

Efficient solutions of implicit discretizations of the immersed boundary method

Jordan Fisher

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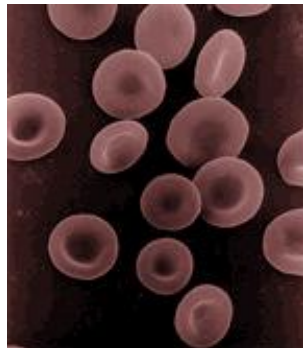
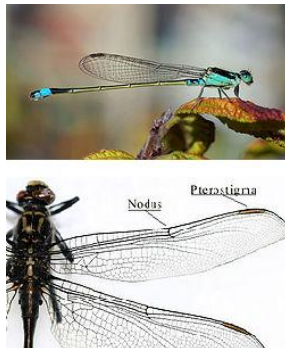
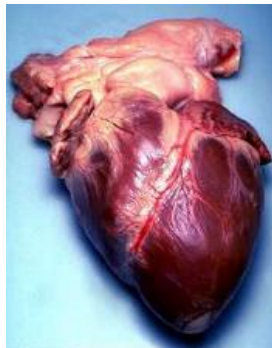
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The standard Navier-Stokes equation for incompressible flow with an arbitrary forcing function \mathbf{f}

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$$\mathbf{f}(\mathbf{x}, t) = \int_B \mathbf{F}(\mathbf{X}(\cdot, \cdot), s, t) \delta(\mathbf{x} - \mathbf{X}(s, t)) ds \quad (4)$$

Standard discretization: Forward Euler/Backward Euler

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We approximate δ with a discrete function

$$d_h(r) = \begin{cases} \frac{1}{4h} (1 + \cos(\frac{\pi r}{2h})) & \text{if } |r| \leq 2h \\ 0 & \text{otherwise,} \end{cases} \quad (5)$$

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and approximate the spreading and smoothing via

$$(\mathcal{S}_n G)(\mathbf{x}) = \sum_{s \in \mathcal{G}_B} G(s) \delta_h(\mathbf{x} - \mathbf{X}^n(s)) h_B, \quad (6)$$

$$(\mathcal{S}_n^* w)(s) = \sum_{\mathbf{x} \in \mathcal{G}_\Omega} w(\mathbf{x}) \delta_h(\mathbf{x} - \mathbf{X}^n(s)) h^2. \quad (7)$$

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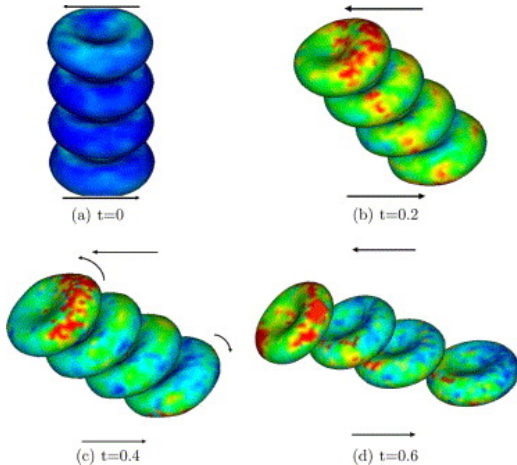
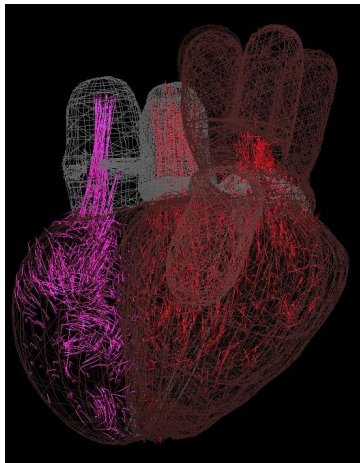
$$\rho \left(\frac{\mathbf{u}^{n+1} - \mathbf{u}^n}{\Delta t} + \mathbf{u}^n \cdot \mathbf{D}_h \mathbf{u}^n \right) = -\mathbf{D}_h p^{n+1} + \mu L_h \mathbf{u}^{n+1} + \mathcal{S}_n \mathcal{A}_{h_B}(\mathbf{X}^n), \quad (9)$$

$$\mathbf{D}_h \cdot \mathbf{u}^{n+1} = 0, \quad (10)$$

$$\frac{\mathbf{X}^{n+1} - \mathbf{X}^n}{\Delta t} = \mathcal{S}_n^* \mathbf{u}^{n+1} \quad (11)$$

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Implicit discretizations!

