# 计算物理第五次作业

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#### Abstract

本次作业学习了Crank-Nicolson 差分格式求解含时微分方程的解法,并运用该方法求解了初 始条件为Gauss函数的扩散方程,并画出了该扩散方程随时间的演化。

# I. 第一个问题

#### A. 问题描述

已知一维扩散方程形式为

$$\begin{split} \frac{\partial u(x,t)}{\partial t} &= D \frac{\partial^2 u(x,t)}{\partial x^2} \\ u(x,t)|_{t=0} &= u_0(x) \\ a_1 u + b_1 \frac{\partial u}{\partial n} &= c_1 \quad (x=0) \\ a_2 u + b_2 \frac{\partial u}{\partial n} &= c_2 \quad (x=a_0) \end{split}$$

例:考虑f=0,D=1,边界条件为 $\phi$ (0) =  $\phi$ (1) = 0。设给定的初始条件是中心在 $x=\frac{1}{2}$ 的一 个Gauss函数,即:

$$\phi(x,t=0) = e^{-20(x-\frac{1}{2})^2} - e^{-20(x-\frac{3}{2})^2} - e^{-20(x+\frac{1}{2})^2}$$

求随后各个时刻的o。

求随后各个时刻的
$$\phi$$
。 若设 $\tau=1+80t$ ,则解析解为: 
$$\phi(x,t)=\tau^{-\frac{1}{2}}[\mathrm{e}^{-20(x-\frac{1}{2})^2}-\mathrm{e}^{-20(x-\frac{3}{2})^2}-\mathrm{e}^{-20(x+\frac{1}{2})^2}]$$
 上例可以用标准的Crank-Nicolson方法求

上例可以用标准的Crank-Nicolson方法求 解。

## 问题分析

Crank-Nicolson是一种数值分析的有限差 分法, 可用于数值求解热方程以及类似形式 的偏微分方程[2]。它在时间方向上是隐式的 二阶方法,可以写成隐式的Runge-Kutta法, 数值稳定。该方法诞生于20世纪,由John Crank与Phyllis Nicolson发展[3]。

可以证明Crank-Nicolson方法对于扩散方 程(以及许多其他方程)是无条件稳定[4]。但 1. 一维扩散方程Crank-Nicolson格式

已知一维扩散方程形式为

$$\begin{cases} \frac{\partial u(x,t)}{\partial t} = D \frac{\partial^2 u(x,t)}{\partial x^2} \\ u(x,t)|_{t=0} = u_0(x) \\ a_1 u + b_1 \frac{\partial u}{\partial n} = c_1 \qquad (x=0) \\ a_2 u + b_2 \frac{\partial u}{\partial n} = c_2 \qquad (x=a_0) \end{cases}$$

$$\begin{cases} x_i = (i-1)h & (i = 1, 2, ..., N) \\ t_k = k\tau & (k = 1, 2, ..., k_{max}) \end{cases}$$

节点处的函数为

$$u(x-i,t_k) = u_i^k$$

是,如果时间步长 $\Delta t$ 乘以热扩散率,再除以空 间步长平方 $\Delta t^2$ 的值过大(根据Von Neumann稳 定性分析,以大于5为准),近似解中将存在 虚假的振荡或衰减。基于这个原因,当要求大 时间步或高空间分辨率的时候,往往会采用数 值精确较差的后向欧拉法进行计算, 这样即可 以保证稳定, 又避免了解的伪振荡。

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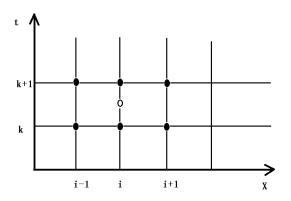


Figure 1: 一维扩散方程的隐式六点差分格式

利用如下差分格式

$$f'(0) \approx \frac{f_1 - f_{-1}}{2h}$$

$$\frac{\partial u}{\partial t} = \frac{u_i^{k+1} - u_i^k}{\tau}$$

$$\frac{\partial^2 u}{\partial x^2} = \frac{1}{2} \left[ \left( \frac{u_{i+1} - 2u_i + u_{i-1}}{h^2} \right)^k + \left( \frac{u_{i+1} - 2u_i + u_{i-1}}{h^2} \right)^{k+1} \right]$$

$$\frac{u_i^{k+1} - u_i^k}{\tau} = \frac{D}{2h^2} \left[ (u_{i+1} - 2u_i + u_{i-1})^k + (u_{i+1} - 2u_i + u_{i-1})^k \right]$$

