

Notes

- **Assignments**
 - ◆ 3.2
 - ◆ 3.27
 - ◆ 3.36
 - ◆ 3.38
 - ◆ 3.50
- **Tutorial problems**
 - ◆ Basic Problems with Answers 3.11
 - ◆ Basic Problems 3.30, 3.37
 - ◆ Advanced Problems 3.49

CT Fourier Series Pairs

$$x(t) = \sum_{k=-\infty}^{+\infty} a_k e^{jk\omega_0 t} = \sum_{k=-\infty}^{+\infty} a_k e^{j2\pi kt/T}$$

Harmonically related

$$a_k = \frac{1}{T} \int_T x(t) e^{-jk\omega_0 t} dt$$

Periodicity Properties of DT Complex Exponentials

- For DT complex exponentials, signals are periodic only when $\omega_0 N = k \cdot 2\pi$, $k = 0, \pm 1, \pm 2, \dots$

$$e^{j\omega_0 n} = e^{j\omega_0 (n+N)} \rightarrow e^{j\omega_0 N} = 1 \rightarrow \omega_0 N = k \cdot 2\pi$$

- For DT complex exponentials, signals with frequencies ω_0 and $\omega_0 + k \cdot 2\pi$ are identical.
$$e^{j(\omega_0 + k \cdot 2\pi)n} = e^{j\omega_0 n} \cdot e^{jk \cdot 2\pi n} = e^{j\omega_0 n}$$
 - We need only consider a frequency interval of length 2π , and in most cases, we use the interval: $0 \leq \omega_0 < 2\pi$, or $-\pi \leq \omega_0 < \pi$

- $e^{j\omega_0 n}$ does ***not*** have a continually increasing rate of oscillation as ω_0 is increased in magnitude.

low-frequency (slowly varying): ω_0 near $0, 2\pi, \dots$, or $2k \cdot \pi$

high-frequency (rapid variation): ω_0 near $\pm \pi, \dots$, or $(2k+1) \cdot \pi$

$$e^{j(2k+1)\pi n} = e^{j\pi n} = (e^{j\pi})^n = (-1)^n$$

DT Fourier Series Representation

Arbitrary periodic DT signal with period N can be written as

$$x[n] = \sum_{k=\langle N \rangle} a_k e^{jk(2\pi/N)n}$$

$\sum_{k=\langle N \rangle}$ = Sum over *any* N consecutive values of k

— This is a *finite* series

$\{a_k\}$ - Fourier (series) coefficients

Frequency component: $\frac{2k\pi}{N}$ $k = 0, 1, 2, \dots, N-1$ or $1, 2, \dots, N$

Why?

Existence

Any DT periodic signal has a Fourier series representation

$$x[n] = \sum_{k=\langle N \rangle} a_k e^{jk\omega_0 n} \quad \forall n$$

\Downarrow

$n=0$

$$x[0] = \sum_{k=\langle N \rangle} a_k$$

$n=1$

$$x[1] = \sum_{k=\langle N \rangle} a_k e^{jk\omega_0}$$

$n=2$

$$x[2] = \sum_{k=\langle N \rangle} a_k e^{j2k\omega_0}$$

\vdots

\vdots

$n=N-1$

$$x[N-1] = \sum_{k=\langle N \rangle} a_k e^{j(N-1)k\omega_0}$$

N equations for N unknowns, a_0, a_1, \dots, a_{N-1}

How to calculate a_k

- Define inner product as

$$\langle e^{jk\omega_0 n} \cdot e^{jm\omega_0 n} \rangle = \frac{1}{N} \sum_{n=0}^{N-1} e^{jk\omega_0 n} e^{-jm\omega_0 n}$$

- We have

$$\begin{aligned} \langle e^{jk\omega_0 n} \cdot e^{jm\omega_0 n} \rangle &= 1 \quad (k = m + Nk') \\ \langle e^{jk\omega_0 n} \cdot e^{jm\omega_0 n} \rangle &= 0 \quad (\textit{Otherwise}) \end{aligned}$$

- $\{e^{jk\omega_0 n} | k = \langle N \rangle\}$ is similar to basis of vector space

• So

$$\begin{aligned} \langle x[n] \cdot e^{jk\omega_0 n} \rangle &= \langle \sum_{m=0}^{N-1} a_m e^{jm\omega_0 n} \cdot e^{jk\omega_0 n} \rangle \\ &= \sum_{m=0}^{N-1} a_m \langle e^{jm\omega_0 n} \cdot e^{jk\omega_0 n} \rangle = a_k \end{aligned}$$

• Hence,

$$a_k = \langle x[n] \cdot e^{jk\omega_0 n} \rangle = \frac{1}{N} \sum_{n=0}^{N-1} x[n] e^{-jk\omega_0 n}$$

DT Fourier Series Pair $\left(\omega_o = \frac{2\pi}{N}\right)$

$$x[n] = \sum_{k=\langle N \rangle} a_k e^{jk\omega_o n} \quad (\text{Synthesis equation})$$

$$a_k = \frac{1}{N} \sum_{n=\langle N \rangle} x[n] e^{-jk\omega_o n} \quad (\text{Analysis equation})$$

Different
from CT
Fourier
series

Cont.

$$x[n] = \sum_{k=\langle N \rangle} a_k e^{jk\omega_0 n} \quad (\text{Synthesis equation})$$

- a_k can be defined for all integers k , and we have $a_{k+N} = a_k$

$$x[n] = a_0 e^{\frac{j0 \times 2\pi}{N}n} + a_1 e^{\frac{j1 \times 2\pi}{N}n} + \dots + a_{N-1} e^{\frac{j(N-1) \times 2\pi}{N}n}$$

$$x[n] = a_1 e^{\frac{j1 \times 2\pi}{N}n} + \dots + a_{N-1} e^{\frac{j(N-1) \times 2\pi}{N}n} + a_N e^{\frac{jN \times 2\pi}{N}n}$$

◆ a_k is periodic w.r.t. k

◆ CT is different

Example #1: Sum of a pair of sinusoids

$$x[n] = \cos(\pi n / 8) + \cos(\pi n / 4 + \pi / 4)$$

— periodic with period $N = ?$

$$x[n] = \frac{1}{2} [e^{j\omega_0 n} + e^{-j\omega_0 n}] + \frac{1}{2} [e^{j\pi/4} e^{j2\omega_0 n} + e^{-j\pi/4} e^{-j2\omega_0 n}]$$

↓

$$a_0 = 0$$

$$a_{15} = a_{-1+16} = a_{-1} = 1/2$$

$$a_1 = 1/2$$

$$a_{-1} = 1/2$$

$$a_{66} = a_{2+4 \times 16} = a_2 = e^{j\pi/4}/2$$

$$a_2 = e^{j\pi/4}/2$$

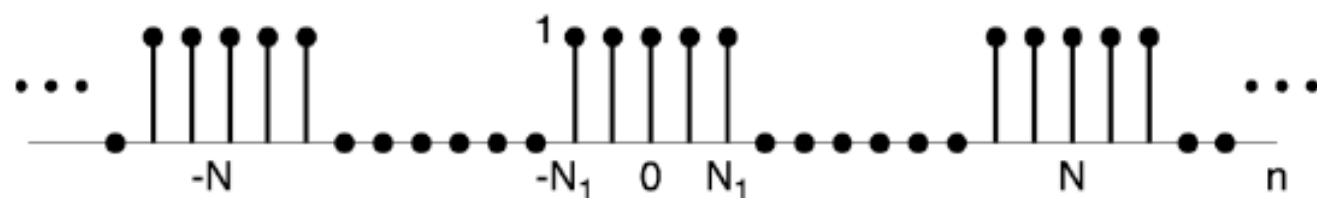
$$a_{-2} = e^{-j\pi/4}/2$$

$$a_3 = 0$$

$$a_{-3} = 0$$

⋮

$$\begin{aligned} \cos(x) &= \operatorname{Re}(e^{jx}) = \frac{1}{2}(e^{jx} + e^{-jx}) \\ \sin(x) &= \operatorname{Im}(e^{jx}) = \frac{1}{2j}(e^{jx} - e^{-jx}) \end{aligned}$$



Period=?