What should we make implicit?

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$$(\mathbf{u}^{n+1}, \mathbf{u}^{n+1})_{\Omega} + (\mathcal{A}_{h_B} \mathbf{X}^{n+1}, \mathbf{X}^{n+1})_B \leq (\mathbf{u}^n, \mathbf{u}^n)_{\Omega} + (\mathcal{A}_{h_B} \mathbf{X}^n, \mathbf{X}^n)_B \quad (12)$$

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Our work has focused on solving this system efficiently

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which we refer to as the fluid solver.

Now we may write

$$\mathbf{u}^{n+1} = \mathcal{L}_h \left[\frac{\Delta t}{\rho} \mathcal{S}_n \mathcal{A}_{h_B}(\mathbf{X}^{n+1}) + \mathbf{u}^n - \Delta t \mathbf{u}^n \cdot \nabla_h \mathbf{u}^n \right] \quad (15)$$

$$\mathbf{X}^{n+1} = \mathbf{X}^n + \Delta t \mathcal{S}_n^* \mathbf{u}^{n+1}. \quad (16)$$

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In fact, $I - \mathcal{M}_n \mathcal{A}_{h_B}$ is positive-definite if \mathcal{A}_{h_B} is negative-semidefinite.

How can we efficiently solve $(I - \mathcal{M}_n \mathcal{A}_{h_B}) \mathbf{X}^{n+1} = \mathbf{b}^n$?

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A matrix representation of $I - \mathcal{M}_n \mathcal{A}_{h_B}$ would greatly speed iterative methods

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The key is an approximation inspired by the continuous equations, as well as a reliance on lookup tables.

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- Slightly more complicated code

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- Faster evaluations of quantities of the form $\mathcal{M}_n \mathbf{F}$, by a factor of 20 to 100. This greatly streamlines iterative methods.

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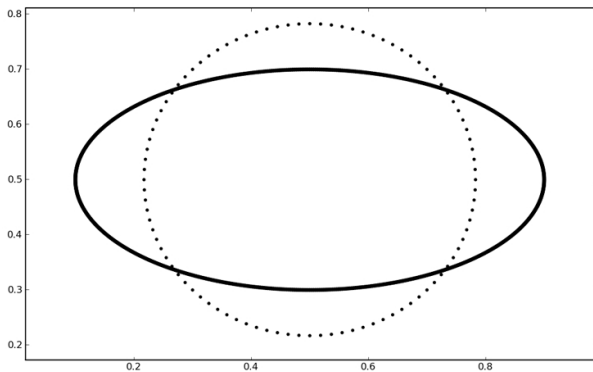
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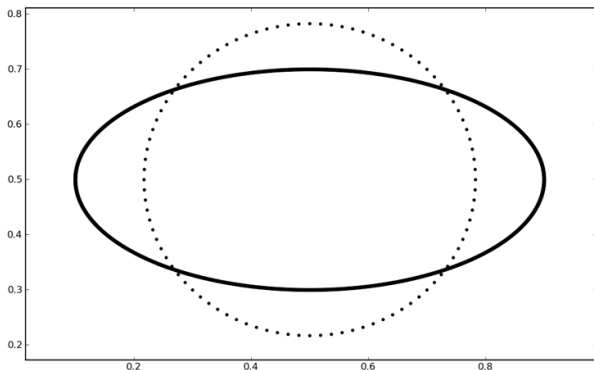
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- **Multigrid!**

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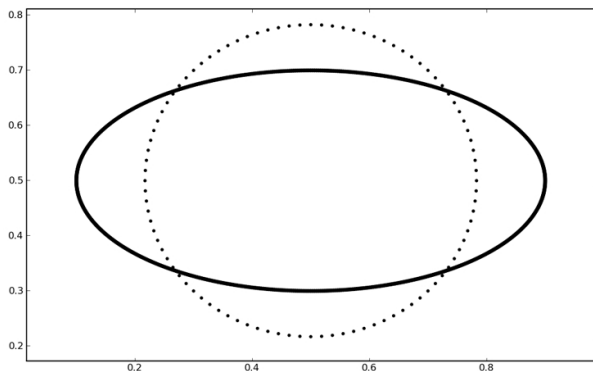


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$$\mathcal{A}_{h_B}(\mathbf{X}) = \sigma \mathbf{X}_{ss} \quad (20)$$

\mathcal{A}_{h_B} is linear, negative-semidefinite, hence we can apply our multigrid.

