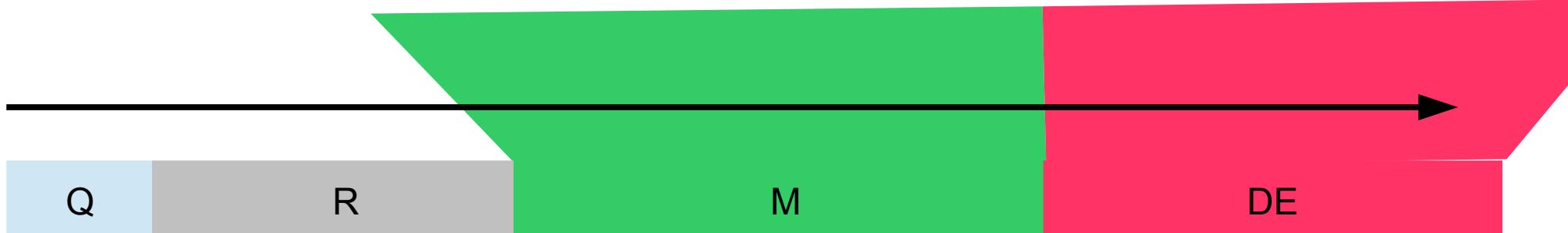
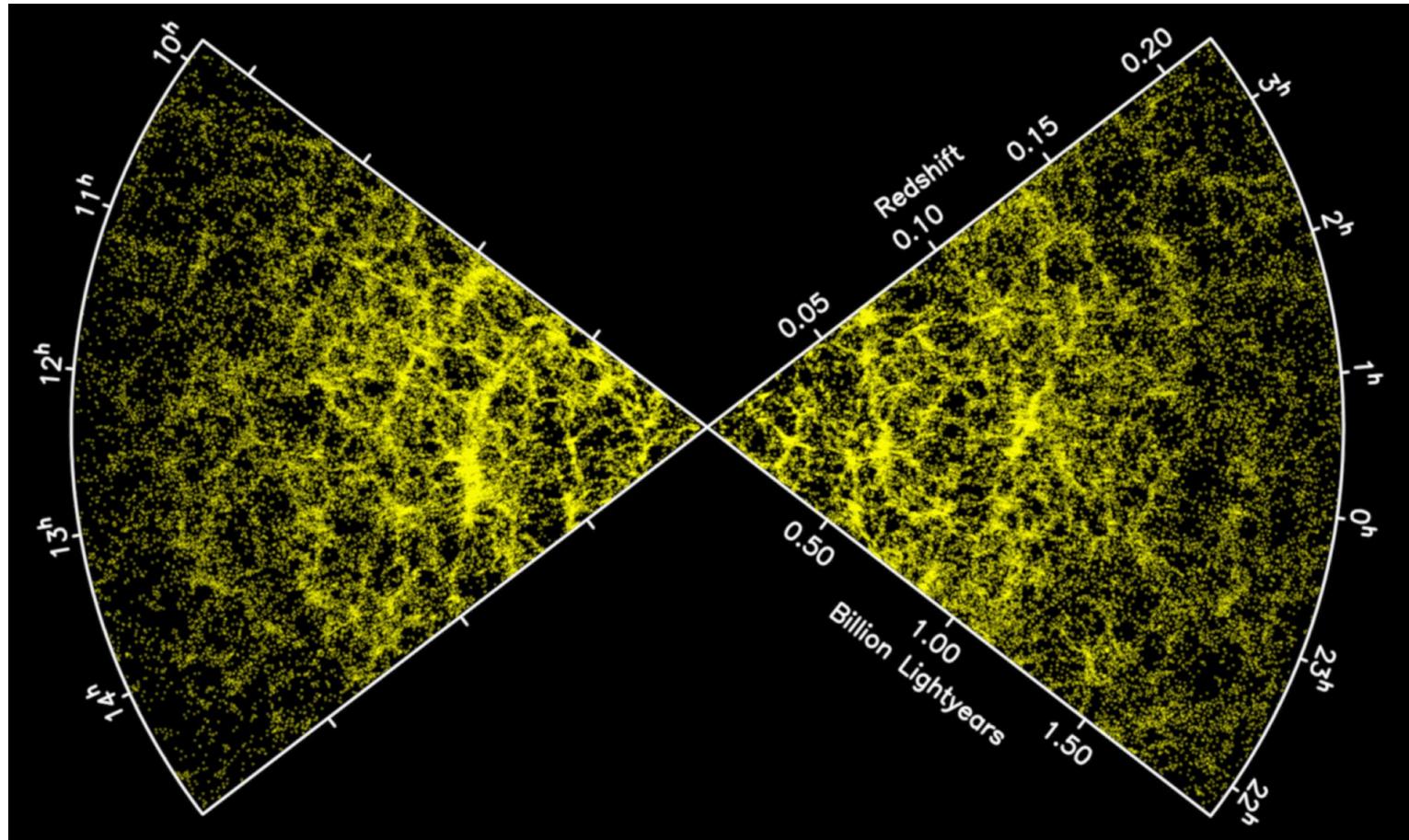
CMB seeds the structure at large scales



Baryonic
Acoustic
Oscillations



The story of “pertuby”

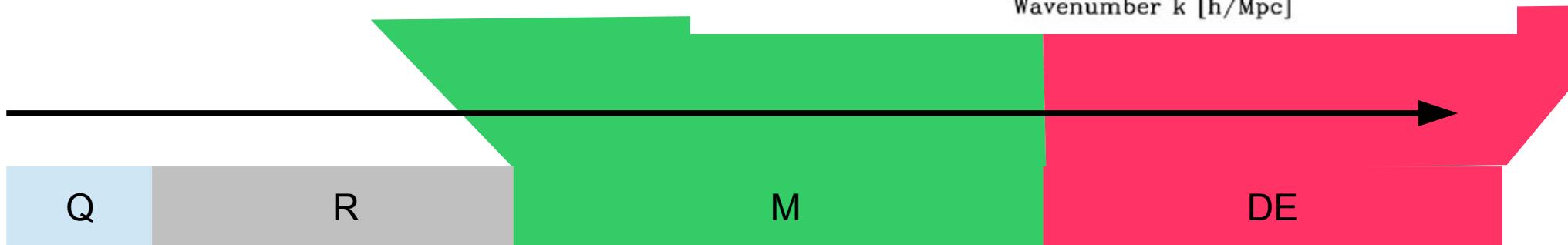
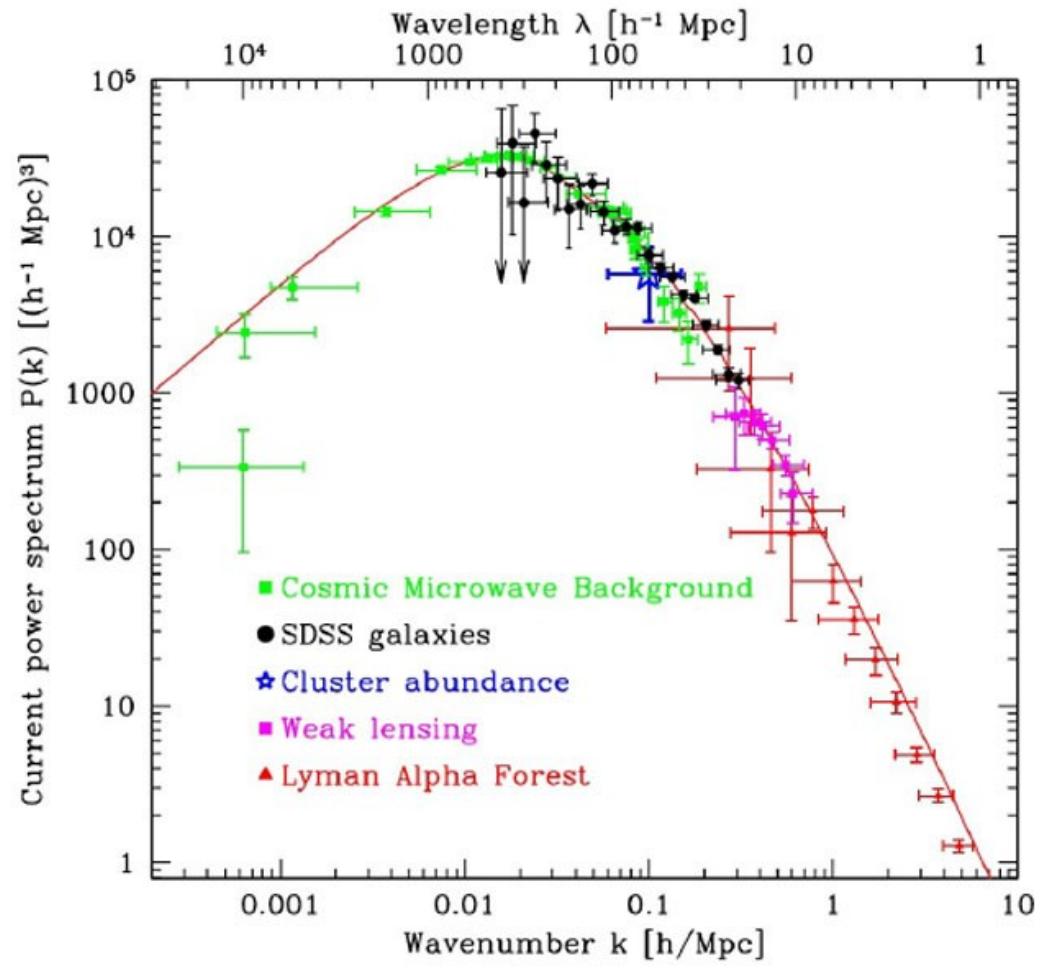


The story of “pertub”

Benefits

1) Volume Vs area (CMB)
*More Resolution!!!