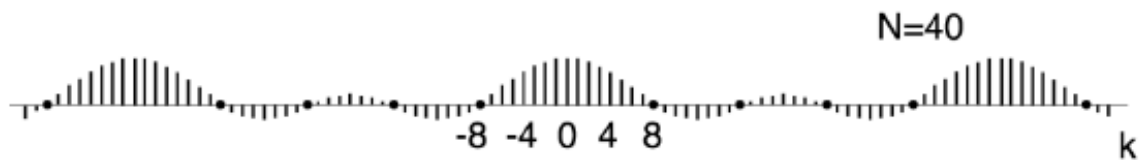
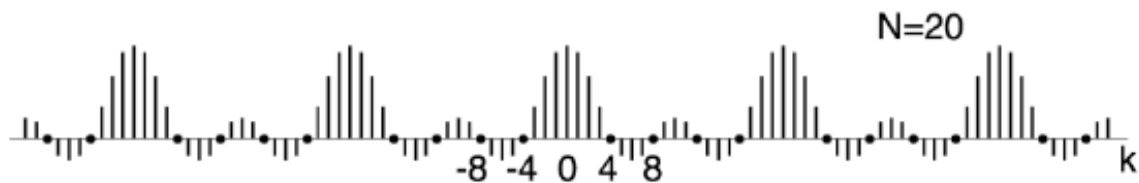
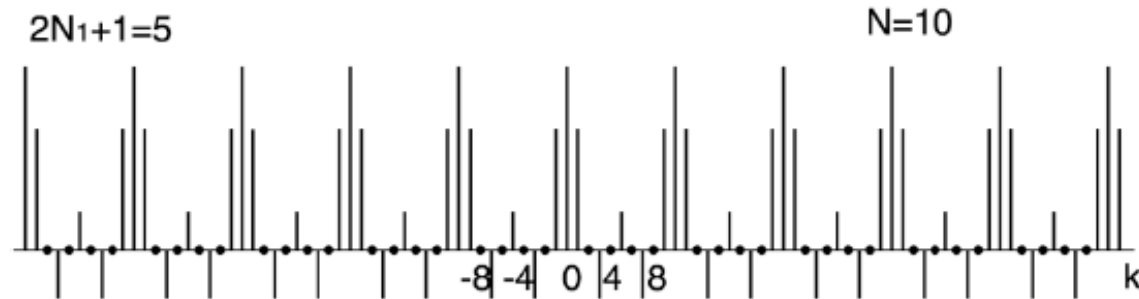
$$a_0 = \frac{1}{N} \sum_{n=-N_1}^{N_1} x[n] = \frac{(2N_1 + 1)}{N} = a_N = a_{-N} = a_{6N} = \dots$$

For $k \neq$ multiple of N :

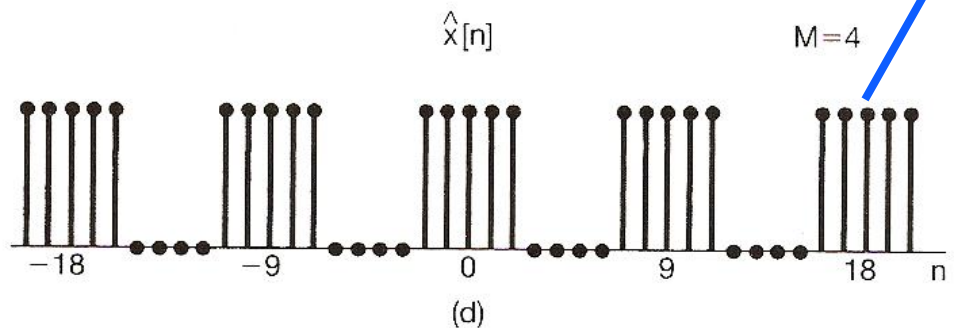
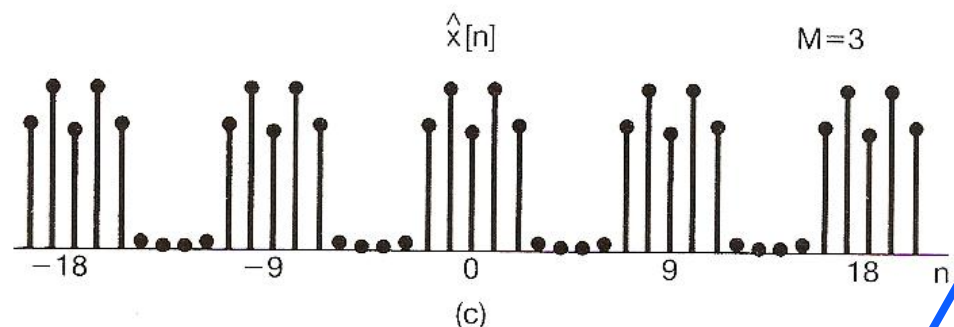
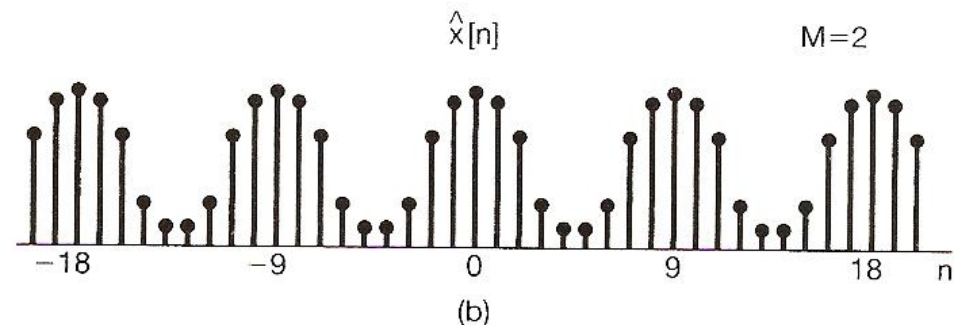
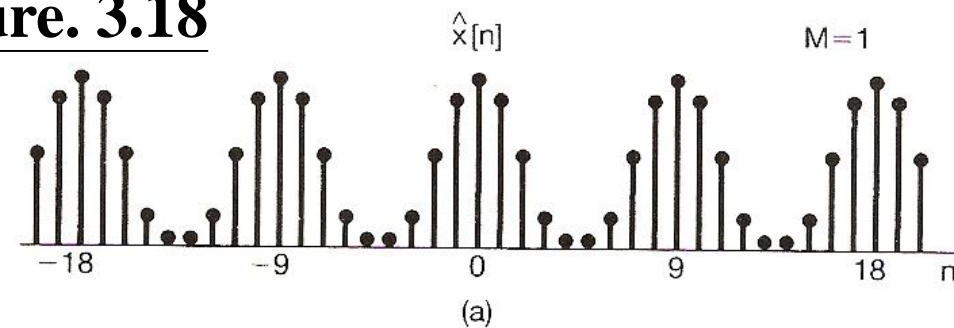
$$\begin{aligned} a_k &= \frac{1}{N} \sum_{n=-N_1}^{N_1} e^{-jk\omega_0 n} \stackrel{n=m-N_1}{=} \frac{1}{N} \sum_{m=0}^{2N_1} e^{-jk\omega_0 (m-N_1)} \\ &= \frac{1}{N} e^{jk\omega_0 N_1} \sum_{m=0}^{2N_1} (e^{-jk\omega_0})^m = \frac{1}{N} e^{jk\omega_0 N_1} \frac{1 - e^{-jk\omega_0 (2N_1 + 1)}}{1 - e^{-jk\omega_0}} \\ &= \frac{1}{N} \frac{\sin[k(N_1 + 1/2)\omega_0]}{\sin(k\omega_0 / 2)} = \frac{1}{N} \frac{\sin[2\pi k(N_1 + 1/2) / N]}{\sin(\pi k / N)} \end{aligned}$$