令 $P=\frac{ au D}{h^2}, P_1=\frac{1}{P}+1, P_2=\frac{1}{P}-1$ ,则可以得到隐式六点差分格式(Crank-Nicolson格式)

$$(-u_{i-1} + 2P_1u_{i+1} - u_{i+1})^{k+1}$$
  
=  $(u_{i-1} + 2P_2u_i + u_{i+1})^k$ 

2. 边界条件的处理

$$\begin{cases} a_1 u + b_1 \frac{\partial u}{\partial n} = c_1 & (x = 0) \\ a_2 u + b_2 \frac{\partial u}{\partial n} = c_2 & (x = a_0) \end{cases}$$

设置虚格点i=0, i=N+1,利用中心差商公式有

$$\begin{cases} a_1 u_1 + b_1 \frac{u_0 - u_2}{2h} = c_1 & (x = 0) \\ a_2 u_n + b_2 \frac{u_{n+1} - u_{n-1}}{2h} = c_2 & (x = a_0) \end{cases}$$

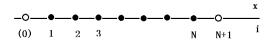


Figure 2: 1维扩散方程的边界条件

解 出 $u_0, u_{n+1}$ 代 入i=1, i=N的Crank-Nicolson差分格式

$$[(b_1P_1 + ha_1) u_1 - b_1u_2]^{k+1}$$

$$= [(b_1P_2 - ha_1) u_1 + b_1u_2]^k + 2hc_1$$

$$[-b_2u_{N-1} + (b_2P_1 + ha_2) u_N]^{k+1}$$

$$= [b_2u_{N-1} + (b_2P_2 - ha_2) u_N]^k + 2hc_2$$

差分方程组及其求解:

将上述差分公式和边界条件结合起来,得 到差分线性方程组,其形式为:

$$AU = R$$

其中,

 $\mathbf{U} = (u_1, u_2, ..., u_n)$ 是未知量组成的矢量。  $\mathbf{R} = (R_1, R_2, ..., R_n)$ 是有前一时刻的u值组成的矢量。

则有

$$R_1 = (b_1 P_2 - ha_1)u_1 + b_1 u_2 + 2hc_1$$

$$R_i = u_{i-1} + 2P_2 u_i + u_{i+1} i = 2, 3, ..., N - 1$$

$$R_N = b_2 u_{N-1} + (b_2 P_2 - ha_2)u_N + 2hc_2$$

系数矩阵A是三对角的

$$\begin{bmatrix} b_1P_1 + ha_1 & -b_1 & 0 & \cdots & 0 & 0 \\ -1 & 2P_1 & -1 & 0 & \cdots & 0 \\ 0 & -1 & 2P_1 & -1 & \ddots & \vdots \\ \vdots & \ddots & \ddots & \ddots & \ddots & \vdots \\ \vdots & \ddots & -1 & 2P_1 & -1 & 0 \\ 0 & \cdots & 0 & -1 & 2P_1 & -1 \\ 0 & 0 & \cdots & 0 & -b_2 & b_2P_1 + ha_2 \end{bmatrix}$$

这是一个三对角问题,应用追赶法即可得到 $u_i^{k+1}$ ,而不需要对矩阵直接求逆。

# C. 计算结果

运行结果见压缩包中的figure文件夹中的动画1dDiffusion.gif,以下只展示截图

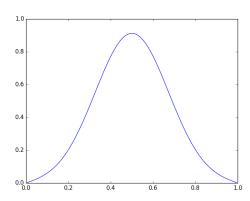


Figure 3: 1维扩散方程瞬时截图

#### 第二个问题 II.

# A. 问题表述

试将第一个问题拓展成二维, 求解二维扩 散方程随时间的演化并画图(初始条件、边界 条件自设, 你可以选择二维Crank-Nicolson方 法,也可以选择简单Eular法,也可以选择上一 章提到的迭代法)。

# B. 问题分析

# 1. 二维扩散方程

将扩散问题延伸到二维的Cartesian网络, 我们有类似的推导:

二维的扩散方程形式为

$$\begin{cases} \frac{\partial u}{\partial t} = D\left(\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2}\right) & \begin{pmatrix} 0 < x < a_0 \\ 0 < y < b_0 \\ 0 < t < t_{max} \end{pmatrix} & \text{边界条件的差分格式} \\ u(x,y,0) = u_0(x,y) & \begin{cases} a_i u + b_i \frac{\partial u}{\partial n} = c_i(y,t) & (1)(2) \text{边界}(i=1,2) \\ a_i u + b_i \frac{\partial u}{\partial n} = c_i(x,t) & (3)(4) \text{边界}(i=3,4) \end{cases}$$

取 $\Delta x = \Delta y = h$ 的正方形网格覆盖x-y平面,并 取 $\Delta t = \tau$ , 节点坐标为

$$\begin{cases} x_i = (i-1)h & (i = 1, 2, \dots, N) \\ y_i = (i-1)h & (j = 1, 2, \dots, M) \\ t_k = k\tau & (k = 1, 2, \dots, k_{max}) \end{cases}$$

节点处的函数为 $u(x_i, y_j, t_k) = u_{ij}^k$ 

代替,将 $\frac{\partial^2 u}{\partial x^2}$ 用k+1时的中心差商代替,将 $\frac{\partial^2 u}{\partial u^2}$ 用k时的中心差商代替,则二维扩散方程 变为

$$\frac{u_{ij}^{k+1} - u_{ij}^k}{\tau} = \frac{D}{h^2} [(u_{i+1,j} - 2u_{ij} + u_{i-1,j})^{k+1} + (u_{i,j+1} - 2u_{ij} + u_{i,j-1})^k]$$

 $\dot{a}(i,j,k+rac{3}{2})$ 点,将 $rac{\partial u}{\partial t}$ 用k+1时中心差商代替,将 $rac{\partial^2 u}{\partial x^2}$ 用k+1时的中心差商代替,将 $rac{\partial^2 u}{\partial y^2}$ 用k+2时的中心差商代替,则二维扩散 方程变为

$$\frac{u_{ij}^{k+2} - u_{ij}^{k+1}}{\tau} = \frac{D}{h^2} [(u_{i+1,j} - 2u_{ij} + u_{i-1,j})^{k+1} + (u_{i,j+1} - 2u_{ij} + u_{i,j-1})^{k+2}]$$

令
$$P=rac{ au D}{h^2}, P_1=rac{1}{P}+1, P_2=rac{1}{P}-1$$
,则有

$$(-u_{i-1,j} + 2P_1u_{ij} - u_{i+1,j})^{k+1}$$

$$= (u_{i,j-1} + 2P_2u_{ij} + u_{i,j+1})^k \qquad \cdots (a)$$

$$(-u_{i-1,j} + 2P_1u_{ij} - u_{i+1,j})^{k+2}$$

$$= (u_{i,j-1} + 2P_2u_{ij} + u_{i+1,j})^{k+1} \qquad \cdots (b)$$

说明·

- 1.求解方法与一维情形类似。
- 2.k为奇数时用(b)式沿y方向计算, k为偶数时 用(a)式沿x方向计算,并且只输出后一结果, 以减少因为时间上的不一致引起的误差。
- 3.该方法是无条件稳定的。

# 2. 边界条件的处理

$$\begin{cases} a_i u + b_i \frac{\partial u}{\partial n} = c_i(y, t) & (1)(2) \dot{\mathcal{D}} \, \Re(i = 1, 2) \\ a_i u + b_i \frac{\partial u}{\partial n} = c_i(x, t) & (3)(4) \dot{\mathcal{D}} \, \Re(i = 3, 4) \end{cases}$$

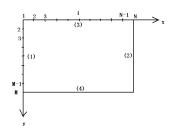


Figure 4: 2维扩散方程边界条件

设置虚节点: 在(1)(2)边界上有

$$\begin{cases}
b_1 u_{0j} = 2hc_1(y_i, t) - 2ha_1 u_{1j} + b_1 u_{2j} & (1) \\
b_2 u_{n+1,j} = 2hc_2(y_i, t) - 2ha_1 u_{nj} + b_2 u_{n-1,j} & (2)
\end{cases}$$

在(3)(4)边界上有

$$\begin{cases}
b_3 u_{i0} = 2hc_3(x_i, t) - 2ha_3 u_{i1} + b_3 u_{i2} \\
b_4 u_{i,m+1} = 2hc_4(x_i, t) - 2ha_4 u_{im} + b_4 u_{i,m-1}
\end{cases} (3)$$