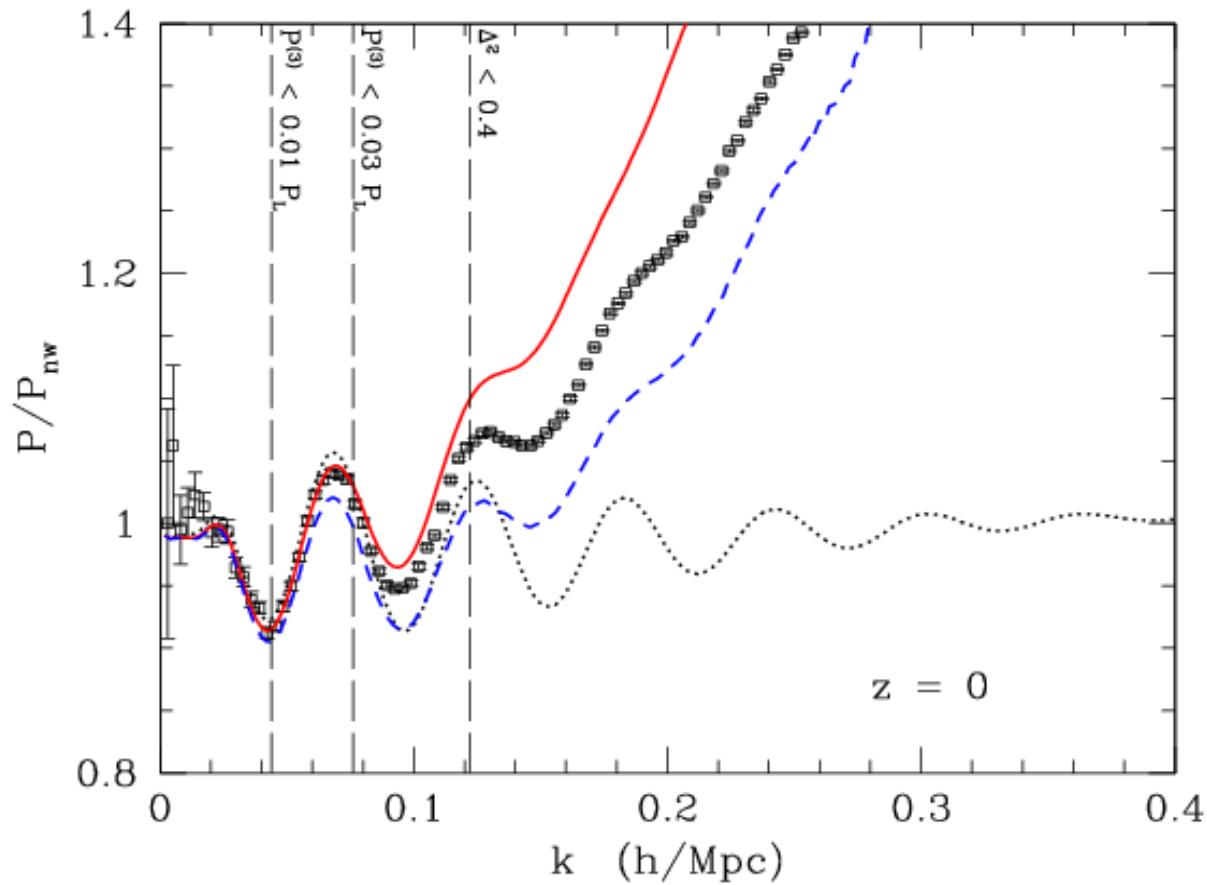
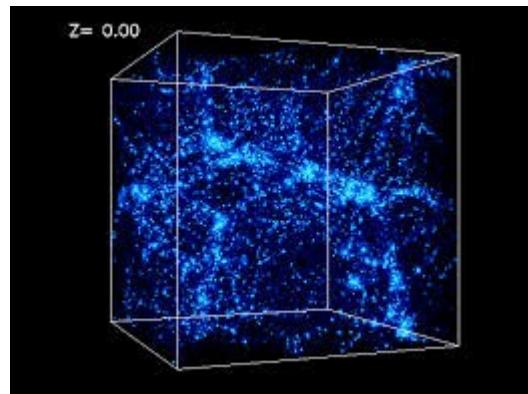
2) Non-linear physics
*More difficult
*Richness



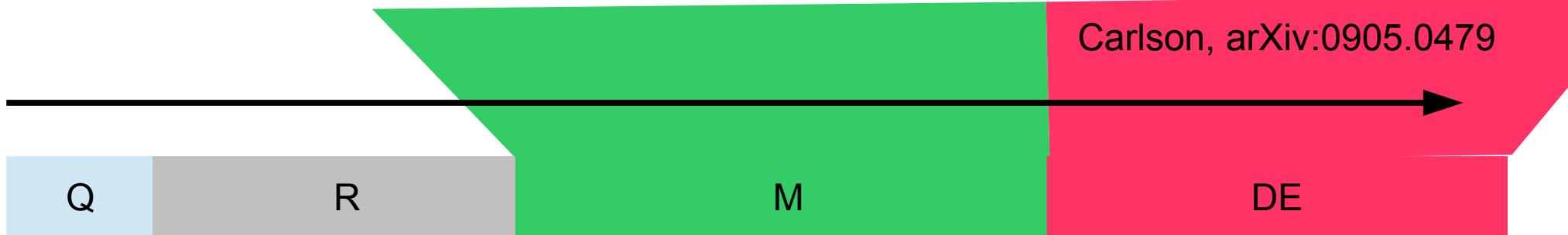
The story of “pertuby”