Comparison of total CPU between an explicit method and our implicit implementation

Table: Elliptical drop relaxation for Navier Stokes. The average CPU time per time-step and the total CPU time up to a simulation time of $T = 0.005$ is given. Δt is the time-step taken and is the maximum allowed while maintaining stability.

N	Explicit			Implicit		
	Δt	Average	Total	Δt	Average	Total
128	$1.95 \cdot 10^{-6}$	0.03	73.35	$1.17 \cdot 10^{-4}$	0.10	4.34
256	$9.76 \cdot 10^{-7}$	0.14	730.21	$7.81 \cdot 10^{-5}$	0.55	35.06
384	$6.51 \cdot 10^{-7}$	0.34	2613.19	$5.21 \cdot 10^{-5}$	1.24	118.98
512	$4.87 \cdot 10^{-7}$	0.63	6491.95	$3.91 \cdot 10^{-5}$	2.34	299.10

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What if \mathcal{A}_{h_B} isn't linear?

- When the Jacobian of \mathcal{A}_{h_B} is negative semidefinite we can solve Newton iterations via multigrid. This works well.
- Unfortunately, multigrid fails if the Jacobian isn't definite
- In certain important cases, we can fix the problem by opting for a fixed-point iteration

Fixed-point iteration

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The most immediate choice for a fixed point iteration is

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Intuitively, the reverse of this iteration should be stable.

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We dub the following the splitting method

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Typically the Jacobian of \mathcal{A}_{h_B} is sparse, so this process is cheap.

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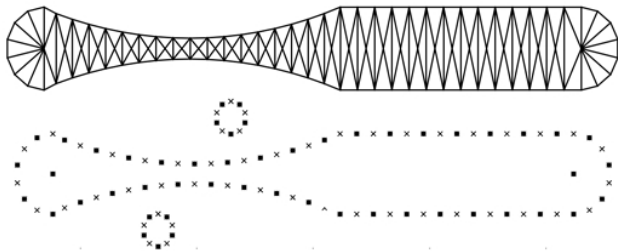
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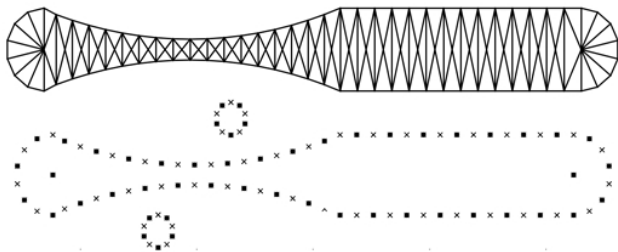
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Provided the eigenvalues of the Jacobian of \mathcal{A}_{h_B} are negative and large enough, (22) converges rapidly.

Application: blood flow past a heart valve



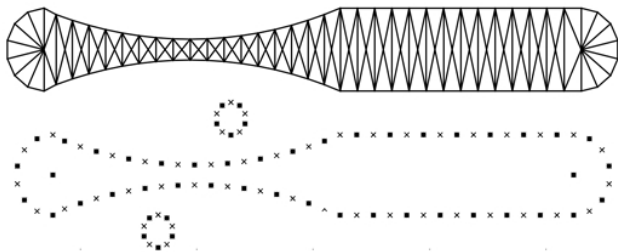
Application: blood flow past a heart valve



The force between two nodes is given by Hooke's law

$$\mathbf{f}_{i,j} = -k_{i,j} \frac{\mathbf{x}_i - \mathbf{x}_j}{|\mathbf{x}_i - \mathbf{x}_j|} (|\mathbf{x}_i - \mathbf{x}_j| - L_{i,j}) \quad (23)$$

Application: blood flow past a heart valve

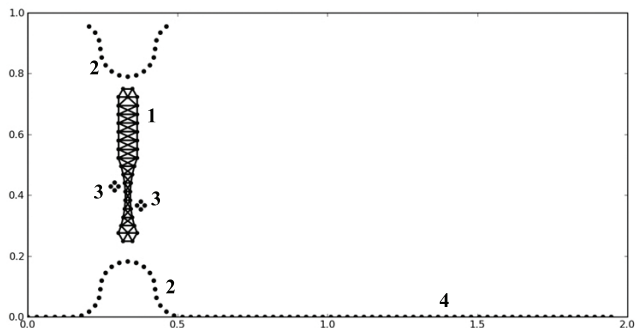


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This is nonlinear when $L_{i,j} \neq 0$!