解出 $u_{i0}, u_{i,m+1}$  结果会是带形矩阵的方程式。

# C. 计算结果

运行结果见压缩包中的figure文件夹中的动画2dDiffusion.mp4,以下只展示截图

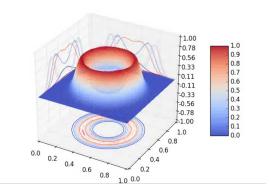


Figure 5: 2维扩散方程动画瞬时截图

# **Appendices**

# A. 使用Crank-Nicolson方法求解一维扩散方程

Here is the program to solve the one dimensional diffusion equation by Crank-Nicolson method in Python programming language.

# Input Python source:

```
2
    This program solves the heat equation
        11 t = 11 xx
 4
    with dirichlet boundary condition
5
        u(0,t) = u(1,t) = 0
    with the Initial Conditions
7
        u(x,0) = \exp(-20*(x-0.5)**2) - \exp(-20*(x-1.5)**2) - \exp(-20*(x+0.5)**2)))
8
    over the domain x = [0, 1]
10
        The program solves the heat equation using a finite difference method
11
    where we use a center difference method in space and Crank-Nicolson in time.
12
13
    import scipy as sc
14
    import scipy.sparse as sparse
15
   import scipy.sparse.linalg
   import numpy as np
17
   import matplotlib.pyplot as plt
    # Number of internal points
19
   N=200
    # Calculate Spatial Step-Size
21
   h=1/(N+1.0)
22
    # Create Temporal Step-Size, TFinal, Number of Time-Steps
23
    k=h/2
24
   TFinal=1
   NumOfTimeSteps=int (TFinal/k)
    # Create grid-points on x axis
    x=np.linspace(0,1,N+2)
28
    x=x[1:-1]
29
    # Initial Conditions
   |u=np.transpose(np.mat(np.exp(-20*(x-0.5)**2)-np.exp(-20*(x-1.5)**2)-np.exp(-20*(x+0.5)**2))|
31
    # Second-Derivative Matrix
    data=np.ones((3,N))
33
    data[1] = -2*data[1]
34
   diags=[-1,0,1]
   D2=sparse.spdiags(data,diags,N,N)/(h**2)
36
    # Identity Matrix
    I=sparse.identity(N)
38
    # Data for each time-step
   data=[]
40
    for i in range(NumOfTimeSteps):
41
        \# Solve the System: (I-k/2*D2) u_new=(I+k/2*D2)*u_old
42
        A = (I - k / 2 * D2)
43
        b = (I + k/2*D2)*u
44
        u=np.transpose(np.mat(sparse.linalg.spsolve(A,b)))
45
        # Save Data
46
        data.append(u)
47
    # Define of the Create Movie Function
48
    def CreateMovie(plotter, numberOfFrames, fps=10):
49
        import os, sys
50
        import matplotlib.pyplot as plt
51
52
        for i in range(numberOfFrames):
53
            plotter(i)
            fname='_tmp%05d.png'%i
```

```
55
56
            plt.savefig(fname)
57
            plt.clf()
58
    # Define the Frame Speed and Movie Length
59
    FPS=60
60
   MovieLength=10
61
    # Function to plot any given Frame
62
    def plotFunction(frame):
63
        plt.plot(x,data[int(NumOfTimeSteps*frame/(FPS*MovieLength))])
64
        plt.axis((0,1,0,1.0))
65
    # Generate the movie
66
    CreateMovie(plotFunction,int(MovieLength*FPS),FPS)
```

#### B. 一维扩散方程动画

Here is the program to merge png pictures together to gif animation in Python programming language.