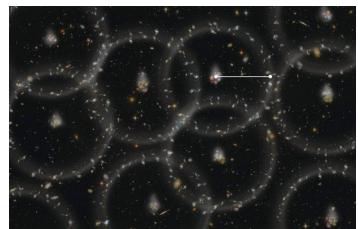
Perturbation theory

- 1) Series does not converge
- 2) Active area of research
- 3) Tool



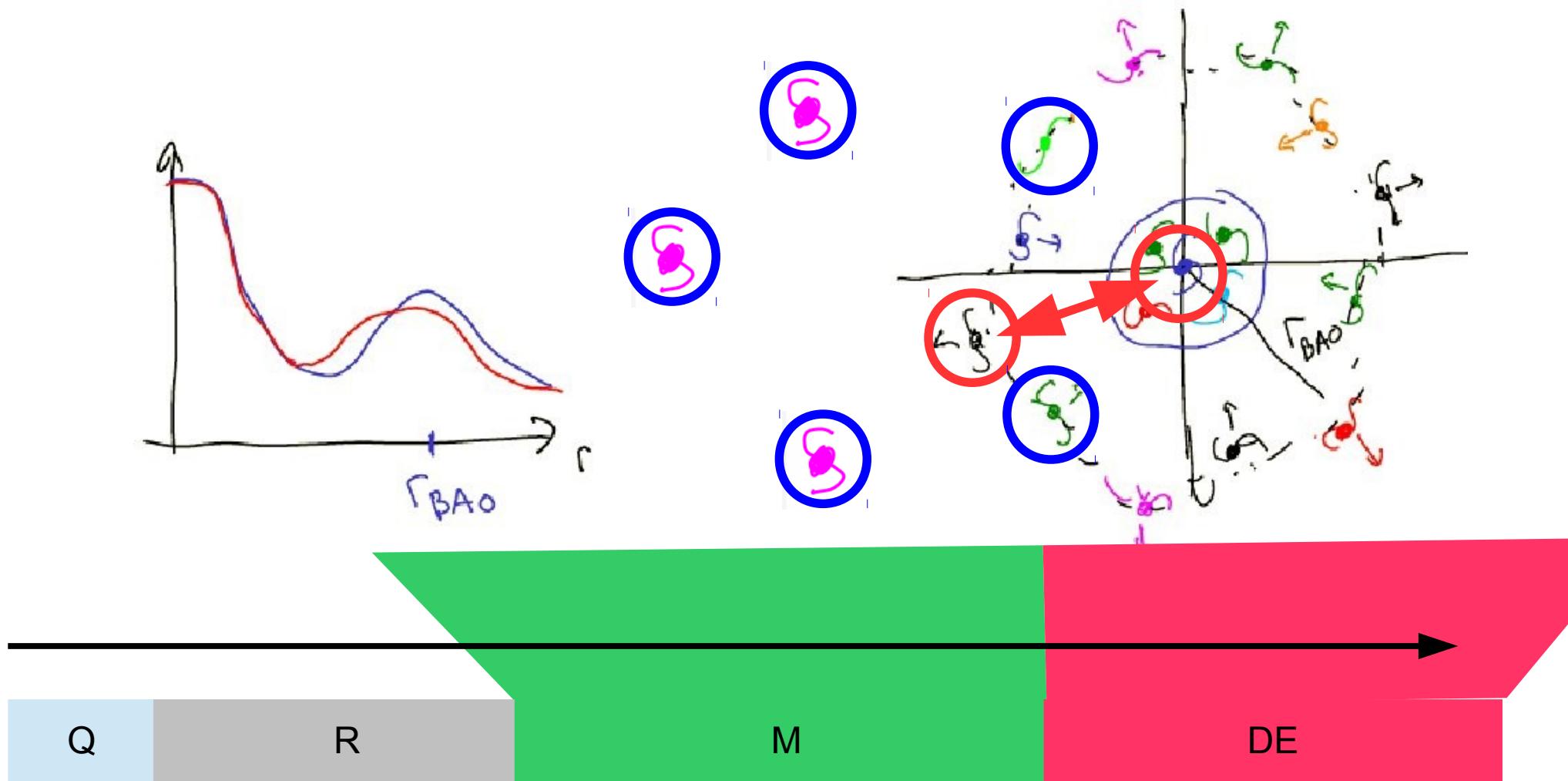
Carlson, arXiv:0905.0479





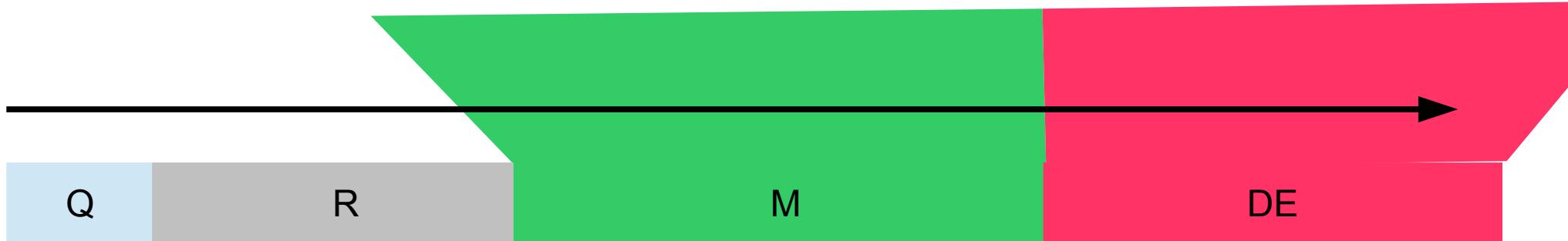
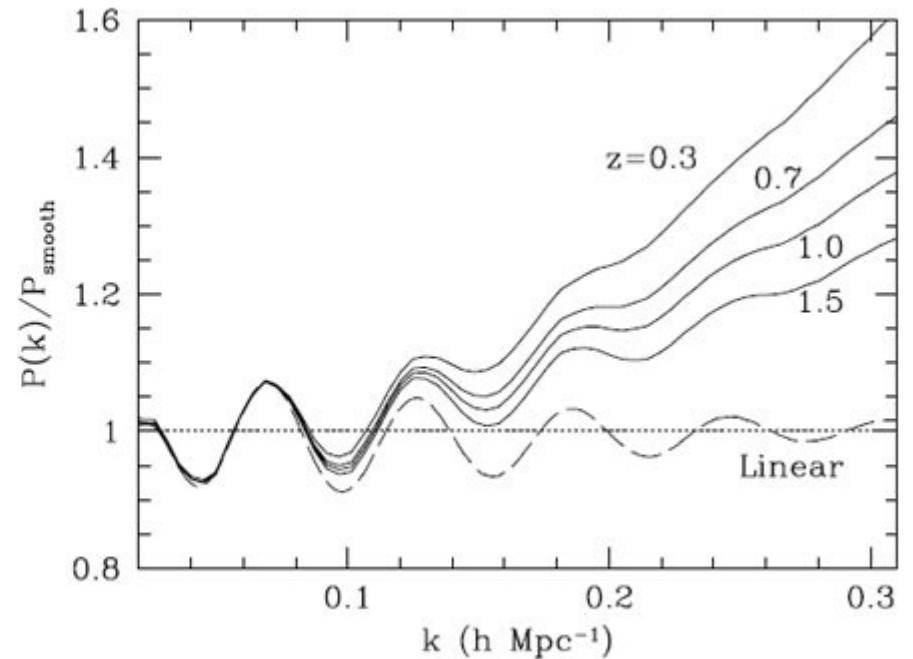
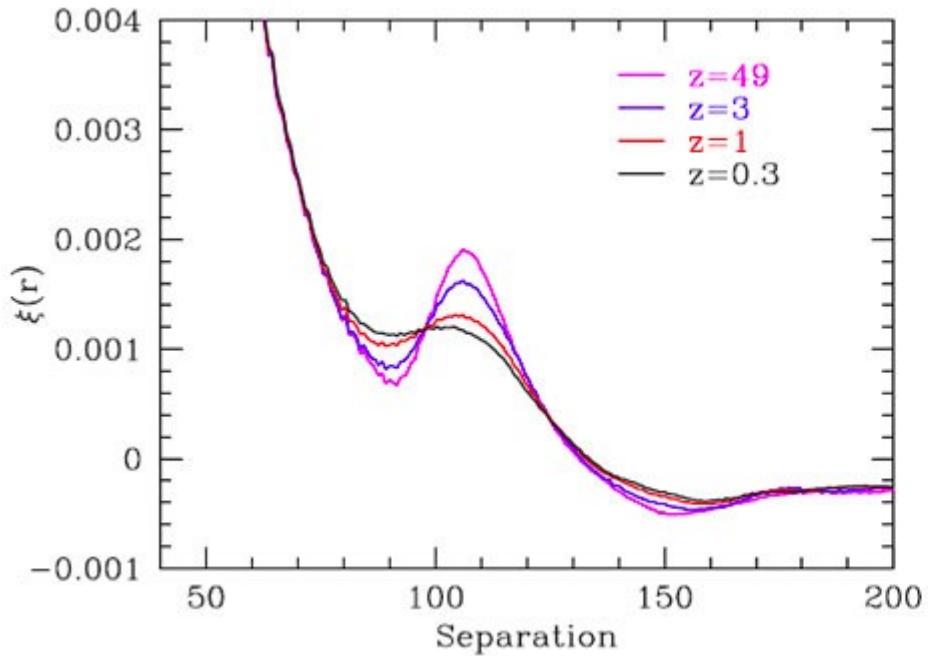
BAO signal in LSS

- 2PCF \rightarrow Non-linear mode evolution
 - It does contain the 3PCF, but *averaged*



BAO signal in LSS

- 2PCF  Non-linear mode evolution



Equations (Eulerian)

Continuity

$$\dot{\delta} + \nabla \cdot [(1 + \delta) \mathbf{v}] = 0,$$

Euler

$$\dot{\mathbf{v}} + (\mathbf{v} \cdot \nabla) \mathbf{v} = -\mathcal{H} \mathbf{v} - \nabla \phi,$$

Poisson

$$\nabla^2 \phi = 4\pi G a^2 \bar{\rho} \delta,$$

(Conformal time and comoving coordinates)

Equations (Eulerian)