DT Square wave (continued)

$$a_k = \frac{1}{N} \frac{\sin[2\pi k(N_1 + 1/2)/N]}{\sin(\pi k/N)}$$



$$N=9, 2N_1+1=5$$



- 1) The same as original DT square wave
- 2) **No** Gibbs phenomenon, and **no** discontinuity

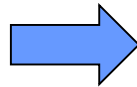
Figure 3.18 Partial sums of eqs. (3.106) and (3.107) for the periodic square wave of Figure 3.16 with $N = 9$ and $2N_1 + 1 = 5$: (a) $M = 1$; (b) $M = 2$; (c) $M = 3$; (d) $M = 4$.

DT Fourier Series - Properties

- Strong similarities between the properties of DT and CT Fourier series [Comparing Table 3.2 to Table 3.1.]

Example: $x[n] \longleftrightarrow a_k$

$e^{jM\omega_0 n} x[n] \longleftrightarrow b_k = ?$



$$x[n]e^{jM\omega_0 n} = \sum_{r=\langle N \rangle} a_r e^{jr\omega_0 n} e^{jM\omega_0 n}$$

$$= \sum_{k=\langle N \rangle}^{k=r+M} a_{k-M} e^{jk\omega_0 n}$$



$$b_k = a_{k-M}$$

Frequency shift
 $jM\omega_0 \rightarrow j(k-M)\omega_0$

Two Important Properties

- Periodic convolution:

Suppose x and y are two periodic signals with common period N , the periodic convolution between x and y is defined as

$$x[n] \circledast y[n] = \sum_{k=\langle N \rangle} x[k]y[n-k]$$

- Suppose $x[n] \rightarrow a_k$ and $y[n] \rightarrow b_k$, then

$$x[n] \circledast y[n] \rightarrow Na_k b_k \quad \text{and} \quad x[n]y[n] \rightarrow a_k \circledast b_k$$

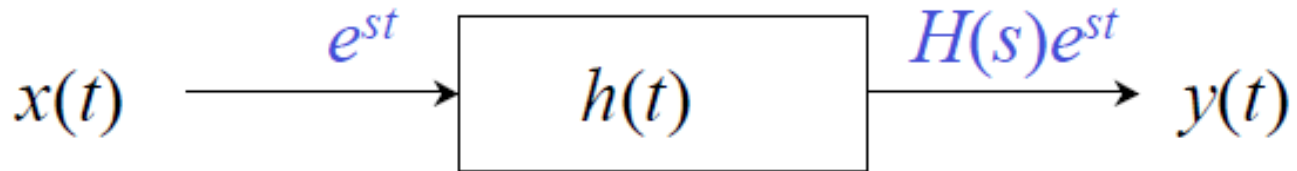
- Parseval's Relation

$$\frac{1}{N} \sum_{n=\langle N \rangle} |x[n]|^2 = \sum_{k=\langle N \rangle} |a_k|^2$$

Frequency Behavior of LTI Systems

System Functions $H(s)$ or $H(z)$

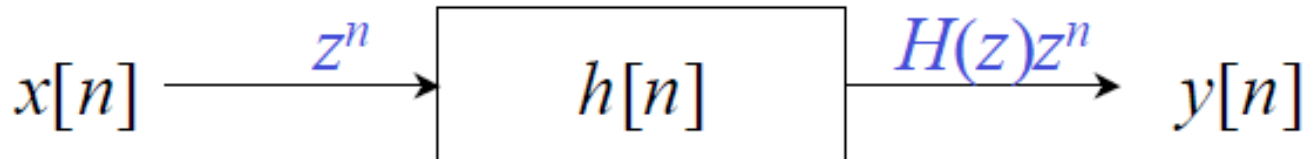
→ **CT:**



$$H(s) = \int_{-\infty}^{\infty} h(t) e^{-st} dt$$

$$x(t) = \sum a_k e^{s_k t} \longrightarrow y(t) = \sum H(s_k) a_k e^{s_k t}$$

→ **DT:**

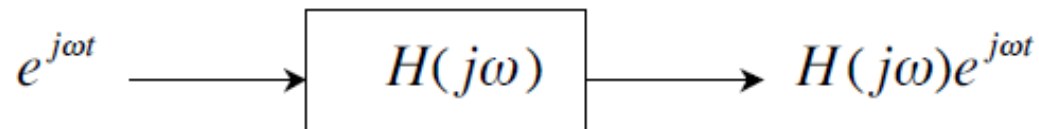


$$H(z) = \sum_{n=-\infty}^{\infty} h[n] z^{-n}$$

$$x[n] = \sum a_k z_k^n \longrightarrow y[n] = \sum H(z_k) a_k z_k^n$$

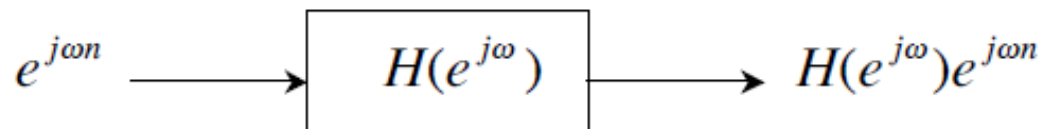
Frequency Response of an LTI System

$$(s = j\omega)$$



CT Frequency response:

$$H(j\omega) = \int_{-\infty}^{+\infty} h(t)e^{-j\omega t} dt$$



$$(z = e^{j\omega})$$

DT Frequency response:

$$H(e^{j\omega}) = \sum_{n=-\infty}^{+\infty} h[n]e^{-j\omega n}$$

Periodic

Fourier Series and LTI Systems

$$x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t} \longrightarrow \boxed{h(t)} \longrightarrow y(t) = \sum_{k=-\infty}^{\infty} H(jk\omega_0) a_k e^{jk\omega_0 t}$$

$$H(j\omega)$$

$$a_k \longrightarrow \underbrace{H(jk\omega_0) a_k}_{\text{"gain"}}$$

$$H(jk\omega_0) = |H(jk\omega_0)| e^{j\angle H(jk\omega_0)},$$

includes both amplitude & phase

$$x[n] = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 n} \longrightarrow \boxed{h[n]} \longrightarrow y[n] = \sum_{k=-\infty}^{\infty} H(e^{jk\omega_0}) a_k e^{jk\omega_0 n}$$