## Input Python source:

```
import matplotlib.pyplot as plt
import os,imageio
images = []
filenames=sorted((fn for fn in os.listdir('.') if fn.endswith('.png')))
for filename in filenames:
    images.append(imageio.imread(filename))
imageio.mimsave('1dDiffusion.gif', images,duration=1/60)
```

## C. 求解二维扩散方程

Here is the program to two dimensional diffusion equation in Python programming language. **Input Python source:** 

```
1
    import scipy as sp
   import time
   from mpl_toolkits.mplot3d import Axes3D
    from matplotlib import cm
    from matplotlib.ticker import LinearLocator, FormatStrFormatter
    import matplotlib.pyplot as plt
7
    import mpl_toolkits.mplot3d.axes3d as p3
    import matplotlib.animation as animation
9
10
    dx = 0.01
11
    dy = 0.01
12
    a = 0.5
13
    timesteps=500
14
    t = 0.0
15
16
    nx = int(1/dx)
17
   ny = int(1/dy)
18
19
    dx2=dx**2
20
    dy2=dy**2
21
22
   dt = dx2*dy2/(2*a*(dx2+dy2))
23
24
    ui = sp.zeros([nx,ny])
25
    u = sp.zeros([nx,ny])
26
   for i in range(nx):
```

```
28
        for j in range(ny):
29
            if (((i*dx-0.5)**2+(j*dy-0.5)**2 <= 0.1)
30
                &((i*dx-0.5)**2+(j*dy-0.5)**2>=0.05)):
31
                    ui[i,j] = 1
32
    def evolve_ts(u, ui):
33
        u[1:-1,1:-1] = ui[1:-1,1:-1] + a*dt*(
34
                         (ui[2:,1:-1]-2*ui[1:-1,1:-1]+ui[:-2,1:-1])/dx2+
35
                         (ui[1:-1,2:]-2*ui[1:-1,1:-1]+ui[1:-1,:-2])/dy2)
36
37
    def data_gen(framenumber, Z, surf):
38
        global u
39
        global ui
40
        evolve_ts(u,ui)
41
        ui[:]=u[:]
42
        Z=ui
43
44
        ax.clear()
45
        plotset()
46
        surf=ax.plot_surface(X,Y,Z,rstride=1,cstride=1,cmap=cm.coolwarm,
47
                            linewidth=0, antialiased=False, alpha=0.7)
48
        return surf,
49
50
    fig=plt.figure()
51
    ax=fig.add_subplot(111, projection='3d')
52
53
    X = sp.arange(0,1,dx)
54
   Y = sp.arange(0,1,dy)
55
   X, Y = sp.meshgrid(X, Y)
56
57
    Z = ui
58
59
    def plotset():
60
        ax.set_xlim3d(0.0,1.0)
61
        ax.set_ylim3d(0.0,1.0)
62
        ax.set_zlim3d(-1.0,1.0)
63
        ax.set_autoscalez_on(False)
64
        ax.zaxis.set_major_locator(LinearLocator(10))
65
        ax.zaxis.set_major_formatter(FormatStrFormatter('%.02f'))
        cset=ax.contour(X,Y,Z,zdir='x',offset=0.0,cmap=cm.coolwarm)
66
67
        cset=ax.contour(X,Y,Z,zdir='y',offset=1.0,cmap=cm.coolwarm)
68
        cset=ax.contour(X,Y,Z,zdir='z',offset=-1.0,cmap=cm.coolwarm)
69
70
    plotset()
71
    surf=ax.plot_surface(X,Y,Z,rstride=1,cstride=1,cmap=cm.coolwarm,
72
                            linewidth=0, antialiased=False, alpha=0.7)
73
74
75
    fig.colorbar(surf,shrink=0.5,aspect=5)
76
77
    ani=animation.FuncAnimation(fig,data_gen,fargs=(Z,surf),frames=1000,interval=30,blit=False)
    ani.save("2dDiffusion.mp4", bitrate=1024)
```

# 参考文献

- [1] Tuncer Cebeci. Convective Heat Transfer. Springer. 2002. ISBN 0-9668461-4-1.
- [2] Crank, J.; Nicolson, P. A practical method for numerical evaluation of solutions of partial differential equations of the heat conduction type. Proc. Camb. Phil. Soc. 1947, 43 (1): 50–67. doi:10.1007/BF02127704.
- [3] Thomas, J. W. Numerical Partial Differential Equations: Finite Difference Methods. Texts in Applied Mathematics 22. Berlin, New York: Springer-Verlag. 1995. ISBN 978-0-387-97999-1.. Example 3.3.2 shows that Crank–Nicolson is unconditionally stable when applied to  $u_t = au_{xx}$
- [4] Wilmott, P.; Howison, S.; Dewynne, J. The Mathematics of Financial Derivatives: A Student Introduction. Cambridge Univ. Press. 1995. ISBN 0-521-49789-2.