In Fourier space

$$\boxed{\frac{\partial \delta_{\mathbf{k}}}{\partial \tau}(\tau) + \theta_{\mathbf{k}}(\tau)} = - \int \frac{d^3 k_1}{(2\pi)^3} \int d^3 k_2 \delta^D(k_1 + k_2 - k) \frac{\mathbf{k} \cdot \mathbf{k}}{k_1^2} \theta_{\mathbf{k}_1}(\tau), \delta_{\mathbf{k}_2}(\tau),$$

$$\boxed{\frac{\partial \theta_{\mathbf{k}}}{\partial \tau}(\tau) + \mathcal{H}(\tau) \theta_{\mathbf{k}}(\tau) + \frac{3}{2} \mathcal{H}^2(\tau) \Omega_m(\tau) \delta_{\mathbf{k}}(\tau)}$$

$$= - \int \frac{d^3 k_1}{(2\pi)^3} \int d^3 k_2 \delta^D(k_1 + k_2 - k) \frac{k^2 (\mathbf{k}_1 \cdot \mathbf{k})}{2k_1^2 k_2^2} \theta_{\mathbf{k}_1}(\tau) \theta_{\mathbf{k}_2}(\tau),$$

Linear Piece

Non-linear Piece

Equations (Eulerian)

To first order in perts

$$\boxed{\frac{\partial \delta_k}{\partial \tau}(\tau) + \theta_k(\tau)} = 0$$

$$\boxed{\frac{\partial \theta_k}{\partial \tau}(\tau) + \mathcal{H}(\tau)\theta_k(\tau) + \frac{3}{2}\mathcal{H}^2(\tau)\Omega_m(\tau)\delta_k(\tau)} = 0$$

$\delta_k^{(1)}$ = Decaying mode + Growing mode

Standard Perturbation Theory (SPT)

At higher orders

$$\delta_n(\mathbf{k}) = \int d^3\mathbf{q}_1 \dots \int d^3\mathbf{q}_n \delta_D(\mathbf{k} - \mathbf{q}_{1\dots n}) F_n(\mathbf{q}_1, \dots, \mathbf{q}_n) \delta_1(\mathbf{q}_1) \dots \delta_1(\mathbf{q}_n)$$



Obey simple recursion relations for some backgrounds

(Bernardeau et al, Phys.Rept. 367 (2002) 1-248)

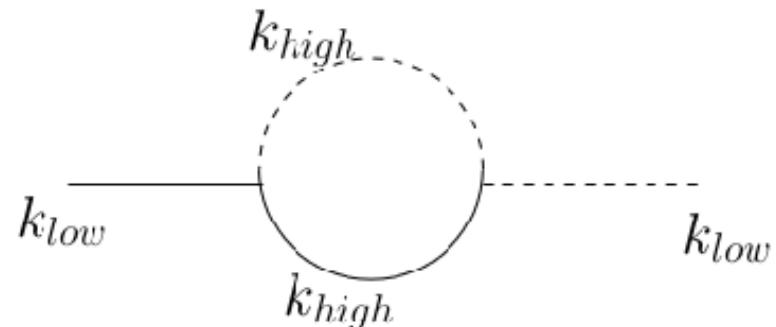
Standard Perturbation Theory (SPT)

FIRST ISSUE (UV)

$$\langle \delta_k^{(2)} \delta_k^{(2)} \rangle \sim \int d^3 k' \langle \delta_{k-k'}^{(1)} \delta_{k-k'}^{(1)} \rangle \langle \delta_{k'}^{(1)} \delta_{k'}^{(1)} \rangle$$

Notice that

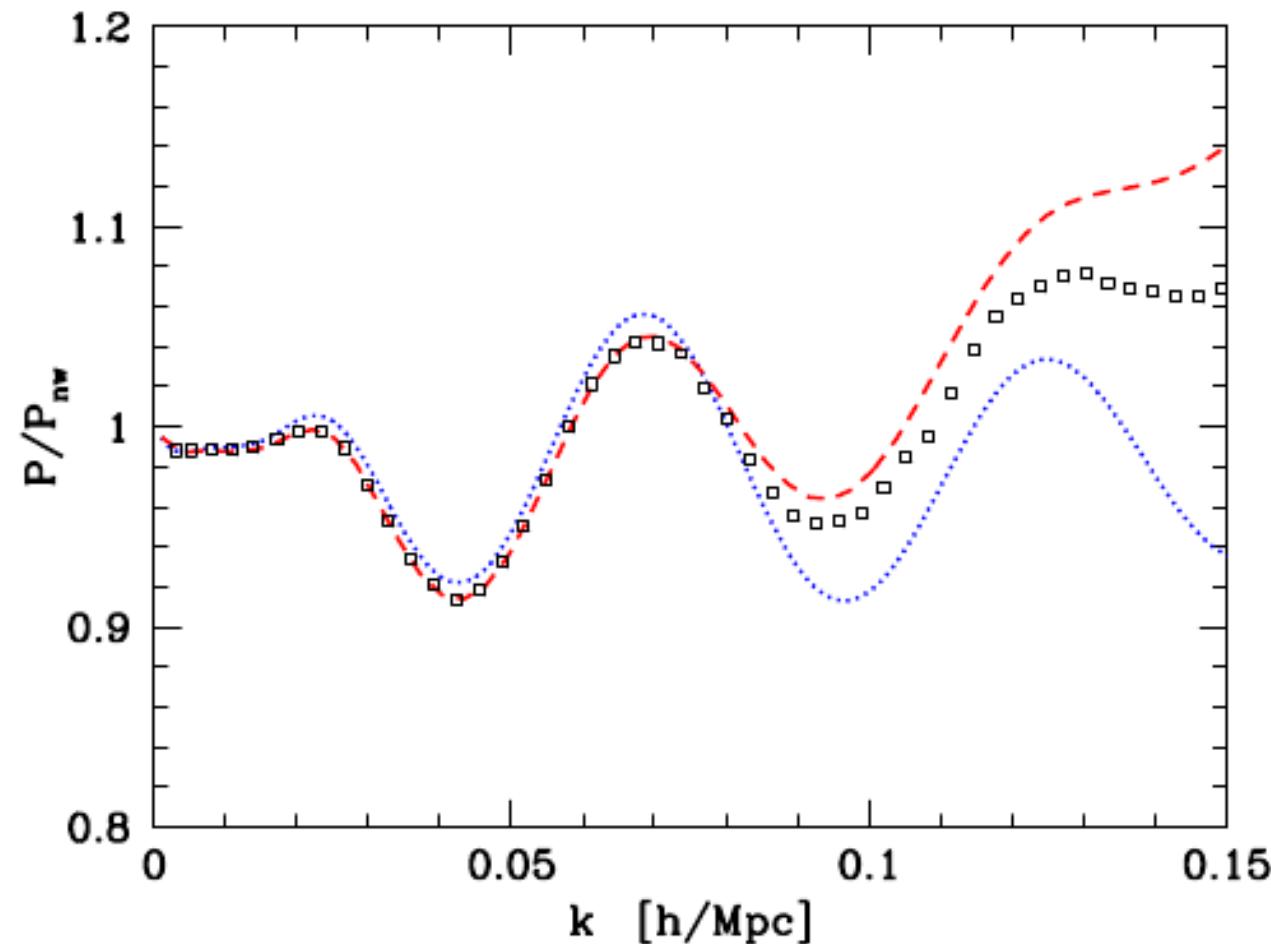
$$\delta \sim \frac{k}{k_{NL}} \gg 1 \quad \text{for} \quad k \gg k_{NL}$$



Perturbation theory breaks down!!!

Standard Perturbation Theory (SPT)

How does it fit our Universe (with BAO)?



Bright Future

- DESI
- LSST
- DES
- HETDEX
- WFIRST
- EUCLID
- Others.....

Bright Future

- DESI
- LSST
- DES
- HETDEX
- WFIRST
- EUCLID
- Others.....