$$H(e^{j\omega})$$

$$a_k \longrightarrow \underbrace{H(e^{jk\omega_0}) a_k}_{\text{"gain"}}$$

$$H(e^{jk\omega_0}) = |H(e^{jk\omega_0})| e^{j\angle H(e^{jk\omega_0})},$$

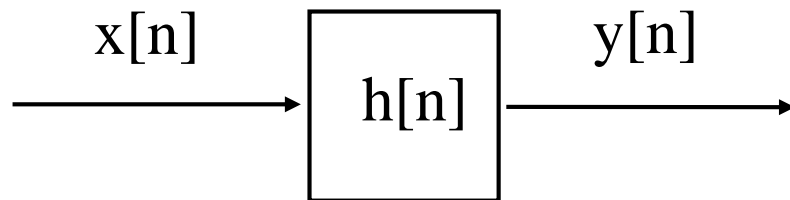
includes both amplitude & phase

The effect of the LTI system is to modify each a_k through multiplication by the value of the frequency response at the corresponding frequency.

Example 3.17

$$h[n] = \alpha^n u[n] \quad , \quad |\alpha| < 1$$

$$x[n] = \cos\left(\frac{2\pi n}{N}\right) = \frac{1}{2}e^{j(\frac{2\pi}{N})n} + \frac{1}{2}e^{-j(\frac{2\pi}{N})n}$$



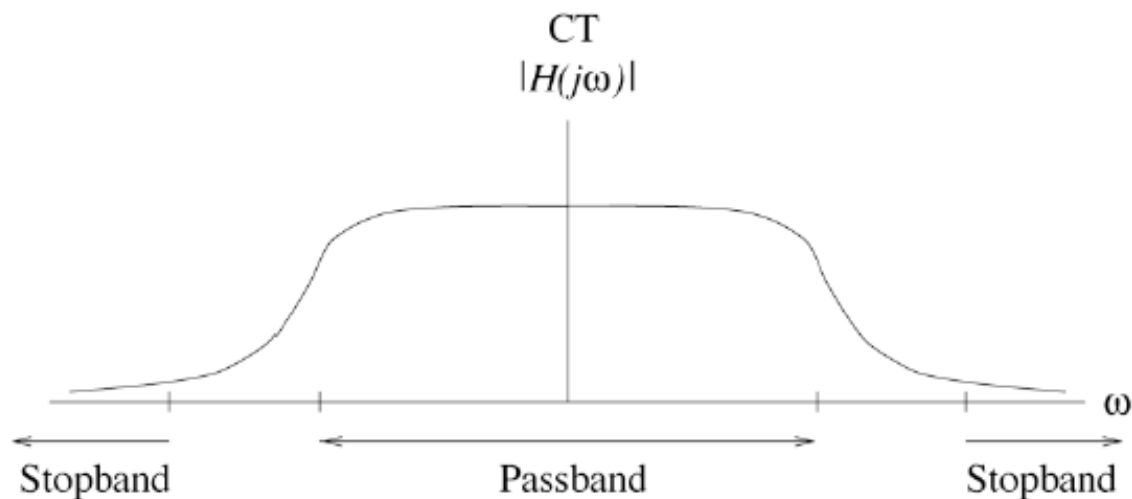
$$H(e^{j\omega}) = \sum_{n=0}^{\infty} \alpha^n e^{-j\omega n} = \frac{1}{1 - \alpha e^{-j\omega}}$$

$$\begin{aligned} y[n] &= \frac{1}{2} H\left(e^{j\frac{2\pi}{N}}\right) e^{j(\frac{2\pi}{N})n} + \frac{1}{2} H\left(e^{-j\frac{2\pi}{N}}\right) e^{-j(\frac{2\pi}{N})n} \\ &= r \cos\left(\frac{2\pi n}{N} + \theta\right) \end{aligned}$$

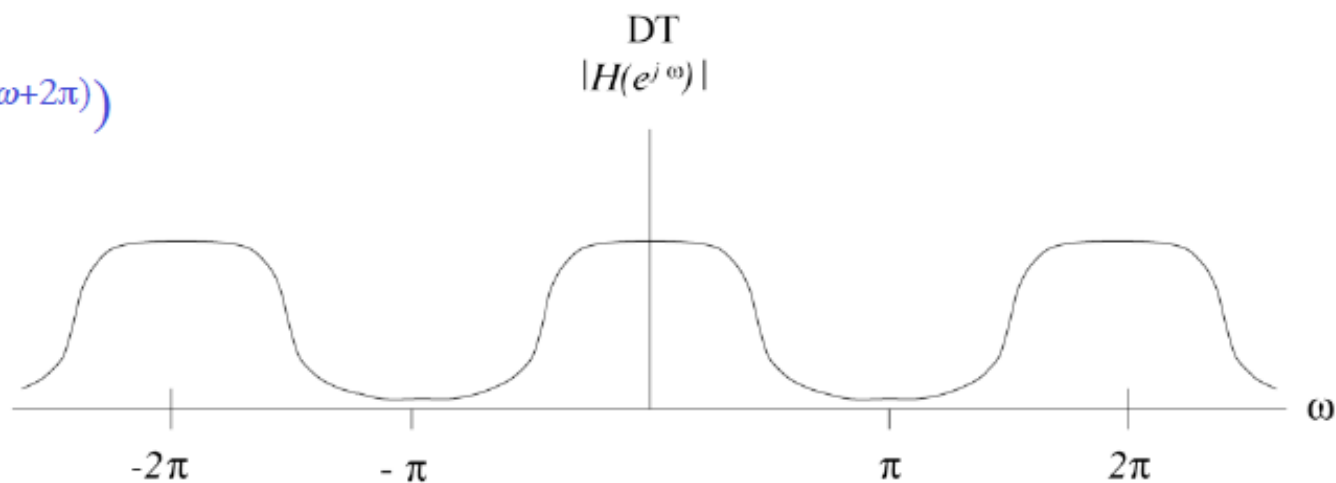
$$\text{where } re^{j\theta} = \frac{1}{1 - \alpha e^{-j\frac{2\pi}{N}}}$$

Lowpass Filter

Lowpass Filters:
Only show
amplitude here.

