

Tutorial 2

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Problem 1

- (i) $\frac{dx^\mu}{d\tau}$ is a tensor for all kinds of transformations. $\frac{dx'^\mu}{d\tau} = \frac{d}{d\tau} \left(x^\nu \frac{dx'^\mu}{dx^\nu} \right) = \frac{dx'^\mu}{dx^\nu} \frac{dx^\nu}{d\tau}$, and thus we get a contravariant tensor
- (ii) $\frac{d^2x^\mu}{d\tau^2}$ is also a tensor because it is the derivative of a tensor (Use the property proved in (i) by replacing x^μ with $\frac{dx^\mu}{d\tau}$). This tensor is also contravariant.
- (iii) $\partial_\mu \phi$ is a tensor. $\partial'_\mu \phi = \frac{\partial \phi}{\partial x'_\nu \frac{dx^\mu}{dx^\nu}} = \frac{dx_\nu}{dx'_\mu} \frac{\partial \phi}{\partial x_\nu} = \frac{dx_\nu}{dx'_\mu} \partial_\nu \phi$, making it a covariant tensor.
- (iv) This is only a tensor when special relativistic constraints are present. We know that the covariant derivative of the tensor F , $F_{\mu\nu;\rho} = \partial_\rho F_{\mu\nu} - \Gamma_{\nu\rho}^\sigma F_{\mu\sigma} - \Gamma_{\mu\rho}^\sigma F_{\nu\sigma}$ is a tensor itself. However, $\Gamma_{\nu\rho}^\sigma F_{\mu\sigma}$, $\Gamma_{\mu\rho}^\sigma F_{\nu\sigma}$ are not tensors in the general case as proved in class, so $\partial_\rho F_{\mu\nu}$ is not. However, in special relativity the Christoffel symbols are zero, so the remaining portion of the equation must be a tensor too.

Problem 2

We first calculate the velocity and then γ_v of the particle in the frame:

$$\begin{aligned}
v^2 &= \dot{x}^2 + \dot{y}^2 + \dot{z}^2 \\
&= (a + b\omega \cos \omega t)^2 + (-b\omega \sin \omega t)^2 \\
&= a^2 + b^2\omega^2 + 2ab\omega \cos \omega t \\
\therefore \gamma_v &= \frac{1}{\sqrt{1 - \frac{a^2 + b^2\omega^2 + 2ab\omega \cos \omega t}{c^2}}}
\end{aligned}$$

From this, we get $w^\mu = \frac{1}{\sqrt{1 - \frac{a^2 + b^2\omega^2 + 2ab\omega \cos \omega t}{c^2}}} \begin{bmatrix} a + b\omega \cos \omega t \\ -b\omega \sin \omega t \\ 0 \end{bmatrix}$