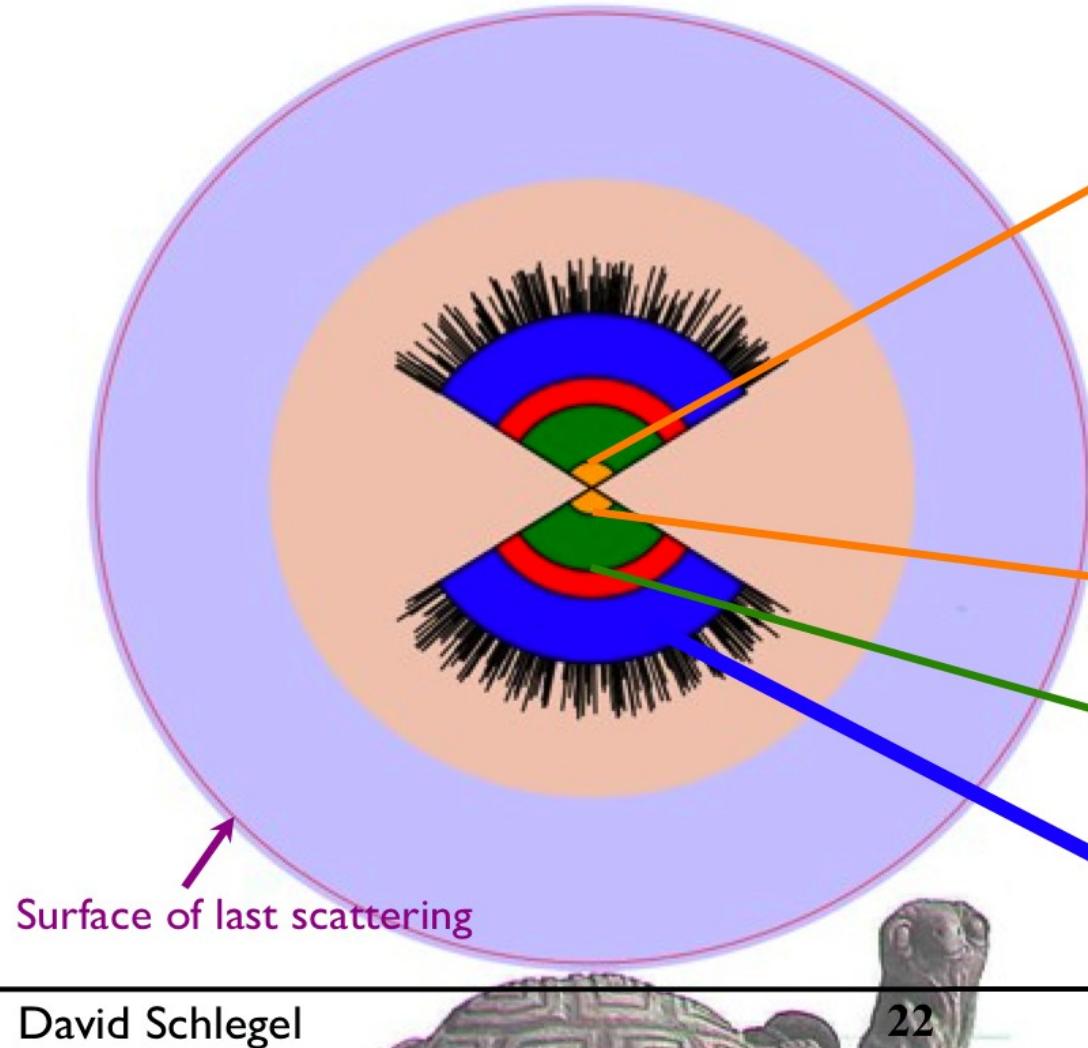
DESI MEXICO

- 10 researchers
 - 4 UNAM ([A. de la Macorra](#) & M. Vargas in IF, O. Valenzuela in IA, M. Alcubierre in CN)
 - 3 DCI-UG: Luis Ureña, **Gustavo Niz**, Alma González
 - 1 Cinvestav (T. Matos)
 - 1 ININ (J. Cervantes, A. Avilés)
- “[Proyecto fronteras de la ciencia](#)” CONACYT
- Cosmological simulations
- Data analysis
- LCDM-alternative models

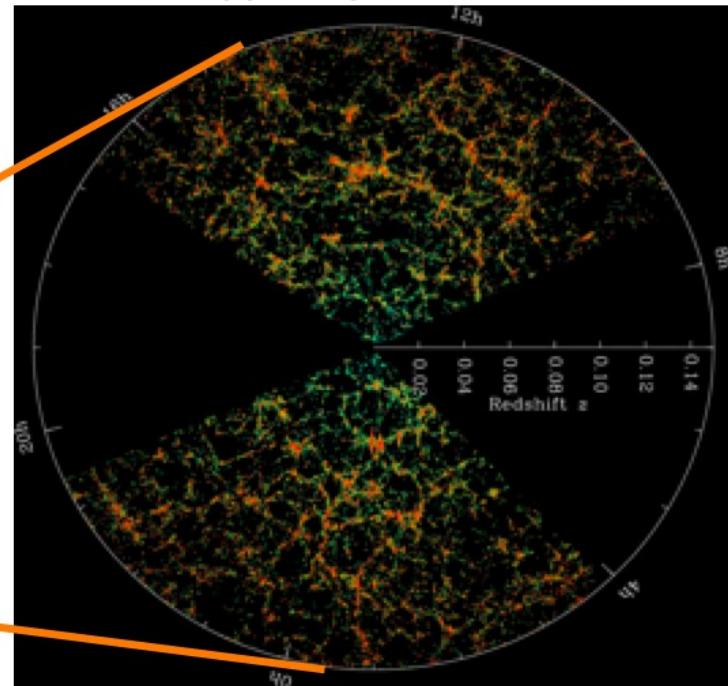
Dark Energy Spectroscopic Instrument

Sensitivity to new physics scales as volume surveys -- # of modes

Our observable Universe



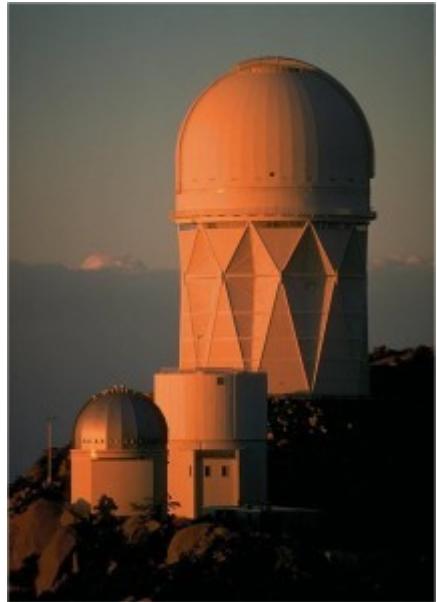
Volume mapped by SDSS + SDSS-II



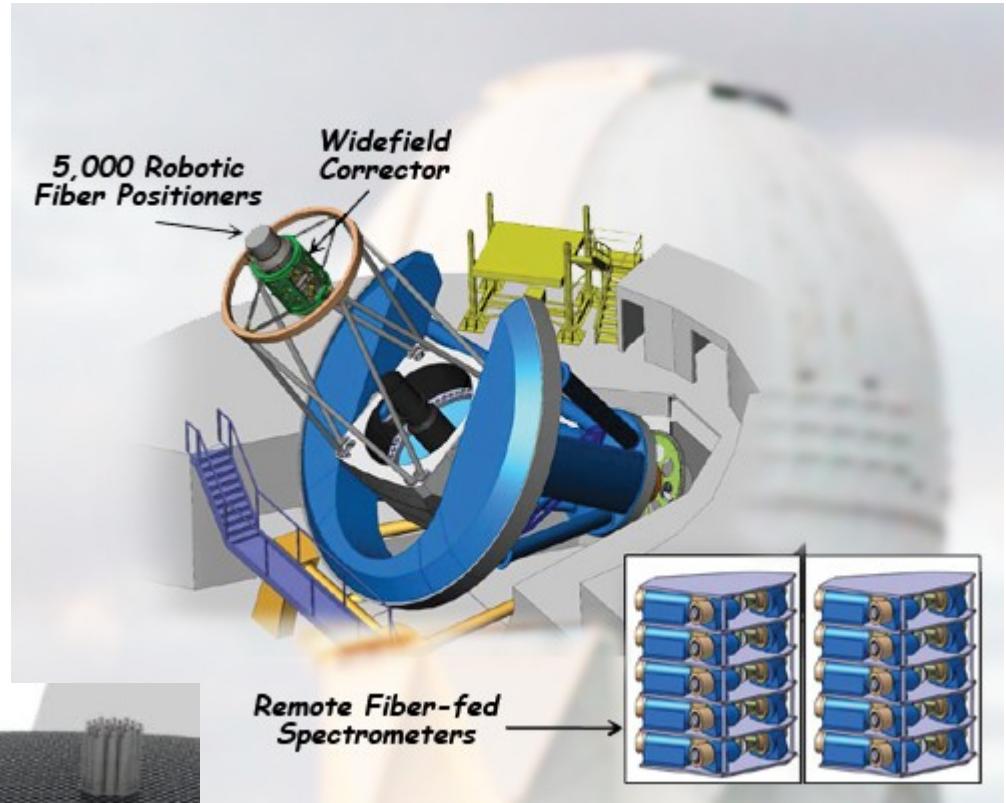
Volume to be mapped by SDSS-III/BOSS
(ca. 2015)