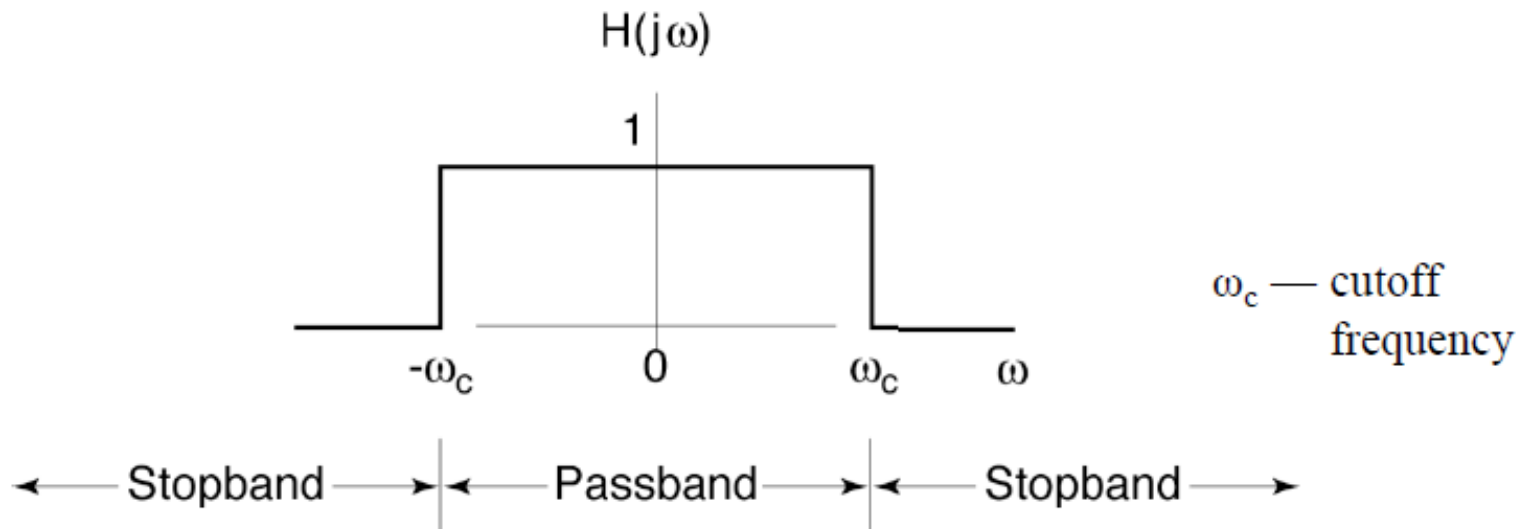


Note for DT:
 $H(e^{j\omega}) = H(e^{j(\omega+2\pi)})$

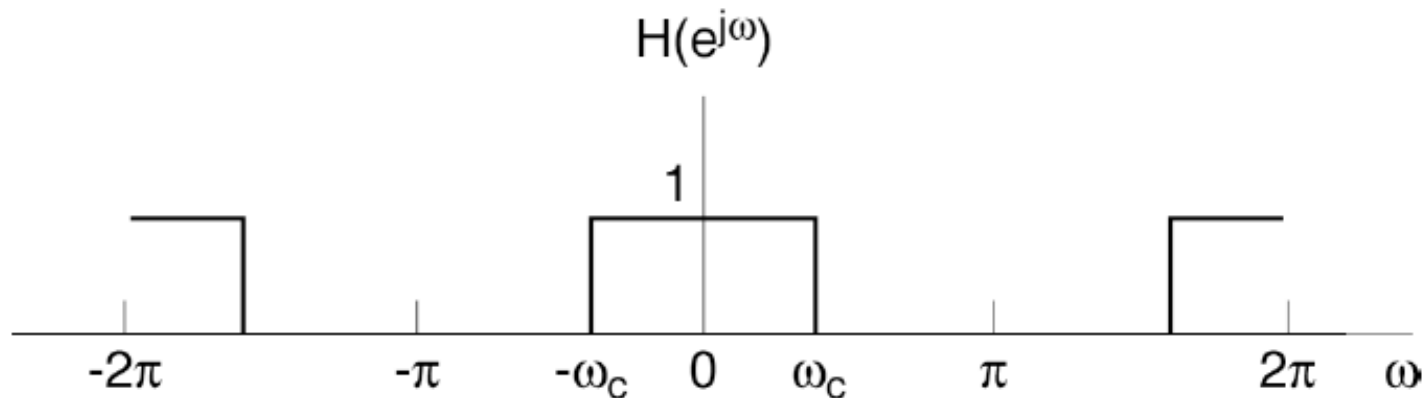


Ideal Lowpass Filter

CT



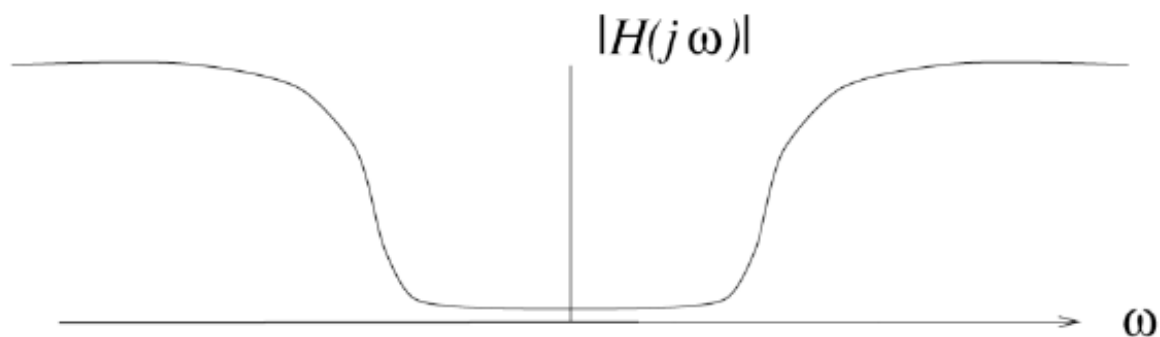
DT



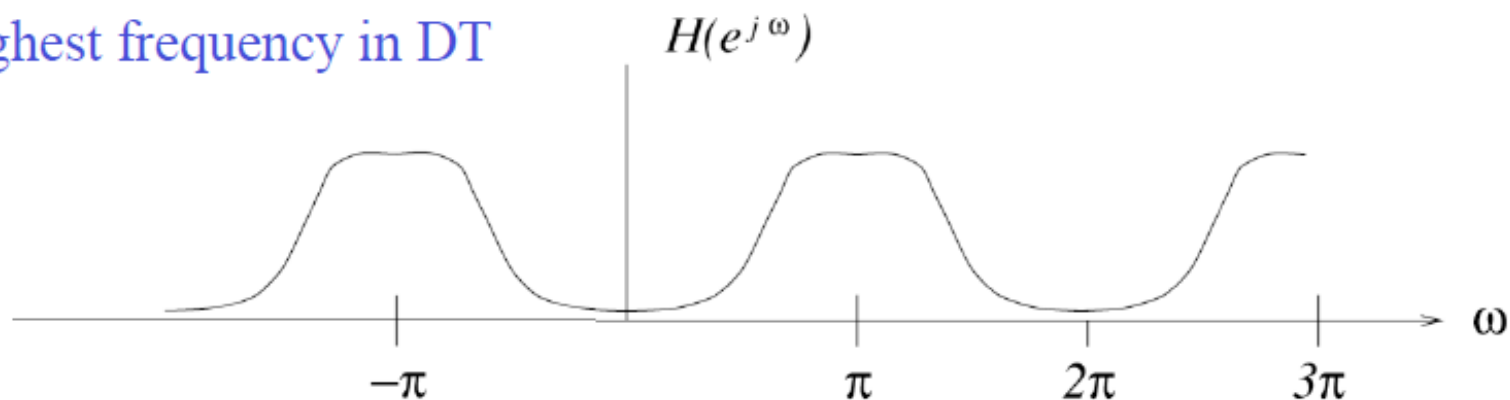
Note: $|H| = 1$ and $\angle H = 0$ for the ideal filters in the passbands,
no need for the phase plot.

