

## HW 10

## Problem [1]

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(a)

$$C_{el} = \frac{\pi^2}{2} \left( \frac{k_B T}{\varepsilon_1} \right) n k_B \quad (\text{A&M 2.81})$$

$$= \frac{\pi^2}{2} \left( \frac{1.38 \times 10^{-23} \times 4}{3.24 \times 1.6 \times 10^{-19}} \right) \times \left( 2.65 \times 10^{22} \#/\text{cm}^3 \right) \times 1.38 \times 10^{-23} \text{ J/K}$$

$$= 1.92 \times 10^{-4} \text{ J/K.cm}^3$$

$$(N_a, \varepsilon_1 = 3.24 \text{ eV})$$

$$(b) C_{ph} = 234 \left( \frac{T}{\Theta_D} \right)^3 n k_B$$

$$= 234 \times \left( \frac{4k}{150k} \right)^3 \times (2.65 \times 10^{22} \#/\text{cm}^3) \times 1.38 \times 10^{-23} \text{ J/K}$$

$$= 1.62 \times 10^{-3} \text{ J/K.cm}^3$$

(c)

$$\text{When } C_{el} = C_{ph} \Rightarrow \frac{\pi^2}{2} \left( \frac{k_B T}{\varepsilon_1} \right) n k_B = 234 \left( \frac{T}{\Theta_D} \right)^3 n k_B$$

$$\Rightarrow \underbrace{\frac{\pi^2}{2}}_{234} \underbrace{\frac{k_B \Theta_D^3}{\varepsilon_1}}_{= 1} = T = 1.38 \text{ K}$$

$$\text{Ans: } 1.38 \text{ K}$$

## Problem [2]

Total energy of electrons is

$$u = \int_{-\infty}^{\infty} \varepsilon g(\varepsilon) f(\varepsilon) d\varepsilon$$

$$\approx U_0 + \int_{\varepsilon_-}^{\varepsilon + k_B T} (\varepsilon - \varepsilon_F) g(\varepsilon) f(\varepsilon) d\varepsilon$$

$$= U_0 + \int_{\varepsilon_- - k_B T}^{\varepsilon + k_B T} \alpha |\varepsilon - \varepsilon_F| e^{-\frac{\varepsilon - \mu}{k_B T}} + 1$$

$$(g(\varepsilon) \propto |\varepsilon - \varepsilon_F| := \alpha |\varepsilon - \varepsilon_F|)$$

$$\approx U_0 + \alpha (k_B T)^3 \int_{-1}^1 \frac{s |s| ds}{e^s + 1} \quad s := \frac{\varepsilon - \varepsilon_F}{k_B T}$$

$$:= U_0 + AT^3$$

$$(A := \text{const.})$$

$\mu \approx \varepsilon_F$  at low temp.

$$\therefore C_{el} := \frac{\partial u}{\partial T} \propto AT^2 \propto T^2$$

### Problem [3]

Total carrier density is fixed:

$$n = \int_0^\infty g(\varepsilon) f(\varepsilon) d\varepsilon$$

$$= \int_0^\infty \frac{m}{\pi \hbar^2} \frac{1}{e^{\frac{\varepsilon - \mu}{k_B T}} + 1} d\varepsilon$$

$$= \frac{m k_B T}{\pi \hbar^2} \int_{-\frac{\mu}{k_B T}}^\infty \frac{ds}{e^s + 1} \quad s := \frac{\varepsilon - \mu}{k_B T}$$

$$= \frac{m k_B T}{\pi \hbar^2} \int_{-\frac{\mu}{k_B T}}^\infty \frac{e^{-s} ds}{e^{-s} + 1}$$

$$= \frac{m k_B T}{\pi \hbar^2} \int_{-\frac{\mu}{k_B T}}^\infty \frac{e^{-s} ds}{e^{-s} + 1}$$

$$= \frac{m k_B T}{\pi \hbar^2} \ln \left| e^{\frac{\mu}{k_B T}} + 1 \right|$$

$$\Rightarrow \mu = k_B T \ln \left| e^{\frac{n \pi \hbar^2}{m k_B T}} - 1 \right|$$

$$\mu(T \rightarrow 0) = \lim_{T \rightarrow 0} k_B T \ln \left| e^{\frac{n \pi \hbar^2}{m k_B T}} - 1 \right|$$

$$= \lim_{T \rightarrow 0} k_B T \cdot \frac{n \pi \hbar^2}{m k_B T} = \frac{n \pi \hbar^2}{m} = \varepsilon_F$$

Problem [4] A&M 23.2

$$\begin{aligned}
 a) g(\omega) &= \sum_s \int \frac{d\vec{k}}{(2\pi)^3} g(\omega - \omega_s(\vec{k})) \\
 &= \sum_s \frac{1}{3} \int \frac{d\Omega}{4\pi} \int d\vec{k} \frac{4\pi k^2}{(2\pi)^3} \cdot 3 \cdot g(\omega - \omega_s(\vec{k})) \\
 &= \sum_s \frac{1}{3} \int \frac{d\Omega}{4\pi} \int d\vec{k} \frac{3k^2}{2\pi} g(\omega - \omega_s(\vec{k}))
 \end{aligned}$$