Dark Energy Spectroscopic Instrument

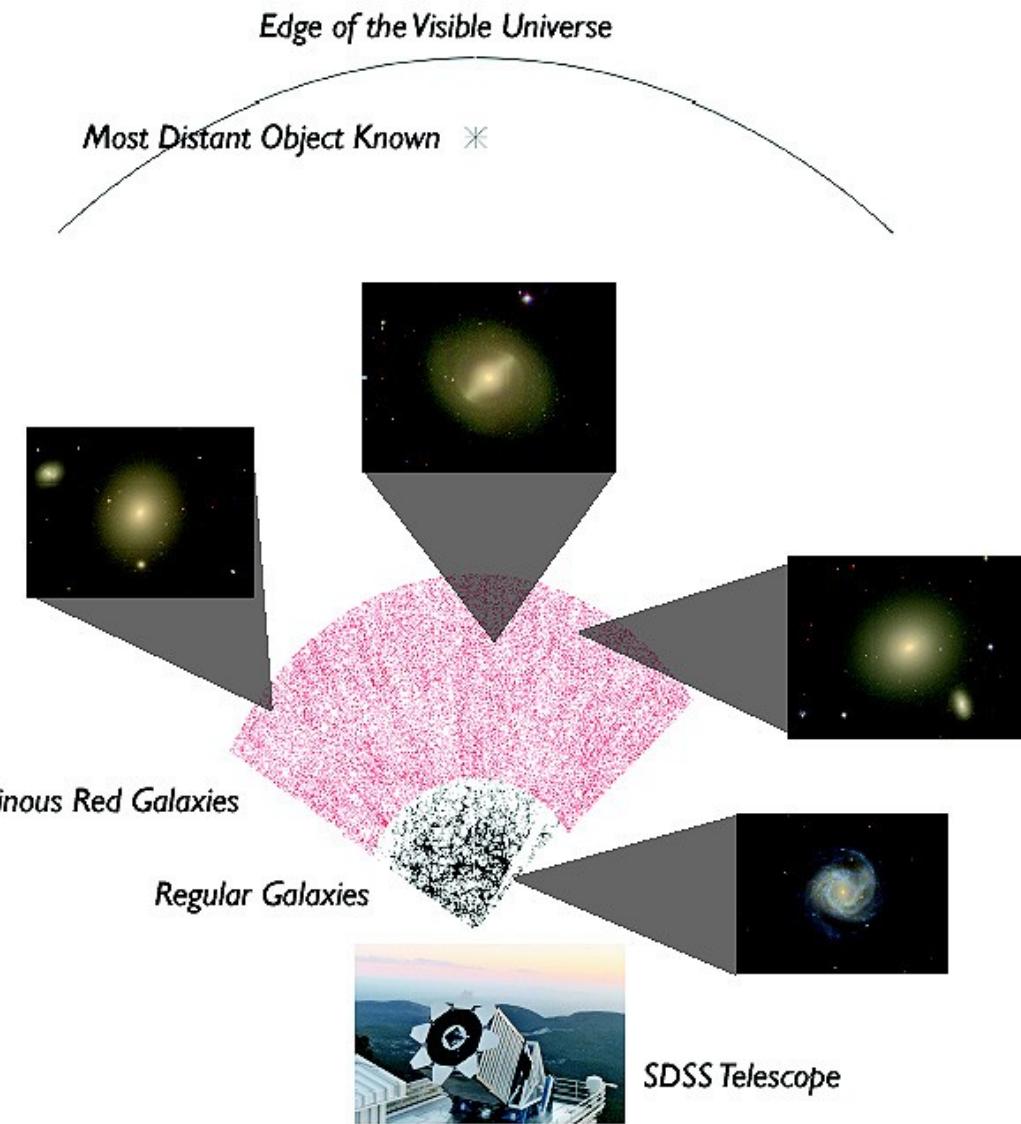


Kitt Peak (Mayall 4m)



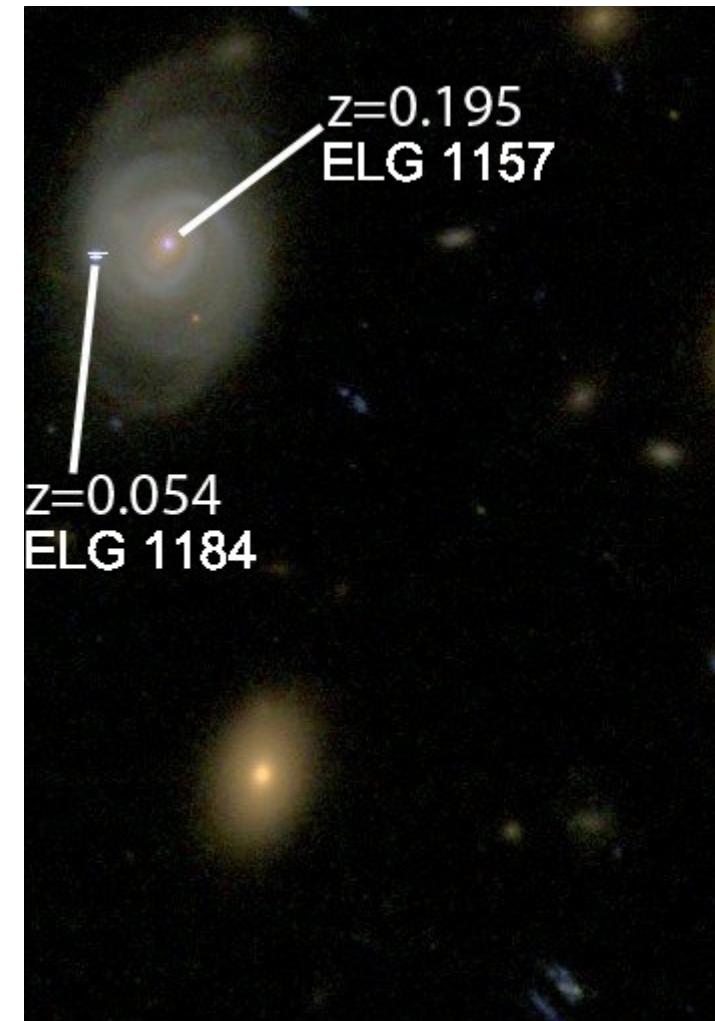
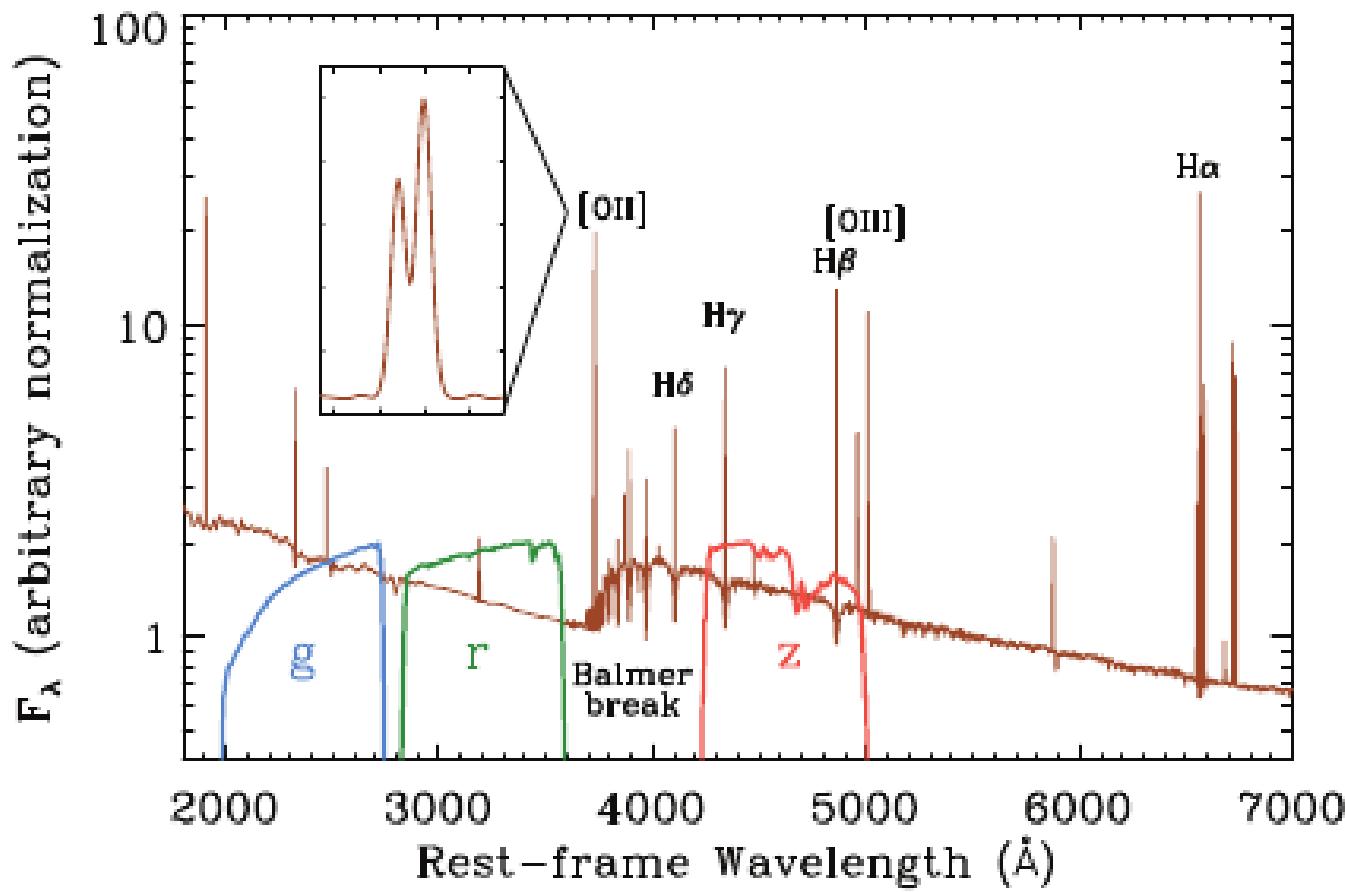
DESI

