

EP 307 Assignment 2

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Problem 1

$$\begin{aligned}\langle e_1 | e_1 \rangle &= \langle e_2 | e_2 \rangle = 2 \\ \langle e_3 | e_3 \rangle &= 4 \\ \langle e_1 | e_2 \rangle &= i\sqrt{2} \\ \langle e_1 | e_3 \rangle &= 1 + i \\ \langle e_2 | e_3 \rangle &= 2\end{aligned}$$

First, we normalize the states to get:

$$\begin{aligned}\langle a_1 | a_1 \rangle &= \langle a_2 | a_2 \rangle = 1 \\ \langle e_3 | a_3 \rangle &= 1 \\ \langle a_1 | a_2 \rangle &= i\frac{1}{\sqrt{2}} \\ \langle a_1 | a_3 \rangle &= \frac{1+i}{2\sqrt{2}} \\ \langle a_2 | a_3 \rangle &= \frac{1}{\sqrt{2}}\end{aligned}$$

Taking $\hat{p}_{a_i} = |a_i\rangle \langle a_i|$ to be the projection operator with respect to a_i , by Gram-Schmidt orthogonalization procedure,

$$\begin{aligned}
|\phi_1\rangle &= |a_1\rangle \\
|\phi_2\rangle &= |a_2\rangle - \hat{p}_{a_1} |a_2\rangle \\
|\phi_3\rangle &= |a_3\rangle - \hat{p}_{a_1} |a_3\rangle - \hat{p}_{a_2} |a_3\rangle
\end{aligned}$$

Thus

$$\begin{aligned}
|\phi_1\rangle &= |a_1\rangle \\
|\phi_2\rangle &= |a_2\rangle - |a_1\rangle \langle a_1 | a_2 \rangle \\
|\phi_3\rangle &= |a_3\rangle - |a_1\rangle \langle a_1 | a_3 \rangle - |a_2\rangle \langle a_2 | a_3 \rangle
\end{aligned}$$

and

$$\begin{aligned}
|\phi_1\rangle &= |a_1\rangle \\
|\phi_2\rangle &= |a_2\rangle + \frac{1}{\sqrt{2}i} |a_1\rangle \\
|\phi_3\rangle &= |a_3\rangle - \frac{1+i}{2\sqrt{2}} |a_1\rangle - \frac{1}{\sqrt{2}} |a_2\rangle
\end{aligned}$$

We can calculate

$$\begin{aligned}
\langle \phi_1 | \phi_1 \rangle &= 1 \\
\langle \phi_2 | \phi_2 \rangle &= 1 - \frac{1}{2} = \frac{1}{2} \\
\langle \phi_3 | \phi_3 \rangle &= 1 + \frac{(1-i)(1+i)}{8} + \frac{1}{2} = \frac{7}{4}
\end{aligned}$$

and from here we get an orthonormal basis

$$\begin{aligned}
|\psi_1\rangle &= |a_1\rangle \\
|\psi_2\rangle &= \sqrt{2} \left(|a_2\rangle + \frac{1}{\sqrt{2}i} |a_1\rangle \right) \\
|\psi_3\rangle &= \frac{2}{\sqrt{7}} \left(|a_3\rangle - \frac{1+i}{2\sqrt{2}} |a_1\rangle - \frac{1}{\sqrt{2}} |a_2\rangle \right)
\end{aligned}$$

Rewriting in terms of the original vectors

$$\begin{aligned} |\psi_1\rangle &= \frac{1}{\sqrt{2}} |e_1\rangle \\ |\psi_2\rangle &= |e_2\rangle + \frac{1}{\sqrt{2}i} |e_1\rangle \\ |\psi_3\rangle &= \frac{4}{\sqrt{7}} \left(|e_3\rangle - \frac{1+i}{2\sqrt{2}} |e_1\rangle - \frac{1}{\sqrt{2}} |e_2\rangle \right) \end{aligned}$$

$|\psi_1\rangle, |\psi_2\rangle, |\psi_3\rangle$ are our orthonormal basis vectors.

Problem 2

First, we normalize them by taking integrals in $[-1, 1]$, and get $c_0 = \frac{1}{\sqrt{2}}, c_1 =$

$$\sqrt{\frac{3}{2}}, c_2 = \sqrt{\frac{5}{2}}$$

Applying the orthogonalization process,

$$\begin{aligned} |\phi_1\rangle &= |\psi_1\rangle \\ |\phi_2\rangle &= |\psi_2\rangle - |\psi_1\rangle \langle \psi_1 | \psi_2 \rangle \\ |\phi_3\rangle &= |\psi_3\rangle - |\psi_1\rangle \langle \psi_1 | \psi_3 \rangle - |\psi_2\rangle \langle \psi_2 | \psi_3 \rangle \end{aligned}$$

calculating the inner products, we get

$$\begin{aligned} |\phi_1\rangle &= |\psi_1\rangle \\ |\phi_2\rangle &= |\psi_2\rangle - 0 |\psi_1\rangle \\ |\phi_3\rangle &= |\psi_3\rangle - \frac{\sqrt{5}}{2} \frac{2}{3} |\psi_1\rangle - 0 |\psi_2\rangle \end{aligned}$$

which evaluates to

$$\begin{aligned}\phi_1 &= \frac{1}{\sqrt{2}} \\ \phi_2 &= \sqrt{\frac{3}{2}}x \\ \phi_3 &= \sqrt{\frac{5}{2}}\left(x^2 - \frac{1}{3}\right)\end{aligned}$$