Application: blood flow past a heart valve



- 1. Valve
- 2. Cushions
- 3. Hinges
- 4. Artery wall

Application: blood flow past a heart valve

A horizontal flow is induced by adding a a time dependent component to the force field:

$$v_{flow}(t) = \begin{cases} 60000t(0.1 - t) & t < 0.1 \\ -60000(t - 0.1)(0.2 - t) & t \geq 0.1 \end{cases} . \quad (24)$$

Application: blood flow past a heart valve

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Application: blood flow past a heart valve

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\mathcal{A}_{h_B} isn't invertible! It is neither 1-1 nor onto.

\mathcal{A}_{h_B} can't produce translational or rotational forces.

\mathcal{A}_{h_B} is insensitive to translation of its input

Factoring out components outside the image of \mathcal{A}_{h_B}

We define vectors

$$\mathbf{v}_I^1 = \begin{cases} (1, 0) & \text{if } I \in \mathcal{G}_V, \\ (0, 0) & \text{otherwise,} \end{cases} \quad (25)$$

$$\mathbf{v}_I^2 = \begin{cases} (0, 1) & \text{if } I \in \mathcal{G}_V, \\ (0, 0) & \text{otherwise,} \end{cases} \quad (26)$$

$$\mathbf{v}_I^3 = \begin{cases} (-Y_I + y_c, X_I - x_c) & \text{if } I \in \mathcal{G}_V, \\ (0, 0) & \text{otherwise,} \end{cases} \quad (27)$$

where \mathcal{G}_V is the collection of all indices I such that \mathbf{X}_I is a fiber point belonging to the valve, and (x_c, y_c) is some fixed point.

Factoring out components outside the image of \mathcal{A}_{h_B}

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$$\left\{ \begin{array}{l} \mathbf{X}^{n+1,k+1,0} = \mathbf{X}^{n+1,k}, \\ J(\mathbf{X}^{n+1,k+1,l})(\mathbf{X}^{n+1,k+1,l+1} - \mathbf{X}^{n+1,k+1,l}) \\ \quad = P(\mathbf{F}^{k+1} - \mathcal{A}_{h_B}(\mathbf{X}^{n+1,k+1,l})), \end{array} \right. \quad (28)$$

where

$$P(\mathbf{F}) = \mathbf{F} - \sum_{j=1}^3 \mathbf{V}^j(\mathbf{X}^{n+1,k})(\mathbf{F}, \mathbf{V}^j(\mathbf{X}^{n+1,k}))_B \quad (29)$$

Reintroducing translation and rotation

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There are simple, fully explicit, and stable procedures for calculating the translation and rotation.

Comparison of total CPU time for the heart valve simulation

Table: Heart valve simulation. The average CPU time per time-step and the total CPU time up to a simulation time of $T = 0.5$ is given. For the explicit method Δt is the time-step taken and is the maximum allowed while maintaining stability. For the implicit method Δt is held constant, and is $O(h)$ or smaller.

N	Δt	Explicit		Δt	Implicit	
		Average	Total		Average	Total
128	$1.15 \cdot 10^{-7}$	0.024	105089*	0.0025	2.001	118
256	$2.89 \cdot 10^{-8}$	0.109	1884358*	0.0025	7.640	402
384	$1.29 \cdot 10^{-8}$	0.249	9658428*	0.0025	18.735	949
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Current research

Current research is focused on extending the methodology to 3D. The main problems are

- The matrix \mathcal{M}_n has $O(N^4)$ entries, so can't be used directly
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- Without a good smoother we can't implement a multigrid

Current lines of attack include

- Looking at alternative smoothers, potentially involving banded components of \mathcal{M}_n
- Looking at good preconditioners for non-multigrid solvers, such as GMRES
- Modifying a Fast Multipole Method to implement Gauss-Seidel in $O(N_b)$ steps