Luminous Red Galaxies (LRG)



DESI

Emission Line Galaxies (ELG)

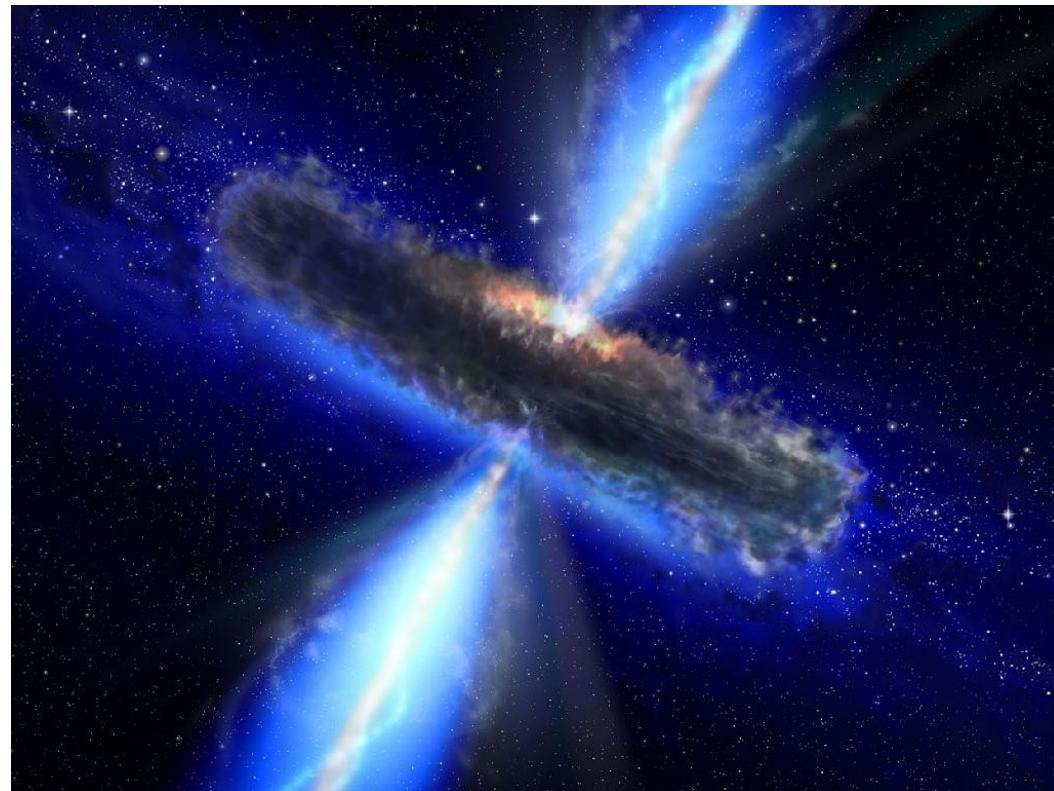


DESI

Quasars

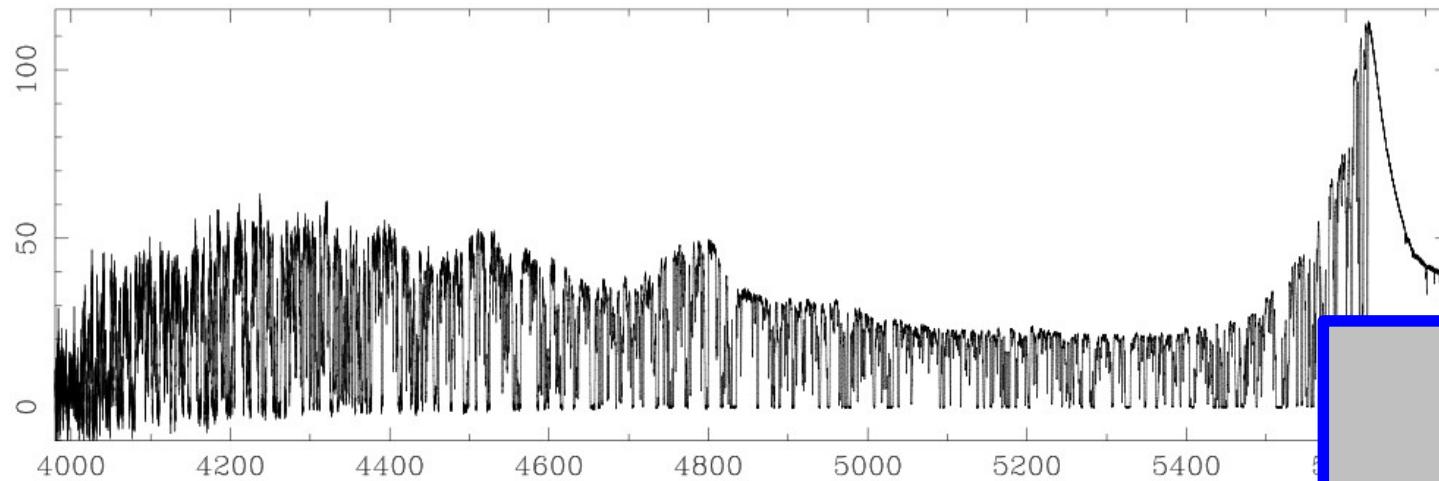
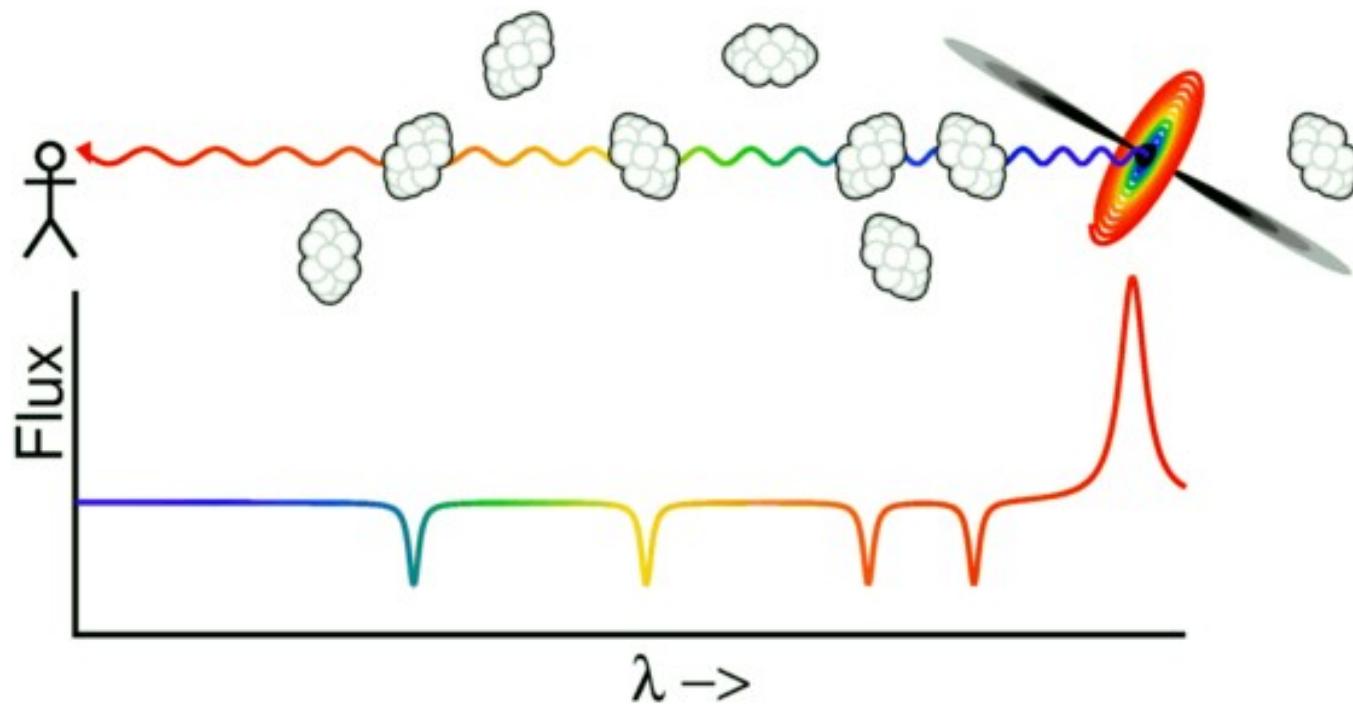