Problem 3

$$\hat{B}g(x) = g(-x)$$

The operator is linear, since $\hat{B}(\alpha f(x) + \beta g(x)) = \alpha f(-x) + \beta g(-x) = \alpha \hat{B}f(x) + \beta \hat{B}g(x)$

It also is Hermitian, as

$$\begin{aligned}\langle f | (\hat{B} |g\rangle) &= \int_{-\infty}^{\infty} \bar{f}(x) \hat{B}g(x) dx \\ &= \int_{-\infty}^{\infty} \bar{f}(x) g(-x) dx \\ &= \int_{-\infty}^{\infty} \bar{f}(-x) g(x) dx \\ &= \int_{-\infty}^{\infty} \bar{\hat{B}f(x)} g(x) dx \\ &= \int_{-\infty}^{\infty} \overline{\hat{B}f(x)} g(x) dx \\ &= \langle \hat{B}f | g \rangle \\ \therefore \langle f | \hat{B}g \rangle &= \langle \hat{B}f | g \rangle\end{aligned}$$

To find eigenvalues, $\hat{B}g(x) = bg(x) = g(-x)$

Since $bg(x) = g(-x)$ and $b g(-x) = g(x)$, $b = \pm 1$ (neglecting the trivial $b = 0$ solution).

Thus, we have eigenvectors:

$$\begin{aligned} \sum_n a_n x^{2n} & \quad (\text{even function}) & \quad \text{for eigenvalue } b = 1 \\ \sum_n b_n x^{2n+1} & \quad (\text{odd function}) & \quad \text{for eigenvalue } b = -1 \end{aligned}$$

To calculate the commutator,

$$\begin{aligned} [B, \hat{x}^n]f &= \hat{B}\hat{x}^n f(x) - \hat{x}^n \hat{B}f(x) \\ &= \hat{B}x^n f(x) - \hat{x}^n f(-x) \\ &= (-x)^n f(-x) - x^n f(-x) \\ &= f(-x)((-x)^n - x^n) \end{aligned}$$

For this to be zero for all states f , n must be even. Thus the commutator is zero only for even n

Problem 4

$$\hat{\Omega} = |\psi\rangle \langle \phi|$$

For it to be Hermitian, $\hat{\Omega}^\dagger = \Omega$. Applying it to the state $|\Psi\rangle = |\psi\rangle + |\phi\rangle$,

$$\begin{aligned} |\psi\rangle \langle \phi| \Psi &= |\phi\rangle \langle \psi| \Psi \\ |\psi\rangle \langle \phi| \psi + |\phi\rangle \langle \psi| \psi &= |\psi\rangle \langle \phi| \phi + |\phi\rangle \langle \psi| \phi \\ (\langle \phi| \psi\rangle - \langle \phi| \phi\rangle) |\psi\rangle &= (\langle \psi| \phi\rangle - \langle \psi| \psi\rangle) |\phi\rangle \end{aligned}$$

For this to hold, either $|\psi\rangle = c|\phi\rangle$ or $(\langle \phi| \psi\rangle - \langle \phi| \phi\rangle) = 0 = (\langle \psi| \phi\rangle - \langle \psi| \psi\rangle)$

In the latter case, we get $\langle \psi| \psi\rangle = \langle \phi| \psi\rangle = \langle \phi| \phi\rangle$, which can only happen if the two states are equivalent.

Thus condition for the operator to be hermitian is that $|\psi\rangle = c|\phi\rangle$.

For it to be a projection operator, $c = \frac{1}{\langle \phi| \phi\rangle}$

1 Problem 5

$$\hat{B}\psi(x) = \int_{-\infty}^x x' \psi(x') dx'$$

The eigenvalue problem $\hat{B}\psi = \lambda\psi$ thus becomes $\int_{-\infty}^x x'\psi(x')dx' = \lambda\psi(x)$. Differentiating with respect to x , this is

$$\begin{aligned}
 x\psi(x) &= \lambda\psi'(x) \\
 \implies \frac{\psi'(x)}{\psi(x)} &= \frac{x}{\lambda} \\
 \implies \int_{\psi(0)}^{\psi} \frac{\psi'(x)}{\psi(x)} &= \int_0^x \frac{x}{\lambda} \\
 \implies \log \frac{\psi(x)}{\psi(0)} &= \frac{1}{\lambda} \frac{x^2}{2} \\
 \implies \psi(x) &= \psi(0)e^{\frac{x^2}{2\lambda}}
 \end{aligned}$$

This is an acceptable (square-integrable) wavefunction only when λ is negative.