Highpass Filters

CT



DT



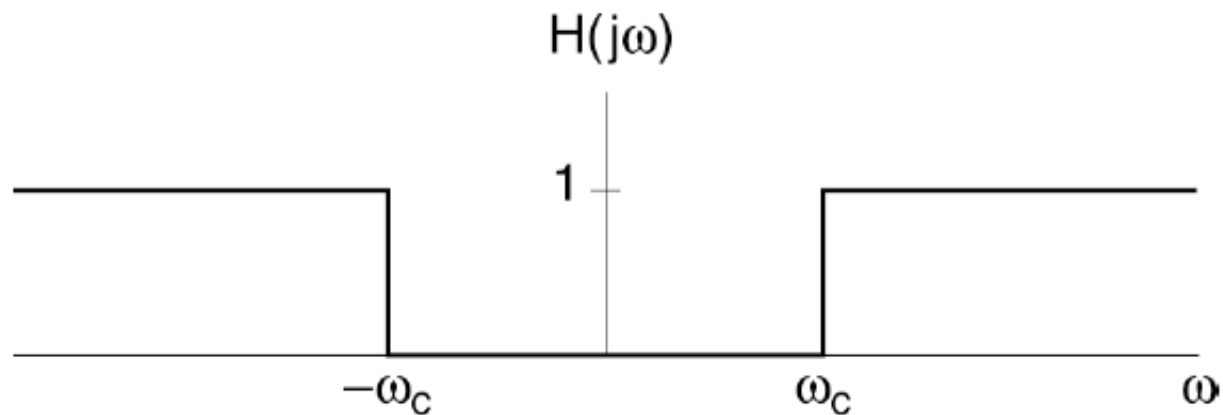
Remember:

$$(-1)^n = e^{j\pi n}$$

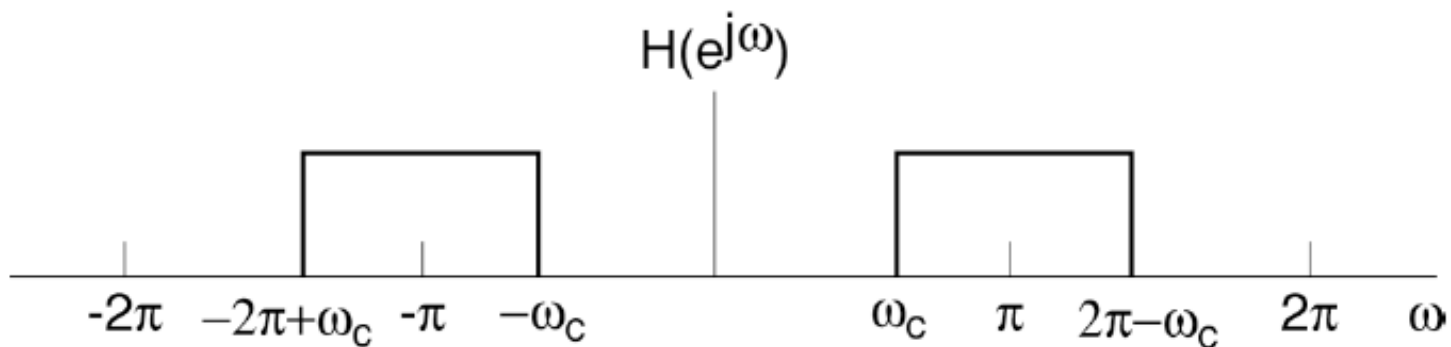
— π = highest frequency in DT

Ideal Highpass Filter

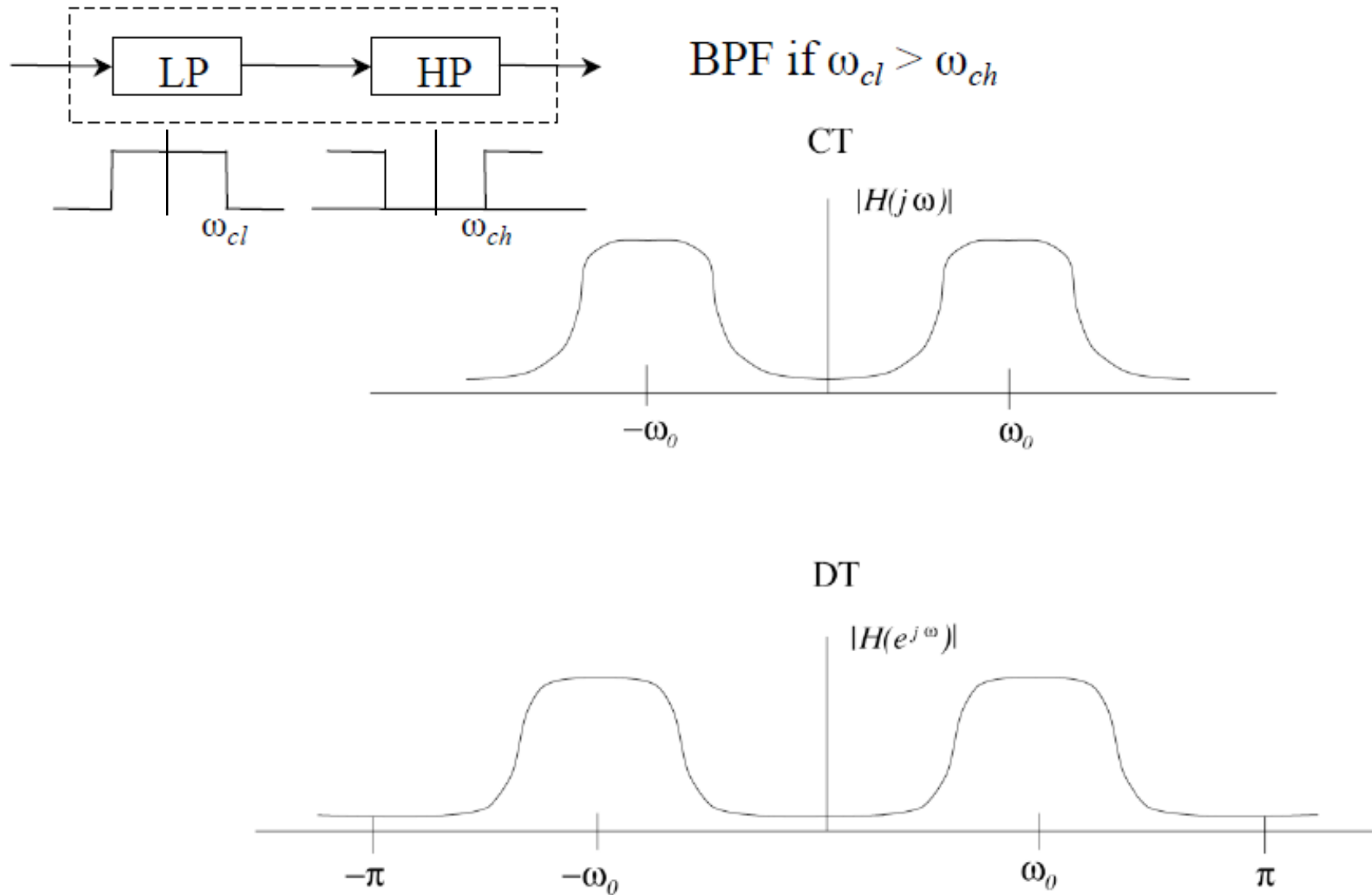
CT



DT

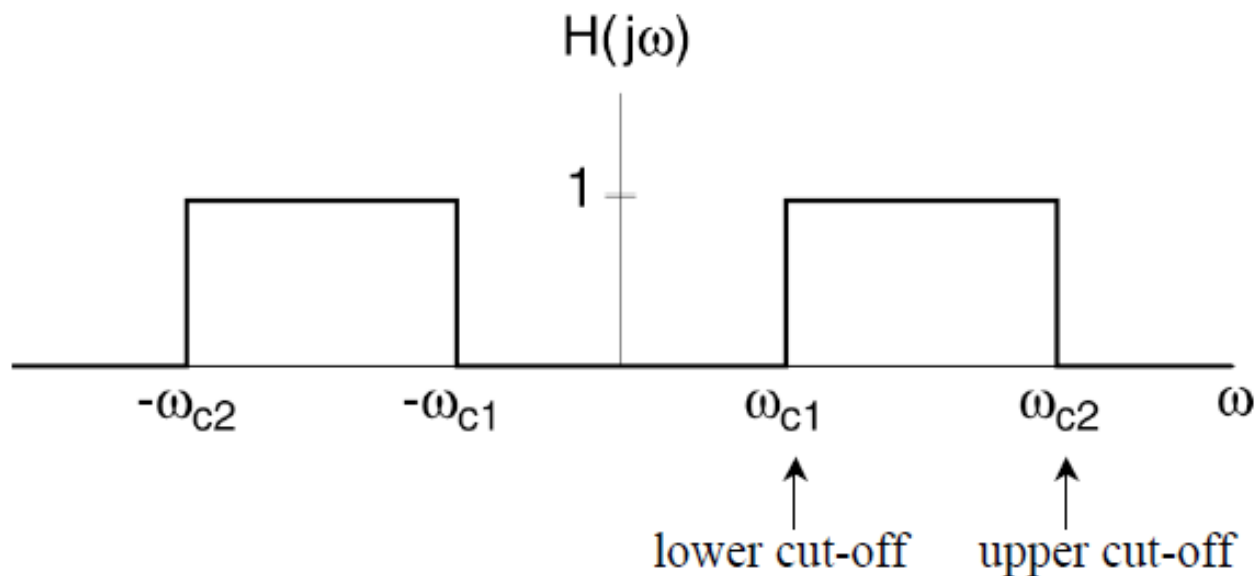


Bandpass Filters

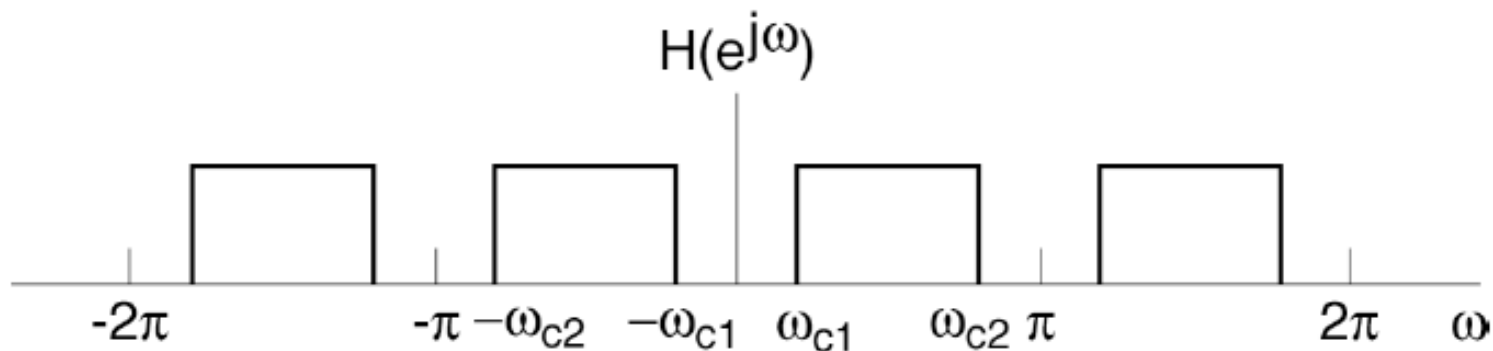


Ideal Bandpass Filter

CT



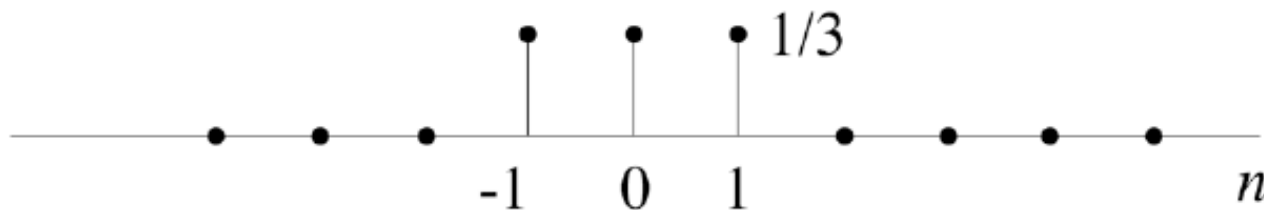
DT



Example: DT Averager/Smoothing

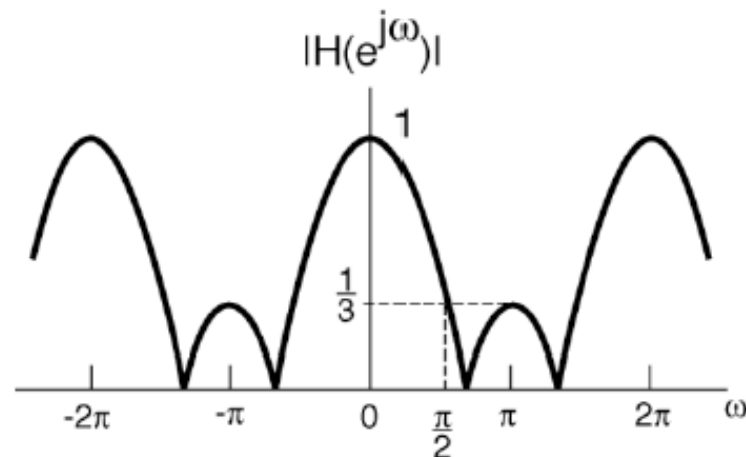
$$y[n] = \frac{1}{3}\{x[n-1] + x[n] + x[n+1]\}$$

$$h[n] = \frac{1}{3}\{\delta[n-1] + \delta[n] + \delta[n+1]\}$$



Frequency response:

$$H(e^{j\omega}) = \sum_{n=-\infty}^{\infty} h[n]e^{-j\omega n} = \frac{1}{3}[e^{-j\omega} + 1 + e^{j\omega}] = \frac{1}{3} + \frac{2}{3}\cos\omega$$