ULAS J1120+0641
($z \sim 7$)



DESI

Lyman-alfa forest



DESI



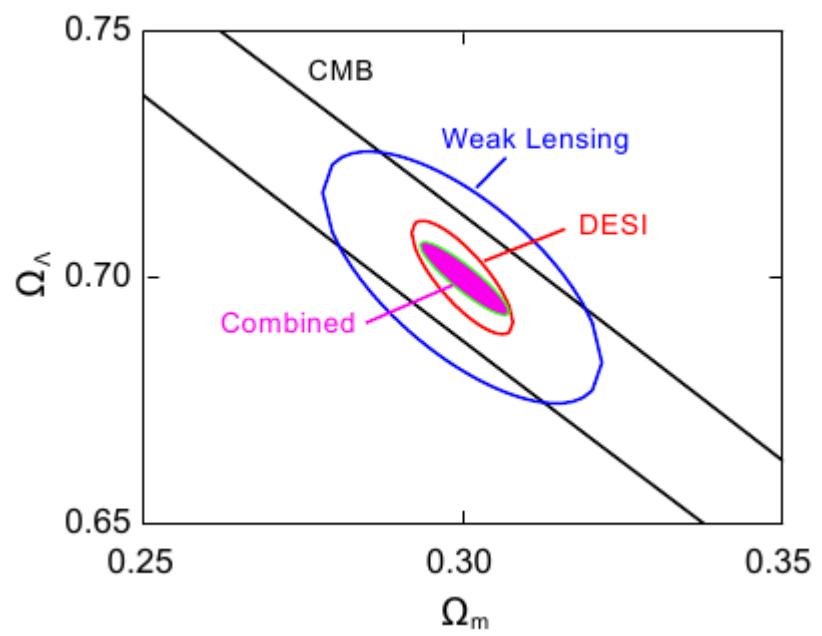
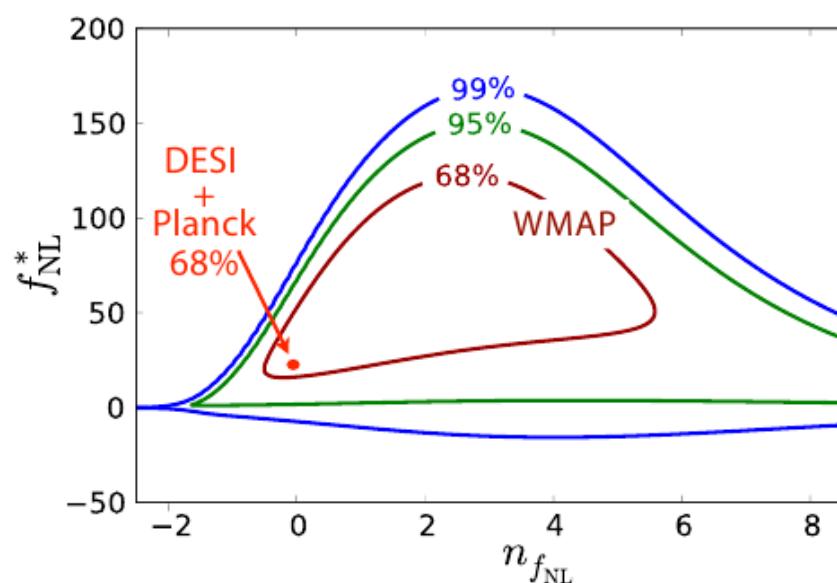
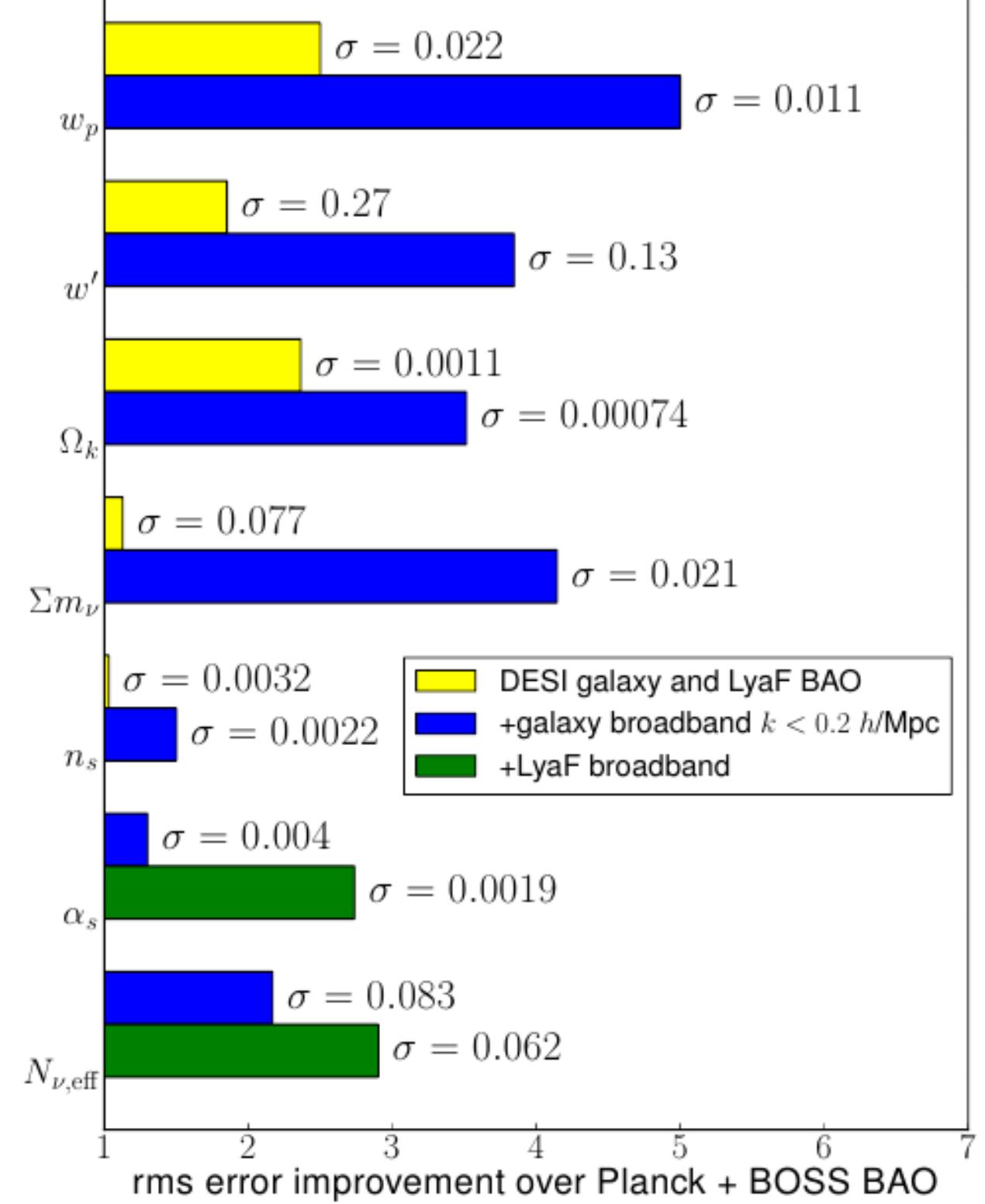
Dark Energy Spectroscopic Instrument

Goal: to map

- > 18 millions of ELG
- > 4 millions of LRG
- > 2.5 millions of quasars
- > 0.5 millions of Ly-alfa forest

Precision < 1%

DESI



DESI

Conclusions

- Observables in the Universe (CMB & LSS) are a door to the Early Universe physics and the dark sector (dark matter & **dark energy**)
- High precision cosmology, such as **DESI**, will be the leading observationally experiment in the next decade



New perspective into the Universe

Cosmic collider

Universe
evolution

Dark energy
Dark matter
(Modified) gravity

Distortion