For low-freq, the integrand of  $d\vec{k}$  can be taken approximately to  $k_D \sqrt{\Omega}$

$$\begin{aligned}
 \therefore g(\omega) &\approx \sum_s \frac{1}{3} \int \frac{d\Omega}{4\pi} \int_0^{k_D} dk \frac{3k^2}{2\pi^2} g(\omega - k c_s(k)) \\
 (\text{Where } \omega_s(\vec{k}) &:= k \cdot c_s(k)) \\
 &= \sum_s \frac{1}{3} \int \frac{d\Omega}{4\pi} \frac{3}{2\pi^2} \cdot \left( \frac{\omega}{c_s(k)^3} \right) \\
 &= \frac{3}{2\pi^2} \omega^2 \sum_s \frac{1}{3} \int \frac{d\Omega}{4\pi} \frac{1}{c_s(k)^3} \\
 &= \frac{3}{2\pi^2} \frac{\omega^2}{C^3}
 \end{aligned}$$

$$\text{Where } \frac{1}{C^3} = \sum_s \frac{1}{3} \int \frac{d\Omega}{4\pi} \frac{1}{c_s(k)^3} \quad (\text{A&M 23.18})$$

$\therefore$  Gives the same form of result as (23.36)

(b) In  $d$ -dim crystal,

$$\int \frac{d\vec{k}}{(2\pi)^d} = \int \frac{d\Omega}{\alpha(d)} \int \frac{\alpha(d) k^{d-1} dk}{(2\pi)^d}$$

For low-freq, ( $\alpha(d)$  is some constant)

$$\begin{aligned} g(\omega) &\approx \sum_S \frac{1}{3} \int \frac{d\Omega}{\alpha(d)} \int_0^\infty \frac{k_B 3\alpha(d)}{(2\pi)^d} k^{d-1} g(\omega - k c_s(\vec{k})) \\ &= \sum_S \frac{1}{3} \int \frac{d\Omega}{\alpha(d)} \frac{3\alpha(d)}{(2\pi)^d} \cdot \frac{\omega^{d-1}}{c_S(\vec{k})^d} \\ &= \frac{3\alpha(d)}{(2\pi)^d} \cdot \frac{\omega^{d-1}}{C^d} \propto \omega^{d-1} \end{aligned}$$

( $C$  is the average phase velocity)

(c)

$$u \approx \int_0^\infty dw \cdot \frac{\hbar w g(w)}{e^{\frac{\hbar w}{k_B T}} - 1}$$

$$= A \int_0^\infty dw \frac{w^d}{e^{\frac{\hbar w}{k_B T}} - 1} \quad (\because g(w) \propto w^{d-1})$$

$$= \left(\frac{k_B T}{\hbar}\right)^{d+1} A \int_0^\infty \frac{x^d}{e^x - 1} \quad (A, B \text{ are just some constants})$$

$$= B T^{d+1}$$

$$\Rightarrow C_V = \left(\frac{\partial u}{\partial T}\right)_V = (d+1) B T^d \propto T^d$$

(d)

Assume that  $w = a k^v$

Then we can reformulate the  $g(w)$  as:

$$\begin{aligned} g(w) &= \frac{3}{(2\pi)^d} \int_0^{k_D} \frac{\alpha(d) k^{\frac{d-1}{2}} dk}{(2\pi)^d} \cdot \delta(w - a k^v) \\ &= \frac{3\alpha(d)}{(2\pi)^d} \int_0^{k_D} k^{\frac{d-1}{2}} dk \cdot \delta(w - a k^v) \\ &= \frac{3\alpha(d)}{(2\pi)^d} \left[ \left( \frac{w}{a} \right)^{\frac{1}{v}} \right]^{\frac{d-1}{2}} \frac{1}{(av)} \left( \frac{w}{a} \right)^{\frac{v-1}{v}} \\ &= \frac{3\alpha(d)}{(2\pi)^d} \frac{1}{av} \cdot \left( \frac{w}{a} \right)^{\frac{d-v}{v}} \mathcal{L} w^{\frac{1}{v}-1} \end{aligned}$$

$$\therefore C_V = \frac{\partial}{\partial T} \Big|_V \int_0^\infty \frac{g(w) \frac{\hbar w}{k_B T} dw}{e^{\frac{\hbar w}{k_B T}} - 1}$$

$$\begin{aligned} &= \frac{\partial}{\partial T} \Big|_V A \int_0^\infty \frac{w^{\frac{d}{v}} dw}{e^{\frac{\hbar w}{k_B T}} - 1} \\ &= \frac{\partial}{\partial T} \Big|_V A \left( \frac{k_B T}{\hbar} \right)^{\frac{d}{v}+1} \int_0^\infty \frac{x^{\frac{d}{v}} dx}{e^x - 1} \\ &\mathcal{L} \frac{1}{T^{\frac{d}{v}}} \end{aligned}$$