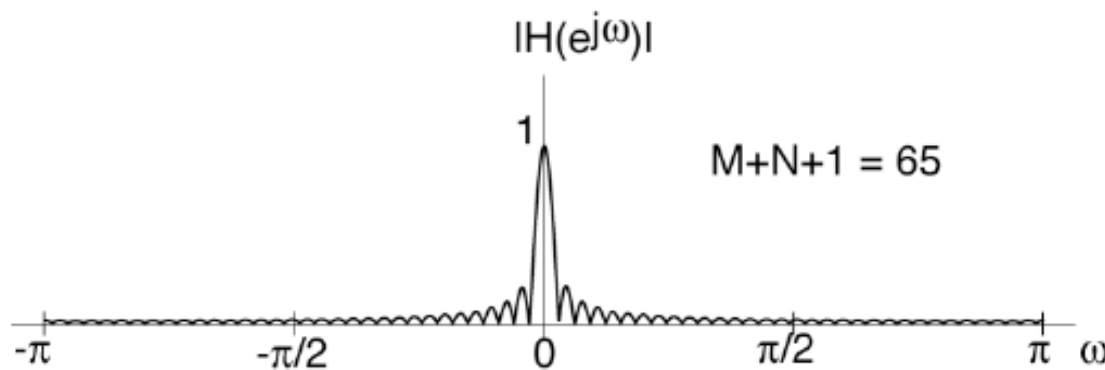
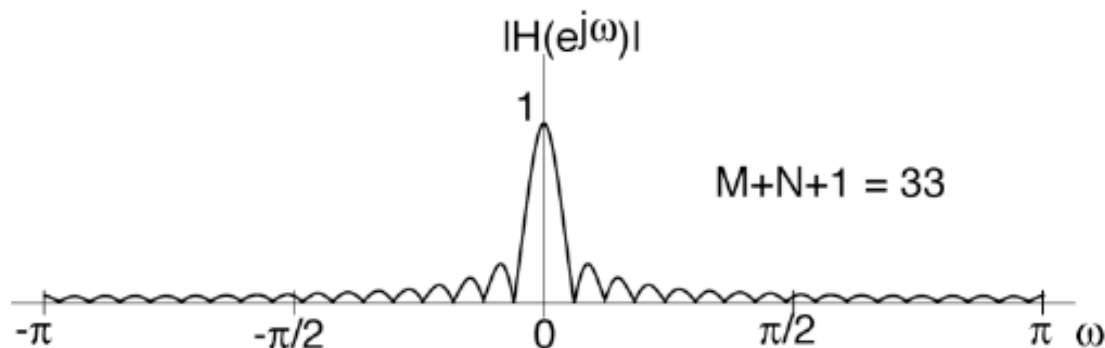
A LPF

Example: Nonrecursive DT (FIR) filters

$$y[n] = \frac{1}{N + M + 1} \sum_{k=-N}^M x[n - k] \longrightarrow h[n] = \frac{1}{N + M + 1} \sum_{k=-N}^M \delta[n - k]$$

Frequency response:

$$H(e^{j\omega}) = \frac{1}{N + M + 1} \sum_{k=-N}^M e^{-jk\omega} = \frac{1}{N + M + 1} e^{j\omega(N-M)/2} \frac{\sin[\omega(M + N + 1) / 2]}{\sin(\omega / 2)}$$



Rolls off at lower ω as $M+N+1$ increases

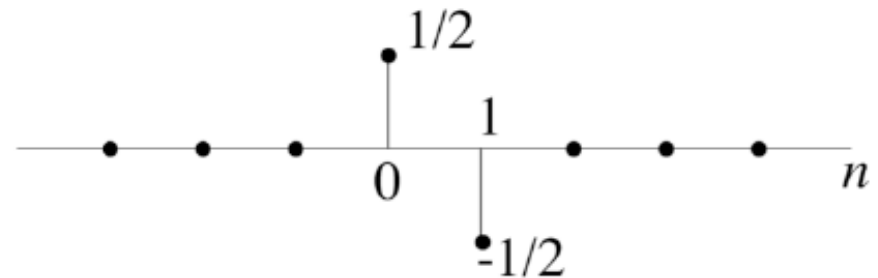
Example

Simple DT “Edge” Detector

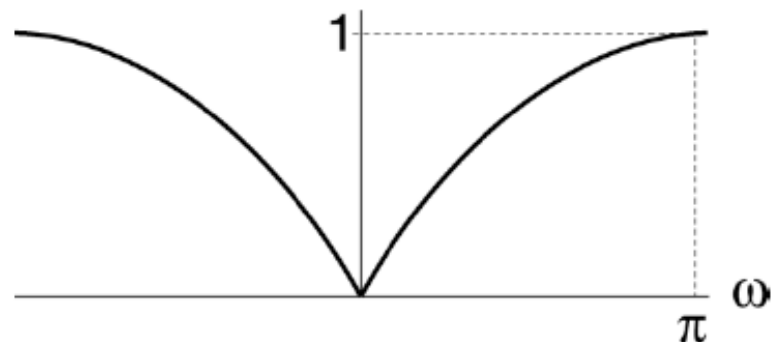
— DT 2-points “differentiator”

$$y[n] = \frac{1}{2}[x[n] - x[n-1]]$$

$$h[n] = \frac{1}{2}[\delta[n] - \delta[n-1]]$$



$|H(e^{j\omega})|$



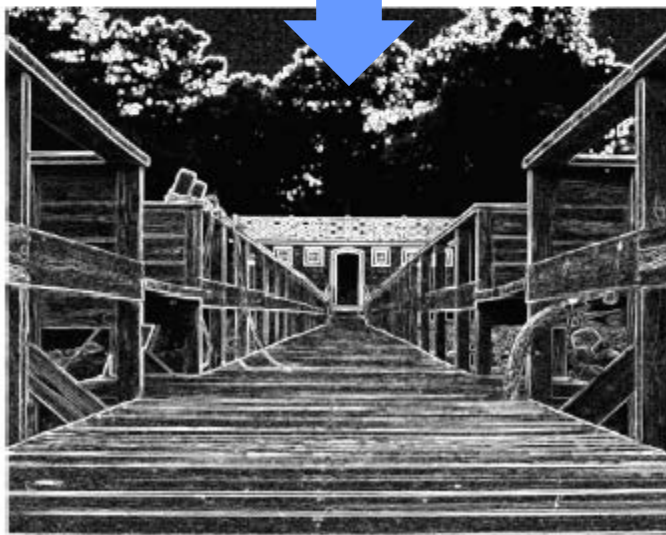
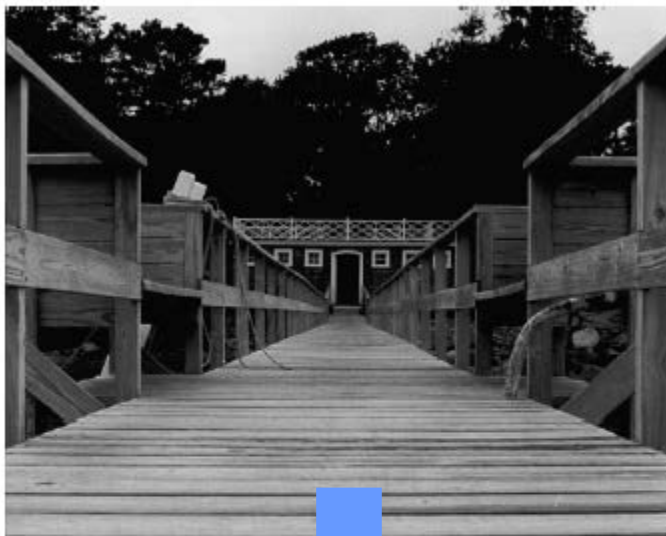
Amplifies high-frequency components

Frequency response:

$$H(e^{j\omega}) = \frac{1}{2}(1 - e^{-j\omega}) = je^{j\omega/2} \sin(\omega / 2)$$

$$|H(e^{j\omega})| = |\sin(\omega / 2)|$$

Edge enhancement using DT differentiator



Summary

- **DT Fourier Series pair**
 - ◆ Understand the difference between CT and DT
- **Frequency response**
 - ◆ How to determine frequency response?
- **Filtering**